Statistical Thermodynamics for Engineers Professor Saptarshi Basu Indian Institute of Science, Bengaluru Lecture 36 The Four Quantum Numbers and Multielectron System

(Refer Slide Time: 00:02)

	9:41 AM Tue 9 Jan く 昭 Q 〇	Ŀ		Untitled No	*** tebook (32) ~	ۍ م		State B
6	i deop im × Untitled N	ictebook. ×			T O O		VIntitled Noteb	
Statistical Thermodynamics for Engineers Lec25			atom anart n: pr det	is spec un nur maipul	e of a hyp ified by ibos, amantum is the elec	Four #		
	9:41 AM. Tue 9 Jan く 昍 Q, []	ŀÔ		Untitled No	•••• tebook (32) ~	ئ ر		A CONTRACT
S	ratóp im × Untitled N	40 Z		VINTELLE NOTEBOOK × U	T	ebooki. × Untified Notebook	× Untitled Noteb	
Ingineel					consular t designa			
Statistical Thermodynamics for Engineers Lec25			X		oobital a			
hermodyn			mi: s	obital	magneti	۲ <u>۲</u>		
Statistical 1 Lec25			te	e '2' (mmber omponent angula	ofthe		
		-	m ₆ ;	Z-10m	ponents	the		
			D PP-	n ma	ponent o montar v oratic a	nomente. enant		

9:41 AM Tue 9 Jan く 昭 Q 囗 ① E / O D & Q @ D E & O O - - - number. ms has two possible values Electronic degeneracy 200 the hydrovsen atom becomes Lec25 $gel = 2m^2$

	9:41 AM					ut S	P 100% 🛄	as les
	< 8		Ω Φ	Untitled Notebook (32) ~	5	÷Ð	× …	
	drop im	× Untit		× Untitled Notebook	ebook	× Untitled	i Noteb 🗔	A.
			6	′ ◇ ∕ & ○ @ ⊇ ⊡ / / ● ● ● •				and Sugar
Statistical Thermodynamics for Engineers Lec25				my has two possible value	2			
ingin				+				
for E				12				
nics				Electronic clegeneracy 200 the hydrosen atom becomes				
ynan				1 K hudaya Ku				
pou				doo the hydroosen atom				
Ther				be iomes				
ical								
Statisti Lec25				gel = 2m²				
Le St			_					
				4 amontun #s.				
				4 anantan -				
							-	
			-					
								5
						1		

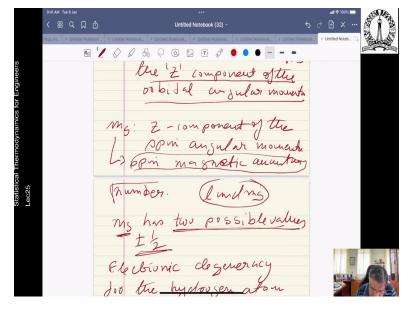
9:41 AM Tue 9 Jan く 韶 Q 囗 ① ቘ 1⁄2 ◊ ◊ ◊ ◊ ◊ @ ⊑ 표 ♂ ● ● ● - - ● MG: Z-component of the ppin angular momente bopin magnetic anautur number. (limiting) ms has two possible values Electrionic degeneracy 200 the hydrosen atom beiomes 10 c 10 mes

< 器 Q 口 位
< 器 Q Д 仓 隆 。

Untitled Notebook (32) ₠ 1⁄ ◊ ◊ ◊ ◊ ◎ ₽ ₸ ₡ ● ● ● - - ● Electionic state of a hydrosen atom is specified by Four amartan numbers n. principal amountain # det en mines the electronic ever sy 1: orbital angular aumter # designated as L' -> oobital angular mo mentur Untitled Notebook (32)

፼ 1/ ◇ / & ○ @ ⊇ 〒 ♂ ● ● ● - - ever 57 1: orbital angular anuten It designated as L'. - > oo bital angular no newtoni me: sobital magnetic anantun muber socifies the '2' component of the or bidal angular mount MG: Z - 10m ponent of the Apri moure. A

Untitled Notebook (32) ፼ 1⁄2 0∕ 0⁄ 0√ 00 12 17 0⁄ 0 0 0 − − − L'. -> oo bital angular momentum M: sobital magnetic anantun muser spaities lie 2' component of the pobidal angular mountu Mg: Z-component of the ppin angular momenter Dopin magnetic anautur number. (limiting) 1-1- 10/11



< 器 Q 口 凸 T ungular momente auntry she rumber honic hydro lep 10mps 155 of 155

So, welcome to lecture number 25 of the Statistical Thermodynamics course. So, as we told in the previous lecture that the electronic state, electronic state of an atom, of a hydrogen atom, hydrogen atom is specified, specified by four quantum numbers. So, of this let us see which quantum number does what, and which is also the principal quantum, number quantum number, and determines the electronic energy, determines electronic energy.

Now, orbital 1 is basically the orbital, angular orbital, angular quantum number designated as l. So, it defines the orbital angular momentum. So, this particular number is basically, this, it is totally for designating the orbital angular momentum. Now, then comes m l which is called the orbital magnetic quantum number, magnetic quantum number. That is, m l. It specifies basically, specifies the z component, the z component of the orbital angular momentum.

Lastly, we have m s, which is the z component of the spin quantum number, of spin angular momentum. So, this is called the spin magnetic number magnetic number, magnetic quantum number. So, already you see that the quantum states are defined by 1 and m l. Now, there are two possible values of the fourth quantum number. So, m s has two possible values, plus and minus half.

So, the electronic degeneracy therefore, electronic degeneracy for the hydrogen atom, the hydrogen atom becomes g electronic is equal to 2 n square. The n is the principal

quantum number. So, basically we already have seen that I and m s was the two quantum numbers that was used to generate the degeneracy. Now, you add the fourth quantum number which is has got two possible values, plus minus half. So, therefore n gets multiplied by this factor 2 because of these two possibilities.

So, this is the four quantum numbers therefore. n, principle quantum number which determines electronic energy, l is the orbital angular quantum number which designates the orbital angular momentum, then m l is basically the orbital magnetic quantum number which specifies, in essence, the z component of the orbital angular momentum. And m s is the z component of the spin angular momentum also called the spin magnetic quantum number.

So, these four quantum numbers, and m s has got two possible values. So, this actually completes the total degeneracy, how to calculate the total degeneracy for the hydrogen atom. So, there are four quantum numbers. And all four are basically required to designate the electronic state. And of course, obviously the energy, energy is dependent on the principle quantum number. So, this completes this, for the hydrogen atom.

(Refer Slide Time: 06:17)

く器へ口合 Untitled Notebook (32) ፼ 1/ ◇ / & ○ @ ⊑ 〒 / ● ● ● - - -Flectronic evens mode for multi election species Steady State Schrödeniger wave Eduction for a generic N ele toon atom with unclease charge 7

< 器 Q 口 凸 T ð yarde t ·D NN e2 9 4 accounts dor the motion of N electrons to account for republice forces among various electrons mking up the atom Untitled Notebook (32) 6100 0 - - -Sh. (t) - T • Steady State Schrödinger wave Eduction for a generic N el for atom with unclease wise the T Zme 24 counts dor montion of

Now we move on to the electronic energy mode for multielectron species. So, electronic energy mode or multielectron species. So, again for doing this, we write a steady state, steady state Schrodinger wave equation, wave equation for a generic N electron atom with nuclear charge Z. So, this is the, this is the equation that we are going to write, z at the origin.

So, minus h bar square divided by 2 m e summation i equal to 1 to N del i squared minus summation I equal to 1 to N, Z e squared divided by 4 pi epslion nought r i plus summation I equal to 1 to N summation j greater than I e square divided by 4 pi epsilon nought r i j. This entire thing is multiplied by psi. Now, the first term that you see over

here, first term accounts for the motion, accounts for the motion of N electrons. This is the Term 1. The second term, this particular term, this term accounts for hydrogen like interaction, interaction between proton and electron.

And last term, this i over here, this particular term you see over here, the last term, the third term is for, to account for repulsive forces, forces, repulsive among the various electrons making. So, these are the. So, this is the steady state Schrodinger wave equation which has got a charge of Z at the center of the origin.

So, this is the total, this is how you write that total equation form. So, first one is to account for the motion, second one is to account for the proton electron reaction, the interactions, and the third is to basically to account for the repulsive forces. So, this entire thing is multiplied by the wave function is equal to this. So, this is the Schrodinger's steady state wave equation, to begin with.

(Refer Slide Time: 11:20)

< 器 Q 口 凸 Hustree-Fock procedure Lacurale solung for many atoms = ZN ; Cyexp(-Non malization wotud

 $\Phi(r_i) = \sum N_i C_j exp(f_j)$ Nop malization custut uped to obtimine final poly qi(ai) -> trial wave function Dependic potros to Schoodinger countron

9:41 AM Tue 9 Jan Untitled Notebook (32) ₠ 1⁄ ◊ ◊ ◊ ◊ ◎ ₽ ₸ ♂ ● ● ● - - -Qi(ri) → trial wave function Dependence polys to Schrodinger constan for independent dections moving in an effective potential (internetice) Coeffed by nuclei and Elle chrons.

9:41 AM Tue 9 Jan く 器 Q 口 ① ⇒ ∂ ∄ × … ፼ 1⁄ ◊ ◊ ◊ ◊ 0 0 1 1 ở ● ● ● - - works well dor 2540 does not wink don 2740 (heavier atoms) S. wave Fountion has operational limitations. and cannot provide accurate electionic evenies dos melti-clection atoms and moleules 158 of 158

 < B Q Q Q Initial Notebook (32) S C B X Initial Notebook (32) Initial Notebook (32)	9:41 AM Tue 9 Jan			atts	≈ 100% 🔤
Depend on noleular and atomic spectrocopy. Each electronic eversy level is denoted by	< 器 Q 口 白	Untitled Notebook (32) ~		∂ 🗜	× …
metti-clection atoms and molecules Depend on nolecular and atomic spectroscopy. Each electronic evensy level is denoted by	rop im × Untitled Notebook	× Untitled Notebook × Untitled Notebook × Untitled Notebook × Untitled Notebook × Untitled	ed Notebook	× Untitle	d Noteb 🗔
metti-cle chion atoms and mole unks Depend on molecular and atomic spectroscopy. Each electronic evensy level is denoted by			ed restebook	× Onumer	a Noteo (1.)
metti-cle chion atoms and mole unks Depend on molecular and atomic spectroscopy. Each electronic evensy level is denoted by	<u>e</u>		-		
Depend on nolealar and ato mic spectrocopy. Each electromic eventsy level is denoted by		a the ale the storus an	d		
Depend on noleular and ato mic spectrocopy. Each electromic evensy level is denoted by					
Each de chumic evensy level is denoted by		moleules			
Each de chumic evensy level is denoted by					
Each de chumic evensy level is denoted by					
Each de chumic evensy level is densted by	_		_		
Each de chumic evensy level is denoted by					
Each de chumic evensy level is densted by		Depend on molecular and	1		
Each de chumic evensy level is denoted by					
Each de chumic evensy level is densted by		atomic spectroscopy.			
Each de bonic evensy level is densted by form symbol					
toch deckomic enersy level is densted by form symbol		0 1 1 1			
level is densted byg form symbol		tack dectoric every			
form symbol					
form symbol		here is denoted by a	-		
A CUT IS MADE		toma and had			
		A com is mod			
					SALL .
					2

Untitled Notebook (32) ፼ 1⁄ ◊ ◊ ◊ ◊ 0 ₽ ₽ ₺ / ● ● ● - - -Quantum the for electrons of a multi-electron atom n=1,2,3,,... n l l=0, 1,2,...,m-1 my m= 0, ±1, ±2,...,±l ms ms= ± ± Depecify a amantan Brato

9:41 AM Tue 9 Jan く 昭 Q 口 ① Untitled Notebook (32) ₠ 1⁄ ◊ ◊ ◊ ◊ ◎ ₽ ₸ ♂ ● ● ● - - -We introduce the so culled configuration of an atom #lelectrons - shell (n=1,2 Sub-shell (l=0,1,2,...) of a multi ele clow poster Shell is poecified by the movenical value of n 1,2,3

く器へ口白 ₠ 1⁄ ◊ ◊ ₭ ♡ ७ छ छ ♂ ♦ ● ● ● Shell is precidied by the muggical value of n 1,2,3 etc Sub-phell is ppecificlly s, P, d, J corresponding to l=0, 1, 2, 3, K -> # gelections but shell (S.P. Untitled Notebook (32) ፼ 1/ ◇ / ୠ ○ @ ⊑ 〒 ♂ ● ● ● - - -Quantum the for electrons of a multi-electron atom n=1,2,3,,.... l=0, 1, 2, ..., n-1 m1=0,±1,±2,...,±l specify a quantan

So, it is kind of not always easy to find out solutions to this. So, there is a process called Hartree-Fock procedure. So, what does this procedure actually do? So, Hartree-Fock procedure, we use this procedure to basically find out accurate solutions, solutions for many atoms, many. So, it is usually done like this.

So, the psi, the wave function, the total wave function is multiplied i equal to 1 to N, and this is represented by some element, phi. So, overall wave function is a product of independent wave functions. And especially this phi i, is given a summation of N i j sum over j C i j exponential j r. So, this is a normalization constant, optimization constant. These are, this and this, these two, these two are variational parameters.

These two are variational parameters used to obtain the final solution, solution. So, these are also called, phi i r i are also called the trial wave functions. So, what the solution, so this trial wave functions represent basically specific solutions to, specific solutions to the Schrodinger's wave equation, Schrodinger equation. Or independent electrons moving in an effective potential.

Trial Solutions, these are specific solutions to the Schrodinger's equation for independent electrons moving in an effective potential, mind the word effective potential, this is which is interactive in nature basically, interactive, created by nucleus, nucleus, by the nuclei, by nuclei, actually, by nuclei and electrons. Nuclei and electrons

So, final solution is obtained by iteratively guessing the values of this above new parameters and so on until we get some convergence. So, you, you, what you do is that basically you guess values and then you see that whether it converges. So, many, so, it is like an iterative kind of a solution, iterative solution. So, it works very well, works well for Z less than equal to around 40.

Does not, does not work or Z greater than 40, that means heavier atoms, it does not work at all. So, it works but similar calculations for larger atoms and for diatomic molecules, it does not really work that way. So, therefore we can infer that Schrodinger wave equation, wave equation, equation has operational limitations, operational limitations, operational limitations.

And cannot provide accurate electronic energies, and cannot provide accurate electronic energies, energies for atoms and molecules or multielectron basically, multielectron molecules. So, these are operational limitations. So, you principally depend on, depend on molecular and atomic spectroscopy, atomic and molecular spectroscopy.

So, each electronic level is also characterized by an associated term symbol from which we can, we can extract the electronic degeneracy which is required for the statistical thermodynamics, statistical mechanical calculations. So, so each electronic, electronic energy level, each electronic energy level, energy level is denoted by a term symbol, so which is required for the stat-thermo calculations that we are going to perform later on. So, we have to develop an appreciation for these term symbols, what they actually mean and how they can be used et cetera et cetera. So, as we saw that how Hartree-Fock orbital procedure can do a lot of stuff, it can give very accurate solutions not only for smaller more compact atoms even for the multielectron kind of a system. Now, let us see that we know now that there are four quantum numbers for the electrons of a multielectron atom.

So, the, so the quantum numbers, let us put it like that, quantum numbers for electrons, of a multielectron, multielectron. So, the first one is n, n equal to, n equal to 1, 2, 3, da, da, da. Then of course, it is 1, 1 equal to 0, 1, 2, dot, dot, dot, n minus 1. Then there is m 1, which we know m 1 is equal to 0, plus minus 1 plus minus 2 dot, dot, dot. plus minus 1. Then that is m s which is equal to, m s equal to plus minus half.

So, these are the four sets of quantum numbers. So, all four sets of quantum numbers must specify, all these should specify a quantum state. Furthermore, say only one electron can occupy any given quantum state. This comes from Exclusion Principle. Therefore each electron of a multielectron atom can be identified by a unique combination of these four quantum numbers. So, there is only a unique combination of these four quantum numbers that can identify an electron in a multielectron kind of a setting.

So, we introduce the so called configuration of an atom, we introduce the so called configuration of an atom which is a convenient fashion, the number of electrons occupying each shell and sub-shell. That means, so this actually says the electrons, the number of electrons occupying each shell, shell is basically n, 1, 2, et cetera and then the corresponding sub-shell which is basically equal to 1 equal to 0, 1, 2, dot, dot, dot, so, for multielectron system, multielectron system.

The shell is specified by the numerical values. So, the shell is specified by the numerical value, the numerical value of n, which is basically 1, 2, 3, et cetera. The sub-shell which is orbital quantum number, is specified by s, p, d, f, s, p corresponding to I equal to 0, 1, 2, 3. So, given this nomenclature, so we have identified that the electron should be specified by n and then the corresponding option in which they are residing.

So, this is given in the format n lk. So, n is the principal quantum number which tells you the shell, 1 is the sub-shell or s, p, d, f, and k is nothing but the number of electrons per sub-shell, number of electrons per sub-shell. So, that is the configuration that we have.

1 AM Tue 9 Jan < 器 Q 口 🖞 6 1 000 O @ P T d Folectio Alonnic sodium Z=11 15225020635 Latt of electrons por subshell is limited by me and ms 2 (2lti ፼ 1/ ◇ / ↓ ○ @ ₽ ₸ / ● ● ● - - -2021+1 2 = 18 Is² 2s² 2p⁶ 3s² 3p⁶ L ptomic configuration Specto Diopic Term Oynubolo Jor Multi-plactic atom

(Refer Slide Time: 23:50)

So, n represents the sub-shell number. As an example, if you take, take atomic sodium, let us say atomic, atomic sodium for which the Z is equal to 11, the configuration for the ground electronic state will be 1s2 2s2 2p6 3s. So, the number of electrons, but however you see the number of electrons, electrons per sub-shell, sub-shell is limited by, limited by ml and ms, which is 2 into 2 l plus 1.

So, for example when you are at s, s is equal to 0, so there are only two electrons that are possible. When I is equal to 1, so that means the number becomes 3. 3 into 2, 6 electrons are possible in that particular sub-shell so on and so forth. So, you understand when s means I equal to 0, so that means if you put I equal to 0 here, so the number of ways this can be counted is 2. So, that is why you have the number 2.

And you cannot have, because if you see, I varies up to n minus 1. So, at n is equal to 1, 1 can only be 0, I can only be 0 when n is equal to 1. Correct. So, that is why we have 1 s 2, because there is no p when, when n is equal to 2, 1 can be 0 and 1, respectively. That means s and p. So, when it is s, when I equal to 0, this is 2, when I is equal to 1, this is actually 3 into 2, 6.

Similarly, you can have 3 s but then we exhaust. So, it is like 2, 2, 6, 10 and 11. So, this is, this is 1, actually. It is what it is. So, in this configuration, in this particular way, you can basically do any kind of, any kind of elements. For example, let us take the example of the Z is equal to say 18. This is Ar. So, what will be the configuration over here? So, first you write 1s2, then 2s2 2p6, out of this, these are common, this is basically common for all of these things.

And then of course you write 2p6 and then what? 2p6, and then what comes after that? Then you have 3s2, n is the same. Now you have like s p d now. So, then it becomes 6 once again. So, this completes the configuration 18. So, you understand, you get a good idea that how this variations is. It basically because of this right and because of this, and because of this right. And of course, this is still the principle quantum number. So, even for a multi-electron system, we can represent how the atomic configuration is.

So, this is how the atomic configuration is, atomic configuration can be written in this succinct kind of a way for this. And then you can do it for n number of these things. So, next class, we are going to do what we call spectroscopic term symbols, term symbols, symbols for multielectron atoms, multielectron. Because these symbols are very important. Some of these symbols you already know from your undergraduate. So, no harm in actually seeing it one more time. Thank you.