

**Convective Heat Transfer**  
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**Lecture – 55**  
**Turbulence – Tutorial**

So, in this particular class, we are going to see about five problems in turbulence which will give you an idea that how to solve different problems these are simple problems which does not require you to memorize a whole lot of things ok, but ok. So, let us look at the first problem.

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Water flows with the velocity  $U_\infty = 0.2$  m/s parallel to a plane wall. The following calculations refer to the position  $x = 6$  m measured downstream from the leading edge. The water properties can be evaluated at  $20^\circ\text{C}$ .

(a) A probe is to be inserted in the viscous sublayer to the position represented by  $y^+ = 2.7$ . Calculate the actual spacing  $y$  (mm) between the probe and the wall.

(b) Calculate the boundary layer thickness  $\delta$ , and compare this value with the estimate based on the assumption that the length  $x$  is covered by turbulent boundary layer flow.

(c) Calculate the heat transfer coefficient averaged over the length  $x$ .

Handwritten notes in red:  $U_b = 0.2 \rightarrow$ ,  $U_\infty = 0.2 \rightarrow$ ,  $y^+ = 2.7$ ,  $y = 2.7$ ,  $20^\circ\text{C}$ ,  $x = 6 \text{ m}$ ,  $\delta$ ,  $y^+$ ,  $y$ ,  $v_{SL}$ .

The first problem what does it say is that the water flows if you look at the ppt. So, the water flows with a velocity of  $U_\infty$  equal to 0.2 meter per second parallel to the plane wall, ok. So, the following calculations refer to the position  $x$  equal to 6 meter measured downstream from the leading edge the water properties can be evaluated at 20 degree Celsius.

So, is basically water which is flowing over a flat plate parallel to a plain wall this is water  $U_\infty$  is about 0.2 meter per second, calculations are carried at  $x$  equal to 6 meter, from the leading edge. So, this is the leading edge. So, water properties can be evaluated at 20 degree Celsius. So, what we are trying to do is that we are trying to insert a probe in the viscous sub layer to the position represented by  $y^+$  equal to 2.7, if you

recall the definition  $y$  plus equal to 2.7. Calculate the actual spacing  $y$  between the probe and the wall. So, you have placed the probe at a distance which is given as  $y$  plus equal to 2.7. We want to know the physical distance all right, we want to know the physical distance and this probe is placed in the what we call the viscous sub layer the VSL, ok. So, that is the first problem that what is the actual spacing.

Next part is calculate the boundary layer thickness and compare the value based on the assumption that  $x$  is covered by turbulent boundary layer flow and calculate the heat transfer coefficient which is averaged over the length  $x$ . So, this is the problem that we are dealing with here. So, let us see how this problem can be solved.

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$$\text{Q1: properties of water at } 20^\circ\text{C}$$

$$\rho = 1 \text{ gm/cc} \quad c_p = 4.18 \text{ kJ/kg-K} \quad \gamma = 0.01 \frac{\text{m}^2}{\text{s}}$$

$$k = 0.59 \frac{\text{W}}{\text{mK}} \quad Pr = 7.07$$

$$a) \quad y = y + \gamma \left( \frac{\tau_w}{\rho} \right)^{-\frac{1}{2}}$$

$$\frac{\tau_w}{\rho} = 0.0276 u_s^2 Re_x^{-\frac{1}{2}} \quad \rightarrow Re_x = \frac{u_s x}{\gamma} = \frac{20 \frac{\text{cm}}{\text{s}} \cdot 600 \text{ cm}}{0.01 \frac{\text{cm}^2}{\text{s}}} = 1.2 \times 10^6 \text{ (turbulent)}$$

$$\therefore \frac{\tau_w}{\rho} = 0.0018 u_s^2$$

$$\therefore \left( \frac{\tau_w}{\rho} \right)^{\frac{1}{2}} = 0.0424 \times 20 \frac{\text{cm}}{\text{s}} = 0.849 \frac{\text{cm}}{\text{s}}$$

$$\therefore y = y + \gamma \left( \frac{\tau_w}{\rho} \right)^{-\frac{1}{2}} = 2.7 \times 0.01 \frac{\text{cm}^2}{\text{s}} \frac{\text{s}}{0.849 \text{ cm}} = 0.3 \text{ mm}$$

So, question one ok. So, we can first take the properties of water of water at 20 degree Celsius. So, rho is 1 gram per cc cp is equal to 4.1 one eight kilo joules per kg k you have to take care of the units whatever unit you use you have to just use them properly this is the kinematic viscosity  $\gamma$  is 0.01 watt per meter Kelvin and Prandtl number is about 7.07 that is a Prandtl number for water. So, in order to calculate  $y$ , so, the first part of the problem is to calculate  $y$ . So, we know  $y$  is given plus  $y$  plus into gamma tau wall divided by rho raised to the power of minus half all right, this is the definition that we know from the during the course of our last lectures.

So, you have to know what about this we already know this is 2.7, and we already know the value of this all right it is given there. So, what will be the value of tau wall by rho ?

So, tau wall by rho therefore, is given from our correlations is 2.296 into U infinity square Reynolds number x to the power of minus one fifth in this particular term Reynolds number x is given as U infinity into x by gamma.

So, therefore, if you start substituting U infinity is about 20 centimeter. So, let us put it at centimeter per second, 600 centimeter is the wall length. So, we are converting it into one particular unit and gamma is basically the kinematic viscosity which is 0.1 centimeters square and second on the top. So, this gives you about 1.2 into 10 to the power of 6 which is basically a turbulent flow all right it is basically a turbulent flow. So, therefore, your tau wall the rho is given as 0.0018 ok, if you convert all the terms U infinity squared. So, therefore, tau wall rho raised to the power of half is about 0.4024 into 20 centimeter per second which is 0.849 centimeter per second.

So, therefore, your y is equal to y plus gamma tau wall by rho raised to the power of half. So, this is actually given as 2.7 in to 0.01 centimeter square per second, second divided by 0.849 centimeter. So, this gives you 0.3 millimeter. So, the physical distance y where you have inserted the probe is about 0.3 mm which is about 300 microns. 100 micron is the diameter of your hair, so, it is about 3 times the thickness of your hair where this probe is inserted just to give you a feel.

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Handwritten notes on a whiteboard showing calculations for boundary layer thickness and heat transfer coefficient.

b) BL thickness  $\delta$

$$\delta = 0.37 Re_x^{-1/5} = 0.37 \times 6m (1.2 \times 10^6)^{-1/5} = 13.5 \text{ cm.}$$

$\delta = 13.5 \text{ cm.}$

$\delta_{Lam} = 4.92 \times Re_x^{-1/2} = 4.92 \times 6m (1.2 \times 10^6)^{-1/2} = 2.7 \text{ cm.}$

Laminar BL is much thinner than the real TBL.

c)  $Nu_x = 0.0296 Pr^{1/3} Re_x^{4/5}$

$$\overline{Nu}_x = 0.037 Pr^{1/3} Re_x^{4/5} = 0.037 (7.07)^{1/3} (1.2 \times 10^6)^{4/5} = 5184$$

$\overline{h} = \overline{Nu}_x \frac{k}{x} = 5184 \times 0.59 \frac{1}{6m} = 509 \frac{W}{m^2K}$

So, similarly now the boundary layer thickness which is delta ok, can be evaluated as this is 0.37 if you recall Reynolds number x to the power of minus one fifth. So, this is 0.37

into 6 meter into 1.2 into 10 to the power of 6 raised to the power of minus one fifth, this gives you about 13.5 centimeter. So, that is the thickness. So, you can put this in perspective the total thickness of this boundary layer is about 13.5 centimeter quite a bit.

Now, in the laminar region, so, if you calculate the laminar region for this same thing if you recall it was  $4.92 \times 10^{-1/2}$  into Reynolds number  $x$  to the power of minus half. So, it is a much much stronger dependence so,  $4.92 \times 6 \times 1.2 \times 10^{-1/2}$  to the power of 6 raised to the power of half. So, it is basically becomes 2.7 centimeter because Reynolds number is in the denominator all right. So, it is as much big it is the same number, but it has been raised to the power of half that is raised to the power of one fifth. So, basically the boundary layer thickness blows up all right because it is a smaller number in the denominator effectively.

So, you can compare one is 13.5, one is 2.7, ok. So, that is the extent that the laminar boundary layer is much thinner boundary layer is much thinner then the real then the real turbulent boundary layer. So, if you use the laminar boundary layer correlations you will be way off that is the whole point that we are trying to make.

So, the Nusselt number on the other hand is given as Nusselt number is given as  $0.0296 \text{ Prandtl number}^{1/3} \text{ Reynolds number}^{4/5}$ .

So, Nusselt number  $\bar{x}$  is  $0.037 \text{ Prandtl number}^{1/3} \text{ Reynolds number}^{4/5}$   $0.037 \times 7.07^{1/3} \times 1.2 \times 10^{4/5}$  to the power of 6 to the power of 4 by 5 this gives you about 5184 Prandtl number. What is your heat transfer coefficient  $\bar{h}$  is equal to Nusselt number  $\bar{x}$  divided by  $k \times x$ , all right. So, this gives you about 5184, into 0.59 remember we calculated  $k$  earlier, divided by  $1/6$  meter this actually if you take care of the units it would be 509 watt per meter square Kelvin, got it. So, that is 509 that is what you get from your from your relation.

So, this completes the first problem which shows that how to for example, use existing correlations and find out these are not correlations per say, but these were problems that we relations that we established earlier and now we are able to use it for solving a real problem ok.

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Consider the heat transfer in boundary layer flow from an isothermal wall  $T_0$  to a constant temperature stream  $(U_\infty, T_\infty)$ . The leading laminar section of the boundary layer has a length comparable with the length of the trailing turbulent section; consequently, the heat flux averaged over the entire wall length  $L$  is influenced by both sections. Derive a formula for the  $L$ -averaged Nusselt number, assuming that the laminar-turbulent transition is located at a point  $x$  (between  $x = 0$  and  $x = L$ ) where  $xU_\infty/\nu = 3.5 \times 10^5$ .

So, if you look at the next problem now so, you consider the heat transfer in the boundary layer from an isothermal wall to a constant temperature  $T$  stream all right the leading laminar section of the boundary layer. So, it has got a laminar section and then it has got a turbulent section has a length comparable with the length of the trailing turbulent section.

So, both are kind of comparable you cannot neglect one for the other consequently the heat flux averaged over the entire wall length is influenced by both sections. So, our idea is to derive a formula for the  $L$  averaged Nusselt number assuming that laminar turbulent transition is located at a point  $x$  which is between  $x$  equal to 0 and  $x$  equal to  $L$  and where  $x$  into  $U$  infinity divided by gamma is 3.5 into 10 to the power of 5, ok.

So, that is the problem. So, it has got a laminar section and it has got a turbulent section. So, now, let us attack this problem.

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Definition of L-averaged Nusselt number tells us

$$Nu_{0-L} = \frac{q_{0-L} L}{k \Delta T}$$

$$q_{0-L} = \frac{1}{L} \left[ \int_0^{x_{tr}} q''_{Lam} dx + \int_{x_{tr}}^L q''_{turb} dx \right]$$

In the laminar section,

$$q''_{Lam} (x) = 0.332 \frac{k \Delta T}{x} Pr^{1/3} Re_x^{1/2}$$

Hence,

$$\int_0^{x_{tr}} q''_{Lam} dx = 0.664 k \Delta T Pr^{1/3} Re_{x_{tr}}^{1/2}$$

The diagram shows a horizontal plate with a coordinate system  $x$  starting from the leading edge. A vertical line marks the transition point  $x_{tr}$ . To the left of  $x_{tr}$ , the flow is labeled 'laminar' with a velocity profile  $u_{laminar}$  shown as a smooth curve. To the right of  $x_{tr}$ , the flow is labeled 'turbulent' with a velocity profile  $u_{turbulent}$  shown as a flatter curve. The ambient fluid is indicated by arrows pointing towards the plate.

Question 2, ok. So, the definition definition of L averaged Nusselt number tells us Nusselt number 0 to L k into delta T. So, if you look at the problem now there is a section in which it is kind of like this and then you have this particular thing. So, up to this is basically the x transition that point and so, this part is your laminar and beyond this it is basically turbulent all right, beyond this it is basically turbulent.

So, your  $q_{0-L}$  double prime 1 over L is basically a sum total from 0 to x transition this is  $q_{laminar}$  into dx plus 0 to L minus x t r  $q_{turbulent}$  into dx, ok. So, in the laminar section when you write as  $q_{laminar}$  for Prandtl number greater than 1 at any point x, if you recall your old stuff it will be this is an isothermal plate, so, this is  $0.332 k$  into delta T by x Prandtl number one third Reynolds number half, ok. So, hence 0 to x t r laminar dx gives you  $0.664 k$  delta T Prandtl number one third, Reynolds number up to the transition is about half ok. So, for the turbulent section we do the same.

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So we do same for turbulent section,

$$h''_{turb}(x) = \rho C_p U_{\infty} \Delta T St_x$$

$\rho C_p U_{\infty} \Delta T St_x$  (Colburn)

$$0.0296 Re_x^{-1/5} \text{ (Prandtl)}$$

$$q''_{turb} dx = 0.037 \rho C_p U_{\infty} \Delta T Pr^{-2/3} \left(\frac{U_{\infty}}{\gamma}\right)^{-1/5} (L-x)^{4/5}$$

$$Nu_{0-L} = \frac{q''_{0-L} L}{k \Delta T} = Pr^{1/3} \left[ 0.664 Re_{x,trans}^{1/2} + 0.037 (Re_L - Re_{x,trans})^{4/5} \right]$$

Do the same for the turbulent section for turbulence section so,  $q''_{turb}$  double prime  $x$   $\rho C_p U_{\infty} \Delta T$  Stanton number  $x$ , ok. So, now this particular expression if you look at the different types of expressions for this ok, you have of course, there is the Prandtl two third half  $C f x$  thing which is once again this is given as 0.0296 into Reynolds to the power of minus one fifth is Prandtl if you recall this is basically Colburn once again if you recall ok. So, this is the not this whole thing, but just a standard number part ok. So,  $0$  to  $L$  minus  $x$   $T r q''_{turb} dx$  is equal to  $0.037 \rho C_p U_{\infty} \Delta T Pr^{-2/3} U_{\infty} \gamma^{-1/5} (L-x)^{4/5}$  ok.

So, in conclusion your Nusselt number from  $0$  to  $L$  is basically given as  $q''_{0-L}$  double prime  $L$  by  $k$  into  $\Delta T$  which will be Prandtl one third  $0.664$  Reynolds transition to the power of half plus  $0.037$  Reynolds number  $L$  minus Reynolds number  $x$   $t r$  to the power of  $4$  fifth, all right. So, this gives you the total expression for your Nusselt number.

So, it is basically an addition of two quantities basically one is the under Reynolds number part one is the laminar part and one is the turbulence part. So, that is a very simple way of and we have used two additional relationships one is the  $C f x$  which is Prandtl and one is basically the Stanton number relationship which is from Colburns analogy, all right.

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- (a) Show that the following general relationships exist among the Stanton, Nusselt, Reynolds, Péclet, and Prandtl numbers in boundary layer flow:

$$St_x = \frac{Nu_x}{Re_x Pr} = \frac{Nu_x}{Pe_x}$$

- (b) Show that the Colburn analogy (7.78) also applies to the laminar section of the boundary layer near an isothermal wall if the fluid has a Prandtl number in the range  $Pr \geq 0.5$ .

So, let us look at the next problem. This is this problem is even simpler. So, you can see that the general relationship that exist between Reynolds number and Prandtl number is given by this this we already did multiple times. So, it will be a nice thing to work it out once again and show that the Colburn analogy also applies to the laminar section of the boundary layer for an isothermal plate if the fluid has got a Prandtl number which is greater than 0.5. So, far has been we have been doing it for turbulent flow only let us look at it that whether Colburn analogy is also applicable for laminar flow.

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The image shows a handwritten derivation in a digital note-taking application. The derivation starts with the definition of the Stanton number  $St_x = \frac{h_x}{\rho c_p u_x}$  and the Nusselt number  $Nu_x = \frac{h_x x}{k}$ . It then shows the relationship  $St_x = \frac{Nu_x}{Re_x Pr}$  and  $St_x = \frac{Nu_x}{Pe_x}$ . The derivation then uses the Colburn analogy for laminar flow,  $\frac{1}{2} C_{f,x} = 0.332 Re_x^{-1/2}$ , and the Nusselt number formula for laminar flow,  $Nu_x = 0.332 Re_x^{1/2} Pr^{1/3}$ . The final result is  $St_x Pr^{1/3} = \frac{1}{2} C_{f,x}$ , which is noted as being applicable for laminar flow too.

Q3:

a)  $St_x = \frac{h_x}{\rho c_p u_x} = \frac{h_x \frac{x}{k}}{\rho c_p u_x \frac{x}{k}} = \frac{Nu_x \frac{k}{x}}{\rho c_p u_x \frac{x}{k}} = \frac{Nu_x}{Re_x Pr} = \frac{Nu_x}{Pe_x}$  (1)

b)  $\frac{1}{2} C_{f,x} = 0.332 Re_x^{-1/2}$  (2)  
 $Re_x = \frac{u_x x}{\nu}$   
 Using Eq. (1) and  $Nu_x$  formula. ( $Pr > 0.5$ )  
 $St_x Pr^{1/3} = \frac{Nu_x}{Re_x Pr} = \frac{0.332 Re_x^{1/2} Pr^{1/3}}{Re_x Pr} = \frac{0.332 Pr^{1/3}}{Re_x^{1/2} Pr} = 0.332 Re_x^{-1/2} Pr^{-2/3}$   
 $St_x Pr^{1/3} = \frac{1}{2} C_{f,x}$  → applies for laminar flow too.



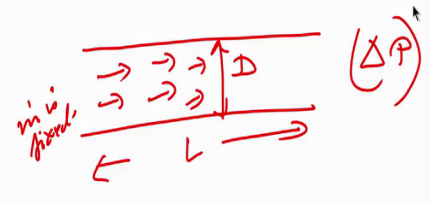
So, question 3 this is easy enough question. So, your Stanton number if you write it from the definition is  $h x$  by  $\rho C_p U_\infty$ . So, if you do it  $h x$  divided by  $\rho C_p U_\infty$   $x$  by  $k$  into  $k$  by  $x$  ok. So, this will be Nusselt number  $x$  into  $\alpha$  by  $U_\infty$  into  $x$  which will be Nusselt number  $x$  into  $\gamma$  by  $U_\infty$  into  $x$  divided by  $\alpha$  by  $\gamma$ . So, this will further lead to Nusselt number by Peclet number which will give you Nusselt number divided by Reynolds number into Prandtl number it is very simple math. So, if you show it this is how this thing was done just starting from the basic definition of your Stanton number.

So, the b part of the question. So, this is the a part ok. So, the Colburn analogy is  $1$  by  $C_f x$  is equal to  $0.332$  to establish the Colburn analogy for the laminar, let us write the  $C_f x$  term also for the for the laminar boundary layer where your Reynolds number  $x$  is given as  $U_\infty$  into  $x$  by  $\gamma$  all right. So, the left side can be estimated of this Colburn analogy can be estimated using this first form of the equation which is basically what we did here, ok. So, this is the second ok. So, using equation – I and the  $N U x$  formula formula greater for Prandtl number greater than  $0.5$  we get Stanton number  $x$  Prandtl number two third is equal to Nusselt number  $x$  Reynolds number  $x$  Prandtl number Prandtl number two third.

So, now, we can substitute for the  $C_f x$  part. So, this is Prandtl number one third Reynolds number half divided by Reynolds number into Prandtl number Prandtl number two third. So, where using the Stanton number expression, ok. So, this once again gives rise to  $0.332$  Reynolds number to the power of minus half see in other words this clearly shows that Stanton number Prandtl number to the power of two third is equal to half  $C_f x$  this is applicable also for laminar applies for laminar boundary layer too interesting revelation all right, it is applicable for laminar boundary layer too, got it, ok.

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Consider the flow of a fluid through a tube of fixed diameter  $D$  and length  $L$ . The mass flow rate  $\dot{m}$  is also fixed. The only change that may occur is the switch from laminar to turbulent flow because the Reynolds number  $Re$  happens to be in the vicinity of 2000. In either regime, the flow is fully developed. Calculate the change in the pumping power required as the laminar flow is replaced by turbulent flow.



The diagram shows a horizontal tube of length  $L$  and diameter  $D$ . Arrows indicate flow from left to right. A note on the left says "m dot is fixed". A pressure drop symbol  $(\Delta P)$  is shown on the right. The text above the diagram is underlined in red.

Now, so, that takes care of the third problem. The fourth problem; so, there is a fluid that is flowing through a tube which has got a fixed diameter of about  $D$  there is a fluid flow through the tube it has got a diameter of about  $D$  and length is about  $L$   $\dot{m}$  is fixed the only change that can happen from laminar to turbulent flow because Reynolds number happens to be in the vicinity of 2000, there is for internal fluid flow that is the transition Reynolds number in either regime the flow is fully developed.

So, this is an important assumption that is given. Calculate the change in pumping power required as the laminar flow is replaced by the turbulent flow. Pumping power is basically given by the how much pressure that you need to put and correspondingly how much mass you are pumping because of that ok. So, let us look at this problem, it is a simple problem.

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Q4: pumping power is proportional to the product  $m \dot{\Delta} P$

$$P = \frac{1}{\rho} m \dot{\Delta} P$$

$$\therefore \Delta P = f \frac{4L}{D} \frac{1}{2} \rho U^2$$

$\left( \begin{matrix} L \\ D \\ m \end{matrix} \right)$  unchanged  
 $U$ : mean velocity also does not change.

$$U = \frac{m}{\frac{\pi}{4} \rho D^2}$$

$$Re_D = \frac{UD}{\nu}$$

$$\frac{P_{turb}}{P_{lam}} = \frac{f_{turb}}{f_{lam}} = \frac{0.079 Re_D^{-1/4}}{16 Re_D^{-3/4}} = 0.00494 Re_D^{3/4}$$

$Re_D \sim 2000$

$$\frac{P_{turb}}{P_{laminar}} = 1.48$$

So, question 4, so, the pumping power is proportional basically proportional to the product  $m \dot{\Delta} P$ . So,  $P$  is basically equal to  $\frac{1}{\rho} m \dot{\Delta} P$ , where your  $\Delta P$  is the one that actually is the, is a key player because it is  $f$  by  $4L$  by  $D$  into half  $\rho U^2$  this we know, where  $f$  is the friction factor.

Now, all these other terms this  $L$ ,  $D$ ,  $U$  and the  $m \dot{\Delta}$  these are kind of unchanged all right, these are unchanged, ok. So, the mean velocity also does not change  $U$  which is the mean velocity also does not change ok. So,  $U$  is basically what mean velocity is  $m \dot{\Delta}$  by  $\pi$  by  $4 \rho D^2$  where your Reynolds number  $Re_D$  is given as  $UD$  by  $\nu$  all right. So, the only thing that can change is the friction factor all right, ok.

So, the pumping power turbulent divided by the pumping power laminar is basically given as  $f$  of the turbulence which is basically the friction factor and  $f$  of laminar because rest of the things are kind of the same. So, for the turbulent friction factor it is given as  $0.079$  to Reynolds number  $Re_D$  to the power of minus one fourth, if you recall your notes and for laminar it is  $16$  by  $Re_D$ .

So, the total thing that we get is  $0.00494$  into Reynolds number  $Re_D$  to the power of  $3/4$ . If the Reynolds number is of the order of roughly about  $2000$ , ok. The  $P_{turbulent}$  by  $P_{laminar}$  these are not pressure these are basically the pumping power is about  $1.48$ . So, the pumping power of the turbulent flow or the pumping power to maintain the turbulent flow is about  $1.48$  than the corresponding laminar counterpart all right.

So, this is one of the most important another interesting short problem and we have done.

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Water is being heated in a straight pipe with an inside diameter of 2.5 cm. The heat flux is uniform,  $q_w'' = 10^4 \text{ W/m}^2$ , and the flow and temperature fields are fully developed. The local difference between the wall temperature and the mean temperature of the stream is  $4^\circ\text{C}$ . Calculate the mass flow rate of the water stream, and verify that the flow is turbulent. Evaluate the properties of water at  $20^\circ\text{C}$ .

*h* : heat transfer coefficient  
Nusselt number  $\rightarrow$  radiation *but we have already established*  
Reynolds Number  $\rightarrow$   $U = ??$   
mass flow rate.  
 $\dot{m} = \frac{\pi}{4} \rho D^2 U$   
 $\dot{m} = 1.281 \text{ kg/s}$

The last problem in this particular series is that water is being heated in a straight pipe with an inside diameter of 2.5 centimeter. The heat flux is uniform and the flow and the temperature fields are fully developed. The local difference between the wall temperature and the mean temperature is about 4 degree Celsius. Calculate the mass flow rate of the water stream and verify that the flow is turbulent evaluate the water properties at about 20 degree Celsius, all right.

So, in order to do this ok, so, I am not going to solve this problem in total, but what I can give you are some hints. So, that you can work it out it is not a homework problem per say, but this is some problem that will kind of give you and we will post the solution a little later ok, but you should know how this problem actually works. So, in order to calculate, ok, I am showing you the steps  $h$  needs to be calculated first which is the heat transfer coefficient that needs to be calculated first, then you need to calculate the Nusselt number, the Reynolds number and lastly the mass flow rate.

So, Nusselt number, Reynolds number and the heat transfer coefficient all can be calculated using the standard correlations and the data that is given over there. So, once you calculate the Reynolds number you can calculate  $U$  and from  $U$  you can calculate what will be going to be the mass flow rate and check what will be the velocity of  $U$ .

From  $U$  you can find out what will be the mass flow rate because  $\dot{m}$  is nothing, but one forth  $\rho D^2 U$ .

So, you can once you know  $U$  you can calculate. So, for knowing  $U$  you need to calculate the Reynolds number all right and for knowing the Reynolds number you need to know the heat transfer coefficient and a Nusselt number. Nusselt number and Reynolds number are connected through the correlation or the correlation by the relation that we established that we have already established. Do this problem and try to see if you get a good match I can give you the answer  $\dot{m}$  is about 0.281 kg per second. Try to see if you can match that ok.

So, with this we finish basically our turbulence small session in which we have shown you how to work out certain problems some numerical problems, and it will also give you an idea how to use the relationships why and how they should be used. Also give you a good idea about how the turbulence physics is in the limited time that was possible and what turbulent convection is all about, ok.

So, in the next class, we will look at some of the other problems in the series which is basically going to deal with droplets, sessile droplets and normal contact free droplet us and try to see that how we one should evolve a methodology for solving those problems.

Thank you.