

Introduction to Mechanical Vibration
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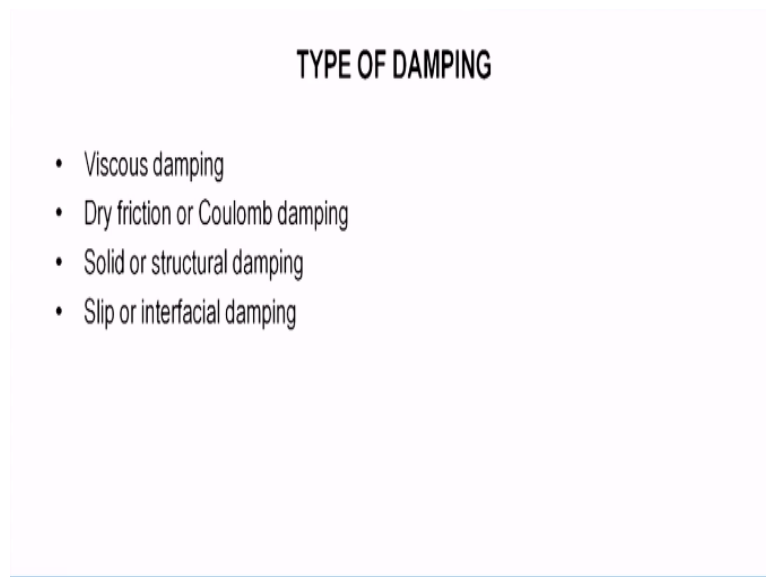
Lecture - 08
Damped Free Vibration

Welcome to the lecture on Damped Free Vibration. So in the previous lecture we discussed the free vibration of single degree of freedom system having mass and stiffness. There was not a damping element so there was not a dissipation of energy in the system. So the vibration of the system was maintained forever. There was no reduction in the amplitude of the vibration.

Now if we introduce a damping element in the system means we have M we have K now we have the C the damping element in the system then what will happen like in the case of the free undamped vibration we had the vibration of the system with the natural frequency ω_n now will the damped system will vibrate with the same frequency as ω_n or different, will the amplitude of the vibration remain constant or it will change with the time, will it decay with the time, will it increase with the time.

So all these questions that we will discuss in this lecture.

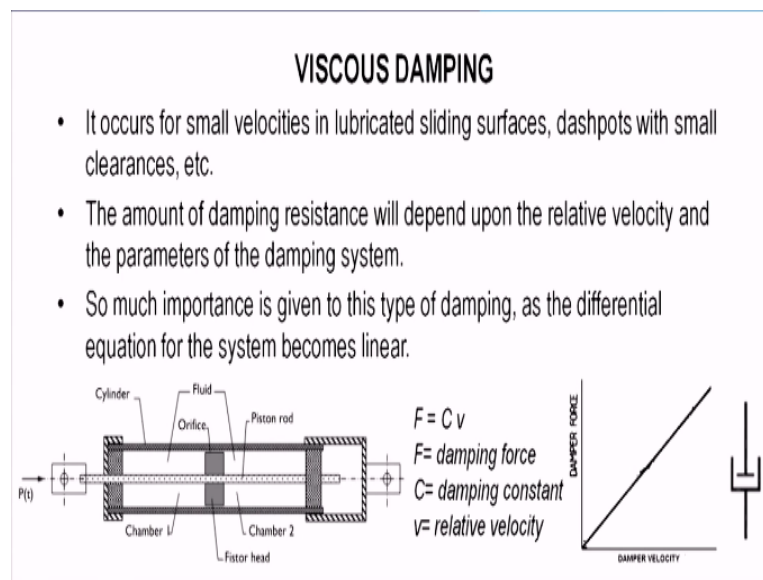
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So we will discuss type of damping because when we want to discuss the damped free vibrations there are several types of damping the most common is the viscous damping

because this is very important and simple so it is used vastly then we have the dry friction or Coulomb damping then solid or structural damping then slip or interfacial damping. So here viscous damping, viscous damping occurs when a piston in certain fluid that has some viscosity. The piston moves in the fluid and the fluid is allowed to pass through some orifice so there is the small velocity of the movement of the piston.

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So this kind of damping we can see that the force and velocity relationship that is linear. Here is the force $F = C v$ where F is the damping force C is the damping constant and v is the relative velocity. Here we can see one damper and if some force is applied here is the piston and it is moving and here is the orifice through which the fluid is passing so this kind of damper is called the viscous damper because they are based on the principle of viscous damping.

Why we give this so important to this kind of damping because here the force viscous this damping force is directly proportional to the velocity so $\dot{x}$. So if we use this put this force in the equations differential equation of motion of the system that makes the equations differential equation linear so that is why we this damping is more important. Several times we use the equivalent of the other type of damping are equivalent viscous damping for the other type of damping just for this reason.

We represent this damping C like this symbol then another type of damping is friction dry friction damping or Coulomb damping. So whenever a mass is sliding on some surface there is friction between the two surfaces and due to the friction there is some energy dissipation

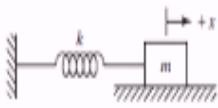
because the basic role of the damping is the energy dissipation so like in the viscous damping it was the energy dissipation in the fluid and energy dissipation is usually results into heat and sound.

So here when a mass is sliding on a surface due to the friction there is energy dissipation into heat. So when there is dry friction here the damping resistance is constant because it depends on μ and therefore it is independent of the rubbing velocity.

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DRY FRICTION OR COULOMB DAMPING

- It occurs when two machine parts rub against each other, dry or unlubricated
- The damping resistance is constant and independent of the rubbing velocity

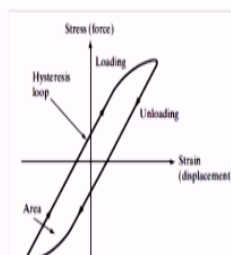


So here we can see we have mass that may slide on the surface and due to the friction between the two surface there could be the energy dissipation and this kind of mechanism is called Coulomb Damping.

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SOLID OR STRUCTURAL DAMPING

- It is due to the internal friction of the molecules
- The stress-strain diagram forms a hysteresis loop, the area of which represents the energy dissipated due to molecular friction per cycle per unit volume
- The size of the loop depends upon the material, frequency and the amount of dynamic stress



Now we have another type of damping that is called Solid or Structural Damping. It is also called a Material Damping or Hysteresis Damping. So there is various name to this damping. So we can see the stress-strain diagram and stress strain diagram when there is a material that has this kind of damping. The stress strain diagram may form a hysteresis loop this loop that is the area there is some area between stress and strain and the area shows the energy dissipated in one cycle of this loop.

So this energy shows the energy that is due to the friction in the molecules so there are the molecular deformation in the material. So therefore this damping depends on material, the frequency and amount of the dynamic stress.

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SLIP OR INTERFACIAL DAMPING

- Energy is dissipated by microscopic slip on the interfaces of machine parts in contact under fluctuating loads
- Microscopic slip also occurs on the interfaces of machine elements forming various types of joints
- The amount of damping depends amongst other things upon the surface roughness of the mating parts, the contact pressure, and the amplitude of vibration.
- It is a non-linear type of damping

There is a one more type of damping that is Interfacial Damping or Slip Damping so we have several mechanical parts and there is roughness between these parts and when they come into contact they meet. There is some micro level of interfaces and rubbing and that makes the energy loss, energy dissipation. For example, there is some joints like bolted joints and some machine parts so this damping depends on the roughness of the surfaces the contact pressure and the amplitude of the vibration.

Now in this lecture we will discuss, we will take the viscous damping as we said that in the taking the viscous damping there is certain advantage that it makes the differential equation linear and so it is easier to solve these kinds of equations. Therefore, we start with a single degree of freedom system having viscous damping. So we can see here we have a spring damper and mass system and again it is a free vibration case.

So due to the free vibration we will give some initial disturbance to the system and then we will isolate this mass to make a free-body diagram means we have to show the forces that is acting on the body and then we have to apply the Newton's law of motion to write the differential equation.

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The slide contains handwritten notes and diagrams for a mass-spring-damper system. On the left, a diagram shows a mass m attached to a wall by a spring with constant k and a damper with coefficient c . The displacement is x and velocity is $\dot{x}$. Below this, the differential equation is derived: $\Sigma F = m\ddot{x}$, $-kx - c\dot{x} = m\ddot{x}$, leading to $m\ddot{x} + c\dot{x} + kx = 0$. The characteristic equation is $s^2 + \frac{c}{m}s + \frac{k}{m} = 0$. The roots are $s_{1,2} = \frac{-c \pm \sqrt{c^2 - 4mk}}{2m}$. The slide then branches into three cases: 1. Real and distinct roots: $x = C_1 e^{s_1 t} + C_2 e^{s_2 t}$. 2. Equal and real roots: $x = (C_1 + C_2 t) e^{st}$. 3. Complex roots: $\alpha \pm i\beta$, $x = e^{\alpha t} (C_1 \cos \beta t + C_2 \sin \beta t)$. On the right, the 'critical damping' case is detailed: $c^2 - 4mk = 0$, $c = 2\sqrt{km} = 2m\omega_n$. The damping factor $\xi = \frac{c}{c_c}$ is introduced, with $\frac{c}{m} = \frac{c_c}{m} = 2\omega_n$. The roots are $s = -\xi\omega_n \pm \omega_n \sqrt{\xi^2 - 1}$, leading to $s_1 = [-\xi + \sqrt{\xi^2 - 1}] \omega_n$ and $s_2 = [-\xi - \sqrt{\xi^2 - 1}] \omega_n$.

So let us do it so we take the single degree of freedom system so this is the system having spring damper and this is mas. So we disturb the system so we put here X now we make a free-body diagram so this is the mas M and because we have pulled this system X and it has the velocity X dot so the forces. So let us assume this is the direction of X double dot and because we have pulled the system with X we have spring force KX and this is X dot.

So we will have here CX dot because when we pull it the damper or spring will apply force opposite to this pull. So this is our system. Now we have to write the equations so here we have sigma F = MX double dot. So sigma F we have two forces spring force and damping force so here it is - KX because their direction is opposite to the direction of X double dot that is why we are giving negative so - KX and - CX dot = MX double dot.

So we can write MX double dot + CX dot + KX = 0 and this is the equation of motion of free vibration of this system. Now you see if we put C = 0 means if we remove the damping our equation results MX double dot + KX = 0. So it is the same as undamped free vibration case. Now we have to solve this equation and to solve this kind of differential equation we assume a solution X = E power ST.

So we have assumed one solution $X = e^{st}$ and we put this into this equation 1 so put this $X = e^{st}$. So $\dot{X} = s e^{st}$ and $\ddot{X} = s^2 e^{st}$. So we put these values in this equation so $M s^2 e^{st} + C s e^{st} + K e^{st} = 0$. So here we can take e^{st} outside so $M s^2 + C s + K = 0$.

So because this is exponential so this cannot be zero. So e^{st} cannot be zero that is why this term should be zero. So $M s^2 + C s + K = 0$. Now this is a quadratic equation in s so we can write the roots it will have two roots. So let us write the roots so $s = s_1$ and $s_2 = \frac{-C \pm \sqrt{C^2 - 4MK}}{2M}$. So $-B$, B is C so it is $-C \pm \sqrt{C^2 - 4MK}$, A is M and C is K upon $2A$.

So $2A$ is M now here we can write so $\frac{-C \pm \sqrt{C^2 - 4MK}}{2M}$. Now the solution of this differential equation if we have roots so roots these roots s_1 and s_2 they can be real and distinct. So we will have solution $X = C_1 e^{s_1 t} + C_2 e^{s_2 t}$. If they are equal and real so we will have solution $X = (C_1 + C_2 t) e^{st}$.

So because s is s_1 or s_2 anyone if they are complex roots so we have if complex roots so that is let us say $\alpha \pm i\beta$. So this is if these kind of complex roots we have $s_1 = \alpha + i\beta$ $s_2 = \alpha - i\beta$. So we have $X = e^{\alpha t} [C_1 \cos \beta t + C_2 \sin \beta t]$. So these are the solutions depending on the different conditions of the roots so roots could be either real and distinct or real and equal or they could be a complex pair.

So in these three conditions the solution will be given by these equations. So now let us so this is our equation now we define certain terms so let us see. Now we define critical damping so critical damping comes when we have the equal and real roots so equal and real roots means this term is zero. So $C^2 - 4MK = 0$. So here we have and this C is CC so we call it $CC^2 - 4MK = 0$.

So $CC = \sqrt{4MK} = 2\sqrt{KM}$ or we can write it 2 into M into ω_n because ω_n is $\sqrt{K/M}$ and M will cancel out. So the critical damping is the damping that we obtain when our roots are equal and real it means this part is zero so this under root part is zero so we get $CC = 2\sqrt{KM}$ or $2M\omega_n$. Another thing we define that is damping

factor.

So damping factor $\zeta = C \text{ by } CC$. So C is the actual damping value and CC is the critical damping value. So the ratio of this is $C \text{ by } CC$. Now we write this equation of S in this term so S equal to so first we write $C \text{ by } M$ so $C \text{ by } M = C \text{ by } CC \text{ into } CC \text{ by } M$. So ζ into $CC \text{ by } M$ $CC \text{ by } M$ is $2 \omega N$ so $= 2 \zeta \omega N$. Now we write here we can write here $- C \text{ by } 2M + - 1 \text{ by } 2 \text{ root } C \text{ by } M \text{ whole square} - 4 K \text{ by } M$ because here M will be $M \text{ square}$ when it will go inside.

So $4K \text{ by } M$ $K \text{ by } M$ is $4 K \text{ by } M$ and $K \text{ by } M = \text{nothing}$, but $\omega N \text{ square}$ because ωN is $\text{root } K \text{ by } M$. So now we can write these values in this equation so S equal to $- C \text{ by } M$ is $2 \zeta \omega N$ and it is $\text{by } 2$ so it is $- \zeta \omega N + - 1 \text{ by } 2 \text{ under root } C \text{ by } M \text{ square}$. So $C \text{ by } M \text{ square}$ is $2 \zeta \omega N \text{ square} - 4 \text{ into } \omega N \text{ square}$.

So here we will get $- \zeta \omega N + - 1 \text{ by } 2$ and this is 4 and this is 4 so 2 will come out and it will cancel out and $\omega N \text{ square}$ ωN will cancel out so this will be cancelled out $\omega N \text{ root } \zeta \text{ square} - 1$. So here we can write $= - \zeta + - \text{under root } \zeta \text{ square} - 1 \omega N$. So these are our roots so our $S_1 = - \zeta + \text{root } \zeta \text{ square} - 1 \omega N$ and $S_2 = - \zeta - \text{root } \zeta \text{ square} - 1 \text{ okay}$.

So these are the two roots S_1 and S_2 that we have obtained now depending on conditions we have to fit these in this three conditions. We have to see in which condition these roots are coming okay. So now I remove this part.

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$\zeta < 1$
 $s_1 = (-\zeta + i\sqrt{1-\zeta^2})\omega_n$
 $s_2 = (-\zeta - i\sqrt{1-\zeta^2})\omega_n$
 $\alpha = -\zeta\omega_n$
 $\beta = \sqrt{1-\zeta^2}\omega_n$
 $x = e^{-\zeta\omega_n t} [C_1 \cos \sqrt{1-\zeta^2}\omega_n t + C_2 \sin \sqrt{1-\zeta^2}\omega_n t]$

$\zeta = 1$
 $s_1 = s_2 = -\zeta\omega_n$
 equal and real roots
 $x = (C_1 + C_2 t) e^{-\zeta\omega_n t}$

$\zeta > 1$
 $s_1 = [-\zeta + \sqrt{\zeta^2 - 1}]\omega_n$
 $s_2 = [-\zeta - \sqrt{\zeta^2 - 1}]\omega_n$
 real and distinct
 $x = C_1 e^{[-\zeta + \sqrt{\zeta^2 - 1}]\omega_n t} + C_2 e^{[-\zeta - \sqrt{\zeta^2 - 1}]\omega_n t}$

So now here if we have $\zeta = 1$. So if we have $\zeta = 1$, so this is 0, this is 0 we will have $s_1 = s_2 = -\zeta$. So it means we have equal and real roots so we will have $x = C_1 + C_2 T$. E power $s_1 T$ is $-\zeta$ into T , okay, for $\zeta = 1$. So if ζ is greater than 1 so ζ is greater than 1 we have $s_1 = -\zeta + \sqrt{\zeta^2 - 1}\omega_n$ and $s_2 = -\zeta - \sqrt{\zeta^2 - 1}\omega_n$. So here we see that because ζ is greater than 1 so this quantity is greater than 0 this is greater than 0.

So these s_1 and s_2 are real and distinct. So we will have this equation so we will have $x = C_1 e^{s_1 T} + C_2 e^{s_2 T}$. So s_1 is $-\zeta + \sqrt{\zeta^2 - 1}\omega_n$ into T + $C_2 e^{s_2 T}$. So s_2 is $-\zeta - \sqrt{\zeta^2 - 1}\omega_n$ into T . So it is $-\zeta$ into ω_n so here it will be into ω_n . Here into ω_n and here is also into ω_n so we have $-\zeta + \sqrt{\zeta^2 - 1}\omega_n$ into T + $C_2 e^{[-\zeta - \sqrt{\zeta^2 - 1}\omega_n] T}$.

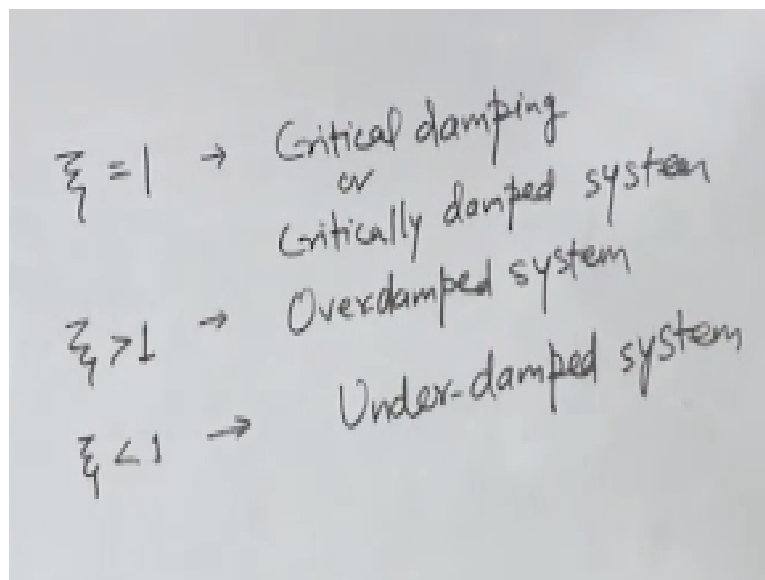
So here we have the condition when ζ is greater than 1. Now another case we will have so let us have the other case. So when ζ is less than 1. So we take a case when ζ is less than 1. So when ζ is less than 1 we will have $s_1 = -\zeta + i\sqrt{1-\zeta^2}\omega_n$ because we can write here is $1 - \zeta^2$. So $\sqrt{1-\zeta^2}$ so $\sqrt{1-\zeta^2}$ can be 1 into ω_n and $s_2 = -\zeta - i\sqrt{1-\zeta^2}\omega_n$.

Now here we see that we are getting a complex pair root. So here if we compare from this $\alpha + i\beta$ we have $\alpha = -\zeta\omega_n$ and $\beta = \sqrt{1-\zeta^2}\omega_n$. So we can write the solution $x = e^{-\zeta\omega_n T} [C_1 \cos \sqrt{1-\zeta^2}\omega_n T + C_2 \sin \sqrt{1-\zeta^2}\omega_n T]$.

So here we have reached to three conditions based on the damping factor ζ . So $\zeta = 1$ so $\zeta = 1$ we get equal and real roots and the solution, solution means the response of the system that is $X = C_1 + C_2 T$ into E power $-\zeta \omega N T$. When ζ is greater than 1 we have got the two real and distinct roots and the response of the system $X = C_1 E^{\text{power } \omega N - \zeta + \sqrt{\zeta^2 - 1} T} + C_2 E^{\text{power } \omega N - \zeta - \sqrt{\zeta^2 - 1} T}$.

However, for the case when the damping factor ζ is less than 1 we got the complex pair roots and for complex pair roots we have $-\zeta \pm j \sqrt{1 - \zeta^2} \omega N$ and we have α , β and we write $X = E^{-\zeta \omega N T}$ and $C_1 \cos \sqrt{1 - \zeta^2} \omega N T + C_2 \sin \sqrt{1 - \zeta^2} \omega N T$. So these three conditions lead to us the three different types of damp systems.

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So these three different type of damp system like $\zeta = 1$ that is called critical damping or critically damped system. If ζ is greater than 1, we call it over damped system or overdamping of the system. When ζ is less than 1, we call it under damped system or underdamping of the system. So we see that in these three different conditions of damping that is critical damping over damping and underdamping the response of the system is different.

So because from the expressions of the response X we can see that their response is governed by different equations and therefore we have if you want to study a system we have to decide that what the damping value we want I mean we want our system to keep in the under

damped region, over damped region or critical damp region because all these different regions has their own characteristics and that we will discuss in detail one by one in the next lectures. So thank you for your attention and see you in the next lecture. Thank you.