

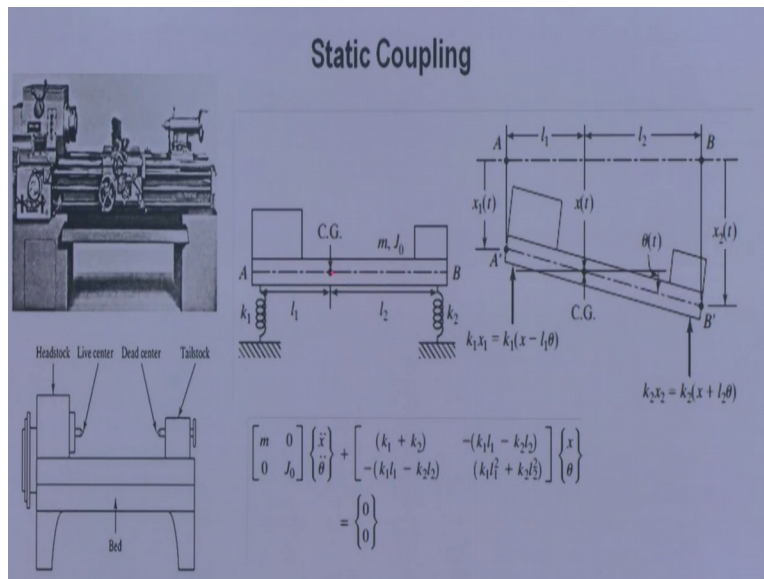
Introduction to Mechanical Vibration
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Lecture - 40
Static and Dynamic Coupling

So welcome to the lecture on multi-degree of freedom system. So today we will discuss the concept of Static and dynamic coupling. So as we know that we can have the - we have defined the degree of freedom so n degree of freedom system if it is a multi-degree of freedom system, so it needs an independent coordinate to represent its configuration at any instant.

Generally, we take these coordinates from the equilibrium position, for example we take from center of mass we can select any other set of degree of freedom.

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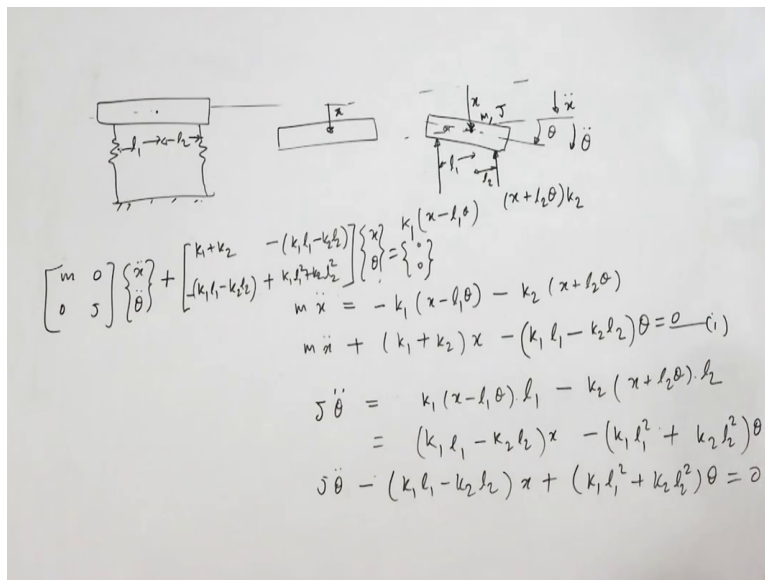
For example, if we have there is a lathe machine and there is a model of this machine, so this machine it has Headstock, Tailstock there is a Bed and this column so we can assume a model of this system and so here there is the elastic stiffness of these columns there is this - this bed okay and we have the center of mass or center of gravity of the system here, so we can because the motion of the system can be in linear as well as in angular directions.

So we can select a degree of freedom passing through center of mass so we can take one the motion of if here x t that shows the motion of the center of mass and the angle θ the rotation angular motion of this system, we can take another set of degrees of freedom that is x_1 t and x_2 t these are another set of the degree of freedom that we can that can represent the system's configuration at any instant.

Also it is possible that we select a degree of freedom that is not on CG but some other point and from there we can get the configuration in terms of these this point, so we can assume some point and its motion in linear and angular motion so we can take that set of coordinates. So all these - these sets they are known as generalized coordinates. Because these coordinates are it depends that they are our choice they are not some property of a system.

They are our selections we select these coordinates and we so for example if we select this coordinate so we have here the center of gravity and its motion x t and then θ t , we take the angle with respect to this axis so what happens.

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So it is a system like this so we have this system and this system is moving here, here with x t and again it is rotating here with θ t , so this is θ and this is x , now if we have this springs k_1 and k_2 and this distances here this is l_1 , l_2 here so the forces acting at this point or this

because the total deflection of this spring is $x - l_1 \theta$, because this is θ and here this is θ , so $x - l_1 \theta$ and into k_1 so that is the force that is acting on this.

Now on this - this spring so the net displacement of this spring is so it is $x + l_2 \theta$ so it is $x + l_2 \theta$ and multiplied by k_2 , so that is the force on this system and the mass of this system is m and the moment of inertia is J , so we assume that the mass of this system is m and the moment of inertia J at the, so center of mass, so now we can write the equation of motion of this system so $m \ddot{x}$ equal to forces.

So these are the, so this is \ddot{x} and that is - so this is opposite so it is $-k_1 x - k_1 l_1 \theta$ and this is again upon it, so $-k_2 x + k_2 l_2 \theta$, so we can rewrite this equation $m \ddot{x} + k_1 x - k_2 x - k_1 l_1 \theta + k_2 l_2 \theta$ and this is equation number one, similarly we for the another degree of freedom that is the rotation or angular motion, so we take, so we have moment of inertia J and $\ddot{\theta}$.

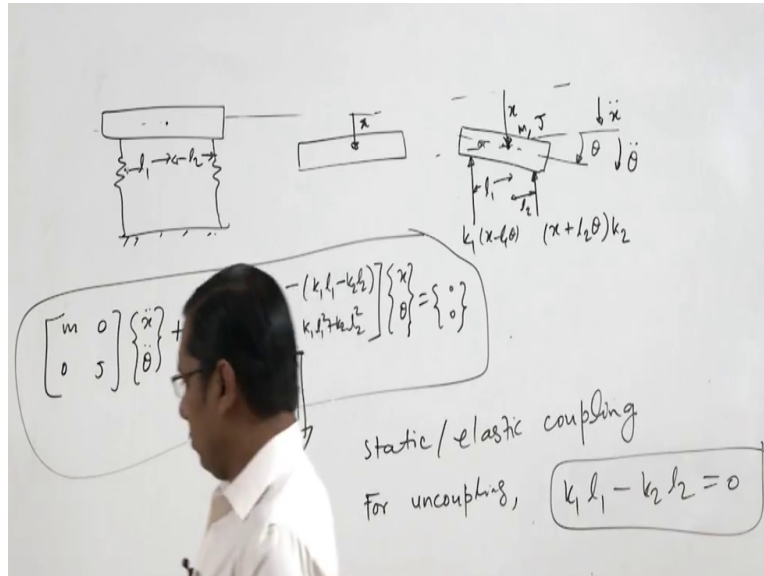
So now we have we take moment about this point, so we have the this is the force and that distance l_1 so this is distance l_1 and this is l_2 , so this is giving a moment about this point, okay, and this $\ddot{\theta}$ is positive as a clockwise and this is a moment in the same direction so it is $k_1 x - k_1 l_1 \theta$ into l_1 and this is in the opposite so it is $-k_2 x + k_2 l_2 \theta$ into l_2 . Now we can rewrite so it is $k_1 l_1 - k_2 l_2 x$ and $-k_1 l_1^2 \theta$, so it is $-k_1 l_1^2 \theta$.

And this is $-k_2 l_2^2 \theta$ and if we take this it is like this, so $J \ddot{\theta} = J \ddot{\theta} - k_1 l_1^2 \theta - k_2 l_2^2 \theta + k_1 l_1^2 \theta + k_2 l_2^2 \theta = 0$, so this is equation number two. Now these equations one and two we can assemble in form of a matrix. So if these equations one and two we assemble so it is $m \ddot{x}$ and $J \ddot{\theta}$ and that is equal to, + we have another matrix so $k_1 + k_2$ and $-k_1 l_1 + k_2 l_2$ so it is $+k_1 l_1 \theta$.

And this is -, so it is - and it is - and this is $-k_1 l_1 - k_2 l_2$ and this is $-k_1 l_1^2 - k_2 l_2^2 + k_1 l_1^2 + k_2 l_2^2$ this is x and θ that is $= 0$ and 0 . So we have arranged so we have this equation, first equation so $k_1 + k_2 x$ so $m \ddot{x}$ and this is 0 , so $+k_1 + k_2 x - k_1 l_1 - k_2 l_2 \theta = 0 = 0$ and here $J \ddot{\theta} - k_1 l_1^2 - k_2 l_2^2 \theta + k_1 l_1^2 \theta + k_2 l_2^2 \theta = 0$.

So now we can have only the matrix form. So now we see that from this expression this matrix expression.

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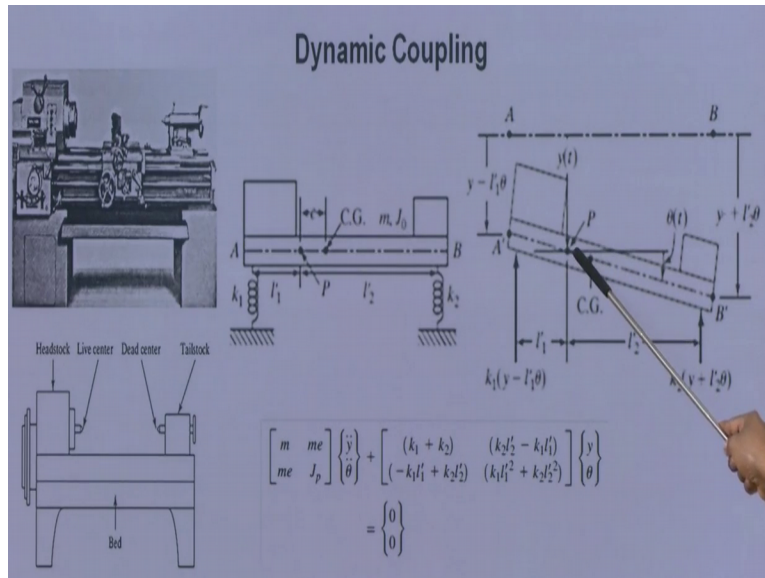


So it is - so this is the matrix expression and from this matrix expression we see that in this equation if we write when the equation in the first equation that is the linear motion degree of freedom we have also this theta term and therefore the system is coupled and similarly, in the angular motion there is the x because - k₁ l₁ - k₂ l₂ x, so this is however there is no x double dot there is no theta double dot, okay.

But here x and theta they are, there is so when there is such a situation so we called it static or elastic coupling - elastic coupling - coupling, but it is not this system is only static coupling not the dynamic coupling there is no any coupling of the acceleration terms however if you want to uncouple this system we have to make these terms zero, because if it is zero the theta component will be zero in equation one and this is zero then x component will be zero in the theta equation.

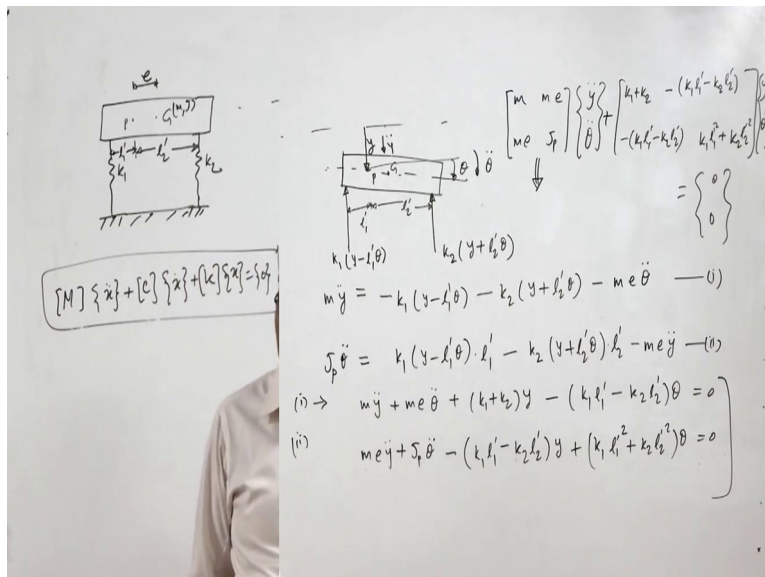
So therefore for uncoupling k₁ l₁ - k₂ l₂ = 0 this is the condition. Now we take the another example I mean another set of degree of the coordinates and another set of degree of freedoms or coordinates - general coordinates.

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So here we can see we have now here is a point P and we are, we have selected the y t that is the linear motion of this point and angular motion theta as a generalized co-ordinate and now we will study this system.

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So now we have a system like this so the same system here is G and we take a point P that is at a distance here the distance is e and here we have this springs and here we have this spring that is k1 and k2. now we take this distance here this is l1 dash and this is l2 dash, so from point P to this spring okay, there is l1 dash and this is the l2 dash, now again we have to, we have displaced this system so linearly as well as angular motion.

So we have here y t no here we have P and this P so we have P this is G and here is M and J the mass M and J so mass M is center this is center of mass and J is moment of inertia about this point, now we have this y and this is θ - θ and here we have this is spring forces so now we have this points l_1 dash and this is l_2 dash, so we can write now we should understand that when we take it here.

Now, we write equation of motion so $m \ddot{y} =$, so here it is $k_1 y - l_1$ dash into θ and here it is $k_2 y + l_2$ dash into θ , so because we are here y - because it is y - and this is $y + l_2$ dash θ , so $m \ddot{y}$ so exalation here is this is y double dot and this is θ double dot, so $m \ddot{y} =$, so this is opposite - $k_1 y - l_1$ dash $\theta - k_2 y + l_2$ dash θ . However, there is θ double dot.

So there is the inertia force that is m times $e \theta$ double dot, so it is $m y$ and inertia of force is opposite so - $m e \theta$ double dot, so that is equation number one, similarly for the angular degree of freedom so J_P , J_P is - J_P is the moment of inertia about the point P , okay, so J_P in to $\theta =$, so now about this point θ double dot is in clockwise so this is giving clockwise, so it is $k_1 y - l_1$ dash θ into l_1 dash.

However, it is giving opposite so - is $k_2 y + l_2$ dash θ into l_2 dash again there will be a inertia moment due to the degree of freedom y double dot, so that moment will be $m y$ double dot the inertia force and into e , because this distance is, the distance between these two points e , so the moment is - $m e y$ double dot, because $m y$ double dot in to e , now we can the equation two. So equation one we can write again as $m \ddot{y} + m e \theta$ double dot.

Then y terms, so we have $k_1 y$ and $k_2 y$ so it is $k_1 + k_2 y$ and here $k_1 l_1$ dash θ and this is $k_2 - k_2 l_2$ dash θ , so it is - $k_1 l_1$ dash - $k_2 l_2$ dash $\theta = 0$. similarly, equation 2 can be written as $m e y$ double dot + $J_P \theta$ double dot then y term, so it is - $k_1 l_1$ dash and here $k_2 l_2$ dash and that is + so it is - $k_2 l_2$ dash y and then θ . So θ term is here $k_1 l_1$ dash square so it is + $k_1 l_1$ dash square and this is $k_2 l_2$ dash square with - so + $k_2 l_2$ dash square and it is $\theta = 0$.

So now again we can write these two equations in matrix form, so we can write them as $m, m_e, J, \ddot{\theta} = -k_1 \theta - k_2 \theta$ and here $k_1^2 + k_2^2$ and this is θ and equal to 0. So we have written the equation in form of a matrix, now we see that now for this matrix there is no that is not diagonal.

So there is the of diagonal terms are non-zero and we can see the effect in the equation that the exalation terms are coming in the equation each equation so here in equation of $\ddot{y} - y$ motion there is $\ddot{\theta}$ term and in $\ddot{\theta}$ there is \ddot{y} , so this is called coupling mass coupling or that is also dynamic coupling. And this is as we have already discussed this is the static coupling.

So here we have static and dynamic both coupling static coupling from here due to the stiffness matrix and the dynamic coupling due to, here the mass matrix. Now if you take a general system that is having the mass damping and stiffness matrix, so if we have system like $M\ddot{x} + C\dot{x} + kx = 0$, so this is a system. So if the stiffness matrix is not diagonal then the system has elastic or static coupling.

However, if the damping matrix is not diagonal it has a damping or velocity coupling, if the mass matrix is not diagonal the system has mass or inertia coupling. These two couplings due to damping and mass matrix that is called the dynamic coupling and this is the static coupling. So here we discussed the static and dynamic coupling. So we stop our lectures and I wish you all the best for this course and the examinations, thank you.