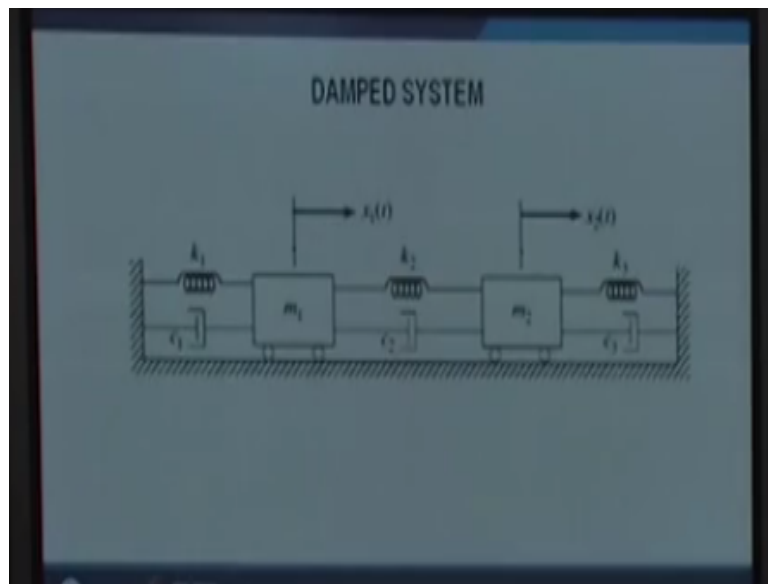


Introduction to Mechanical Vibration
Prof. Anil Kumar
Department of Mechanical and Industrial Engineering
Indian Institute of Technology – Roorkee

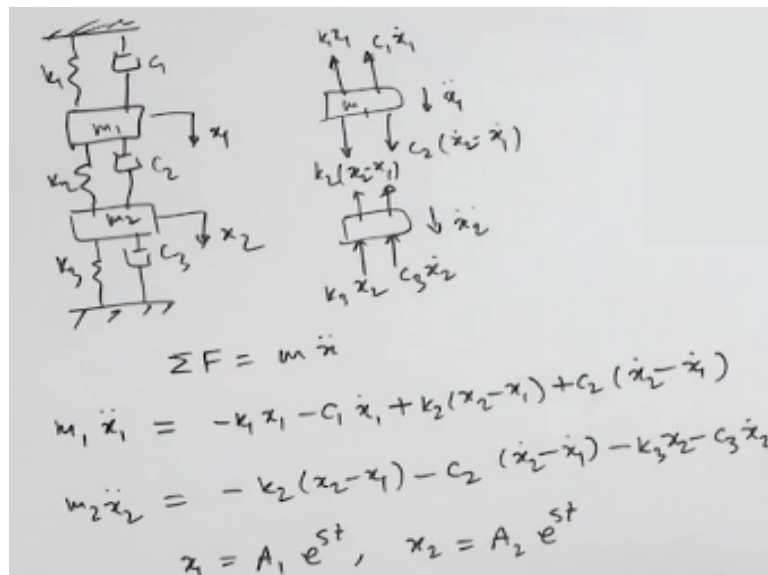
Lecture - 29
Damped Free Vibration

So welcome to the lecture on two degree of freedom systems. So we will consider now the damped free vibration of two degree of freedom systems.

(Refer Slide Time: 00:32)



(Refer Slide Time: 00:38)



So we have such a system that is two degree of freedom, so this is the mass m_1 , m_2 ; this is k_1 , k_2 , k_3 ; c_1 , c_2 , c_3 and this is the x_1 and x_2 , the two degree of freedoms. Now to study these systems what we have to do, we have to first write the equations of motion so we make

the free body diagram and here one more force is added in compare to the previous analysis that is the force due to the damping element.

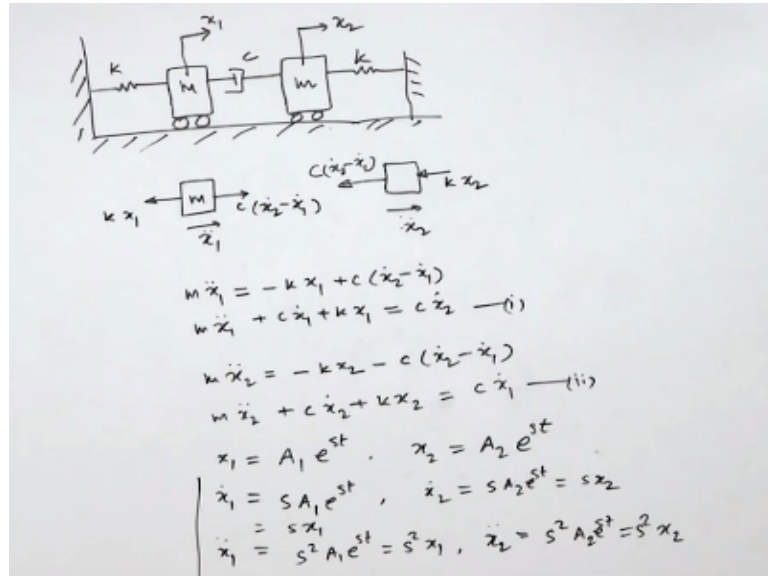
So damping element c_1, c_2, c_3 ; they will act certain forces in these equations and so we will write the equations and that we can write. So if this is m_1 , we have the spring force here, damping force here okay and then similarly here we have these forces $k_2 x_2 - x_1$, $c_2 \dot{x}_2 - \dot{x}_1$ and here $k_1 x_1$ and $c_1 \dot{x}_1$. Similarly, this mass will have same forces in the opposite directions and here these forces $a_3 x_2$ and $c_3 \dot{x}_2$.

So and we take this like this. So these are the forces on the free body diagram and when we apply the $\Sigma f = m \ddot{x}$ for both the cases, we can write the equations of motion. So here for example $m_1 \ddot{x}_1 = -k_1 x_1 - c_1 \dot{x}_1 + k_2 x_2 - x_1 + c_2 \dot{x}_2 - \dot{x}_1$. Similarly, for mass m_2 , so $m_2 \ddot{x}_2 = -k_2 x_2 - x_1 - c_2 \dot{x}_2 - \dot{x}_1 - k_3 x_2 - c_3 \dot{x}_2$.

So this is the equation of motion and to solve these equations of motion, we will assume some solution like we will assume $x_1 = A_1 e^{st}$ and $x_2 = A_2 e^{st}$. So when we put here where A_1 and A_2 , they are the amplitude so their ratio will define the mode shape and A_1, A_2 they are constant, s is also a constant and when we put in these equations, we will find s .

And we solve for the s that is the root of that characteristic equation and from there we will find the final solutions with the help of the initial conditions. So rather than solving on this system, we take some simple damped system to demonstrate.

(Refer Slide Time: 05:23)



So we have this system. So here is this mass, here is a damper and there is a second mass and this is the rigid. So there is no friction between the mass and the surface and this is k and here is again one spring, so this is also k and this is c , this is m and so here it is x_1 and this is x_2 . So this is the system and now we have to write the equations of motion for this simple system. So this is the first free body diagram of this mass m first mass.

So we will have here this force kx_1 or if we pull here, it will be kx_1 and then this is the damper and if we assume $x_2 > x_1$ so this damper force will pull it $c \dot{x}_2 - \dot{x}_1$ and this is x_1 . Now here the second mass, we will have the same force here $c \dot{x}_2 - \dot{x}_1$ opposite direction and this is spring because when it will compress, this will apply a force k times x_2 and here it is x_2 .

So we can write this equation of motion for this system so here $m\ddot{x}_1 = -kx_1 + c\dot{x}_2 - \dot{x}_1$. So here $m\ddot{x}_1 + c\dot{x}_1 + kx_1 = c\dot{x}_2$ so this is equation one. Now for the second mass, we will have this so we have to write the equation of motion $M\ddot{x}_2 = -kx_2 - c\dot{x}_2 + \dot{x}_1$ so $M\ddot{x}_2 + c\dot{x}_2 + kx_2 = \dot{x}_1$.

So because this is in this direction and they are opposite forces so they will be having - sign. So here $M\ddot{x}_2 + c\dot{x}_2 + kx_2 = c\dot{x}_1$. So this is equation number two. Now as I said that we have to assume the motion $x_1 = A_1 e^{st}$ and $x_2 = A_2 e^{st}$. So here why we are taking such a form because s is a complex, it could be complex number and due to the introduction of damping the motion, the roots could be complex.

So that is why we take this kind of solution and so now we put this here. So what is x_1 dot is s into $A_1 e^{st}$ and equal to $s x_1$ and here x_2 dot = $s A_2 e^{st}$ so that is equal to $s x_2$. Similarly, x_1 double dot = $s^2 A_1 e^{st}$ so that is $s^2 x_1$ and x_2 double dot = $s^2 A_2 e^{st}$, so that is $s^2 x_2$ and x_2 double dot = $s^2 A_2 e^{st}$, so that is $s^2 x_2$.

(Refer Slide Time: 11:36)

$$\begin{aligned}
 m s^2 x_1 + c s x_1 + k x_1 &= c s x_2 \\
 (m s^2 + c s + k) A_1 e^{st} &= c s A_2 e^{st} \\
 [(m s^2 + c s + k) A_1 - c s A_2] e^{st} &= 0 \\
 (m s^2 + c s + k) A_1 - c s A_2 &= 0 \\
 \frac{A_1}{A_2} &= \frac{c s}{m s^2 + c s + k} \quad \text{--- (iii)}
 \end{aligned}$$

$$\begin{aligned}
 m s^2 x_2 + c s x_2 + k x_2 &= c s x_1 \\
 (m s^2 + c s + k) x_2 &= c s x_1 \\
 (m s^2 + c s + k) A_2 e^{st} &= c s A_1 e^{st} \\
 [(m s^2 + c s + k) A_2 - c s A_1] e^{st} &= 0 \\
 (m s^2 + c s + k) A_2 - c s A_1 &= 0 \\
 \frac{A_1}{A_2} &= \frac{m s^2 + c s + k}{c s} \quad \text{--- (iv)}
 \end{aligned}$$

So now we put these values in equation one and two so what do we get, so $m x_1$ double dot is s^2 and $m s^2$ into $x_1 + c s x_1 + k x_1 = c s x_2$. Now $m s^2 + c s + k x_1$ and x_1 is $A_1 e^{st}$ and equal to $c s$ and x_2 is $A_2 e^{st}$. So of course we can have because e^{st} if we do like $m s^2 + c s + k A_1 - c s A_2 e^{st} = 0$.

Of course e^{st} exponential it cannot be 0, so we have to write $m s^2 + c s + k A_1 - c s A_2 = 0$ and so we can write here A_1 by $A_2 = c s$ by $m s^2 + c s + k$ and this is equation number three. So here we have obtained the amplitude ratio as a function of s . Now we put these values in equation number two, so we get $m s^2 x_2 + c s x_2 + k x_2 = c s x_1$. So x_2 double dot is $s^2 x_2 + c s x_2 + k x_2 = c s x_1$.

So here $m s^2 + c s + k$ and $x_2 = c s x_1$ so here $m s^2 + c s + k$. Now x_2 is $A_2 e^{st}$, so $A_2 e^{st} = c s$, x_1 is $A_1 e^{st}$. So again here $m s^2 + c s + k A_2 - c s A_1 e^{st} = 0$. Because e^{st} cannot be 0, we have to write $m s^2 + c s + k A_2 - c s A_1 = 0$. So we can get A_1 by $A_2 = m s^2 + c s + k$ by $c s$. So we get this amplitude ratio A_1 by A_2 .

(Refer Slide Time: 16:20)

$$\begin{aligned}
& \frac{cs}{ms^2 + cs + k} \xrightarrow{\text{cross}} \frac{ms^2 + cs + k}{cs} \\
& (ms^2 + cs + k)^2 = (cs)^2 \\
& (ms^2 + cs + k)^2 - (cs)^2 = 0 \\
& (ms^2 + cs + k + cs)(ms^2 + cs + k - cs) = 0 \\
& (ms^2 + 2cs + k)(ms^2 + k) = 0 \\
& ms^2 + k = 0 \\
& ms^2 + 2cs + k = 0 \\
& s = \frac{-k}{m} = j^2 \frac{k}{m} \\
& s_{1,2} = \pm j \sqrt{\frac{k}{m}} \\
& \omega_n = \sqrt{\frac{k}{m}} \\
& s_{1,2} = \pm j \omega_n
\end{aligned}$$

Now what we can do, we can equate these two equations three and four because both are the ratios of A1 by A2 and so therefore we find we do like cs by ms square + cs + k = ms square + cs + k upon cs. Now we cross multiply and so we get ms square + cs + k square because these are the same terms so they will be square equal to cs square. So we can write ms square + cs + k square - cs square = 0.

Now we can because a square + b square = a + b into a - b. So we can use this formula so ms square + cs + k + cs and second factor is ms square + cs + k - cs and this is = 0. So what we get ms square + 2cs + k and here cs will be cancel out so we will have ms square + k = 0. So it means that ms square + k = 0 and ms square + 2cs + k = 0.

So now from these two equations, we can get the values of s. So when ms square + k = 0 so s square = -k by m and that is - complex number j square because - is = j square. So s = j root k by m + - and root k by m is if we say that omega n = root k by m, so we can write s1, 2 = + - j omega n.

(Refer Slide Time: 19:54)

$$\begin{aligned}
 ms^2 + 2cs + k &= 0 \\
 s_{3,4} &= \frac{-2c \pm \sqrt{4c^2 - 4mk}}{2m} \\
 &= -\frac{c}{m} \pm \sqrt{\left(\frac{c}{m}\right)^2 - \frac{k}{m}} \\
 &= -\frac{c}{m} \pm \sqrt{\left(\frac{c}{m}\right)^2 - \frac{k}{m}} \\
 \frac{c}{m} &= 2\zeta\omega_n \\
 s_{3,4} &= -2\zeta\omega_n \pm \sqrt{(2\zeta\omega_n)^2 - \omega_n^2} \\
 &= \left[-2\zeta \pm \sqrt{(2\zeta)^2 - 1} \right] \omega_n \\
 &= \left[-2\zeta \pm j\sqrt{1 - (2\zeta)^2} \right] \omega_n \quad \left(2\zeta < 1 \right. \\
 &\quad \left. \zeta < \frac{1}{2} \right)
 \end{aligned}$$

Now what about this equation? So if we write $ms^2 + 2cs + k = 0$ so we have $s = -b$ so $-b$ is $-2c \pm \sqrt{b^2 - 4mk}$ so b^2 is $4c^2 - 4mk$ so 2 into A is m and c is k by $2A$ so 2 into A is m . And so we will have this is third and fourth root of s where we got here two roots of course this is negative $-j$ root k by m and this is $+j$ root k by m and here 3 by 4.

So this we can write $-c$ by $m \pm \sqrt{c^2 - km}$ so we take m inside so we will get c by m whole square $-$. So here $2mk$ by $4m^2$ so $2km$ by $4m^2$. So here it is 2 and m is cancel out, $-4ac$ so it is $b^2 - 4ac$ so it is 4 into a into c . so it is 4 here it is 4 , so 4 4 cancel out, we will get k by m . So $-c$ by $m \pm \sqrt{c^2 - km}$, c by m whole square $- k$ by m .

Now c by $m = 2\zeta\omega_n$ so we have $s_{3,4} = -2\zeta\omega_n \pm \sqrt{(2\zeta\omega_n)^2 - \omega_n^2}$ and k by m is ω_n^2 . So we will get $-2\zeta \pm \sqrt{(2\zeta)^2 - 1}$ and ω_n . Now we assume that this part is such that $1 - (2\zeta)^2$ is positive. So 2ζ , it is less than 1 so it means ζ is less than half so damping factor is less than half so we are taking basically the under damp system.

So here we will have $-2\zeta \pm j\sqrt{1 - (2\zeta)^2}$ so we take it outside so it is root $-j$ under root $1 - (2\zeta)^2$ whole square and this is ω_n . Okay so here we have obtained the roots, now what will be the solution of this. So here we see that we have the natural frequencies here so they are because they are the complex roots they are in pair.

(Refer Slide Time: 24:27)

$$\begin{aligned}
 x_1 &= A_{11} \sin(\omega_{n1}t + \phi_1) + A_{12} e^{-2\zeta\omega_{n1}t} \sin(\sqrt{1-4\zeta^2}\omega_{n1}t + \phi_2) \\
 x_2 &= A_{21} \sin(\omega_{n1}t + \phi_3) + A_{22} e^{-2\zeta\omega_{n1}t} \sin(\sqrt{1-4\zeta^2}\omega_{n1}t + \phi_4)
 \end{aligned}$$

$A_{11}, \phi_1, A_{12}, \phi_2 -$
 $A_{21}, \phi_3, A_{22}, \phi_4 -$

And so we will have a solution like $x_1 = A_{11} \sin \omega_{n1} t + \phi_1$ because this root will give us a solution like this, harmonic one and the other one due to the second frequency we will get the solution like $A_{12} \text{ exponential} - 2 \zeta \omega_{n1} t$ into $\sin \sqrt{1 - 2 \zeta^2} \omega_{n1} t + \phi_2$. Okay so this you can understand from because this is complex root.

So we will have some exponential free decay vibration as we discussed to see the damped response of the single degree of freedom system. So this is similar, but only here the damping is not zeta, but it is 2 zeta so it is $1 - 2 \zeta^2$ is $4 \zeta^2$ and here again it is 2 zeta. Similarly, x_2 will be $A_{21} \sin \omega_{n1} t + \phi_3 + A_{22} \text{ exponential}$.

So here it is $\omega_{n1} - 2 \zeta \omega_{n1} t$ into $\sin \sqrt{1 - 4 \zeta^2} \omega_{n1} t + \phi_4$. So here we can see these are the free vibration response of the system and the values of this constants that is $A_{11} \phi_1$; $A_{12} \phi_2$ or $A_{21} \phi_3$; $A_{22} \phi_4$. They will all depends on the initial conditions.

Now if the system is given an initial condition equal disturbance so if initially we displace what the system with equal amplitude like 5 mm both then the system will vibrate purely with the first natural frequency because of course from here we get the mode shape corresponding to the first natural frequency and mode shape corresponding to second frequency. So when if the equal initial disturbance to both the masses, the system will vibrate with the first natural frequency.

And so only the first term that is $\sin \omega_n t + \phi_1$ will be there for both the equations. However, if we are giving completely opposite initial conditions, then the system will vibrate with the damped free vibration. So the second part will be there. Because if we put here this s_{12} , we will get A_1 by $A_2 = +1$ and when we put these values, we will get A_1 by $A_2 = -1$.

So the system will be in case of ω_n it will be vibrating as a damped free vibration. Okay so that is how we can study a damped vibration of two degree of freedom system. So in case of first natural frequency, damper is ineffective so there is no role of damper because both the system they are vibrating as pure sin wave okay, but in other case damper is effective with a damping factor of 2ζ okay. So I thank you for this lecture and see you in the next lecture.