

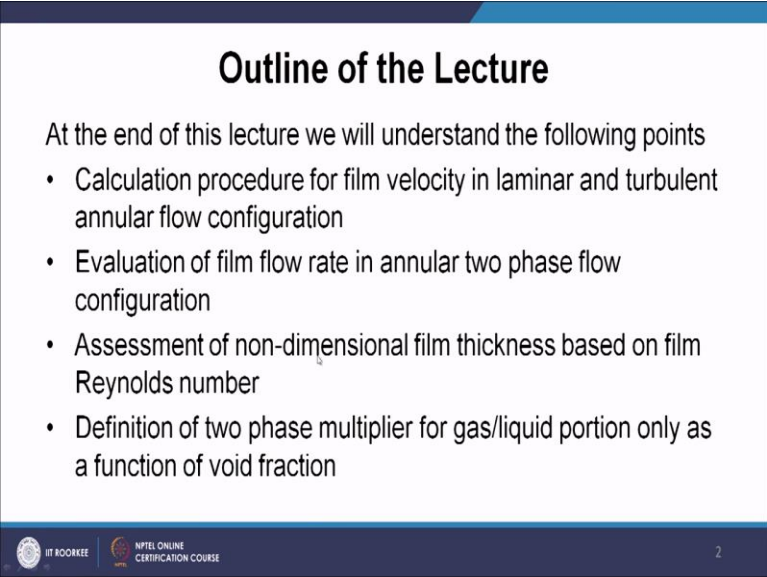
**Two phase flow and heat transfer**  
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**Lecture No.08**  
**Annular Flow**

Hello, welcome in the eighth lecture of the course Two Phase Flow and Heat Transfer. In this course you will be understanding annular flow models. Annular flow model is a typical flow pattern in gas liquid Two Phase Flow where we will be finding out liquid film is added with the pipe and in the core of the tube you will be finding out gaseous phase.

So in this lecture we will be stressing about the calculation of film velocity in both laminar and turbulent regime. We will be evaluating film flow rate in the Two Phase Flow configuration, annular Two Phase Flow configuration.

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**Outline of the Lecture**

At the end of this lecture we will understand the following points

- Calculation procedure for film velocity in laminar and turbulent annular flow configuration
- Evaluation of film flow rate in annular two phase flow configuration
- Assessment of non-dimensional film thickness based on film Reynolds number
- Definition of two phase multiplier for gas/liquid portion only as a function of void fraction

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We will be assessing non-dimensional film thickness based on film Reynolds number. As well as we will be defining 2 phase multiplier for gas or liquid portion only as a function of void fraction. So first let us see that how annular flow configuration can be shown here in pictorial view. Here I have shown you a pipe inside which annular flow is occurring. You see inside this, we are having first film added with wall and inside the core we are having gaseous phase okay.

Now here, you can see in the film we can consider a small element like this where the small element will be getting shear stress  $\tau$  from the adjacent liquid layers.

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### Annular Flow

For film flow:  $\tau = \tau_i \frac{r_i}{r} - \frac{1}{2} \left( \rho_f g + \frac{dP}{dz} \right) \left( \frac{r^2 - r_i^2}{r} \right)$

If film flow is laminar:  $\tau = \mu_f \frac{du_f}{dy}$

Here,  $y = d - r$ ,  $dy = dr$

Hence,  $u_f = \frac{1}{\mu_f} \left\{ \left[ \tau_i r_i + \frac{1}{2} \left( \frac{dP}{dz} + \rho_f g \right) r_i^2 \right] \ln \frac{D}{2r} - \frac{1}{4} \left( \rho_f g + \frac{dP}{dz} \right) \left( \frac{D^2}{4} - r^2 \right) \right\}$

For turbulent film flow:  $\tau = (\mu_f + \mu_t) \frac{du_f}{dy}$

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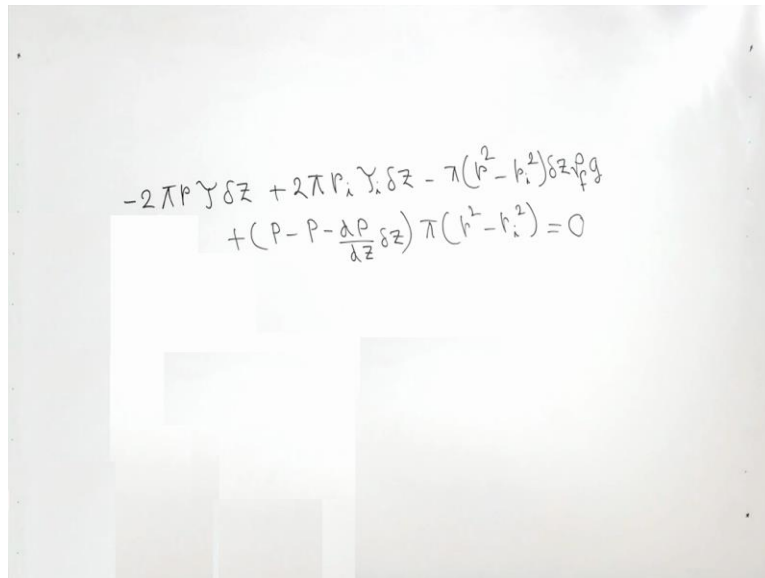
And in the interface it will be getting the shear stress  $\tau_i$  in the opposite direction. Now near the wall this shear stress  $\tau$  is being converted to  $\tau_w$  which is nothing but the wall shear stress okay. And if you think about this fluid element the fluid element is having pressure from the bottom as  $P$  and from the top it is  $P+dP$  okay.

This element length we have considered over here  $\delta z$ . Now if you try to have a force balance then you will be finding out that we are having over here the shear stress at the interface. The shear stress at the outer layer of this fluid element as well as we are having some weight of the fluid element acting in the downward direction.

Apart from that you will be having also the pressure forces  $P+dP$  and  $P$  over here okay. So if you write down the force balance equation over here, you can see the first term whatever I have written this is nothing but  $2\pi r \delta z$  that is the area whatever through which this  $\tau$  or shear stress from the other liquid layer is being acted okay. So this is in the downward direction because  $\tau$  is in the downward direction okay.

In a similar fashion here we have written the shear stress, what it is getting the liquid layer is getting from the gaseous interface which is in the upward direction.

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$$-2\pi r \tau \delta z + 2\pi r_i \tau_i \delta z - \pi(r^2 - r_i^2) \delta z \rho_f g + (P - P - \frac{dP}{dz} \delta z) \pi(r^2 - r_i^2) = 0$$

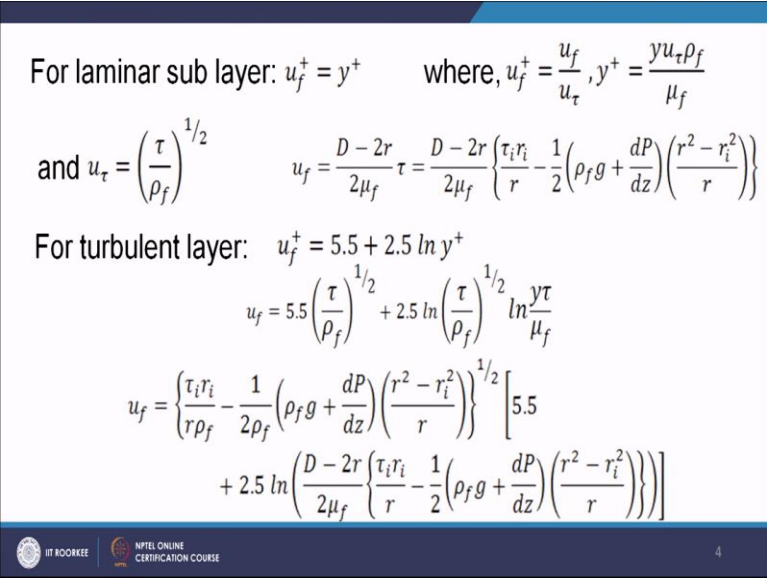
So that is why here you are having a +sign the area will be once again  $2\pi r \cdot dz$  and the magnitude of shears interfacial shears stress is nothing but  $\tau_i$ . So these 2 are coming from the shear stresses and this is actually your interfacial shear stress. Now as I have told that we are having over here some sort of weight of the fluid element. So the weight can be found out by volume \*density. So for the volume  $\pi \cdot r^2 - r_i^2 \cdot dz$ .

So this is the volume multiply by density of the liquid  $\rho_f$  and the  $g$  which is acting in the downward direction. And from the pressure force you see here we have pressure  $P$  from the downward side and from the upward side  $P+dP$ . So this will be  $P - (P+dP)$ . Now this  $dP$ , if you consider for unit length  $dz$ , so we can write down  $P - dP / dz \cdot dz$  so this  $dz$  cancels out. But for facilitating this calculation of  $dP/dz$ , I have written like this.

This pressure is acting on the area  $\pi \cdot r^2 - r_i^2$  which is the bottom area or you can say the top area for the liquid film. So from here if you simplify this one and try to calculate the value of the  $\tau$  then you will be getting the  $\tau$ . As  $\tau = \tau_i \cdot r_i / r - (1/2 \cdot \rho_f \cdot g + dP/dz)$ ,  $dP / dz$  comes from here multiplied by  $r^2 - r_i^2 / r$ .

A simple calculation from these equations, if you try to figure out what is the value of this  $\tau$ , you will be getting in this equation after cancellations. Now we are having different types of a film flows. So for film flow that I have showed this is a equation. Incase of laminar film flow, we can write down  $\tau = \mu_f \cdot du_f / dy$  okay. So this further can be calculated by some transformation, the transformations are like this okay.

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For laminar sub layer:  $u_f^+ = y^+$       where,  $u_f^+ = \frac{u_f}{u_\tau}$ ,  $y^+ = \frac{y u_\tau \rho_f}{\mu_f}$

and  $u_\tau = \left( \frac{\tau}{\rho_f} \right)^{1/2}$        $u_f = \frac{D-2r}{2\mu_f} \tau = \frac{D-2r}{2\mu_f} \left\{ \frac{\tau_i r_i}{r} - \frac{1}{2} \left( \rho_f g + \frac{dP}{dz} \right) \left( \frac{r^2 - r_i^2}{r} \right) \right\}$

For turbulent layer:  $u_f^+ = 5.5 + 2.5 \ln y^+$

$u_f = 5.5 \left( \frac{\tau}{\rho_f} \right)^{1/2} + 2.5 \ln \left( \frac{\tau}{\rho_f} \right)^{1/2} \ln \frac{y \tau}{\mu_f}$

$u_f = \left\{ \frac{\tau_i r_i}{r \rho_f} - \frac{1}{2 \rho_f} \left( \rho_f g + \frac{dP}{dz} \right) \left( \frac{r^2 - r_i^2}{r} \right) \right\}^{1/2} \left[ 5.5 + 2.5 \ln \left( \frac{D-2r}{2\mu_f} \left\{ \frac{\tau_i r_i}{r} - \frac{1}{2} \left( \rho_f g + \frac{dP}{dz} \right) \left( \frac{r^2 - r_i^2}{r} \right) \right\} \right) \right]$

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So let us try to see what happens in case of turbulent film flow. So in case of turbulent film flow, we know that there are several layers we have to consider. First one is laminar sub layer then we will be having buffer layer and we can also have turbulent layer. So I have given here the example of 2 layers only, laminar sub layer and the extreme one turbulent layer. So we know from fluid mechanics that in the laminar sub layer  $u_f^+ = y^+$  where  $u_f^+$  is nothing but  $u_f / u_\tau$  and  $u_\tau$  can be written as shear stress by density of the fluid to the power  $1/2$ .

And  $y^+$  is a characteristic length you can write down that one as  $y \cdot u_\tau \rho_f / \mu_f$  right. So what you can do just in a similar fashion if you write down that what is  $u_f^+ = y^+$  over here using this equations inside this then you will be finding out  $u_f$  can be written in the form of  $\tau$  okay. And that comes as  $u_f = (D-2r) / (2 \mu_f \cdot \tau)$  okay. Already we have expressed what  $\tau$  is over here for the film flow okay.

So we can replace the Tau over here and write down  $u_f$  for the turbulent, laminar sub layer zone  $d - 2r / 2 \mu_f \tau_i / r - (1/2 * \rho_f g + (dP/dz) * (r^2 - r_i^2 / r))$  okay. Similar thing can be also done in case of turbulent layer. For turbulent layer we know  $u_f^+ = 5.5 + 2.5 * \ln(y^+)$ . This is typically coming from fluid mechanics. So what we can do in a similar fashion as we have done in laminar, laminar sub layer, we can also perform this one for turbulent layer.

We can find out that  $u_f$  can be written as  $5.5 \tau / \rho_f$  to the power  $1/2 + 2.5 \ln(\tau / \rho_f$  to the power  $1/2) \ln(y \tau / \mu_f)$ . So once again what we can do  $\tau$ , we can replace from the previous equation what we have found out for the film flow and finally obtain the film velocity in this fashion okay.

So for both the regimes laminar and turbulent we have shown how film velocity can be obtained. This is very, very important for an annular flow regime okay. Next as we have found out the flow velocity, laminar film flow velocity, it will be very critical to find out what is the film flow rate. So what we have done over here  $w_{FF}$  which is the film flow rate we have found out by integrating 0 to  $\delta$   $\pi D dy u_f \rho_f$ .

So this is basically  $\pi D \int_0^\delta u_f \rho_f dy$  so the limit for  $y$  is from 0 to  $\delta$  okay.

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### Film flow rate (assuming $\delta$ is small)


For laminar flow:

$$w_{fF} = \int_0^\delta (\pi D dy u_f) \rho_f = \pi D \rho_f \int_0^\delta \frac{y \tau}{\mu_f} dy = \frac{\pi D \rho_f \tau \delta^2}{2 \mu_f}$$

For turbulent flow:

$$w_{fF} = \pi D \rho_f \int_0^\delta (u_f dy) = \pi D \int_0^{\delta^+} \mu_f u_f^+ dy^+$$

Where,  $\delta^+ = \frac{\delta u_\tau \rho_f}{\mu_f}$



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So you can write down this one for laminar domain that this  $u_f$  will be written as  $y\tau_w/\mu_f$  okay. And if you do this derivation of the integration from 0 to  $\delta$ , you will be getting the final expression for the laminar film flow is  $\pi D \rho_f \tau_w \delta^3 / 6 \mu_f$  where  $\tau_w$  is the shear stress we have already derived. Similar thing can be done for turbulent flow.

In case of turbulent flow it will be  $\pi D \rho_f \int_0^\delta u_f dy$  and here  $u_f$  will be replaced by  $u_f^+$  okay and  $dy$  will be replaced by  $dy^+$  at the same time your limit will be transforming from  $\delta$  to  $\delta^+$  okay. Already we have said what is the relationship between  $y$  and  $y^+$ . Here  $\delta^+$  will be  $\delta * u\tau_w * \rho_f / \mu_f$  okay. So once you know the expression for this  $u_f^+$  depending on which layer you are in, you will be finding out the integration value and you can get the film flow rate for the turbulent regime okay.

Next let us see some other extend if in the previous case we have shown the film was moving up okay. So let us say in this case we are taking you are having falling film so the film is actually falling down and the gas is moving up.

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### Falling film flow

Obtained at very low gas phase velocities

$$\frac{dP}{dz} = -\rho_g g \text{ and } \tau_i \approx 0$$

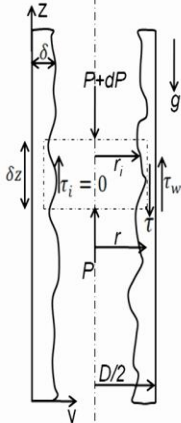
Hence,  $\tau = -\frac{1}{2}(\rho_f - \rho_g)g \left( \frac{r^2 - r_i^2}{r} \right)$



As,  $\delta \ll D$ ;  $r + r_i \approx 2r$

$$\tau = -(\rho_f - \rho_g)g(r - r_i)$$

Again,  $r - r_i \approx \delta - y$

$$\tau = -(\rho_f - \rho_g)g(\delta - y)$$





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So in this case what we can do, we can typically find out similar figure okay whatever we have shown in the last 1 only thing that this  $\tau_w$  is in the upward direction because the film is now falling down okay. Now let us try to see that what happens for the gaseous velocities. So if we

obtain if you consider that very low gas velocities then we can write down  $dP/dz = -\rho_f g$  okay. And in this case if you are having very low gas velocity, you can assume that at the interface there is no shear stress okay. So  $\tau_i$  is nearly equivalent to 0 okay.

So  $\tau_i$  is nearly equivalent to 0 over here okay. So if you apply the previous equation whatever we have derived over here for  $\tau_i$  and make this  $\tau_i = 0$ , you will be getting something like this  $\tau = -\frac{1}{2}(\rho_f - \rho_g)g(r^2 - r_i^2)/r$  okay. Now let us try to convert this one in easier forms. So what will be doing, we will be taking assumption that film thickness is very small compare to the pipe diameter.

So  $\delta$  is less than  $d$ . Under this assumption, we can write down  $r + r_i$ . So  $r$  was the arbitrary cross sectional radius and  $r_i$  was the interfacial radius that will be more or less equivalent to  $2r$  because  $\delta$  is very small. So under this we can write down if you see over here  $(r^2 - r_i^2)$  so that will be  $(r + r_i)(r - r_i)$ . So  $r + r_i$  we can write down as  $2r$  so that  $r$  and this  $r$  will be canceling out. Finally, we will be getting  $-\frac{1}{2}(\rho_f - \rho_g)g(r - r_i)$  okay.

Once again we can write down  $(r - r_i)$  is nearly equivalent to  $(\delta - y)$  remember  $y$  is being calculated from the wall so  $(r - r_i)$  will be  $(\delta - y)$ . So what we can do  $(r - r_i)$  we can replace in terms of  $(\delta - y)$  and final expression for the falling film will be coming as  $\tau = -(\rho_f - \rho_g)g(\delta - y)$  right. So this is the shear stress expression for falling film flow. Next let us try to understand if this is laminar flow then what happens.

In a typical fashion, if you replace this  $\tau$  as  $\mu \frac{du}{dy}$  and once again integrate with respect to  $y$  then will be getting  $u = -\frac{g}{\mu}(\rho_f - \rho_g)\frac{(\delta - y)^2}{2}$  okay.

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$$\text{For laminar flow: } u_f = -\frac{g}{\mu_f}(\rho_f - \rho_g)\left(\delta y - \frac{y^2}{2}\right)$$

$$\text{Mass flow rate per unit width: } \Gamma = -\rho_f \int_0^\delta (u_f dy) = \frac{\rho_f g (\rho_f - \rho_g) \delta^3}{\mu_f 3}$$

$$\text{Film Reynolds number: } Re_\Gamma = \frac{4\delta u_f \rho_f}{\mu_f} = \frac{4\Gamma}{\mu_f}$$

$$\text{Dimensionless film thickness: } \delta^* = \frac{\delta g^{1/3} (\rho_f - \rho_g)^{1/3} \rho_f^{1/3}}{\mu_f^{2/3}}$$

$$\text{For laminar flow: } \delta^* = \left(\frac{3}{4} Re_\Gamma\right)^{1/3} = 0.909 Re_\Gamma^{1/3}$$

$$\text{For turbulent flow: } \delta^* = 0.115 Re_\Gamma^{0.6}$$

So if you go further and try to find out the the mass flow rate per unit width, so gamma can be written as -  $\rho_f \int_0^\delta u_f dy$  okay. So this is the mass flow rate per unit width. So you can write down this  $u_f$  as this function and once again integrate from 0 to delta. You will be finding out the final expression as  $\rho_f g (\rho_f - \rho_g) / \mu_f \delta^3 / 3$ . So you see the mass flow rate per unit width is actually dependent on the delta cube okay.

Going in this fashion what we can do, we can define a typical non-dimensional number which is called film Reynolds number. So here we are defining film Reynolds number as  $Re_\Gamma$  okay. So  $Re_\Gamma$  will be dependent on the film thickness. So the length scale will be typically the film thickness. So what we have done, we have define the film Reynolds number as  $4 \delta u_f \rho_f / \mu_f$  okay.

So once you get the expression for this mass flow rate per unit width and you know  $u_f$  already we have seen over here so if we plug both this  $u_f$  and gamma over here. We can write down this film Reynolds number as a function of gamma and  $u_f$  only okay. So we find out film Reynolds number will be nothing but  $4 \Gamma / \mu_f$  where gamma, is the mass flow rate per unit width okay.

Then we try to find out what is the dimensionless film thickness and non-dimension film thickness. So that will be delta star. We can write down delta star as  $\delta^* = \delta g^{1/3} (\rho_f - \rho_g)^{1/3} \rho_f^{1/3} / \mu_f^{2/3}$

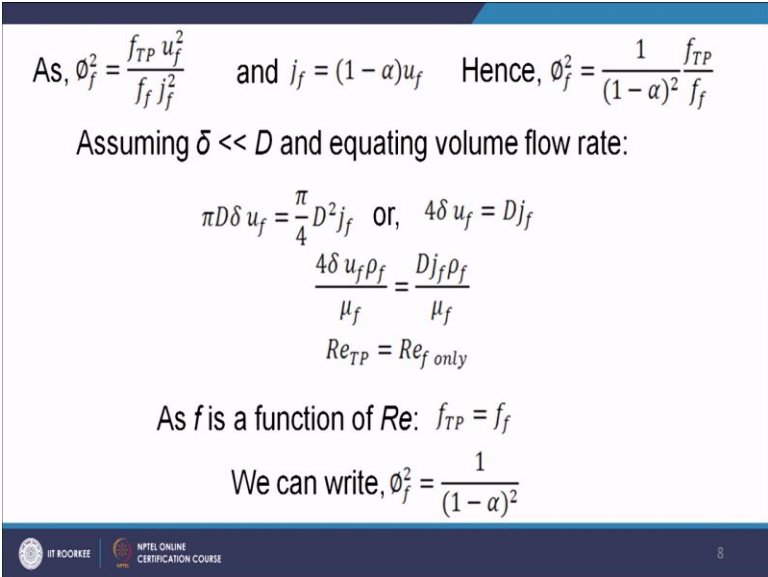


( $\rho_f - \rho_g$ ) to the power  $1/3$  \*  $\rho_f$  to the power  $1/3$  /  $\mu_f$  to the power  $2/3$  okay. Okay with this definition what we can find out, we will be getting relationship between these  $\Delta$  star and  $Re_{\gamma}$  which is nothing but film Reynolds number and little bit of derivation between these 2 can be proving you that  $\Delta$  star is becoming  $3/4 Re_{\gamma}$  to the power  $1/3$ .

Please try to practice this one; this is not very difficult task only replacement of 1 term in the another one will be giving you the correct expression. So we can find out the  $\Delta$  star for laminar film flow comes out to be  $3/4 Re_{\gamma}$  to the power  $1/3$ . So if you do  $3/4$  to the power  $1/3$  so  $\Delta$  star typically becomes  $.909 Re_{\gamma}$  to the power  $1/3$ .

Follow the same procedure for turbulent layer. We will be getting some another expression. The expression comes out to be  $\Delta$  star =  $.115 Re_{\gamma}$  to the power  $.6$  okay. Please practice this turbulent one. It you will be getting the same expression.

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As,  $\phi_f^2 = \frac{f_{TP} u_f^2}{f_f j_f^2}$  and  $j_f = (1 - \alpha)u_f$  Hence,  $\phi_f^2 = \frac{1}{(1 - \alpha)^2} \frac{f_{TP}}{f_f}$

Assuming  $\delta \ll D$  and equating volume flow rate:

$$\pi D \delta u_f = \frac{\pi}{4} D^2 j_f \quad \text{or,} \quad 4 \delta u_f = D j_f$$

$$\frac{4 \delta u_f \rho_f}{\mu_f} = \frac{D j_f \rho_f}{\mu_f}$$

$$Re_{TP} = Re_{f \text{ only}}$$

As  $f$  is a function of  $Re$ :  $f_{TP} = f_f$

We can write,  $\phi_f^2 = \frac{1}{(1 - \alpha)^2}$

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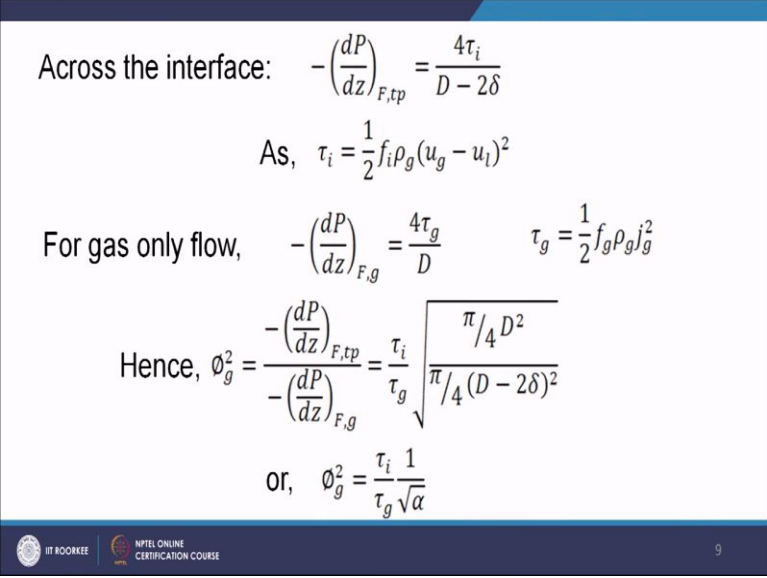
Okay, next let us try to find out what happens for the friction factors. So if you see Two Phase multiplier for fluid portion only so here.  $\phi_f^2$  you know that  $f_{TP}$  Two Phase friction factor by  $f_f$  fluid portion, friction factor \*  $u_f^2 / j_f^2$ , this we have defined in the drift flux chapter and then we know that  $j_f$  can be written as  $(1 - \alpha) u_f$ .

So I can replace this  $j_f$  and  $u_f$  in terms of  $(1 - \alpha)$ . So here you will be getting  $\phi_f^2$  becomes  $(1 - \alpha)$  whole square using this equation over here. So we will be getting  $\phi_f^2$  is  $(1 - \alpha)$  whole square  $f_{TP}/f_f$  okay. Now if you assume that  $\delta$  is very thin, so the film thickness is very small compared to the pipe diameter then I can write down that the volume for the film is  $\pi D \delta u_f$ .

On the other hand if you try to compare the superficial velocity of the film the volume comes out to be  $\pi/4 D^2 j_f$ . So if you equate this 2 then we will be getting a relationship between  $u_f$  and  $j_f$  in this fashion  $4 \delta u_f = D j_f$  right. Little bit of multiplication and division and it will be giving you  $4 \delta u_f \rho_f / \mu_f = D j_f \rho_f / \mu_f$ . Now these are nothing but Reynolds number expression. This is for Two Phase and this is for fluid portion only.

So you find out Reynolds number for Two Phase and Reynolds number for fluid portion only they are same. Assuming  $\delta$  is very small than  $d$  okay. So if Reynolds number is the same obviously friction factors will be the same, so you can write down  $\phi_f^2 = 1/(1 - \alpha)$  whole square. From this equation if we cancel both these terms considering the  $\delta$  less than  $d$  okay.

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Across the interface:  $-\left(\frac{dP}{dz}\right)_{F,TP} = \frac{4\tau_i}{D - 2\delta}$

As,  $\tau_i = \frac{1}{2} f_i \rho_g (u_g - u_l)^2$

For gas only flow,  $-\left(\frac{dP}{dz}\right)_{F,g} = \frac{4\tau_g}{D}$        $\tau_g = \frac{1}{2} f_g \rho_g j_g^2$

Hence,  $\phi_g^2 = \frac{-\left(\frac{dP}{dz}\right)_{F,TP}}{-\left(\frac{dP}{dz}\right)_{F,g}} = \frac{\tau_i}{\tau_g} \sqrt{\frac{\pi/4 D^2}{\pi/4 (D - 2\delta)^2}}$

or,  $\phi_g^2 = \frac{\tau_i}{\tau_g} \frac{1}{\sqrt{\alpha}}$

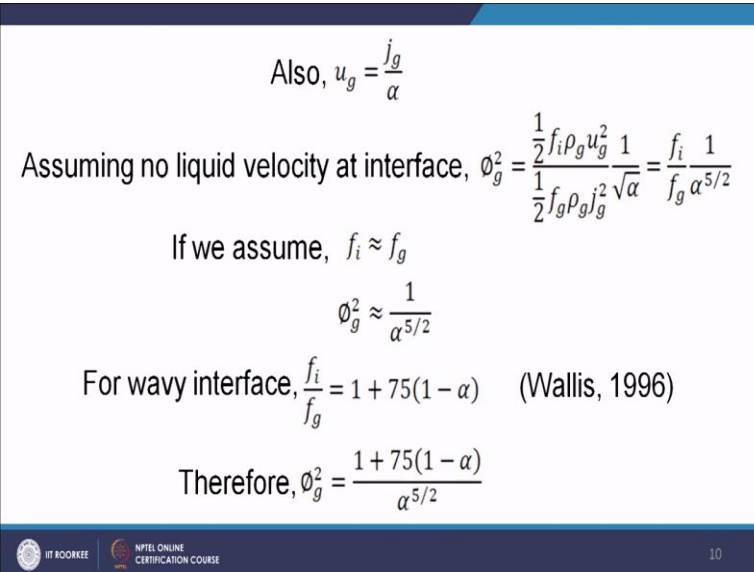
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Now let us try to see the gas portion only. So for the Two Phase you can write down across the interface  $-dP/dz$  friction factor  $= 4 \tau_i / (d - 2 \delta)$  okay. And  $\tau_i$  can be written as this. We

have already explained earlier  $\frac{1}{2} f_i \rho_g u_g^2$  and for gas portion only  $dP/dz$  friction factor for gas is nothing but  $4 \tau_{aug}/D$  and  $\tau_{aug}$  will be  $\frac{1}{2}(f_g \rho_g u_g^2)$ . Now if we plug all these things that means this  $dP/dz$  Two Phase friction factor  $dP/dz$  is gas only fraction factor, gas portion friction factor in the Two Phase multiplier  $\phi_g^2$  then we will be finding out we are getting  $\tau_{ai} / \tau_{aug}$  root over of  $\pi/4 D^2 / \pi/4 (D - 2\delta)^2$ .

Now this gives me  $1/\sqrt{\alpha}$  and  $\tau_{ai} / \tau_{aug}$  comes over here. So we find out that Two Phase gas portion only friction factor is multiplier is actually  $\tau_{ai} / \tau_{aug} * 1/\sqrt{\alpha}$ .

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Also,  $u_g = \frac{j_g}{\alpha}$

Assuming no liquid velocity at interface,  $\phi_g^2 = \frac{\frac{1}{2} f_i \rho_g u_g^2}{\frac{1}{2} f_g \rho_g u_g^2} \frac{1}{\sqrt{\alpha}} = \frac{f_i}{f_g} \frac{1}{\alpha^{5/2}}$

If we assume,  $f_i \approx f_g$

$$\phi_g^2 \approx \frac{1}{\alpha^{5/2}}$$

For wavy interface,  $\frac{f_i}{f_g} = 1 + 75(1 - \alpha)$  (Wallis, 1996)

$$\text{Therefore, } \phi_g^2 = \frac{1 + 75(1 - \alpha)}{\alpha^{5/2}}$$

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Now let us try to find out how root alpha can be calculated. So if you see if we give this  $u_g = j_g / \alpha$ . Then we can write down this  $\tau_{ai} / \tau_{aj}$  from here. Okay over here like this considering no liquid velocity. So in this expression you see here you are having  $u_l$ . So if you cancel this one then you are getting  $\frac{1}{2} f_i \rho_g u_g^2$  and this one is  $\frac{1}{2} f_g \rho_g u_g^2$ .


So you will be finding out relationship between  $u_g$  and  $j_g$  in this fashion. So if we put all this  $\tau_{ai}$  and  $\tau_{aug}$  in this expression okay with assumption of no liquid velocity, we will be finally leading to  $f_i / f_g * 1/\alpha$  to the power  $5/2$ . Now there are several options for this ratio of the friction factors. So if we first assume that both the friction factors are same interfacial friction factor and gas only friction factor then we will be getting that gas portion multiplier is nothing but  $1/\alpha$  to the power  $5/2$  okay.


Now Wallis, he has given another opinion in case of this equal friction factor he has given that at for the wavy interface  $f_i / f_g = 1 + 75(1 - \alpha)$ . If you hook into this  $f_i / f_g$  in this equation then you will be getting  $\phi_g^2 = 1 + 75(1 - \alpha) / \alpha$  to the power of  $5/2$  right.

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### Summary

- In this lecture we have derived the expression for shear stress inside liquid film in annular flow
- We have formulated film velocity for both laminar and turbulent situations
- Using falling film analysis expressed non-dimensional film thickness as a function of film Reynolds number
- At last we have calculated frictional pressure drop by knowing information of average void fraction in annular flow

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So let us summarize this lecture. So in this lecture we have derived the expression for shear stress inside a liquid film in the annular flow okay. We have formulated the film velocity for both laminar and turbulent regime okay.

Using falling film analysis we have expressed the non-dimensional film thickness as a function of the film Reynolds number and at the end we have calculated the friction factors okay or Two Phase multipliers for gas portion and liquid portion only as a function of your void fraction okay for consideration of the annular flow. Right, so these we have summarized in this lecture let us find out some questions to test our understanding.

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## Test your understanding ?

1. Mention assumption for falling film theory
  - a.  $\tau_i \approx 0$
  - b.  $\delta \ll D$
  - c.  $\rho_f \gg \rho_g$
  - d.  $\delta \approx 0$
2.  $\phi_g^2 \approx 1/\alpha^{5/2}$  is valid for
  - a. Low gas velocity
  - b. Low relative velocity
  - c. Zero liquid velocity
  - d. Downward liquid velocity
3. For  $\delta \ll D$  mention the correct relationship
  - a.  $Re_{TP} = Re_f \text{ only}$
  - b.  $Re_{TP} = Re_g \text{ only}$
  - c.  $Re_{TP} = Re_f$
  - d.  $Re_{TP} = Re_g$



So we are having 3 questions over here first one. Mention assumption for falling film theory. So we are having 4 options over here  $\tau_i$  nearly equals to 0,  $\delta \ll D$ ,  $\rho_f \gg \rho_g$  and finally  $\delta \approx 0$ . So I think all of you got which one is the correct answer. So you can get first 1 and second 1  $\tau_i$  nearly equals to 0 and  $\delta \ll D$  is actually the correct answer okay. Then this we have already discussed during the derivation in this lecture okay.

Then second question  $\phi_g^2 \approx 1/\alpha^{5/2}$  is nearly equal to  $1/\alpha^{5/2}$  is valid for. 4 options you are having. Low gas velocity, low relative velocity, zero liquid velocity and finally downward liquid velocity. Answer is obvious for zero liquid velocity. So without using zero liquid velocity ( $u_g - u_l$ ) whole square will not be converting into  $u_g$  square. You cannot cancel this things okay.

Then the third question for  $\delta \ll D$ , mention the correct relationships we are having Reynolds number for Two Phase = Reynolds number for fluid portion only, Reynolds number for Two Phase = Reynolds number for gas phase only, Reynolds number for Two Phase = Reynolds number of fluid portion, Reynolds number of Two Phase = Reynolds number of gas portion okay.

Obviously the answer is part a Reynolds number of Two Phase = Reynolds number of fluid portion okay. So with this we will be concluding this lecture. Thank you.