

# Fundamentals of Combustion

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## Lecture - 55

### Droplet evaporation and combustion - Part 2

#### Equilibrium under steady evaporation of liquid droplet and droplet combustion

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**Equilibrium under Steady Evaporation**


Equating equations (A) and (B):

$$\dot{m} = \frac{4\pi\lambda r_s}{c_{pg}} \ln(1 + B_T) = 4\pi r_s \rho D_{FA} \ln(1 + B_Y)$$

This implies that the transfer numbers based on mass transfer and heat transfer may be related as,

$$(1 + B_Y) = (1 + B_T)^{Le}$$

Here,  $Le = \lambda / (\rho D_{FA} c_{pg})$ . When the molecular weight of the fuel is almost the same as that of the ambient fluid, the value of Le is close to unity. If the molecular weight of the vapor is higher, then Le is greater than unity.



So, that is what we are trying to do now equate the equations A and B, which is the mass evaporation rate is nothing but so from the heat transfer  $4\pi r_s \lambda / c_{pg} \ln(1 + B_T)$ . Here  $4\pi r_s \rho D_{FA} \ln(1 + B_Y)$ . When you equate this, you can see that there, there is no assumption made till now and the main assumption made is the properties are constant.

$$\dot{m} = \frac{4\pi\lambda r_s}{c_{pg}} \ln(1 + B_T) = 4\pi r_s \rho D_{FA} \ln(1 + B_Y)$$

We have not made any other assumption like Lewis number = 1 or anything till now. Now, when you equate this, you can see that the  $4\pi r_s$  cancels and we can write. So, bring this to the left-hand side. So,  $\lambda / c_{pg}$ , So, here  $c_{pg} \rho D_{FA}$ .

So, this parameter taken to the left side, this is actually equal to the Lewis number. So, we can write this equation as  $(1 + B_Y) = (1 + B_T)^{Le}$  which is nothing but  $\lambda / c_{pg} \rho D_{FA}$ .

So, if you assume Lewis number = 1, if Lewis number = 1, we can say  $B_T = B_Y$ ; but we have not made assumptions still. So, without making that assumption, we are saying that individually


the properties are constant. When the properties are individually taken as constant; that means, that they can be evaluated at some average temperatures. But we have not made the assumption that Lewis number is 1. So, this assumption has not been made.

So, this appears here in this term, we can see that when you want to compare the transfer numbers, we have to compare this like this  $(1 + B_Y) = (1 + B_T)^{Le}$ .

So, we can calculate, Lewis number may not be 1. See when it will be 1? If the molecular weight of the fuel is almost same as that of the ambient fluid. There are only binary fluids here. If that is the case, then Lewis number will be close to 1 or else it will not be 1.

So, that assumption has not been made. We can keep this and try to work out. So, use this equation, you know what is  $B_Y$ .  $B_Y$  is  $Y_{FS} - Y_{F,\infty} / 1 - Y_{FS}$  and  $B_T$ , you know  $c_{pg}(T_\infty - T_s)/h_{fg}$ . So, when you substitute this, you can find  $Y_{FS}$  as a function of  $T_s$  or other way also. So, that is what we are going to do this.

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### Equilibrium under Steady Evaporation

Mass fraction of the fuel at the interface may be evaluated from:


$$Y_{FS} = 1 - \frac{1 - Y_{F\infty}}{\left(1 + \frac{c_{pg}(T_\infty - T_s)}{h_{fg}}\right)^{Le}} \quad (C)$$

$T_\infty \rightarrow p_{sat}(T_s)$   
 $T_s \rightarrow p_{sat}(T_s)/h_{fg}$   
 $x_F \rightarrow Y_{FS}$

It is clear that when  $T_\infty$  tends to infinity or  $h_{fg}$  tends to zero,  $Y_{FS}$  tends to unity. Since the mass fraction of fuel vapor at the droplet surface is a function of  $T_s$ , values of  $T_s$  and  $Y_{FS}$  may be evaluated by solving equation (C) with a thermodynamic relation such as **Clausius - Clapeyron equation**, written using  $h_{fg}$ , fuel gas constant and  $T_{boil}$ .

$$X_{FS} = \frac{p_{sat}}{p} = \exp\left(-\frac{h_{fg}}{R} \left[\frac{1}{T} - \frac{1}{T_{boil}}\right]\right) \quad (D)$$

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So, using this, mass fraction of the fuel at the interface can be written like this. So,  $Y_{FS}$  is written in terms of a  $T_s$  here. So,  $Y_{FS} = (1 - Y_{F,\infty}) / (1 + B_T)^{Le}$ . So, that is what we are trying to do. So, you can calculate Lewis number based upon the properties.

So, now, if you want to get the value of  $T_s$ , you can also write this equation by keeping  $T_s$  as unknown. So, this is the equation C, which actually couples the heat and mass transfer. So, that is what we want. So, once you know what will be the equilibrium temperature.

$$Y_{FS} = 1 - \frac{1 - Y_{F\infty}}{\left(1 + \frac{c_{pg}(T_\infty - T_s)}{h_{fg}}\right)^{Le}} \quad (C)$$

So, you know the value of  $T_\infty$ , you know the value of  $hc_{pg}$  and Lewis number, then  $h_{fg}$  at this, you can take very close to 0. So, this you can even neglect it. So, take it to 0,  $Y_{F,\infty}$ . So,  $T_s$  is an unknown here. So, what we can do is what is the equilibrium temperature for a given  $T_\infty$ , that is a particular temperature which will be  $T_s$  and that will obey this equation.


Do you understand? So, when you try to equate these two, there is a particular temperature  $T_s$  which will give a value of the mole fraction. See for example,  $T_s$  you get the saturation pressure at this  $T_s$ , from this  $T_s$ . So, once you know the  $T_s$  value, you can get the saturation pressure at the  $T_s$ . So, this will give you the mole fraction of fuel which is saturation pressure at  $T_s/p$ , total pressure  $p$ .

Now, mole fraction will lead us to mass fraction and once mass fraction is known, you get this value. So, there is a only particular value of  $Y_{FS}$ , when you have a particular  $T_s$  so that will obey this equation. So, that is what the equation is.

So, this equation C is very important and that illustrates that the balance between the heat and mass transfer and that will allow us to calculate the equilibrium surface temperature and the corresponding fuel mass fraction at the interface.

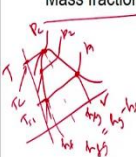
Now, we can see that if  $T_\infty \rightarrow \infty$ ; so, very large value. What happens here then? This denominator will be 0 here; so,  $Y_{FS} \rightarrow 1$ .

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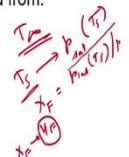


### Equilibrium under Steady Evaporation


Mass fraction of the fuel at the interface may be evaluated from:




$$Y_{FS} = 1 - \frac{1 - Y_{F,\infty}}{\left(1 + \frac{c_{pg}(T_\infty - T_s)}{h_{fg}}\right)^{Le}} \quad (C)$$



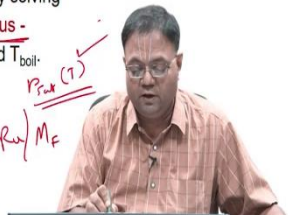
It is clear that when  $T_\infty$  tends to infinity or  $h_{fg}$  tends to zero,  $Y_{FS}$  tends to unity. Since the mass fraction of fuel vapor at the droplet surface is a function of  $T_s$ , values of  $T_s$  and  $Y_{FS}$  may be evaluated by solving equation (C) with a thermodynamic relation such as Clausius - Clapeyron equation, written using  $h_{fg}$ , fuel gas constant and  $T_{boil}$ .



$$X_{FS} = \frac{p_{sat}}{p} = \exp\left(-\frac{h_{fg}}{R}\left[\frac{1}{T_s} - \frac{1}{T_{boil}}\right]\right) \quad (D)$$



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Now, if you take  $h_{fg} \rightarrow 0$  that will happen at critical condition. So, when the pressures are high and you reach the critical temperature and critical pressure conditions, then  $h_{fg} \rightarrow 0$ . So, we know this.

So, you take any phase diagram, so this is a critical point, where the  $h_{fg}$  is 0. So, all other points, you can see based on your pressure say  $p_1$  etcetera you have these regimes here. So, this is  $p_2$  and so on and this is the critical isobar  $p_c$ , you will see that is a critical point. So, the corresponding temperature being  $T_c$ , where  $h_{fg} = 0$ . So, now, in these regimes, where the pressures are lower than the critical pressures etcetera there is a definite saturation temperature say  $T_{sat}$ .

Now, here if this is say corresponding  $h$ , this is the vapor state. So, this is  $h_g$  and this is liquid; so, this is  $h_f$ , then there is a difference in this. So,  $h_{fg}$  will be equal to  $h_g - h_f$ . Now, at the critical point, you will see that  $h_f = h_g$ ; that means,  $h_{fg} \rightarrow 0$ , when you reach the critical point and  $h_{fg} \rightarrow 0$ , then 1 plus this will be  $\infty$ . So, this denominator will be infinity. So, again  $Y_F$  will tend to 1.

So,  $Y_{FS} \rightarrow 1$  under two conditions;  $T_\infty \rightarrow \infty$  or  $h_{fg} \rightarrow 0$ . So, we do not go to this critical point. So, when you are in the equilibrium regime, there should be a definite value of  $h_{fg}$ ; positive value of  $h_{fg}$  should be prevailed. So, that is very important thing. So, this is first thing you get from this relationship.

Second one is since the mass fraction of the fuel vapor at the droplet surface is a function of  $T_s$ ,  $Y_{FS}$  is a function of  $T_s$ , the values of this may be evaluated by solving the equation C. And we need one more equation that will relate the  $Y_{FS}$  and  $T_s$ . So, that is called thermodynamic relation, which will relate see that is what I have written here.

$$X_{FS} = \frac{p_{sat}}{p} = \exp\left(-\frac{h_{fg}}{R} \left[\frac{1}{T} - \frac{1}{T_{boil}}\right]\right) \quad (D)$$

Once you know the value of  $T_s$  or at least guess a value of  $T_s$ , you will be able to find the saturation pressure  $p_{sat}$  at that value of  $T_s$ . So, that relationship is called the thermodynamic relation, one of them is the Clausius, there are several relations; one of them which is actually applicable for ideal gases is Clausius - Clapeyron equation.

Clausius Clapeyron equation, we will actually one problem worked example with this. So, now, this is the Clausius Clapeyron equation in which you can see that the mole fraction of the fuel at the interface is nothing but the saturation pressure by the total pressure that is the pressure ratio. Partial pressure equal to exponent of the latent heat vaporization of the fuel divided by the specific gas constant.

So, this is nothing but  $R_u$  divided by the molecular weight of the fuel, specific gas constant of fuel into  $1/T$ ; any temperature. So, this  $p_{sat}$  at any temperature  $T$ , I am going to get that  $T$  unknown as a known temperature, but I will get the value of  $p_{sat}$ . Now, boiling point of the fuel

also is there. So, we see that if temperature  $T = T_{\text{boil}}$ , what happens? This is 0; so, this will be 1; that means,  $X_{\text{FS}}$  will be equal to 1.

So, as  $T$ , this is surface temperature; say this is surface temperature and this when surface temperature is increased very close to boiling point, as I told you it may not reach the boiling point.

So, it tends to boiling point, then  $X_{\text{FS}} \rightarrow 1$ . So, this is the Clausius Clapeyron equation thermodynamic equation relationship which connects the saturation pressure and temperature; unique relationship between the saturation pressure and the temperature. So, that is the equation. So, this and this if you use this C and this D, if you use, you can get the values of the  $Y_{\text{FS}}$  and  $T_s$ .

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*Evaporation Scenarios*

**Interface Quantities**


The equilibrium temperature,  $T_s$ , for various values of the ambient temperature, may be determined using the below procedure:

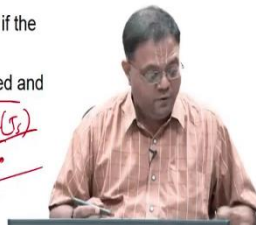
1. First a value of  $T_s$  is assumed.
2. Using the thermodynamic relation (equation (D)), the saturation pressure at  $T_s$  is evaluated.
3. The mole- and mass-fractions of the fuel at the droplet surface are calculated.
4. The values of  $B_T$  and  $B_Y$  are calculated.
5. Substitute these values in Equation (C) and determine if the equation is satisfied.
6. If Equation (C) is not satisfied, the value of  $T_s$  is adjusted and steps 2 to 5 are repeated until Equation (C) is satisfied.

$M_s = X_s M_F + (1 - X_s) M_A$

$X_{\text{FS}} = \frac{p_{\text{sat}}(T_s)}{p}$

*Handwritten notes:*  
 $T_\infty = 350 < 370$   
 $T_\infty = 400 > 370$   
 $T_\infty = 350$   
 $T_b = 370$   
 $T_s = 360$   
 $T_s = 370$   
 $T_s = 360$





The procedure for calculating the equilibrium temperature is as follows. So, you fix the ambient temperature. At any given ambient temperature, we can do this. So, let us do this. First you assume the value, you do not know the value of  $T_s$ . So, let us say, you take any fuel.

For example, say n-heptane the boiling point of this will be say 371 K. So, now, I can assume  $T_s$  at the surface can be say 360 K. This is assumption made because I know that it will never reach the boiling point.

So, this is boiling point of n-heptane is 371 K. So, this  $T_s$  will be close to boiling point based upon the ambient temperature. See for example, now ambient temperature  $T_\infty$  is say only 350 K. Now, there are scenarios you have to understand this guess is very important.

So, let us take. So, there are three temperatures I have; so,  $T_\infty$ ,  $T_{\text{boil}}$  and  $T_s$ . Now, out of the three temperatures, this 350 K, If I will put for this and my boiling point is 370 K; then, the surface temperature cannot exceed this  $T_\infty$ .

So, surface temperature can be guessed as say 340. So, definitely it should be less than the  $T_\infty$ , that condition we have to understand. So,  $T_s$  will be less than the  $T_\infty$ . Evaporation, please understand in the evaporation scenario. I fix the  $T_\infty$ ; now  $T_\infty$  is fixed actually less than the boiling point.

Boiling point is 370 and 350 is the ambient temperature I am fixing. So, evaporation will take place, but the surface temperature will not exceed the boiling point. So, it will not even go close to that because ambient temperature itself is less than the boiling point.

Now, if I say a scenario where the  $T_\infty$  is taken as 400 K which is actually greater than 370 which is the boiling point. Then, I may reach the temperature close to, I cannot exceed boiling point.

Please understand, I may go to 360 K here. So, please understand that  $T_s$  can be either less than the ambient temperature  $T_\infty$ , if the  $T_\infty$  itself is less than the boiling point or it may be less than the boiling point, if the ambient temperature is more than the boiling point.

So, this is the important thing you should remember to guess the value of or assume the value of  $T_s$ . So, if the situation is 350 K for  $T_\infty$ , I cannot guess  $T_s$  as 360 K because I cannot say that the boiling point is 370 K, I am taking temperature less than that; that is not possible.

So, either seeing the scenario, if  $T_\infty$  itself is less than boiling point, you have to take a temperature of the surface as it will be actually less than the  $T_\infty$  which is less than the boiling point or otherwise, if  $T_\infty$  exceeds the boiling point, then my surface temperature cannot exceed the boiling point.

So, it should be less than the boiling point. So, that is it. So, you cannot say that now ambient temperature is more, so the droplet surface temperature can exceed that boiling point to reach that, it is not possible.

So, keep that in mind, these two situations and first value of  $T_s$ ,  $T_s$  is assumed based upon this. Then, use the equation D that is the Clausius Clapeyron equation here. Clausius Clapeyron equation is equation D thermodynamic relationship and find the saturation pressure  $p_{sat}$  at this  $T_s$ . Now,  $p_{sat}$  is known. So, find the mole fraction.

So, once you know equation D is used again. So,  $p_{sat}$  is calculated or  $X_{FS}$  is calculated straight away; the mole fraction of the fuel. So, now you know the mole fraction. So, you have  $X_{FS}$  is calculated as  $p_{sat}/p$  at  $T_s/p$ . So, now, you know this. So, what is the molecular weight of the mixture? Molecular weight at the interface will be equal to  $X_{FS}$  into molecular weight of the fuel plus 1 minus  $X_{FS}$  into molecular weight of the ambient gas A.

Now, you know this now find  $Y_F$ ;  $Y_{FS}$ , molecular weight of the mixture is known and find  $Y_{FS}$ . So, that is what I have told that mole and mass fraction of the fuel at the droplet surface are

calculated. So, straight away the Clausius Clapeyron equation will give you the mole fraction; calculate the molecular weight of the mixture and calculate the mass fraction. So, that is done. Now, you know the mass fraction of the fuel, go back, calculate the values here. So, this  $B_T$  and  $B_Y$ , you can calculate and use this equation and see whether this equation is satisfied; correct? So, calculate the values of  $B_T$  and  $B_Y$  and see whether equation C is satisfied or not because we have guessed the value of  $T_s$ , that will give a value of  $Y_{FS}$ . But if you calculate  $Y_{FS}$  using equation C, that may not be the correct value what you get here.

So, this may not satisfy. So, that is a unique combination of  $T_s$  and  $Y_{FS}$  which will satisfy the equation C. So, that is not going to be there here. So, you will not get the value here. If you do not get this, then your assumption of  $T_s$  is not correct. Now, what you do is the  $T_s$  value is adjusted or iterated and the steps 2 to 5 are repeated to get the convergence of equation C, that is equation C should be finally satisfied.

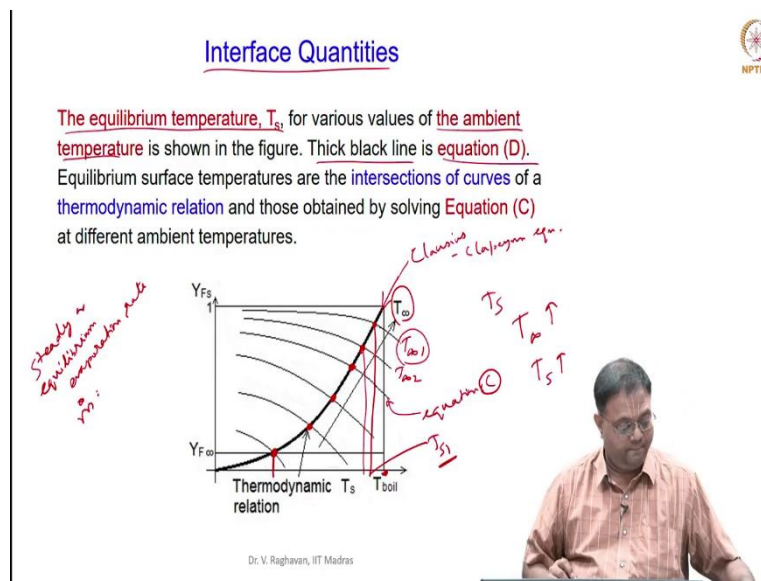
So, let us say for heptane itself, the second case let us take  $T_\infty$  is 400 K and  $T$  boiling point is 370 K. Now, your  $T_s$  can be say 360 K. I am just giving some idea; 360. So, as ambient temperature increases, your  $T_s$  will increase.

So, this iterative procedure to calculate the value of  $T_s$  for corresponding  $Y_{FS}$  values. So, that is the equilibrium condition at which; so, once you get the  $T_s$  and  $Y_{FS}$  value, then you calculate using any equation. Like for example, this; any one of this, you can use to calculate the mass evaporation rate. So, accurate value of  $B_T$ , if you want to calculate  $B_T$ , you will have to know the value of say  $c_{pg} T_\infty - T_s$ .

So,  $T_s$  value should be known to you calculate  $B_T$ . If you calculate  $B_Y$ , you know this value, again  $Y_{FS}$  is required for this. So, once you calculate the correct equilibrium values of  $T_s$  and  $Y_{FS}$ , you can get the mass evaporation rate using one of these equations. So, this is what the procedure which is used to calculate the interface quantities that you have to guess the value of  $T_s$  and understand, that I will repeat once again.

When the ambient temperature is less than the boiling point, your  $T_s$  should be less than the ambient temperature; when the ambient temperature is more than the boiling point, your  $T_s$  should be less than the boiling point. So, it cannot exceed the boiling point.

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Now, again this is the curve what we have basically seen I can do the exercise for several ambient temperatures  $T_\infty$ . So, you can see that the equilibrium temperature  $T_s$  for various ambient temperature is shown in this graph. So, these lines; so, this is say,  $T_{\infty 1}$ , this is  $T_{\infty 2}$  and so on.

So, several  $T_\infty$  values I get and I plot this. Now, what is this curve? This curve is nothing but the equation C which is this equation; so, equation C. So, when I try to calculate  $Y_{FS}$  as a function of  $T_s$  etcetera you get this equation. So, this is  $Y_{FS}$  as a function of  $T_s$ , I plot this. So, I will vary the  $T_s$  and try to get this curve  $Y_{FS}$ .

Now, this dark line, this dark line represents the equation D; thick black line, this is the Clausius Clapeyron equation. Now, what are the equilibrium temperatures? These intersections are the equilibrium temperatures. Graphically, these are the equilibrium temperatures. So, you can see that the equilibrium temperature  $T_s$  values, these are  $T_s$ . So, this is say corresponding to; so, for example, this is  $T_a, T_{\infty 1}$ . Now, this will be the corresponding  $T_{s1}$ .

So,  $T_{s1}$  corresponding value, this  $T_{s1}$  is the equilibrium value for the  $T_{\infty 1}$ , then similarly  $T_{\infty 2}$  corresponding  $T_{s2}$  and so on. So, this will be the corresponding equilibrium temperature. The intersection of these curves with this Clausius Clapeyron equation, it will be the  $T_s$ .

Now, you can see that as  $T_\infty$  increases the  $T_s$  also increases; but please understand that it cannot exceed the boiling point.

So, always I have considered the temperatures basically you can see that always the  $T_s$  is less than the boiling point,  $T_\infty$  itself can be less than boiling point in some scenario. If that is the case, then the  $T_s$  cannot exceed  $T_\infty$  itself. So, there are two things which you have to understand in this.



Graphically I am trying to show you the procedure, what you do I have illustrated here like assuming the value of  $T_s$ , you calculate the Clausius Clapeyron equation, calculate the  $p_{sat}$  and find the mole and mass fractions, then calculate the transfer numbers, then use the equation verify the validity of equation C.

If not, adjust the value of  $T_s$  and iterate. Graphically, equation C are several curves based upon the  $T_\infty$  value, fixed  $T_\infty$  value you get curves for equation C and those curves should intersect the equation D which is the Clausius Clapeyron equation.

Corresponding to that is intersection for a given  $T_\infty$ , we will get the value of  $T_s$ . Now, this is about the evaporation, where you try to find the equilibrium values of the surface quantities like  $T_s$  and  $Y_{F,\infty}$ . Once you get that, you will get the equilibrium evaporation rate or steady evaporation rate  $m$  dot.

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Spherical Droplet Combustion

When an ignition source is provided to an evaporating droplet, a diffusion flame is established around the droplet, at a certain radius  $r_f$ . This is schematically depicted in the figure. Fuel from the droplet surface ( $r = r_s$ ) reaches the flame zone by convection and diffusion, where it is consumed. Similarly, oxygen from the ambient is transported towards the flame zone and it is consumed almost completely in the flame zone.

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So, let us move onto droplet combustion. Again, we will assume the problem to be spherically symmetric. So that means, the droplet is spherical and a flame surrounds a droplet. So, what happen if you just ignite the droplet which is evaporating? Then, a diffusion flame is established around the droplet.

So, there is a flame which is established around the droplet and this is the flame. So, this is diffusion flame. Now, you can see that this is also a spherical, I will say spherical again here. So, droplet is spherical, the flame which is surrounding droplet is also spherical. The radius of the droplet at any time instant say  $r_s$  and the flame surrounds the droplet and its radius say  $r_f$ . So, the ignition source on evaporating droplet will create a flame over the droplet and the flame sustains at radius  $r_f$ . So, this is the scenario at a given time instant. So, what happens is fuel

from this droplet surface goes to the flame and the ambient air from the ambient or the oxidizer comes to this and products are formed here.

So, that is what the diffusion flame nature is here. So, you can see this when this burns it reduces the droplet diameter. So, what happens is the same type of this. So, if you say  $d^2$  versus time in this spherical coordinate and if it is a steady burning, steady droplet combustion I say steady burning problem, you can see that there will be a constant rate; the rate at which this burn is a constant.

So, whatever assumption we have made for the evaporation will be there, we can also make few assumptions here. So, when you talk about the structure of this flame it is like this. So, you can see this is the droplet surface  $r_s$  and the flame radius  $r_f$ . Please understand that this is a very thin flame basically, but I have given a finite thickness here.

In practice, there will be some finite thickness. So, I have exaggerated and shown the thickness here. So, what happens is you can see that the fuel from this mass fraction of the surface  $Y_{FS}$  goes and consume it, it is fully consumed at the flame zone. Similarly, oxygen from the ambient is consumed at the flame zone like this.

So, fuel and oxygen are consumed, where the fuel and oxygen are consumed you can see that the product has the maximum mass fraction there and the product actually diffuses in this zone. This is called a zone between the flame and the surface. So, this zone, I call as inner zone between the flame and the surface of the droplet.

The outer zone is obviously the ambient, from the flame to the ambient, it is called outer zone. The product also is transported to the ambient. So, the product which is formed at the flame zone, maximum value is diffused towards the inner zone and transported towards the outer zone.

So, oxygen is actually transported towards this zone and you can see that the oxygen cannot penetrate the inner zone, fuel cannot go to the outer zone because its consumed within the flame zone.

Temperature also is maximum because the reaction is completed in the flame and you can see the temperature decreases towards what is called  $T_s$ , surface temperature. Again, please understand like in steady evaporation problem, a steady combustion problem that has a definite value of  $T_s$  and  $Y_{FS}$ , equilibrium values.

Again, we can say that the equilibrium prevails. So, we will make the assumptions more clear in the next slides. Now, once the temperature is maximum here, it actually decreases towards the surface to reach the equilibrium temperature  $T_s$  and towards the ambient, it again decreases to reach the ambient temperature  $T_\infty$  at a far-off radius.

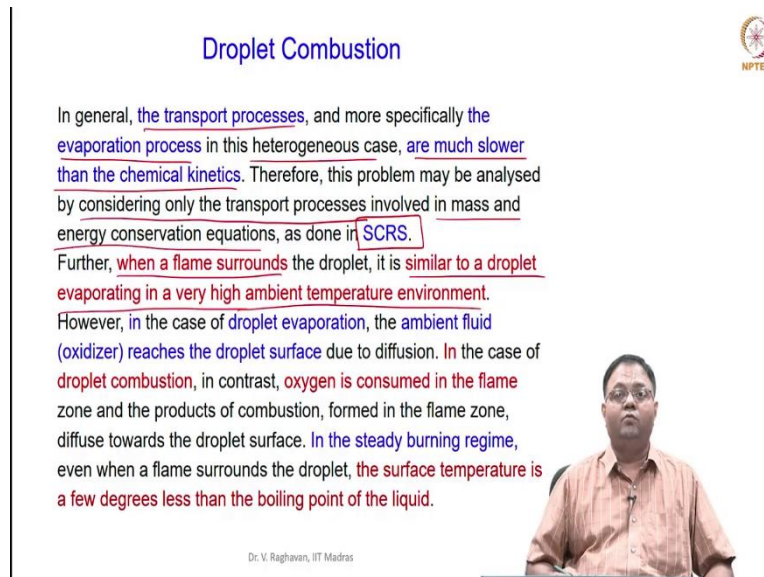
So, this is the structure of the droplet flame and the radial structure profiles. So, fuel from the droplet surface  $r = r_s$  basically reaches the flame zone by both convection and diffusion. Fick's law, we have seen that. So, this is like actually the flame separates the part.

For example, the ambient temperature will not have any influence now. In the previous case, evaporation case, only the ambient temperature actually was the driving force for evaporation. In this case, the flame itself is there which has higher temperature. Now, you can see that the surface temperature can be larger than the ambient temperature.

In the previous case, what we told was the surface temperature cannot be larger than the ambient temperature. Then, if the ambient temperature is less than the boiling point, then the surface temperature cannot exceed the ambient temperature; if the ambient temperature is more than the boiling point, then the surface temperature cannot exceed the boiling point.

So, in this scenario, we can see that there is no influence of the ambient temperature because of a high temperature, this is the flame temperature  $T_f$  that high temperature zone is there in between the droplet surface and the ambient. Similarly, the oxygen from the ambient is transferred towards the flame and it is consumed completely in the flame zone. So, this is the typical structure.

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**Droplet Combustion**


In general, the transport processes, and more specifically the evaporation process in this heterogeneous case, are much slower than the chemical kinetics. Therefore, this problem may be analysed by considering only the transport processes involved in mass and energy conservation equations, as done in SCRS.

Further, when a flame surrounds the droplet, it is similar to a droplet evaporating in a very high ambient temperature environment.

However, in the case of droplet evaporation, the ambient fluid (oxidizer) reaches the droplet surface due to diffusion. In the case of droplet combustion, in contrast, oxygen is consumed in the flame zone and the products of combustion, formed in the flame zone, diffuse towards the droplet surface. In the steady burning regime, even when a flame surrounds the droplet, the surface temperature is a few degrees less than the boiling point of the liquid.

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Now, like in any diffusion flame the transport processes are going to control the combustion zone; How fast the fuel is diffused; how fast the ambient air comes to the flame zone etcetera will control the flame zone location etcetera. But if you see in this case, additionally the evaporation process also is there; like the rate at which there is heat and mass transfer occurring to the droplet surface that also will be.

So, this is heterogeneous in nature because this has to cause a phase change basically and this process is slower, these processes are much slower than the chemical kinetics. So, that means that I do not need chemical kinetics at all to analyze a droplet combustion problem, especially when it is surrounded by a diffusion flame like this.

So, this problem can be analyzed by considering only transport process in mass and energy conservation equations, as done in the simple chemical reacting system SCRS. We already discussed that.

So, we need not invoke any chemical kinetics here. So, that is one of the things. Then, just some characteristics, we will see. When a flame surrounds the droplet, this is similar to the scenario of droplet evaporating in a high ambient temperature environment.