

Fundamentals of Combustion
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Lecture - 32
Characteristics of combustion flame and detonation Part 4
Estimation of detonation velocity and Worked examples


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Chapman – Jouguet Points

Since any real process going from state 1 to state 2 must satisfy both the Rayleigh line equation (equation 5) and the Rankine-Hugoniot relation (equation 8), the points between B and C on the Hugoniot curve are unrealizable because no valid Rayleigh line can be drawn between point A and any point between B and C.

For the upper branch of the Hugoniot, there is a limiting Rayleigh line that is just tangent to the Hugoniot curve at the point D. This point is called the upper Chapman – Jouguet point. It can be shown that Mach number ($M = 1$ at the point D). Points above D are difficult to observe in lab-scale experiments.

Similarly, there is a limiting Rayleigh line, just tangent to the lower branch of the Hugoniot at the point E. This is called the lower Chapman – Jouguet point.

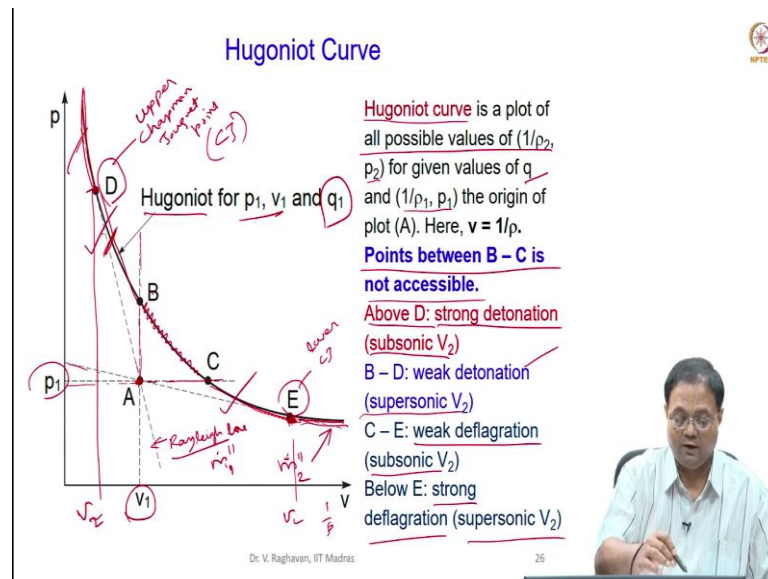


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And you can see some discussions I have written here. So, since any real process going from state 1 to state 2 must satisfy the Rayleigh line equation 5 and the Rankine Hugoniot relation, both should be satisfied. So, the points between B and C on the Hugoniot curve are not realizable because no valid Rayleigh line can be drawn. So, this is the first inference.

Then, the for the upper branch of the Hugoniot there is a limiting Rayleigh line. After that the Rayleigh line cannot be drawn as a tangent basically. So, that is a tangent. So, the limiting line and this is tangent to the Hugoniot curve at the point D and this is called upper Chapman - Jouguet point or say we can also say, Chapman - Jouguet point.

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So, this point D here is called upper Chapman - Jouguet. So, CJ, I will say CJ point upper CJ point and similarly, here it is called lower CJ point; Chapman - Jouguet point. So, it can be shown. Now, I can show that Mach number will be 1 at this point D upper Chapman - Jouguet point, Mach number will be 1. Points above D are difficult to observe in the lab scale experiments.

Theoretically, it may be fine; we can solve the equation and get, but it is not practically possible to achieve this in the lab scale experiments. In some conditions, you may get it, some big engines etcetera you may get it; but not possible. Similarly, there is a limiting Rayleigh line in the lower portion which will give the tangent to the Hugoniot which is point E, we call lower CJ point or Chapman - Jouguet point.

Here, so this point E. So, this is accessible C and below E, it is not accessible. It is not practical; theoretically, you can get the value, but it is not practical; the strong deflagration and supersonic exit velocities.

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Burnt Gas Velocity

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Consider the velocity of the burnt gases V_2 in the laboratory frame, where the flame (combustion wave) is moving with a velocity of V_w . Velocity of unburned gases $V_1 = V_w - u_1 = 0$, and velocity of the combustion wave, $V_w = u_1$ as shown in the figure. The velocity V_2 is expressed as $u_1 - u_2$.

$$V_2 = -(u_2 - u_1) = -(\dot{m}''^2) \left(\frac{1}{\rho_2} - \frac{1}{\rho_1} \right)$$

At point D, since $v_2 < v_1$, $u_1 > u_2$, V_2 is positive. For detonation, burned gases follow the combustion wave. For deflagration, at E, since $v_2 > v_1$, $u_1 < u_2$, V_2 is negative. For deflagration, burned gases flow in the direction opposite to that of the flame.

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Now, let us see about the burnt gas velocity. So, for that let the wave be moving, I will come out of the wave and see the wave propagating to the mixture. So, wave propagates at V_w . So that means, you can say $V_w = u_1$. So, when the wave is at rest the velocity of the unburned gas approaching the wave is u_1 ; so, $V_w = u_1$. So, that is known. Now, what is V_2 , we are going to find it.

So, now V_1 , this is V_1 , V_w , u_1 all the things we can write here. So, let us say. So, when I am in the laboratory frame of reference, what happens is the combustion wave is moving with the velocity of V_w . Velocity of unburned gas, V_1 is see this u_1 was the velocity at which the unburned gas approached wave. Now, if the wave is propagating unburned gas should be at rest.

So, V_1 is the velocity of unburned gas which should be at rest, when the wave propagates over that, so that is equal to 0. Now, $V_w = u_1$ that is known. So, either the unburned gas should approach the wave at velocity of V_1 , u_1 or the wave should propagate at the same velocity towards the unburned gas. Now, the velocity V_2 is $u_1 - u_2$. So, that is the thing here; $u_1 - u_2$. So, we can say $V_2 = u_1 - u_2$ or $-(u_2 - u_1)$. Now, I know the expression for $u_2 - u_1$ that is what we have derived that and kept.

$$V_2 = -(u_2 - u_1) = -(\dot{m}''^2) \left(\frac{1}{\rho_2} - \frac{1}{\rho_1} \right)$$

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Hugoniot Relation


Expression for $u_2 - u_1$ is written as:

$$(u_2 - u_1) = (\dot{m}''') \left(\frac{1}{\rho_2} - \frac{1}{\rho_1} \right)$$

For detonation, $\rho_2 > \rho_1$, and the reverse is true for a deflagration. Thus, the sign of $u_2 - u_1$ is negative for a detonation and positive for a deflagration. An expression for $(u_2)^2 - (u_1)^2$ is obtained as,

$$u_1^2 - u_2^2 = (p_2 - p_1) \left(\frac{1}{\rho_2} + \frac{1}{\rho_1} \right)$$

Substituting the above equation in the energy equation (3), **Hugoniot relation** is obtained.




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So, the $u_2 - u_1$ which we derived here, that can be used here now $-(\dot{m}''')^2 (1/\rho_2 - 1/\rho_1)$. So, now, just to see the direction only we are doing this calculation. So, at point D, $V_2 < V_1$. Here, point D; so, this is V_1 . So, this is V_1 p_1 . Point D, $V_2 < V_1$. So, when $V_2 < V_1$, this is V_2 and this is V_1 ; when $V_2 < V_1$, this is negative. Negative of negative will be positive. So, V_2 is positive. So, for detonation, the burned gas follows the combustion wave; when the wave moves in the positive x direction, the burned gas also follow the wave, that we have to understand.

For deflagration at E, point E let us take. Point E, you can see that the V_2 ; this is V_2 now, $V_1 < V_2$. $V_2 > V_1$. So, $u_1 < u_2$ or V_2 is negative. So, for deflagration, burned gas flow in the direction opposite to that of the flame. This is what we have to understand.

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Estimation of Detonation Velocity

When detonation occurs, the unburned gas enters the combustion wave at supersonic velocity (u_0). Therefore, $u_1 = u_0$. **At the upper CJ point (D), the velocity is sonic. Above D, velocity is subsonic and in B-D it is supersonic.** The continuity equation is written as,

approaching velocity
 $\rho_1 u_1 = \rho_2 c_2 \Rightarrow u_1 = \frac{\rho_2}{\rho_1} c_2 = \frac{\rho_2}{\rho_1} (\gamma R_2 T_2)^{1/2}$ *speed of sound* $c = \sqrt{\gamma R T}$
 $\rho_1 T_1 = \rho_2 T_2$

Relating ρ_2/ρ_1 and T_2 to state 1 or other known quantities, and neglecting pressure at state 1 as compared to state 2 ($p_2 \gg p_1$),

$$\frac{\rho_1 u_1^2}{\rho_2 u_2^2} - \frac{p_2}{\rho_2 u_2^2} = 1 \Rightarrow \frac{\rho_2}{\rho_1} = 1 + \frac{p_2}{\rho_2 u_2^2}; \because u_2 = c_2 = \sqrt{\gamma R_2 T_2}$$

$$\Rightarrow \frac{\rho_2}{\rho_1} = 1 + \frac{p_2}{\rho_2 \gamma R_2 T_2}; \because p_2 = \rho_2 R_2 T_2 = \frac{\rho_2}{\rho_1} \frac{\gamma + 1}{\gamma}$$

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Now, estimation of detonation velocity. So, at the upper Chapman - Jouguet point, CJ point velocity is sonic. Above D, the velocity is subsonic and between B and D, the velocity is supersonic. When I say velocity it is the approaching velocity, please understand this is the approaching velocity; not the burnt gas velocity. Burnt gas velocity is different, we have already seen that. This is approaching velocity

So, when we want to do this, the approaching velocity should be in between the B and D and is supersonic. So, the continuity equation, I can write as $\rho_1 u_1$, the approaching velocity is u_1 equal to $\rho_2 c_2$; so, now, $u_1 = (\rho_2/\rho_1)c_2$. What is c_2 ? c_2 is the speed of sound which is nothing but $\sqrt{\gamma R T}$; $c = \sqrt{\gamma R T}$. So, $c_2 = \sqrt{\gamma R_2 T_2}$. So, this is the expression I get.

$$u_1 = \frac{\rho_2}{\rho_1} c_2 = \frac{\rho_2}{\rho_1} (\gamma R_2 T_2)^{1/2}$$

Now, I can relate this ρ_2/ρ_1 and T_2 to state 1 quantities or other quantities. Now, further I make important assumption that $p_2 \gg p_1$. So, when I say the momentum equation $p_1 + 1/2 \rho_1 u_1^2 = p_2 + 1/2 \rho_2 u_2^2$ for example. So, now, $p_1 = 0$ because I can neglect p_1 and consider the value of p_2 .

So, p_1 is very low. So, I write this equation. Please understand that, I write this when I use this momentum equation, I can write $\rho_1 u_1^2 / \rho_2 u_2^2 - p_2/\rho_2 u_2^2$. Just divide throughout by $\rho_2 u_2^2$. So, I have to put 1 here; 1 and 2 here.

So, momentum equation. $p_1 + 1/2 \rho_1 u_1^2 = p_2 + 1/2 \rho_2 u_2^2$, by this assumption I say $p_1 = 0$.

So, I can write this. From that, I can get the value of ρ_2/ρ_1 as this

Now, $u_2 = c_2$; so, using that I can write ρ_2/ρ_1 . Then, using the definition of the c_p etcetera plus the equation of state $p = \rho RT$, I can get a simple formula for ρ_2/ρ_1 as $(\gamma+1)/\gamma$.

$$\frac{\rho_2}{\rho_1} = \frac{\gamma + 1}{\gamma}$$

So, you can just use it very simple derivation here.

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Estimation of Detonation Velocity

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Solving the energy conservation equation for T_2 , eliminating u_1 using the continuity equation, substituting $c_2 = u_2$, using the density ratio as obtained above, the following expression is obtained.

$$c_p T_2 + \frac{1}{2} u_2^2 = c_p T_1 + \frac{1}{2} u_1^2 + q$$

$$T_2 = T_1 + \frac{u_1^2 - u_2^2}{2c_p} + \frac{q}{c_p} = T_1 + \frac{q}{c_p} + \frac{\gamma R_2 T_2}{2c_p} \left[\left(\frac{\gamma + 1}{\gamma} \right)^2 - 1 \right]$$

Solving for T_2 , and using the relationship between γ , c_p and R , the following relation is obtained.

$\gamma_1 = \frac{c_p}{c_v}$
 $\gamma_2 = \frac{c_p}{c_v}$
 $\gamma_1 = \gamma_2$

$T_2 = \frac{2\gamma^2}{\gamma + 1} \left(T_1 + \frac{q}{c_p} \right)$

$\frac{\gamma R_2}{\gamma - 1} = c_p$
 $p = \rho RT$

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So, now similarly T_2 , I want a relationship for ρ_2/ρ_1 which I have got now; similarly, I want a relationship for T_2 . So, for that what I do? Solving the energy conservation equation for T_2 , eliminating u_1 using the continuity equation substituting $c_2 = u_2$, I can get this and also, using the density ratio. So, much thing I do here. I am not deriving it. So, that is what the summary is.

So, what I do here is I know that $c_p T_2 + 1/2 u_2^2 = c_p T_1 + 1/2 u_1^2$. So, that I try to now take value of T_2 . So, $T_2 = T_1$ divide throughout by c_p , $u_1^2 + q$ is there, $(u_1^2 - u_2^2)/2c_p + q/c_p$. It is only energy equation, simple.

Now, this you retain, we will try to write this γRT . So, $u_2 = c_2 = \sqrt{\gamma R_2 T_2}$ that you substitute you get this equation. Now, you know that $\gamma R_2/(\gamma - 1) = c_p$; similarly, $p = \rho RT$. So, use this equation; finally, this equation you get.

$$T_2 = T_1 + \frac{u_1^2 - u_2^2}{2c_p} + \frac{q}{c_p}$$

$$= T_1 + \frac{q}{c_p} + \frac{\gamma R_2 T_2}{2c_p} \left[\left(\frac{\gamma + 1}{\gamma} \right)^2 - 1 \right]$$

So, now this we have taken to this side and do some algebra, you get a simple equation for T_2 as $2\gamma^2/(\gamma-1)(T_1 + q/c_p)$. Plus understand that this is the simple equation because we have assumed c_p to be constant.

$$T_2 = \frac{2\gamma^2}{\gamma+1} \left(T_1 + \frac{q}{c_p} \right)$$

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Estimation of Detonation Velocity

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Substituting the density ratio and expression for T_2 in the expression for u_1 , detonation velocity is obtained as,

$$u_D = u_1 = \left[2(\gamma + 1)\gamma R_2 \left(T_1 + \frac{q}{c_p} \right) \right]^{1/2}$$

The above expression for detonation velocity is approximate because of the assumptions invoked on constant physical properties and due to the assumption that $p_2 \gg p_1$.

If the assumption of constant and equal specific heats is relaxed, a more accurate, but still approximate, expressions for the state-2 temperature and the detonation velocity can be derived.

$c_{p,1} = c_{p,2}$



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So, now you use this density ratio ρ_2/ρ_1 and T_2 and get the velocity for detonation, that is all. So, detonation velocity is nothing but the approaching velocity $u_1 = [2(\gamma+1)\gamma R_2 (T_1 + q/c_p)]^{1/2}$.

$$u_D = u_1 = \left[2(\gamma + 1)\gamma R_2 \left(T_1 + \frac{q}{c_p} \right) \right]^{1/2}$$

Now, the above expression for detonation velocity is very approximate. First of all, main assumption is p_2 is much more greater than p_1 and also, we have the constant physical properties.

Now, slightly we can relax it. So, what we will do? First assumption say, constant and equal specific heats say $c_{p,1} = c_{p,2}$, let us avoid this. When I do that, you get the expression for state 2 temperature and detonation velocity slightly better; but still it is approximate. I can get slightly better values, but it is still approximate. This is the thing.

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Estimation of Detonation Velocity



Expressions for T_2 and u_D with varying specific heats are:

$$T_2 = \frac{2\gamma_2^2}{\gamma_2 + 1} \left(\frac{c_{p1}}{c_{p2}} T_1 + \frac{q}{c_{p2}} \right)$$

$$u_D = u_1 = \left[2(\gamma_2 + 1)\gamma_2 R_2 \left(\frac{c_{p1}}{c_{p2}} T_1 + \frac{q}{c_{p2}} \right) \right]^{1/2}$$

$$\begin{aligned} c_{p1} &= c_{v1} \gamma_1 \\ c_{p2} &= c_{v2} \gamma_2 \end{aligned}$$

A more exact mathematical description of an one-dimensional detonation wave is available in literature. A numerical solution method for such a formulation is described by Gordon and McBride and implemented in the NASA Chemical Equilibrium Code (CEC).



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So, now I try to put instead of considering $c_{p,1} = c_{p,2}$, now I will consider it is not. So, I will put the ratio of $c_{p,1}/c_{p,2}$ here. So, T_2 is slightly better, see if you see this, you can get γ . There is no γ_1 or γ_2 ; now, I can specify its γ_1 or γ_2 . So, please understand $\gamma_1 = c_{p1}/c_{v1}$; $\gamma_2 = c_{p2}/c_{v2}$.

$$T_2 = \frac{2\gamma_2^2}{\gamma_2 + 1} \left(\frac{c_{p1}}{c_{p2}} T_1 + \frac{q}{c_{p2}} \right)$$

Now, if I assume that $c_{p1} = c_{p2}$ and $c_{v1} = c_{v2}$, then $\gamma_1 = \gamma_2$. These are all assumptions made here to get this. But when I relax that at least this the concept of equal specific heat etcetera for the products and this if we relax, I will say I will take $\gamma_1 = c_{p1}/c_{v1}$ and $\gamma_2 = c_{p2}/c_{v2}$.

Now, you can see that I will retain γ_2 here and the ratio of c_{p1}/c_{p2} etcetera here and get this. But still it is approximate, but slightly better than the previous equation because I am now relaxing the constant property assumption here. So, I get an improved equation for this u_D . So, this you can use. But please understand that more exact mathematical description of a one-dimension detonation wave.

$$u_D = u_1 = \left[2(\gamma_2 + 1)\gamma_2 R_2 \left(\frac{c_{p1}}{c_{p2}} T_1 + \frac{q}{c_{p2}} \right) \right]^{1/2}$$

Please understand one more thing, detonation is not one-dimensional phenomenon, it is 3D phenomenon. However, some assumptions can be made and one-dimensional detonation wave propagation can be adjudged. For that, know you need a numerical solution and NASA Chemical Equilibrium Code, CEC has a module to do this.

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Worked Example



A combustion wave propagates with a mass flux of $3500 \text{ kg/m}^2\text{s}$ through a mixture initially at 298 K and 1 atm . The molecular weight and specific-heat ratio of the mixture (burned and unburned) are 29 kg/kmol and 1.3 , respectively, and the heat release is $3.4 \times 10^6 \text{ J/kg}$. Determine the state of the burned gas and determine in which region this state lies on the Rankine-Hugoniot curve. Also, determine the Mach number of the burned gases.

Solution: Equation (5) is written as, $p_2 = p_1 + (\dot{m})^2(v_1 - v_2)$. Equation (8) is written as,

$$q = \frac{\gamma(p_2v_2 - p_1v_1)}{\gamma - 1} - \frac{1}{2}(p_2 - p_1)(v_1 + v_2)$$

Specific volume at state 1:

$$v_1 = RT_1/p_1 = (8314/29) \times 298/101325 = 0.843 \text{ m}^3/\text{kg}.$$

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Now, at last we will see a worked example. A combustion wave propagates with the mass flux of $3500 \text{ kg/m}^2\text{-s}$ and the initial mixture has a temperature of 298 K and 1 atmospheric pressure.

The molecular weight and the specific heat ratio of the mixture burned and now, both are same. Molecular weight is 29 kg/kmol and specific heat ratio $\gamma=1.3$. So, that is what is given in this.

So, the heat release rate q is given as $3.4 \times 10^6 \text{ J/kg}$. So, now, determine the state of the burned gas and determine in which region, this state lies in the Rankine Hugoniot curve; then, Mach number of the burned gases. So, this is the problem. Now, let us use the equation of Rayleigh line (equation 5) Rayleigh line. So, $p_2 = p_1 + (\dot{m})^2 (v_1 - v_2)$.

Similarly, equation 8 is written like this; equation 8, that is the Rankine Hugoniot curve equation. Now, I want to fix the state. So, what is v_1 ? $p = \rho RT$; so, $v_1 = RT_1/p_1$. So, I get $0.843 \text{ m}^3/\text{kg}$, that is the volume 1; pressure is 1 atmosphere, initial pressure is 1 atmosphere. So, initial state is fixed as 1 atmosphere and specific volume as $0.843 \text{ m}^3/\text{kg}$.

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Worked Example

$v_1 = 0.843 \text{ m}^3/\text{kg}$

$p_2 = 101325 + 12250000(0.843 - v_2) = 10428075 - 12250000v_2$

Heat addition equation is written as,
 $3400000 = 4.333(p_2 v_2 - 85416.98) - 0.5(p_2 - 101325)(0.843 + v_2)$


Substituting for p_2 in the above equation, a quadratic equation in v_2 is
 got as, $-4.7 \times 10^7 (v_2)^2 + 4.52 \times 10^7 (v_2) - 8.122 \times 10^6 = 0$
 $\Rightarrow 4.7(v_2)^2 - 4.52v_2 + 0.8122 = 0$

Solving, $v_2 = [4.52 \pm (4.52 \times 4.52 - 4 \times 4.7 \times 0.8122)^{0.5}] / (2 \times 4.7)$
 $\Rightarrow v_2 = (4.52 \pm 2.2718) / 9.4 = 0.723 \text{ m}^3/\text{kg}$ or $0.239 \text{ m}^3/\text{kg}$.

Corresponding values of p_2 are:
 For $v_2 = 0.723 \text{ m}^3/\text{kg}$, $p_2 = 10428075 - 12250000 \times 0.723 = 1571325 \text{ Pa}$.
 For $v_2 = 0.239 \text{ m}^3/\text{kg}$, $p_2 = 10428075 - 12250000 \times 0.239 = 7500325 \text{ Pa}$.

The velocities are calculated using $m'' = \rho V = V/v \Rightarrow V = m'' \times v$ $\rho = 1/v$
 For $v_2 = 0.723 \text{ m}^3/\text{kg}$, $V_2 = 3500 \times 0.723 = 2530.5 \text{ m/s}$
 For $v_2 = 0.239 \text{ m}^3/\text{kg}$, $V_2 = 3500 \times 0.239 = 836.5 \text{ m/s}$

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Now, p_2 ; p_2 this equation, Rayleigh line equal to p_1 . Here all the values. So, you can see $p_1 + (m'')^2$, $m'' = 3500 \text{ kg/m}^2\text{-s}$, then v_1 and v_2 . So, v_1 , but I do not know v_2 . So, I form an equation for p_2 in terms of v_2 here like this; form an equation for p_2 in terms of v_2

Similarly, heat addition, this equation 8 that is the Rankine Hugoniot equation, I can substitute the values here; but please understand I do not know both p_2 and v_2 , but I know now from this equation, Rayleigh line equation, I know p_2 in terms of v_2 .

So, now substituting for p_2 from the Rayleigh line equation, you get a quadratic equation in v_2 and this is a quadratic equation, when I solve I get 2 volumes v_2 's that is $0.723 \text{ m}^3/\text{kg}$ and $0.239 \text{ m}^3/\text{kg}$.

You can readily see that v_1 was $0.843 \text{ m}^3/\text{kg}$ and we got two v_2 's from the quadratic equation, both are less than that. That means, we are on the detonation curve. So, if you see here this is the initial point, but both v_2 's are less than this now.

So, either I will be here or here somewhere, I am going to be there in this regimes. So, both v_2 's are less than the v_1 , initial v_1 what I got. So, now that is the 2 values I get. Now, for both we will try to fix the state

So, corresponding to the v_2 of 0.723 , p_2 value is got as this is the same 1.57 MPa Similarly, for v_2 of 0.239 , I get 7.5 MPa . So obviously, you can see the higher, we are travelling in that curve to a higher value of pressure.


Now, the velocities are calculated now I know the pressure, let us say state 2 is fixed now. Pressure for corresponding each state is fixed. Now, how will we calculate velocities? I know the mass flux; mass flux, mass flow rate is what? Mass flow rate is ρAV , mass flux is ρV . So, divided by A ; so, ρV . So, this is ρV . But $\rho = 1/v$.

So, capital V is a velocity; so, velocity by specific volume. So, velocity will be equal to mass flux into specific volume. So, for specific volume of point 2 0.723, velocity will be 2530 m/s, state 2 velocity. For this v_2 equal to 0.239, it will be 836.5 m/s. So, I am getting 2 velocities ok; 2530 for the specific volume of 0.723 which is slightly less than less than the initial 0.843.

If it is much less than the initial this v_2 equal 0.239, pressure is very high. For that, I am getting a lower velocity of 836.5 correct.

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Worked Example



Speed of sound is calculated as $a = (\gamma RT)^{0.5}$.

Temperatures at state 2:

For $v_2 = 0.723 \text{ m}^3/\text{kg}$, $T_2 = 1571325 \times 0.723 / (8314/29) = 3963 \text{ K}$.

For $v_2 = 0.239 \text{ m}^3/\text{kg}$, $T_2 = 7500325 \times 0.239 / (8314/29) = 6253 \text{ K}$.

For $v_2 = 0.723 \text{ m}^3/\text{kg}$, $a_2 = [1.3 \times (8314/29) \times 3963]^{0.5} = 1215.3 \text{ m/s}$.

For $v_2 = 0.239 \text{ m}^3/\text{kg}$, $a_2 = [1.3 \times (8314/29) \times 6253]^{0.5} = 1526.6 \text{ m/s}$.


Mach number = $M_2 = V_2/a_2$.

For $v_2 = 0.723 \text{ m}^3/\text{kg}$, $M_2 = 2530.5/1215.3 = 2.08$

For $v_2 = 0.239 \text{ m}^3/\text{kg}$, $M_2 = 836.5/1526.6 = 0.55$

Both specific volumes at state 2 are less than $v_1 = 0.843 \text{ m}^3/\text{kg}$.
 State (7.5 MPa, 0.239 m^3/kg), which has a subsonic velocity, lies in the region above upper CJ point D. State (1.57 MPa, 0.723 m^3/kg), which has a supersonic velocity lies somewhere in the region between points B & D. These states result from very high mass flux.

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Now, I calculate the speed of sound $(\gamma RT)^{1/2}$. For that, I need temperatures. So, for each volume, again I calculate temperature as for 0.723, I get the temperature as 3963 K and for 0.239, I get temperatures of 6253 K.

So, now, 6000 etcetera, I cannot realize in the lab scale, that is what we are trying to make a point here. It is not realizable in the lab scale. So, now from this, I can calculate the value of the speed of sound a_2 for 0.723 m^3/kg , it is 1215 m/s and for 0.239, it is 1526 m/s. So, Mach numbers will be for the specific volume 0.723, it is 2.08. That means supersonic. So, that will be in between the B and D value correct; B and D that portion and this is subsonic; that will be above D. For 0.239, it is above D. So, both specific volumes are less than state 1. So, this state corresponds to subsonic velocity lies in the upper CJ point; above the point D, 7.5 MPa, 0.239.

Subsonic velocity 0.55 that is above that above D and this 1.57 MPa, 0.723 m^3/kg that lies in the region between the point B and D because it has a supersonic velocity. So, you can see that the very high mass flux causes detonation. If you reduce the mass flux etcetera, then you go to the deflagration part. So, I will stop here.