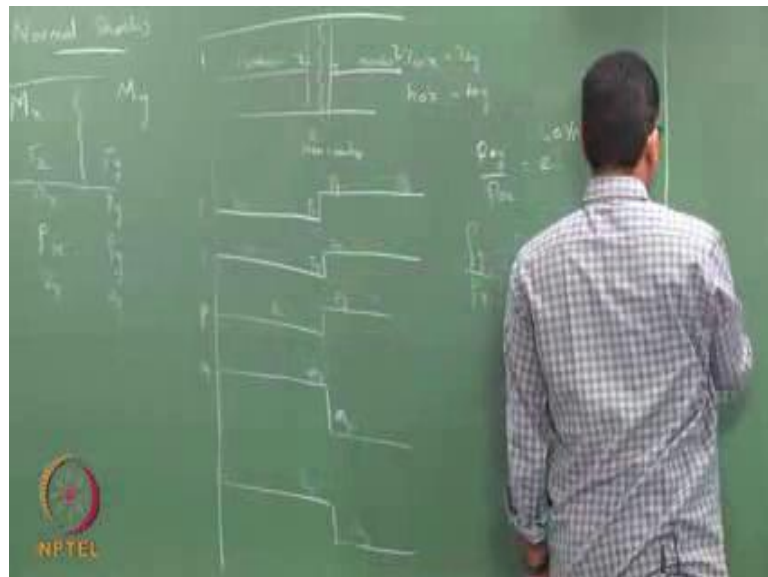


Fundamentals of Gas Dynamics
Dr. A. Sameen
Department of Aerospace Engineering
Indian Institute of Technology, Madras

Week - 08
Lecture – 32
Rankine - Hugoniot equation

So, last class, we were discussing about normal shocks and the variation of properties across the normal shock. So, what we had seen is, if I have a shock, given my value of Mach number before the shock, I know an expression to find the Mach number after the shock.

(Refer Slide Time: 00:24)



We also know how to find the ratio T_x and T_y , likewise, P_x and P_y , and ρ_x and ρ_y ; again, v_x and v_y , and all such quantities, we can get, if we know Mach number before the shock. So, these ratios, we had seen as increasing with, with respect to a , with respect to the shock. So, once it crosses the shock, you have an increase in all these properties, except the stagnation pressure.

So, we have also seen the stagnation temperature to be same, because, there is no heat, and there is no shaft work. So, this is going to be true; but, it is not isentropic process. So, up to this, we had seen, from 1 to x, we have an isentropic flow; and then, suddenly, it encounters a shock, which is a non-isentropic flow; and, from shock to some other

location 2, from y to 2, you have, again, isentropic solution.

So, if I know these properties, isentropic properties up to shock, then, I can find, across the shock, what happens, and from across the shock, after the shock, you again compute everything, using isentropic solution. So, if I plot my density, if this is my location of shock, the density would be some value up to the shock, then, suddenly, there is a jump, and then, increases my... So, this is my ρ_y , and this is my ρ_x . Likewise, your temperature also increases by some amount, which would be here, and here, assuming, assuming the thickness of the shock as very, very small, so, if there is a sudden discontinuity of density and temperature.

Then, you have the pressure also, does the same thing. The Mach number could be large. The moment it encounters the shock, it goes down to subsonic value. So, it goes down to subsonic value. So, I have an m_x here, and m_y here, which is lower than; this is ρ ; this is t ; this is P ; and this is m . And, my stagnation pressure also goes down; P_{0x} , P_{0y} , because it is an isentropic, non-isentropic process, your P_{0y} by P_{0x} is given by $e^{-\Delta s / R}$, where Δs is the entropy change that is happening due to the irreversibility that are associated with your shock.

So, expression for this all, we have seen. We have also seen, how these ratios change with function of m_x . So, you could actually write all these ratios in terms of Mach number before the shock. To consolidate what all we had said in the T S diagram?

(Refer Slide Time: 04:52)



So, from 1 to x, 1 to x, 1 is here; x here, which is an isentropic process; from x to y is a non-isentropic process; but, there is an increase in my temperature. So, I would increase my temperature here. And, from x to 2 is again, y to 2 is again an isentropic process. So, I have 1, goes to x; then, suddenly it encounters a shock; there is an entropy change; and then, it goes to y after the shock; and, from y to 2, you have again an isentropic solution.

Both this have the, all this states have the same stagnation temperature, because, there is no shaft work; there is no heat transfer. We have also seen this, how this would appear in a p v diagram. So, I have 1, goes to x; then, there is a shock; because of the shock, there is a density increase, which means the volume would decrease. So, I have a point somewhere here, which also means, my pressure has increased. So, my y is on a line, isentropic line above this.

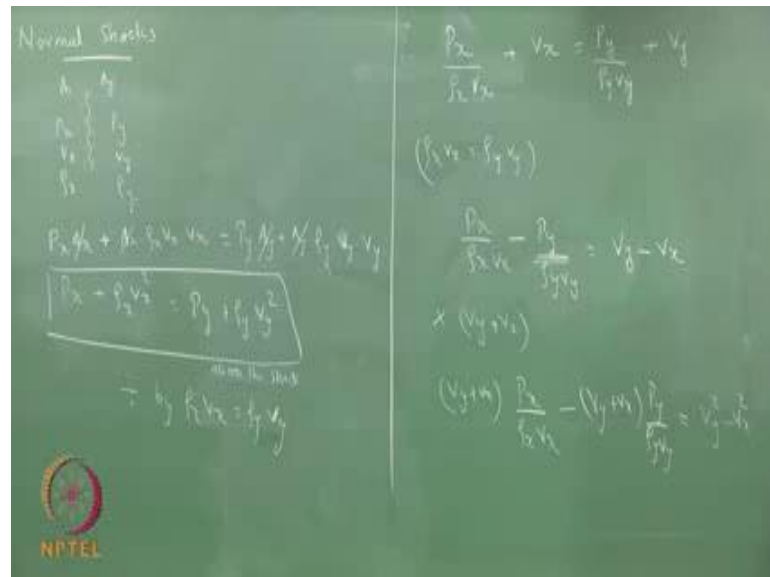
So, this would be my y, which again goes to 2, isentropic way. So, these are isentropic lines, and this is my shock. This would produce a Δs ; Δs is here. This would be my $\Delta \rho$. So, this is a pressure, some pressure here; some pressure here. So, if you look at the stagnation pressure here, this $P_0 2$ is now less than $P_0 1$, according to this relation. So, your $P_0 y$ is always less than your $P_0 x$, if you want an increase in my entropy. Difference is, yes, 1 by Δv . So, this difference is Δv , which is 1 by $\Delta \rho$. So, this is for the shock that is happening here; from 1 to x, it is isentropic; x to y is non isentropic; and, from y to 2 is again isentropic.

So when, if I have a shock in a, shock in a varying area duct, so, up to the shock, I can assume to be isentropic; if there are no heat, and no shaft work, I can assume that to be isentropic. So, we can use all the relation that we have been using for isentropic flow; then, after that, non-isentropic, and then, again, isentropic. So, in the, whether it is a converging nozzle, diverging nozzle, a diffuser, or a c d nozzle, wherever there is a shock, you can assume that, it is, up to that, it is isentropic, and then.

Now, we will try to derive an equation which relates P_1 , P_x and P_y , in terms of ρ_x and ρ_y . This equation is called Rankine - Hugoniot equation. So, there is a jump in the density and pressure, and this jump, and density and pressure are related; and this expression that gives you the density jump in terms of pressure jump, or pressure jump in terms of density jump, is typically called as the Rankine - Hugoniot equation. So, we will derive that, and we will, after that, we will try to do some numerical problems to see how

this works.

(Refer Slide Time: 10:31)



So, like before, I start with the shock here, which has a control volume that is very close by. So, my A_x is same as my A_y ; and, I have a P_x here, P_y here; v_x and v_y ; ρ_x and ρ_y . I apply the momentum equation. I would get $P_x A_x$, plus $A_x \rho_x v_x$ into v_x , which is equal to my $P_y A_y$ plus $A_y \rho_y v_y$, multiplied by v_y .

So, this is your force due to pressure, and this is the kinetic energy. So, this is $m \dot{v}_x$; this is $m \dot{v}_y$. Since A_x and A_y are same, I can cancel out A_y . So, I would end up with P_x plus $\rho_x v_x^2$, equals P_y plus $\rho_y v_y^2$, across the shock.

So, now, our job is to eliminate v_x and v_y , and we end up with Rankine - Hugoniot equation; that is all that is need to be done here. So, for that, we will do some minor algebra. So, I am going to divide by $\rho_x v_x$. So, this would be x that would divide by y . So, I would divide by $\rho_x v_x$, plus v_x , equals P_y by $\rho_y v_y$, which is same as my $\rho_x v_x$ plus v_y . So, my $\rho_x v_x$ is equal to $\rho_y v_y$. So, what I have done here is, I have divided this equation by $\rho_x v_x$, which is also equal to $\rho_y v_y$.

So, I would end up with an equation. Now, I would multiply this equation with v_y plus v_x . So, I would get v_y plus v_x multiplied by P_x by $\rho_x v_x$, minus v_y plus v_x multiplied by P_y by $\rho_y v_y$, equals v_y^2 minus v_x^2 . Now, I expand that term, term on this side.

(Refer Slide Time: 15:25)



So, this would be v_y by v_x into P_x by ρx , plus $v_y v_x$, $v_x v_x$ cancels. So, it would be P_x by ρx minus v_x into v_x by v_y into P_y by ρy minus $P_y \rho y$ equals v_y square minus v_x square. Now, from the continuity equation, which I have written here, my v_x by v_y is also my ρ_y by ρ_x . So, wherever I see this ratio, I substitute that in terms of ρ_y and ρ_x .

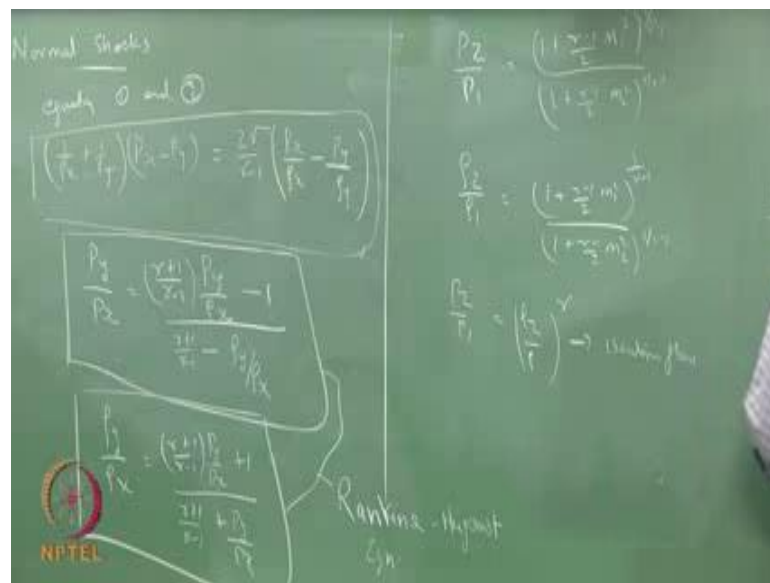
So, my v_y by v_x is now ρ_x by ρ_y , into P_x by ρx , plus P_x by ρx , minus v_x by v_y is ρ_y by ρ_x , into P_y by ρy , minus P_y by ρy , equals v_y square minus v_x square. So, this and this cancels; this and this cancel. So, I can write in this particular form, where I write a P_x into $1 + \rho_x$, plus 1 by ρ_y , minus P_y into $1 + \rho_x$, plus 1 plus 1 by ρ_y , equals v_y square minus v_x square; or, 1 by ρ_x plus 1 by ρ_y equals P_x minus P_y , equals v_y square minus v_x square; yes.

So, this is from the momentum equation, and I have used continuity equation to substitute v_y by v_x as ρ_x by ρ_y . Now, we will use the energy equation. So, this is my equation 1. Now, from the energy equation, I can write v_x square by 2 plus $C_p T_x$ equals v_y square by 2 plus $C_p T_y$, which is nothing, but $h_0 x$ equals $h_0 y$, across the shock, from which I can write, v_x square minus v_y square by 2 equals $C_p T_y$ minus T_x , which is $\gamma R T_y$ minus $\gamma R T_x$, divided by γ minus 1 . So, I had to change to C_p to γR by γ minus 1 . So, $R T$ I can write in terms of p v ; p by ρ equals $R T$. So, I use this particular relation to get, convert that T to pressure and

density ratio. So, my $v_x^2 - v_y^2 = 2(\gamma - 1) \frac{P_x - P_y}{\rho_x}$, into γ ; $R T$ y I substitute as P_y / ρ_y , minus P_x / ρ_x .

So, this is from my energy, which is equation 2. So, what I have done here is, I have taken the momentum equation, I eliminated my velocity in this particular form. Here, I have taken the energy equation, eliminated the temperature, and I have come across this particular form. So, velocity in terms of pressure and density; here again, velocity in terms of, now, equate these two that would give me the Rankine - Hugoniot equation.

(Refer Slide Time: 21:39)



So, equating equation 1 and 2, which I have just now written, this would be $1 + \rho_x v_x^2 - P_x = 1 + \rho_y v_y^2 - P_y$, is minus of the other one. So, it would be $2(\gamma - 1) \frac{P_x - P_y}{\rho_x} = P_x - P_y - \rho_x v_x^2 + \rho_y v_y^2$. So, I have included the minus sign inside the bracket.

Student: (Refer Time: 22:04).

So, this is P_x / ρ_x ; everywhere else its fine. So, I get this equation, which can be reduced to the following form. So, I can write this in terms of P_y and P_x , equals $\frac{\gamma + 1}{\gamma - 1} \frac{\rho_y}{\rho_x} - 1$, the whole divided by $\gamma + 1$ by $\gamma - 1$, minus ρ_y / ρ_x . Or, $\rho_y / \rho_x = \frac{\gamma + 1}{\gamma - 1} \frac{P_y}{P_x} + 1$, divided by $\gamma + 1$ by $\gamma - 1$, plus P_y / P_x , these two equations are called Rankine.

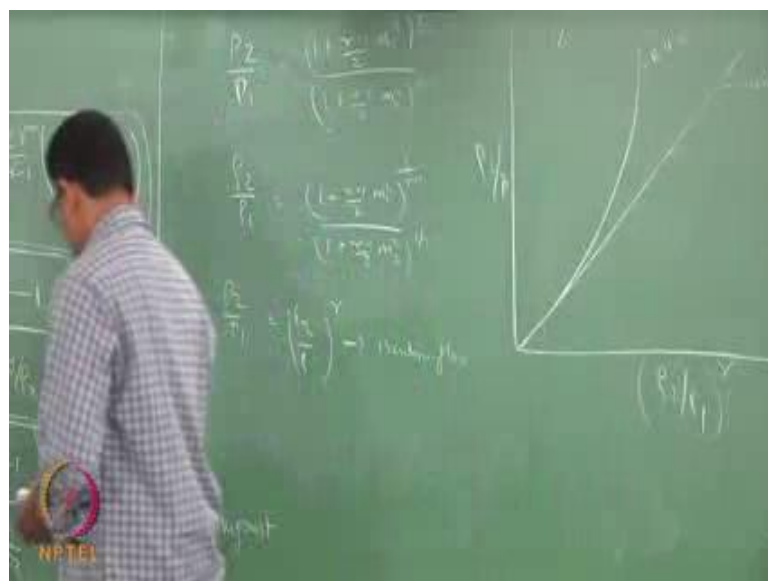
Student: (Refer Time: 22:54).

ρ_2 by ρ_1 ; Rankine - Hugoniot equation; this is pronounced as Hugoniot, Hugoniot equation. Now, here, in this process, we have not assumed anything that is irreversible; all that $x_0 > 1$, $x_0 < x$ equals $x_0 > y$, is the only condition that we had taken, which is no heat, no heat, no shaft work. Then, that includes your $T_0 < x$ equals $T_0 > y$. Then, it is the momentum equation, and continuity equation. So, energy equation and the perfect gas equation, and the momentum equation and the continuity equation is what we have used. So, if it is an isentropic flow, you know how P_1 and P_2 vary with respect to ρ_1 and ρ_2 . So, P_1 and P_2 vary with respect to.

How does it vary? So, if it is an isentropic flow, in order to get P_1 and, P_1 by P_2 in terms of Mach number $1 + \frac{\gamma - 1}{2} M_1^2$, ρ_1 square to the power $\frac{\gamma}{\gamma - 1}$ by $\frac{\gamma - 1}{2} M_2^2$ square to the power $\frac{\gamma}{\gamma - 1}$ by $\frac{\gamma - 1}{2} M_2^2$ square to the power $\frac{\gamma}{\gamma - 1}$. Likewise, your ρ_2 by ρ_1 is $1 + \frac{\gamma - 1}{2} M_1^2$ square to the power $\frac{1}{\gamma - 1}$, right; divided by $1 + \frac{\gamma - 1}{2} M_2^2$ square to the power $\frac{1}{\gamma - 1}$; or, your P_2 by P_1 is ρ_2 by ρ_1 to the power, this to the power $\frac{\gamma}{\gamma - 1}$ is, to the power $\frac{\gamma}{\gamma - 1}$.

So, this is your isentropic flow; because of non-isentropic, you have this relation across the shock. So, if I plot P_1 by P_2 versus ρ_1 by ρ_2 .

(Refer Slide Time: 27:26)

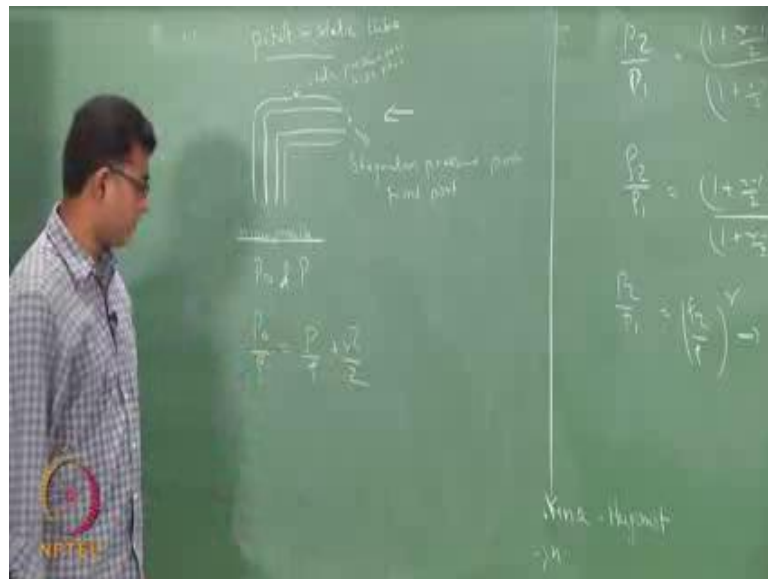


So, let us take ρ_2 by ρ_1 , P_2 by P_1 . If I, instead of y and x , I put 2 and 1, then, I would get something like this, for Rankine – Hugoniot, something like this for isentropic flow; this is Rankine - Hugoniot equation. So, there is a difference between these two ratios, because of the shock; that is what it.

Student: Just in case a straight line.

This will not be a straight line, because, there is an exponent here. So, this line is some curve with, you know, slope changing at every location. Or, what you can do is, you can plot this; and this, if I plot this versus this, this would be a straight line with 45 degrees, which means that, for every P_2 by P_1 , which is equal to this quantity; whereas, in Rankine- Hugoniot, it will be something else. Now, we will see what happens to the scenario, when we have a pitot tube kept in a supersonic flow.

(Refer Slide Time: 29:38)

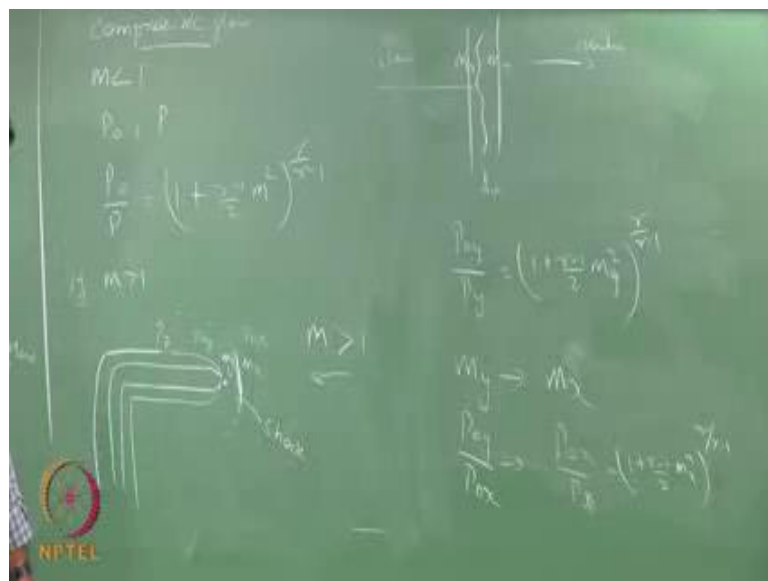


So, pitot tube is something that measures your, a stagnation pressure, at this point, if the flow is in this direction, flow comes, and then, takes, the streamlines take deviation like this. So, there is a stagnation point here. The velocity is about 0 isentropically; hence, you measure your stagnation pressure here. And, somewhere here, if you have a port, that will pressure your static, our static pressure, which is also, which is something different called the pitot static tube; we will discuss that, may be at the end of the class. But now, for the time being, we are going to measure only the stagnation pressure.

So, in an incompressible flow, or, let us discuss pitot static tube right away. So, this, I will modify. So, I have a tube that is going to measure my stagnation pressure. It is also going to measure a static pressure at this surface. So, this is my static pressure port, or, simply, side port. This is my stagnation pressure port, or simply, the front port. So, if I say front port, it essentially measures stagnation pressure, and side port of the pitot static tube measures the static pressure. So, the flow come and goes to a 0 value at this particular point, because, that is a stagnation point there, as you would have discussed in your fluid mechanics course.

So, there is a stagnation point, and there is a static port here. So, I am not going to discuss the pitot, working of the pitot tube, but just these 2 points will be stressed. So, I measure P_0 and P , which is my static. If it is an incompressible flow, I use Bernoulli's equation to find my velocity. So, I pressure my, I measure my stagnation pressure, and the static pressure, I get my velocity of the fluid.

(Refer Slide Time: 33:06)



If the flow is compressible, with the case subsonic flow, m less than 1, then, that particular port will measure your P_0 and P , and you know, the P_0 by P is related to your Mach number, which is $1 + \frac{\gamma - 1}{2} m^2$ to the power $\frac{\gamma}{\gamma - 1}$, from which you can find your Mach number. So, you know, if you measure your stagnation pressure, you measure your static pressure, then, you get your Mach number, which can be converted to your velocity, if you know the temperature.

Now, if m is greater than 1, then, we have a problem, which is the following. So, I have this tube, then, I also measure my static pressure. So, it measures something here; but, the problem here is, I have a flow which is m greater than 1, and it suddenly sees an obstruction. So, what does it do? There is a shock that is generated. So, there will be a shock that is generated. It would be something called a shock; for the time being, let us assume this to be a normal shock.

So, there will be a shock that is generated in front of your measuring device, which is your pitot static tube. Now, if you go back to your description of a shock, to the schematic which I had described sometime back, I have a shock here, m_x , m_y ; and, this is isentropic; this is isentropic; only the shock is non-isentropic. Same thing happens here; the flow comes, and upto this point is isentropic, and after this also is isentropic; only the shock is non - isentropic.

So, you can relate all the properties up to the shock using isentropic relation, and after the shock, and further downstream, using your isentropic relation; this jump, you have to use shock relation. So, I have my m_x coming here, and, I have an m_y here; I have a P_0_x here; I have a P_0_y here. This P_0_y is measured at your pitot tube. So, this point which measures some pressure is your P_0_y ; not P_0_x .

So, now, if you know P_0_y and your P_y , if I know my P_0_y and P_y , I can get my Mach number, m_y square to the power γ by γ minus 1. So, I know my P_0_y ; I know my P_y , from which I find my m_y first. Now, if I know my m_y , I can get my, I can get my m_x . If I know my m_x , I can also get P_0_y by P_0_x . I repeat; I get, I measure P_y and P_0_y , because, there is a shock sitting here. So, whatever I measure in the static, pitot static tube is P_0_y and P_y . The moment I get this, I can get my Mach number after the shock.

The moment I get Mach number after the shock, I have a relation which will tell me what is the Mach number before the shock, I can get a Mach number before the shock. So, I get a Mach number before the shock. If I get a Mach number before the shock, then, I can also get the pressure ratios, stagnation pressure ratios, from which I can find P_0_x . If I know P_0_x , then, I can use P_0_x by P_x relation in terms of Mach number to find my pressure before the shock. So, the static pressure before the shock is also obtained, once I have this. And, once you have the Mach number before the shock, you can find,

eventually, find all the quantities, if you know the temperature at that location; the velocity can be easily found out.

So, the point here is, in supersonic flow, there will be a shock sitting, which would actually make the measurement not as direct as this, but, there will be one more step involved in finding the Mach number.

So, what we have seen is the relation in normal shock, the properties, how the properties change across the normal shock, and how is it useful, when it comes to something like a measurement here, or in a, say, for example, the inner diverging section in the c d nozzle, if there is a shock standing somewhere in the diverging section, the properties across the shock is going to be different, which actually takes it to a, some other Mach number than the designed Mach number. So, all those conditions, we will understand this better by doing some numerical problems.