

Conduction and Convection Heat Transfer
Prof. S.K. Som
Prof. Suman Chakraborty
Department of Mechanical Engineering
Indian Institute of Technology – Kharagpur

Lecture - 39
Energy Conservation Equation – II

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The image shows a handwritten derivation on a green chalkboard. The top equation is:

$$\frac{\partial}{\partial t}(\rho e) + \frac{\partial}{\partial x_j}(\rho u_j e) = \dot{q}'' - \frac{\partial q_j''}{\partial x_j} + \frac{\partial}{\partial x_j}(\tau_{ij} u_i) + b_i u_i$$

An arrow points from the left side of this equation to the next line, which is:

$$\rho \frac{\partial e}{\partial t} + \rho u_j \frac{\partial e}{\partial x_j} + \rho u_j \frac{\partial e}{\partial x_j} = \dot{q}'' - \frac{\partial q_j''}{\partial x_j} + \frac{\partial}{\partial x_j}(\tau_{ij} u_i) + b_i u_i$$

Then, the equation is rearranged to separate the mechanical energy terms:

$$= \rho \left[\frac{\partial e}{\partial t} + u_j \frac{\partial e}{\partial x_j} \right] + e \left[\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x_j}(\rho u_j) \right]$$

The term $e \left[\frac{\partial \rho}{\partial t} + \frac{\partial}{\partial x_j}(\rho u_j) \right]$ is crossed out with a diagonal line and labeled as ϕ (continuity). The final simplified equation is:

$$= \rho \left[\frac{\partial e}{\partial t} + u_j \frac{\partial e}{\partial x_j} \right]$$

In heat transfer we are interested in conservation of thermal energy. So, if you have conservation of total energy how can you get conservation of thermal energy from there. So, in the total energy, you have mechanical energy plus thermal energy. So, you have to write an expression for conservation of mechanical energy, mechanical means kinetic and potential.

So, what essential is this, this is kinetic plus potential plus internal. So, we have to write something in terms of the internal energy only. So, we have to subtract the equation that involves kinetic and potential energy. That is the mechanical energy equation. Now, one has to be careful in one thing. In this equation when we write e actually we are including within it the internal energy, kinetic energy but not potential energy, why?

Because we are already having an expression for work done by body force. So, if we now again write potential energy here, then that will duplicate the same effect. So, either you do not write the work done by body force but put the corresponding potential energy here or you do not put the potential energy here but you put the work done by body force. So, here we have preferred to write the work done by body force.

Therefore, we will not include the potential energy here. Now, work done by potential energy if it is not included, if the energy, if the potential energy is not included that means work done by body force is already included then, this equation will give you the mechanical energy. So, this equation has given you the total energy where you have somewhere the potential energy effect and internal energy and kinetic energy.

We want to get another equation which will give us the mechanical energy, mechanical energy equivalent of the one. So, mechanical energy that means the kinetic and the potential, potential will be in form of work done by body force. So, which equation can give us.

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The image shows a handwritten derivation on a green chalkboard. It starts with the Navier equation: $\rho \left[\frac{\partial u_i}{\partial t} + u_j \frac{\partial u_i}{\partial x_j} \right] = \frac{\partial \tau_{ij}}{\partial x_j} + b_i$. This is then multiplied by u_i to get: $\rho \left[\frac{\partial}{\partial t} \left(\frac{u^2}{2} \right) + u_j \frac{\partial}{\partial x_j} \left(\frac{u^2}{2} \right) \right] = \left(\frac{\partial \tau_{ij}}{\partial x_j} u_i \right) + b_i u_i$. The final result is: $\rho \left[\frac{\partial}{\partial t} \left(\frac{u^2}{2} \right) + u_j \frac{\partial}{\partial x_j} \left(\frac{u^2}{2} \right) \right] = \dot{Q} - \frac{\partial \tau_{ij}}{\partial x_j} u_i + \tau_{ij} \frac{\partial u_i}{\partial x_j}$.

So, Navier equation, what was it? It was an equation of momentum balance. So, in a conservative form, rho... This equation why we start with a Navier equation but not the Navier-Stokes equation is that, we want to generalize it for all types of fluids. So, once we commit ourselves to Navier-Stokes equation, we are restricted to Newtonian and Stokesian fluid.

But when we write the Navier equation or the Cauchy equation whatever we are not bothered about what is the special type of fluid that we are talking about for any fluid you can use this. So, now can you convert this momentum balance equation into an energy equation? Only a small trick you have to play to do that. See, you want to convert it into mechanical energy equation.

That is kinetic energy and potential. Potential will be work done because of the body force. So, how can you convert this to kinetic energy? You multiply this with U_i . So, then it becomes, $U_i \text{ del } U_i \text{ del } T$, that is $\text{del del } T$ of U_i square by 2. U_i square by 2 is kinetic per unit mass. So, you multiply by U_i . So, ρ - Let us say this is equation number one. This is equation number two.

Now you subtract equation number two from equation number one. What you get? You will get an equation that governs the change in internal energy only because this is kinetic, this contribution is due to potential and you will see that when you subtract this e includes what? Internal and kinetic. So, internal and kinetic minus kinetic will become only internal. So, if you subtract, equation two from equation one.

We are using the symbol i for internal energy, the reason is that we have already used u for velocity. So, just to not to confuse with the symbol, whatever symbol u that you have used for internal energy in thermodynamics that same is i . If you subtract from this term, this term, then what will be left? Use the product rule, one time will be cancelled. What will be left is $\tau_{ij} \text{ del } u_i \text{ del } x_j$.

So, see, the change in internal energy, this is the total rate of change of internal energy, the rate of change of internal energy may be due to volumetric heat generation, surface heat transfer but also due to stresses in the flow field. So, this part, this contribution comes directly from fluid mechanics. So, $\tau_{ij} \text{ del } u_i \text{ del } x_j$ physically what does it mean. So, this is like a heat source. That means this will result in heating.

You may ask that, why can't it result in cooling. We will give an answer to that, that it will trivially give rise to heating and not cooling. So, it is something like this. Let us say that you are shearing the flow. If you are shearing the flow then two layers of the fluid are sliding against each other, so as if you are rubbing your palms. So, if you are rubbing your palms, this work is not utilized.

But this work is irreversibly converted into internal energy through viscous dissipation. So, this is called as viscous dissipation where there is an irreversible conversion of work to intermolecular form of energy. So, if you rub your palms you cannot get any work out of that. Simply it will generate heat. So, it is effectively like a heat source. But now the question

always come that well. We are claiming that it is like a heat source, but where is the guarantee that it will always give rise to heating and not cooling.

So, we have to show that whatever is the effect associated with it, that effect associated with it because of viscous part of the stress, so remember that this τ_{ij} also has hydrostatic part of the stress. So, we will isolate the hydrostatic part of the stress and consider only the deviatoric part which is the viscous part and show that the viscous part contribution to this term is always positive. That means it will give rise to heating and not cooling.

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So, let us see the term τ_{ij} . So, we assume homogeneous, isotropic and Newtonian fluid and also Stokesian fluid. So, τ_{ij} is minus $p \delta_{ij}$, plus $\lambda \frac{\partial u_k}{\partial x_k} \delta_{ij}$ plus $\mu \left(\frac{\partial u_i}{\partial x_j} + \frac{\partial u_j}{\partial x_i} \right)$. This we discussed in quite some length. So, $\tau_{ij} \frac{\partial u_i}{\partial x_j}$. What will be the hydrostatic component? Remember δ_{ij} is equal to one when j is equal to i . So, it will become minus p into $\frac{\partial u_i}{\partial x_i}$. What is this?

This is minus p multiplied by $\text{div } \mathbf{v}$. So, this is like the $p \, dv$ work in thermodynamics because the divergence of the velocity vector gives the volumetric change. So, this is analogous to $p \, dv$. So, we will isolate that with the viscous component, remember this λ is equal to minus two third of μ . So, plus λ , what is $\frac{\partial u_k}{\partial x_k}$ multiplied by δ_{ij} , into $\frac{\partial u_i}{\partial x_i}$?

See, $\frac{\partial u_k}{\partial x_k}$ is $\frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2} + \frac{\partial u_3}{\partial x_3}$. δ_{ij} into $\frac{\partial u_i}{\partial x_i}$ is $\frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2} + \frac{\partial u_3}{\partial x_3}$, so

it becomes $\mu \left(\frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2} + \frac{\partial u_3}{\partial x_3} \right)^2$ This term, this into $\mu \frac{\partial u_i}{\partial x_j}$, plus the remaining terms.

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The chalkboard shows the following derivations:

- For $i=1, j=2$ and $i=2, j=1$: $\mu \left[\left(\frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1} \right) \frac{\partial u_1}{\partial x_2} + \left(\frac{\partial u_2}{\partial x_1} + \frac{\partial u_1}{\partial x_2} \right) \frac{\partial u_2}{\partial x_1} \right] \rightarrow \mu \left[\frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1} \right]^2$
- For $i=2, j=3$ and $i=3, j=2$: $\mu \left[\frac{\partial u_2}{\partial x_3} + \frac{\partial u_3}{\partial x_2} \right]^2$
- For $i=3, j=1$ and $i=1, j=3$: $\mu \left[\frac{\partial u_1}{\partial x_3} + \frac{\partial u_3}{\partial x_1} \right]^2$
- For $i=1, j=1$, $i=2, j=2$, and $i=3, j=3$: $2\mu \left[\left(\frac{\partial u_1}{\partial x_1} \right)^2 + \left(\frac{\partial u_2}{\partial x_2} \right)^2 + \left(\frac{\partial u_3}{\partial x_3} \right)^2 \right]$

So, here we will write several terms, first i equal one, j equal to two, i is equal to two, j equal to one. This is one set. Then i equal to 2 k equal to 3, i equal to 3, j equal to 2. This is one set. Then i equal to 3, j equal to 1, i equal to one, j equal to 3. This is another set. Then i is equal to one, j is equal to one, i is equal to 2, j is equal to 2, i is equal to 3, j is equal to 3. Total 9, right? So, first we will consider this set.

So, $\mu \frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2}$ multiplied by $\frac{\partial u_1}{\partial x_1}$. Okay, let me write one term then it will be easier to write. Let me write the full one term. So, this is $\frac{\partial u_1}{\partial x_1}^2 + \frac{\partial u_2}{\partial x_2}^2 + 2 \frac{\partial u_1}{\partial x_2} \frac{\partial u_2}{\partial x_1}$. So, it is like a plus b whole square formula, right? So, this is μ multiplied by... Okay. So, can you physically interpret this. This is like $\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}$, rate of deformation.

So, square of rate of deformation times the viscosity is actually giving rise to a heating effect. Okay? So, how the kinematics of the fluid is getting converted into heat transfer effect, that is something which is not so intuitive but you can see from here. So, this is μ multiplied by... and this is ... Now what will be this. i equal to one, j equal to one means? Right? So, now it is trivially clear that these 3 terms together are positive.

Of course, this term is positive, but these term together with these term whether it is positive negative. So, now it is trivially clear that these three terms together are positive. Of course, this term is positive, but these term together with these term whether it is positive or negative it has to be adjudged, why because lambda is negative. Lambda is not positive; lambda is minus two third of mu where mu is positive.

So, we will try to show that these term together with these term eventually becomes positive. So, let us try to do that exercise.

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The image shows a chalkboard with handwritten mathematical expressions. At the top, there is a term $\frac{2}{3}\mu \left[\frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2} + \frac{\partial u_3}{\partial x_3} \right]^2 + 2\mu \left[\left(\frac{\partial u_1}{\partial x_1} \right)^2 + \left(\frac{\partial u_2}{\partial x_2} \right)^2 + \left(\frac{\partial u_3}{\partial x_3} \right)^2 \right]$. Below this, the expression is expanded to $\frac{2\mu}{3} \left[3 \left(\frac{\partial u_1}{\partial x_1} \right)^2 + 3 \left(\frac{\partial u_2}{\partial x_2} \right)^2 + 3 \left(\frac{\partial u_3}{\partial x_3} \right)^2 - \left(\frac{\partial u_1}{\partial x_1} \right)^2 - \left(\frac{\partial u_2}{\partial x_2} \right)^2 - \left(\frac{\partial u_3}{\partial x_3} \right)^2 - 2 \frac{\partial u_1}{\partial x_1} \frac{\partial u_2}{\partial x_2} - 2 \frac{\partial u_1}{\partial x_1} \frac{\partial u_3}{\partial x_3} - 2 \frac{\partial u_2}{\partial x_2} \frac{\partial u_3}{\partial x_3} \right]$. The final part of the derivation shows $\frac{4\mu}{3} \left[\left(\frac{\partial u_1}{\partial x_1} \right)^2 + \left(\frac{\partial u_2}{\partial x_2} \right)^2 + \left(\frac{\partial u_3}{\partial x_3} \right)^2 - \frac{\partial u_1}{\partial x_1} \frac{\partial u_2}{\partial x_2} - \frac{\partial u_1}{\partial x_1} \frac{\partial u_3}{\partial x_3} - \frac{\partial u_2}{\partial x_2} \frac{\partial u_3}{\partial x_3} \right]$ with a bracket underneath and the word "positive" written below it.

Lambda is equal to minus two third of mu. So, if you take 2 mu as common or 2 mu by 3 as common. So, these essentially becomes like, if you take 2 as common, so 3 minus 1 is 2. If you take 2 as common, then it is 4 mu by 3. So, this is a positive term. Like a square plus b square plus c square minus ab minus bc minus ca. That is half multiplied by a minus b whole square plus b minus c whole square plus c minus a whole square.

So, this is positive, right? So, our conclusion is that, the deviatory component of these will give rise to trivially heating and not cooling because this entire term, this plus whatever we have written, sum total of this is a positive term and that term is known as viscous dissipation. Now task is not yet complete because we have written the equation in terms of internal energy.

But in experiment we do not measure internal energy, we do not have any internal energy meter. We have say thermometer to measure temperature. So, we have to convert the internal

energy based equation to a temperature based equation which will give us a governing equation for temperature which is the so-called energy equation in convection, that we will take up in the next lecture, thank you.