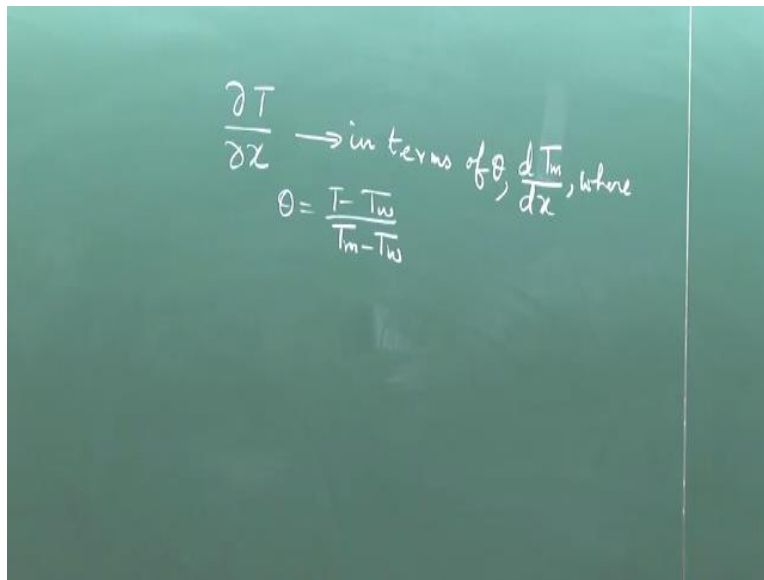


Conduction and Convection Heat Transfer
Prof. S.K. Som
Prof. Suman Chakraborty
Department of Mechanical Engineering
Indian Institute of Technology – Kharagpur

Lecture - 32
Internal Force Convection - II

In the previous lecture, we were discussing about the concept of hydrodynamically and thermally fully developed flow. We discussed about the consequences of hydrodynamic and thermally fully developed flow in terms of expressing a non-dimensional temperature and a non-dimensional velocity. Now, we have seen the term $\delta T \delta x$ can be expressed in terms of θ and $d T_m d x$, where θ is T minus T_{wall} by T_m minus T_{wall} .

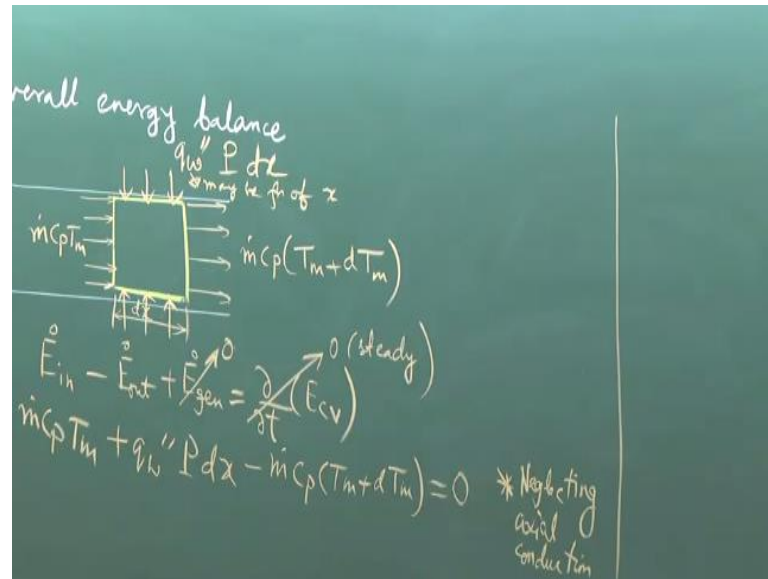
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$$\frac{\partial T}{\partial x} \rightarrow \text{in terms of } \theta, \frac{dT_m}{dx}, \text{ where}$$
$$\theta = \frac{T - T_w}{T_m - T_w}$$

So, for constant wall temperature it is θ into $d T_m d x$ and for constant wall heat flux, it is just $d T_m d x$ without involving θ that much we discussed in the previous class. Now, irrespective of the boundary condition, this parameter is always involved. So, one thing we have been successful in reducing the dimensionality of the problem that when we are expressing $\delta T \delta x$ in terms of $d T_m d x$.

We are converting like a 2-dimensional or a 3-dimensional problem to (0) (02:37) 1-dimensional problem. But the question is how to calculate this dT_m/dx , so for that we will make an overall energy balance.

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So, let us say that this is the channel and we take a control volume like this. Across this control volume, we are going to write an energy balance, so when we are going to write an energy balance, there is some energy that is entering. Let us say that this length is dx . There is some energy that is leaving, so what is the energy that is entering in terms of the bulk mean temperature.

The bulk mean temperature is an equivalent temperature which would have existed uniformly across the cross section to make the flow of the same energy as that of the actual case. So, if T_m is the bulk mean temperature then what is the rate at which energy is transferred across this section, it is $\dot{m} H$, $\dot{m} H$ is $\dot{m} C_p T_m$. We are considering incompressible fluid and let us say this one is $\dot{m} C_p T_m$ plus dT_m .

Then, we have heat transfer at the wall, what is the heat transfer at the wall? If the wall heat flux is q_w'' this may be a constant or it may be a variable. Let us say that the wall heat flux is q_w'' . So, what is the heat transfer rate across the wall. Let us say this is a circular

pipe of radius r . So, if the wall heat flux is q'' , what is the rate of heat transfer? q'' into $2\pi r dx$, right. $2\pi r dx$ is the surface area of the pipe.

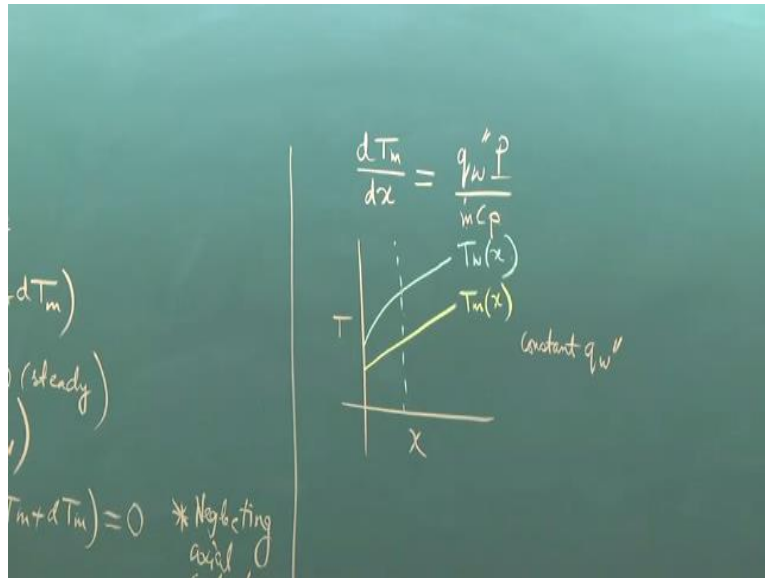
So, $2\pi r dx$ is nothing but perimeter into dx . So, to generalize it for any section, we will write it as q'' into perimeter into dx . So, we can write for this control volume, rate of energy in minus rate of energy out plus rate of energy generated. So, this is zero because it is steady flow and steady state. Rate of energy generation is zero that we are not considering in this problem. If there is some rate of energy generation, we can accommodate this in this formulation very easily.

Rate of energy in is $\dot{m} C_p T_m$ plus $q'' P dx$. What is the rate of energy out $\dot{m} C_p T_m$ plus dT_m . Now, this equation is not perfect, but it has an approximation. My question is what is that approximation that is there in this equation. This is not exactly correct. "Professor - student conversation starts" Yes, no, no, no, it is not taken as constant. "Professor - student conversation ends".

Even if it is a function of x it is true because we have taken a small element over which it will be constant. Q'' , this may be function of x . So, in the axial direction we have considered that there is heat transfer due to fluid flow, but we have not considered that there is heat transfer due to axial conduction. So, we have neglected axial conduction, so ideally, we should have taken a heat flux minus $K dT_m/dx$ here and minus $K dT_m/dx$ at $x + dx$ here.

So, we had not taken any axial conduction. So, we have neglected axial conduction and that is valid in many practical problems where advection is much more dominating than axial conduction. So, this very importantly is valid neglecting axial conduction.

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So now from this equation, you can write $d T_m / d x$ is equal to $q_w'' P$ by $\dot{m} C_p$. Now, can you tell from here that $d T_m / d x$ is constant if wall heat flux is constant, yes or no, yes. Because for steady flow rate is \dot{m} is constant, we assume that the properties are constant then the perimeter of the cross section, it is the constant for the geometry that we are considering. So, if the wall heat flux is constant then the $d T_m / d x$ is constant.

So, we had earlier shown that for the constant wall heat flux, $\Delta T / \Delta x$ is equal to $d T_m / d x$ is equal to $d T_w / d x$ is equal to constant, ok. But not for any general constant wall heat flux. For constant wall heat flux and thermally fully developed flow, ok. So now, we come to an inference that this is the constant for constant wall heat flux not just in the thermally fully developed flow region, but about the entire region.

Because this overall energy balance theory is very general it does not take into account whether it is thermally fully developed or not. It is just the simple energy balance using the definition of bulk mean temperature that is converting the multidimensional problem to 1 dimensional problem, but it does not take into account whether it is thermally fully developed or not. So, the conclusion is that even if the flow is not thermally fully developed for constant wall heat flux $d T_m / d x$ is the constant, but this is still an approximation.

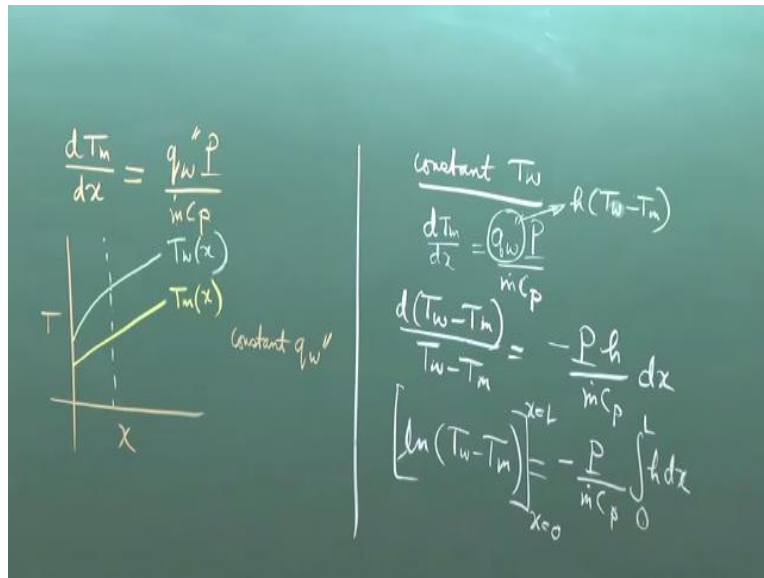
Because it has neglected axial conduction. If you do not neglect axial conduction that is not true. Whereas this is exactly true, does not matter whether axial conduction is there or not, ok. So, we can make a graph of say T versus x for constant wall heat flux. So, when we write plot T versus x , we essentially want to plot basically T_m and T_{wall} because we are now converting it to 1 dimensional problems where the functions of x are T_m and T_{wall} .

So, let us say that T_{wall} is greater than T_m now what is the graph of T_m versus x , it will be a straight line because $d T_m / d x$ is a constant assuming that axial conduction is negligible it is a single straight line throughout otherwise in the thermally developing region this may not be a constant, so it may be a curve and then in the thermally fully developed region, it will be a constant that formula.

So, let us make as sketch of this is T_m versus x assuming that the wall is heated, now what will be the graph of T_{wall} versus x assuming T_{wall} is greater than T_m . So first that graph should be above this if it is heated then how will the graph look like, see $d T_{wall} / d x$ is a constant only if it is a thermally fully developed flow. So, let us say that the flow becomes thermally fully developed from here.

So, up to this it will be a curve beyond this it will be a straight line. What straight line, straight line parallel to this because $d T_m / d x$ is equal to $d T_{wall} / d x$. So, the slopes of these 2 straight lines are equal in the thermally fully developed region. So, this is T_{wall} as a function of x , ok.

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Now, this is the case of constant wall heat flux. Let us study the qualitative behavior of constant wall temperature. So, first we studied quantitatively and then we show it in a plot. So, the case of constant wall temperature. So, dT_m/dx is equal to $q''_w P$ by $\dot{M} C_p$. How can you write q''_w in terms of the heat transfer coefficient h , q''_w equal to what h , h into T_m minus T_w or T_w minus T_m .

So, when you are considering q''_w , look at this figure, here the sign convention is the heat transfer is from the wall to the fluid that means T_w minus T_m , right. If it is the opposite the sign itself will take care. So, in place of q''_w , we will take h into T_w minus T_m , so we can write $d(T_w - T_m) / (T_w - T_m)$ is equal to minus P into h by $\dot{M} C_p$ dx .

So, what we have done is, we have written dT_m/dx as T_w of d/dx of T_m minus T_w because T_w is a constant it does not matter whether we take it inside the derivative or not. It will make no difference because dT_w/dx is 0 for constant T_w , ok and then this minus sign is observed because we have converted T_m minus T_w to T_w minus T_m .

So, it is of the form (18:21) it will be log of this, so $\ln(T_w - T_m)$ is equal to minus P by $\dot{M} C_p$ integral of $h dx$ from x equal to 0 to x equal to L , where L is the total length of

the channel or the pipe. This of course we have to put a definite limit from x equal to 0 to x equal to L . x equal to 0 is the starting and x equal to L is the ending of the channel.

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Handwritten derivation on a chalkboard:

$$\Delta T = T_w - T_m$$

$$\ln \frac{\Delta T_L}{\Delta T_0} = -\frac{PL\bar{h}}{\dot{m}C_p}$$

$$\dot{Q} = \dot{m}C_p(T_{m,x=L} - T_{m,x=0})$$

$$= -\dot{m}C_p \left[(T_w - T_m)_{x=L} - (T_w - T_m)_{x=0} \right]$$

$$\dot{m}C_p = \frac{\dot{Q}}{\Delta T_L - \Delta T_0}$$

$$\ln \frac{\Delta T_L}{\Delta T_0} = \frac{-A\bar{h}(\Delta T_L - \Delta T_0)}{\dot{Q}} \Rightarrow \dot{Q} = A\bar{h} \frac{\Delta T_L - \Delta T_0}{\ln \frac{\Delta T_L}{\Delta T_0}}$$

Graph on the right shows temperature T vs position x . The wall temperature T_w is constant, and the fluid temperature T_m increases from $T_{m,0}$ to $T_{m,L}$. The temperature difference ΔT decreases from ΔT_0 to ΔT_L . The formula is labeled as LMTD (logarithmic mean temp diff).

So, we can write \ln of ΔT_L by ΔT_0 where ΔT is $T_{\text{wall}} - T_m$ is equal to, so what is this, this is the average heat transfer coefficient times L . So, minus P into L into h average by $\dot{m} C_p$. Now what is the total heat transfer rate, \dot{Q} . Let us say you transfer some heat from the wall to the fluid, so how can you measure what heat has been transferred from the wall to the fluid.

Let us say you are doing an experiment, so the situation is that you have some heat let us say you have a heating coil at the wall, you have a heater at the wall, you are transferring heat to the fluid. So how do you measure that what heat actually has been transferred to the fluid. So, you can measure the thermal energy at the inlet and you can measure the thermal energy at the outlet. So, the difference between these 2-thermal energy is the heat that is supplied from the wall.

So, you can say \dot{Q} is nothing but $\dot{m} C_p$ into T_m at x equal to L minus T_m at x equal to 0, right. This is just simple energy balance. So, this heat transfer this is nothing but the change in enthalpy. So, if you write the first law of thermodynamics for the control volume this is what you will get as a heat transfer. There is no (\dot{Q}) (22:18) done in this case. So, all the heat that is supplied is used to change this $\dot{m} C_p$ into T that is \dot{m} dot into h .

Of course, we neglect the changes in kinetic energy and potential energy. So, you can write these as $\dot{m} C_p$, see we can write this because this T_{wall} is the constant so it is just adding and subtracting the same constant from the 2 terms, but why we have done is because this is the ΔT at x equal to L and this is the ΔT at x equal to 0 . So, we can write $\dot{m} C_p$ is equal to \dot{Q} minus \dot{Q} by $\Delta T L$ minus $\Delta T 0$.

So, you can substitute that here and write \ln of $\Delta T L$ by $\Delta T 0$, perimeter into length is what the total surface area, so this A is not cross-sectional area this is the surface area that is $2\pi r$ into L for a circular pipe of length L . So, minus area into h then in place of M this one you can write \dot{Q} by $\Delta T L$ minus $\Delta T 0$. So, \dot{Q} is equal to A into h average into $\Delta T L$ minus $\Delta T 0$ by \ln of $\Delta T L$ by $\Delta T 0$.

So, you can see that normally what is $\dot{Q} A$ into h into ΔT , the temperature difference between the 2 systems across which the heat transfer is taking place here one system is the wall another is the fluid. So, it would have been ideally A into h into T_{wall} minus T_m , but it is not A into h into T_{wall} minus T_m . It is h into these, the reason is that T_{wall} minus T_m is not a constant. It is continuously varying with x .

So, if you make a plot of say T_{wall} so T versus x . So, for constant T_{wall} , so this is constant T_{wall} . For constant T_{wall} , what is T_{wall} as the function of x , this is T_{wall} constant. What about the T_m , see look at this equation $d T_m / dx$, so d of T_{wall} minus T_m by T_{wall} minus T_m is this. So that means these T_{wall} minus T_m where is with e to the power minus x , right. Because \log of these varies with minus x so this T_{wall} minus T_m varies with e to the power minus x .

So, that means T_m has an exponential variation with x . So, for constant wall temperature you get assuming that the wall is again heated, this is T_m as the function of x . This is an exponential curve. Why exponential curve you can look from the analytical expression. So, if you want to say that the rate of heat transfer is equal to A into h into ΔT . Now, the ΔT is different. ΔT at x equal to 0 is this, then it becomes this, it becomes this, it becomes this, like this.

So, ΔT is continuously varying, so you require some equivalent average ΔT . We have shown here by your calculation that average ΔT is not $\frac{\Delta T_L + \Delta T_0}{2}$, not, it is $\frac{\Delta T_L - \Delta T_0}{\ln \Delta T_L / \Delta T_0}$. This is called as LMTD or logarithmic mean temperature difference. so, why do we require a logarithmic mean temperature difference because the temperature difference itself is continuously varying.

So, what would be the logarithmic mean temperature difference, if the temperature difference is constant let us say that T_{wall} is the constant and T_m is the constant, then what would be the logarithmic mean temperature difference or T_{wall} like this T_m like this. This difference is the constant. Then what would be the logarithmic mean temperature difference, no. The logarithmic mean temperature difference physically represents the equivalent temperature difference.

So, if this difference is a constant this constant value itself is the logarithmic mean temperature difference. See if or anything you can do it mathematically. How do you do it mathematically, see if a constant temperature difference is there. Then, it is ΔT_L equal to ΔT_0 . So, it is $\frac{0}{\ln 1}$, $\frac{0}{0}$. You can use $(\frac{0}{0})$ (30:16) rule to find out what is the logarithmic mean temperature difference.

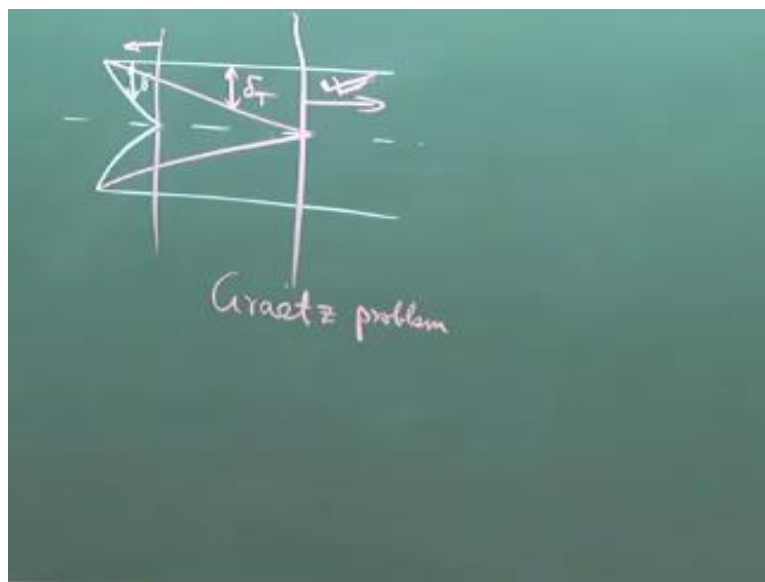
But my question is why should we do that from physical understanding we understand that the logarithmic mean temperature difference is the effective average temperature difference between the 2 fluids. If the 2 fluids temperature are continuously varying, you have to judiciously use that formula, but if the temperature difference itself is the constant then the logarithmic mean temperature difference will be that constant itself because it is the equivalent temperature difference.

So, we can say that to summarize so sometimes in short form this is written as ΔT_{LM} , logarithmic mean temperature difference, so this is just a short form of doing it, but it is the very important parameter and in one of the latter chapters on heat exchange here, what we learn which is a very important topic for practical engineering applications. This terminology will come over and again. So please make a note this very important terminology.

So, we have discussed about the situations of the constant wall heat flux and constant wall temperature somewhat qualitatively that is how the wall temperature and how the bulk mean temperature varies with x , but question is what is the rate of heat transfer that is what is the Nusselt number. So, we need to discuss about what is Nusselt number for thermally fully developed flow.

So, there may be situations when the flow is hydrodynamically fully developed, but it is not thermally fully developed like let me give you an example.

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So, this is hydrodynamic boundary length and this is thermal boundary length. So, if you have a fluid say of value of high Prandtl number, then what will happen, this δ will be significantly more as compared to δT , so the hydrodynamic boundary layer will grow very fast. So, the hydrodynamic boundary layers will merge here and the flow will become hydrodynamically fully developed.

So, from here to here, there is a region when the flow is hydrodynamically fully developed, but not thermally fully developed, right. So, that kind of problem is known as Graetz problem. This kind of problem is actually in the (33:32) advanced level of convective heat transfer. So, we will not come into that here. There is another type of problem where the flow may be both hydrodynamically and thermally developing like this region.

But in this particular course, we will be concentrating on this region where the flow is both hydrodynamically and thermally fully developed and we will evaluate the Nusselt number for that cases.

(Refer Slide Time: 34:58)

Hydrodynamically & thermally fully developed
flow through parallel plate channel

Case 1 Constant q_w''

Energy eq: $u \frac{d^2T}{dy^2} + v \frac{dT}{dy} = \alpha \left(\frac{d^2T}{dx^2} + \frac{d^2T}{dy^2} \right)$

$u \frac{dT_m}{dx} = \alpha \frac{d^2T}{dy^2}$

$\left(\frac{dT}{dx} = \text{const} \right)$

So, hydrodynamically and thermally fully developed flow through parallel plate channel. Case 1, constant wall heat flux. So, what answer can we expect for these types of problems. What are you assured of when you are solving a problem of thermally fully developed flow. We are interested about the Nusselt number. The Nusselt number will be a constant that is what is expected.

So, if we do not come up with the constant Nusselt number that means something is wrong in our analysis or approach. So, the Nusselt number will be a constant and our objective will to evaluate the constant value. So, we have to keep in mind that see what earlier cases we have studied for evaluation of Nusselt number flow over flat plate. The Nusselt number was of the form of Reynolds number to the power something say half into Prandtl number to the power something say half or one third.

This kind of expressions we have seen. So, Reynolds number to the power m into Prandtl number to the power n . In general, for force convection, Nusselt number is actually of the form

Reynolds number to the power m into Prandtl number to the power n , but for hydrodynamically and thermally fully developed flow for internal forced convection that is the special case for Nusselt number becomes a constant.

But in general, if you have a situation, even the situation changes if the flow from laminar becomes turbulent, so then I mean you cannot work out those problems analytically you have to deal with sudden correlations based on experimental studies which engineers used for designing a system with the turbulent flow. I will come to 1 or 2 such expressions, but before that remember here we are assuming laminar flow, so constant wall heat flux.

Now let us write the governing equation, the energy equation. This is our governing equation. Now tell what will be the simplification. So, let us write it with different color may be, so in place of $\Delta T \Delta x$ for constant wall heat flux, we can write dT/dx , which is the constant, whatever this term, v is equal to 0 for hydrodynamically fully developed flow, so this will be 0. What about this term in the right-hand side, $\Delta T \Delta x$, $\Delta T \Delta x$ is what?

$\Delta T \Delta x$ is equal to constant for thermally fully developed flow with constant wall heat flux. Because $\Delta T \Delta x$ is constant, $\Delta^2 T \Delta x^2$ is 0. So, you can write $d^2 T/dx^2$ is equal to $\alpha d^2 T/dy^2$. Now, we will change these variables from T to θ , the nondimensional temperature which does not vary with x anymore for thermally fully developed flow.

Why we want to change the variable, because T is the function of both x and y , but for thermally fully developed flow θ is only a function of y , but not a function of x . So, θ is $T - T_{\text{wall}}$ by $T_m - T_{\text{wall}}$. So, you can write $\Delta T \Delta y$, remember $T_{\text{wall}} - T_m$ is the functions of x , so with respect to y derivative those are like constants.

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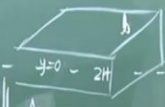
$$\theta = \frac{T - T_w}{T_m - T_w} \Rightarrow \frac{\partial T}{\partial y} = (T_m - T_w) \frac{d\theta}{dy} \quad (\theta \neq \theta(x) \text{ for TFDF})$$

$$\frac{\partial^2 T}{\partial y^2} = (T_m - T_w) \frac{d^2 \theta}{dy^2}$$

$$\frac{dT_m}{dx} = \frac{q_w' P}{\dot{m} C_p} = \frac{q_w'' P}{\rho \bar{u} A C_p}$$

$$\frac{dT_m}{dx} = \frac{q_w''}{h P \bar{u} C_p}$$

$$\frac{u}{u} \frac{q_w''}{h C_p} = \frac{k}{h C_p} (T_m - T_w) \frac{d^2 \theta}{dy^2}$$



$$\frac{A}{P} = \frac{b \cdot 2H}{2(b + 2H)}$$

$$= \frac{2H}{2(1 + \frac{2H}{b})}$$

$$= H \text{ as } \frac{H}{b} \rightarrow 0$$

So, similarly the second derivative and what is $d T_m / dx$, $d T_m / dx$ is what is that expression q double dash P by $\dot{m} C_p$, so q double dash in place of P , let us draw a parallel plate channel. Let us say that the width is b and height is $2H$ with the center line is y equal to 0. So, what is the perimeter, so before that let us write another step. What is \dot{m} , \dot{m} is ρu average into A into C_p by perimeter.

So basically, we have to calculate perimeter by area, so perimeter by area what is perimeter or let us calculate area by perimeter whatever. We will invert that area by perimeter what is that b into $2H$ by 2 into b plus $2H$, right. Now, we have to consider a limit value calculating the area by perimeter what is that limit, limit is that h by b tends to 0 because by definition the parallel plate channel the width is infinitely large.

So, what we will do is we will divide both numerator and denominator by b , so this becomes $2H$ by 2 into 1 plus $2H$ by b , so this will become h as h by b tends to 0, right. So, when h by b tends to 0, this term is zero, this term is 1 and this 2 get cancelled, so it becomes H . So, you can write $d T_m / dx$ is equal to q double prime into perimeter by area is 1 by $h \rho u$ average C_p . So, $u d T_m / dx$ is u by u average.

So, we are writing this equation, this differential equation $u d T_m / dx$ is u by u average into q double prime by $\rho H C_p$ is equal to α , α is K by ρC_p into T_m minus T_w into d

$\frac{d^2\theta}{dy^2} + \frac{u}{\alpha} \frac{h}{H} = 0$, ok. So, then ρC_p gets cancelled. In place of Q_{wall} what we can write Q_{wall} is $h(T_{wall} - T_m)$.

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The chalkboard contains the following equations:

$$\frac{d^2\theta}{dy^2} + \frac{u}{\alpha} \frac{h}{H} = 0$$

$$\bar{y} = \frac{y}{H}$$

$$\frac{d^2\theta}{d\bar{y}^2} + \frac{u}{\alpha} \left(\frac{hH}{K} \right) = 0$$

$$\frac{3}{2}(1 - \bar{y}^2) = f(\bar{y})$$

At $\bar{y} = 0, \frac{d\theta}{d\bar{y}} = 0$
 At $\bar{y} = 1, \theta = 0$

$$\frac{d\theta}{d\bar{y}} = -f(\bar{y}) Nu_H$$

$$\frac{d\theta}{d\bar{y}} = -Nu_H \int f(\bar{y}) d\bar{y} + C_1$$

$$\theta = -Nu_H \int f(\bar{y}) d\bar{y} + C_1 \bar{y} + C_2$$

$$Q = \frac{1}{2} (Nu_H \bar{T})$$

So, we can write $\frac{d^2\theta}{dy^2} + \frac{u}{\alpha} \frac{h}{H} = 0$, right. Now you can non-dimensionalize like this, right. You can define a new parameter \bar{y} as y/H , ok. So, if you define a new parameter \bar{y} as y/H then this will be $\frac{d^2\theta}{d\bar{y}^2} + \frac{u}{\alpha} \frac{hH}{K} = 0$. So, there will be a $1/H^2$ here that H^2 coming here will make it hH/K .

So, in a non-dimensional form this becomes what. This is a Nusselt number based on H and what is this, this is the velocity profile. This is where fluid mechanics comes into the heat transfer calculation. So, the answer how θ will vary with y will depend on what will depend on the velocity profile. Now there are many types of velocity profiles which are possible like there may be a case when molten metal is flowing very slowly.

So, in that case the entire velocity profile may be almost uniform and that is called as the plug flow or slug flow. So, if have a plug flow or a slug flow that means uniform velocity profile then what is $u/u_{average}$ 1, right. So that is a very special case, very simple algebra to deal with. I am not working it out here, but please take it as a homework and complete it by yourself that consider a plug flow with $u/u_{average}$ is equal to 1, then complete the remaining derivation.

I will do the derivation for a more involved scenario when u by u average is not equal to 1, but it is a fully developed pressure driven flow that is Poiseuille flow. So, for a plane Poiseuille flow what is the velocity profile u by u average is equal to $\frac{3}{2} (1 - y^2)$ by 8 square this was derived by Prof. Som when he was discussing about the exact solution of Navier-Stokes equation.

So, you substitute that here, so this is a like from the fully developed flow Navier-Stokes equation you can derive. There is no problem associated with this. So, now what are the boundary conditions, this is y bar, this is non-dimensional y . What are the boundary conditions at y bar equal to 0, what is the boundary condition? y bar is equal to 0 is the center line, what is the boundary condition for θ at the center line.

Top wall and bottom wall are having symmetric boundary condition, so $d\theta/dy$ equal to 0 at y equal to 0, it is the center line symmetry and at nondimensional y is equal to 1 that is the wall, what is θ . Remember the definition of θ , $T - T_{wall}$ by $T_m - T_{wall}$, so θ is 0. So, with this you have $d^2\theta/dy^2$ is equal to some function of y . Let us say this is $f(y)$ bar into Nusselt number with minus sign.

So, $d\theta/dy$ is equal to minus Nusselt number, integral of $f(y)$ dy bar plus some constant C_1 , so when you integrate you will get θ is equal to, so this is let us say $f_1(y)$, right. I have just written it generically. This integration is very easy because this is a simple polynomial, I do not waste time by doing the integration. I want to give you the frame work in which you can just put the numbers and values.

So, you can evaluate the constants C_1 and C_2 by using this 2 boundary conditions, but does this solve the problem. You have an expression for θ is equal to a function of some function f_2 Nusselt number and y . These are all nondimensional y , so make this as y bar, but this does not tell you exquisitely θ as a function of y because you do not still no what is Nusselt number, what is a value of the Nusselt number?

In fact, obtaining that is the objective and remember that although I have written it as a function of Nusselt number and y , actually Nusselt number is separated from the function of y . It is Nusselt number into some function of y . I mean it is not mixed with y . Nusselt number is outside this integration. So, it will come out to be Nusselt number into some function of y . Now you might argue that yes, we have used 2 boundary conditions for a second order differential equation.

So, what makes the situation that still the problem is not completely solved. One of the reasons is that actually this boundary condition $\theta = 0$ at wall is not actually giving you any new information. Because θ by definition is $T - T_w$ by $T_m - T_w$, so at wall θ will always be 0 that is the definition of θ , but it does not depend on whether it is constant wall temperature or constant wall heat flux whatever.

It is just a simple English sentence $T = T_w$ at wall, so nothing more than that. So, that is what that has been utilized here. So, we need to close this problem by putting some additional constraints what is that additional constraint that additional constraint is the definition of bulk mean temperature which we have not yet used. We have not yet used the definition of the bulk mean temperature.

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Handwritten mathematical derivation on a chalkboard:

Define the definition of T_m -

$$T_m = \frac{\rho c_p \int u T dA}{\rho c_p \int u dA} \quad b \times 2H$$

$$= \frac{\int_{-H}^{+H} \frac{u}{\bar{u}} [(T_m - T_w)\theta + T_w] dy}{\int_{-H}^{+H} \frac{u}{\bar{u}} dy} \quad \int_{-H}^{+H} u dy = \bar{u} 2H$$

$$= (T_m - T_w) \int_0^1 f dy + T_w \int_0^1 f dy \rightarrow 1$$

On the right side of the board:

$$\frac{T_m - T_w}{T_m - T_w} = (T_m - T_w) \int_0^1 f dy$$

$$\Rightarrow \int_0^1 f dy = 1$$

where $f = f_2(Nu, y)$

$Nu_H = ?$

So, now use the definition of T_m . The ρC_p is constant for this problem. So, we cancel the ρC_p part. So, we can write what is integral of $u \, dA$ that is u average into A . So, we have integral of u by u average what is dA , dA is so if we have a parallel plate channel like this so what is dA , so at a distance y you take a small strip of a width dy . dA is dy into b , so dA is dy into b and A is b into $2h$. So, integral u by u average in place of p you can write T_m minus T_{wall} into θ plus T_{wall} .


You can change the variable from y to y^* by putting this H within the derivative. So, it will become dy^* . This is the function f and integral from minus H to H is 2 into integral of 0 to H , because it is symmetric in the 2 sides. So, that 2 and this 2 will cancel. So, this will become integral of $f \theta$ into T_m minus $T_{wall} \, dy^*$ plus T_{wall} , right. We have absorbed the y by H within the derivative.

So, if that be the case then can you tell what is integral of $f \, dy^*$, nondimensional, see f is what f is the velocity profile. So, integral of $f \, dy^*$ non-dimensionally it will become 1 from zero to 1 , right. Because integral of $u \, dy$ is u average into the height, so u by u average into integral of dy nondimensional is 1 . So, this is from the condition that u integrally $u \, dy$ is equal to u average into $2H$, so u by u average integral, this will become 1 . So, this is equal to 1 .

So, we can say that T_m minus T_{wall} is equal to T_m minus T_{wall} into integral of $f \theta \, dy^*$ that means integral of $f \theta \, dy^*$ equal to 1 , this is the constraint. So, θ will be now you can write f^2 of Nusselt number and y^* . The functions f and f^2 are completely known, so this will tell you what is Nusselt number based on H . Now normally we can calculate the Nusselt number on the basis of H , but engineers refer to use a length scale which is called as hydraulic diameter.

So, what is hydraulic diameter, if have a circular pipe you have the actual diameter, but you do not have a circular pipe, there is nothing called a physical diameter, but something equivalent to the physical diameter is called as the hydraulic diameter. So, for as circular pipe what is the physical diameter.

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Now, use the definition of T_m - 

$$T_m = \frac{\rho c_p \int u T dA}{\rho c_p \int u dA} \rightarrow b \times 2H$$

$$= \int_{-H}^{+H} \frac{u}{\bar{u}} \left[\frac{(T_m - T_w)\theta + T_w}{2H} \right] dy$$

$$T_m = T_w \int_0^1 f dy + T_m \int_0^1 f dy$$

$$\int_{-H}^{+H} u dy = \bar{u} 2H$$

$A \rightarrow \pi R^2$
 $P \rightarrow 2\pi R$
 $\left(4 \frac{A}{P}\right) = \frac{4R}{2} = D$
 Hydraulic dia,
 $D_h = \frac{4A}{P} = 4H$
 $Nu_{D_h} = Nu_{4H} = 4Nu_H$
 $= \frac{140}{17}$

So, what is the area, area is πr square. What is the perimeter $2\pi r$, so what is area by perimeter? So how can you get diameter from here, multiply by 4, so these becomes diameter, ok. So, 4 into area by perimeter for a circular pipe is the actual diameter. For a noncircular geometry that is some effective diameter which is called hydraulic diameter. So, hydraulic diameter let us write it here is 4 into area by perimeter.

So, what will be the hydraulic diameter for a parallel plate channel the half height is H . Area by perimeter we have already calculated what was that H , so 4 area by perimeter is $4H$. So, the Nusselt number based on hydraulic diameter is the Nusselt number based on $4H$ that is 4 times the Nusselt number based on H and if you evaluate it all these numbers I have given the outline, but if you now evaluate the value this value will come out to be 140 by 17.

So, this is your homework 140 by 17, okay. So, I have given the full frame work, in this frame work just you have to do the integrations, integral A f 1 f 2 all this and once you substitute, you will get the value of this, so we can see that the Nusselt number is the constant for thermally fully developed flow for this example, this is the value of the constant. We will stop here and we will continue in next lecture.