

Computational Fluid Dynamics for Incompressible Flows
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Module 11: Finite Volume Method- 2
Lecture 2

Finite Volume Discretization of Unsteady Convection- Diffusion Equation

Hello everyone, so today we will discretize Unsteady Convection Diffusion Equation using finite volume method. So, in last lecture, we discretize steady convection diffusion equation using finite volume method, same procedure will use here only we need to integrate the governing equation over time as well. So, first let us write the two dimensional convection diffusion equation.

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Finite Volume Discretization of Convection-Diffusion Equation

2-D unsteady convection diffusion equation

$$\frac{\partial(\rho\phi)}{\partial t} + \frac{\partial(\rho u\phi)}{\partial x} + \frac{\partial(\rho v\phi)}{\partial y} = \frac{\partial}{\partial x} \left(\Gamma \frac{\partial\phi}{\partial x} \right) + \frac{\partial}{\partial y} \left(\Gamma \frac{\partial\phi}{\partial y} \right) + S$$

$$\frac{\partial(\rho\phi)}{\partial t} + \nabla \cdot (\rho \vec{u}\phi) = \nabla \cdot (\Gamma \nabla\phi) + S$$

Flux vector,
 $\vec{J} = \rho \vec{u}\phi - \Gamma \nabla\phi$

$$\frac{\partial(\rho\phi)}{\partial t} + \nabla \cdot \vec{J} = S$$

So, 2D unsteady convection diffusion equation. So, in two dimension you can write in conservative form $\frac{\partial(\rho\phi)}{\partial t} + \frac{\partial(\rho u\phi)}{\partial x} + \frac{\partial(\rho v\phi)}{\partial y} = \frac{\partial}{\partial x} \left(\Gamma \frac{\partial\phi}{\partial x} \right) + \frac{\partial}{\partial y} \left(\Gamma \frac{\partial\phi}{\partial y} \right) + S$. Here, Γ is diffusion coefficient and we have written this convection diffusion equation for any general variable ϕ and S is the source term and ρ is fluid density.

So, in vector form if you write, then this equation you will get like divergence of $\rho u \phi$ is equal to divergence of $\gamma \text{grad } \phi$ plus S . So, where u is the velocity vector and this is known, so it is linear equation because u is known. So, u_i plus v_j where you have in Cartesian coordinate. If this is your x and this is your y direction, then velocity u is in the x direction and v in y direction and i is the unit vector in the x direction and j is the unit vector in y direction and grad is $\text{del of del } x_i$ plus $\text{del of del } y_j$. So, this governing equation now will integrate over the main cell P .

So, you can see this is your main cell. So, in this main cell P we will integrate and these are the face centers, small e , small n , small w and small s and these are the neighbors. So, our variables are available, the value of the variable is available only at this node capital P , capital E , capital W , capital N and capital S but we need to interpolate the value if required at the face center.

So, let us introduce the flux vector. So, your flux vector will be J . So, we can write $\rho u \phi$ minus $\gamma \text{grad } \phi$. So, this flux vector now is having two components, one is your convection term, one is diffusion term so, J is like J_c convection term plus J_d where J_c is your $\rho u \phi$ and J_d is your minus $\gamma \text{grad } \phi$.

So, this flux vector if you put it here, then what will be the governing equation? Equation will be $\text{del row } \phi$ by Δt plus divergence of J is equal to S . So, first term represents the temporal term and the second term obviously flux, we have two fluxes one is convection flux and diffusion flux and right hand side term is your source term.

So, in the last class already we have discretize excluding the temporal term. So now, let us integrate this equation over the control volume and over the time.

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Finite Volume Discretization of Convection-Diffusion Equation

$$\int_{t_n}^{t_{n+1}} \int_V \frac{\partial(\rho\phi)}{\partial t} dV dt + \int_{t_n}^{t_{n+1}} \int_V \nabla \cdot \vec{J} dV dt = \int_{t_n}^{t_{n+1}} \int_V S dV dt$$

$$\int_{t_n}^{t_{n+1}} \frac{d(\rho\phi_P \Delta V)}{dt} dt + \int_{t_n}^{t_{n+1}} (\sum_f \vec{J}_f \cdot \vec{S}_f) dt = (\bar{S} \Delta V) \Delta t$$

$$\rho \Delta V (\phi_P^{n+1} - \phi_P^n) + (\sum_f \vec{J}_f \cdot \vec{S}_f) \Delta t = (\bar{S} \Delta V) \Delta t$$

$$\frac{\rho \Delta V}{\Delta t} (\phi_P^{n+1} - \phi_P^n) + (\sum_f \vec{J}_f \cdot \vec{S}_f) = (\bar{S} \Delta V)$$

$$(\sum_f \vec{J}_f \cdot \vec{S}_f)^{n+1} = f_t (\sum_f \vec{J}_f \cdot \vec{S}_f)^{n+1} + (1-f_t) (\sum_f \vec{J}_f \cdot \vec{S}_f)^n$$

$$(\bar{S} \Delta V)^{n+1} = f_t \bar{S}^{n+1} \Delta V + (1-f_t) \bar{S}^n \Delta V$$

$\Delta t = t_{n+1} - t_n$

So, let us write here $\frac{\partial(\rho\phi)}{\partial t}$ by $\frac{d}{dt}$. Now you integrate over time t_n to $t_n + 1$ where Δt is your time step, $t_n + 1$ minus t_n and integral volume, so this is your in this main cell P we are integrating it. So, this is the main cell, this is the main cell. So, here we are integrating $dV dt$. The second term, similarly you can write t_n to $t_n + 1$ volume integral and you have divergence of $J dV dt$ and in the right hand side again the source term you integrate $t_n + 1$ volume integral $S dV dt$.

So, first let us see the temporal term. So, this temporal term if you integrate this term over this control volume. So, what you will get? So, if volume is not changing with time so non deforming control volume if you assume, then you will get $\frac{d}{dt} \int_V \rho\phi dV$. These volume integral dV you can write $\frac{d}{dt} \int_V \rho\phi dV$ and this you can write as $\frac{d(\rho\phi_P)}{dt}$, so you are now taking so, you are averaging it out over the control volume, so you are taking the value of ϕ at the cell center P. So, this is your $\frac{d(\rho\phi_P)}{dt}$ by Δt and the volume of the cell P. So, these now you can write $\frac{d(\rho\phi_P)}{dt} \Delta V$.

So, now this we need to integrate over the time. So, if you write it so, you will get t_n to $t_n + 1$. So, you will get $d(\rho\phi_P)$ and that ΔV is not changing with time, so you can write ΔV so this is your integral, first integral, so these dt and there is a dt so, this

will get or you can write Δt . So, obviously integral will give just $\rho \phi_P \Delta V$ from t_n to t_{n+1} .

Now, the second time so t_n to t_{n+1} . So, now you can see the volume integral, volume integral of this flux term, you can convert it into surface integral using Gauss divergence theorem and after making the assumption that J varies linearly and the average value of J lies at the face center.

So, that if you make the assumptions, then you can write the flux term as summation of over the faces $J_f \cdot S_f$ over all the faces and Δt , so this is the second term and similarly the last term also so, you are taking t_n to t_{n+1} and you are now averaging over the cell and average value you are taking at the cell center. So this is your \bar{S} , at the cell center P you are taking and ΔV and Δt .

So now, you have to integrate over the time, but we will take some, at some time τ and we will take some average between these t_n and t_{n+1} because values are not available in between, only values are available at time t_n and t_{n+1} . So, this will write as, so this will get $\rho \Delta V \phi_P^{n+1} - \phi_P^n$ and this term you will get plus, we are taking $J_f \cdot S_f$ at some time level τ , which is in between t_n and t_{n+1} and the last term also \bar{S} at the cell center P , \bar{S} , so, $\bar{S} \Delta V$ at time level $\tau \Delta t$, here Δt will be there because you are integrating over this, so this is Δ by Δt .

So now, you integrate now, you divide both sides by Δt , so what you will get $\rho \Delta V$ by $\Delta t \phi_P^{n+1} - \phi_P^n$ plus summation of $J_f \cdot S_f$ at τ is equal to $\bar{S} \Delta V$ at τ . So, now this value of fluxes at time τ will take as summation of $J_f \cdot S_f$ at time level τ will take as summation of $J_f \cdot S_f$ with a factor f_t at level $n+1$ and one minus $f_t J_f \cdot S_f$ at time level n .

So f_t lies between 0 and 1, so if f_t is 0 what will we this scheme only at any time level you will give get? So you will get explicit scheme, similarly the source term $\bar{S} \Delta V$ at τ you will take as $f_t \bar{S}^{n+1} \Delta V + (1 - f_t) \bar{S}^n \Delta V$.

Now, let us write the flux term so, already these flux terms using central difference convection scheme you have discretize it in last class, so, let us write down here.

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Finite Volume Discretization of Convection-Diffusion Equation

Use central difference convection scheme
 $\vec{S}_e = A_e \hat{i}$ $\phi_e = \frac{\phi_E + \phi_P}{2}$

$\vec{J} = \rho \vec{u} \phi - \Gamma \nabla \phi$

$$\sum \vec{J}_f \cdot \vec{S}_f = \vec{J}_e \cdot \vec{S}_e + \vec{J}_w \cdot \vec{S}_w + \vec{J}_n \cdot \vec{S}_n + \vec{J}_s \cdot \vec{S}_s$$

$$= \frac{F_e}{2} (\phi_E + \phi_P) - D_e (\phi_E - \phi_P)$$

$$- \frac{F_w}{2} (\phi_P + \phi_W) + D_w (\phi_P - \phi_W)$$

$$+ \frac{F_n}{2} (\phi_N + \phi_P) - D_n (\phi_N - \phi_P)$$

$$- \frac{F_s}{2} (\phi_P + \phi_S) + D_s (\phi_P - \phi_S)$$

$$= \left[(D_e + \frac{F_e}{2}) + (D_w - \frac{F_w}{2}) + (D_n + \frac{F_n}{2}) + (D_s - \frac{F_s}{2}) \right] \phi_P - \left[(D_e - \frac{F_e}{2}) \right] \phi_E$$

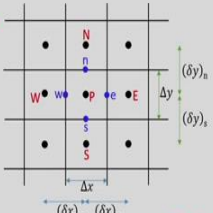
$$+ \left[(D_w + \frac{F_w}{2}) \right] \phi_W + \left[(D_n - \frac{F_n}{2}) \right] \phi_N$$

$$+ \left[(D_s + \frac{F_s}{2}) \right] \phi_S$$

$$= (a_E + a_w + a_n + a_s + F_e - F_w + F_n - F_s) \phi_P$$

$$- [a_E \phi_E + a_w \phi_W + a_n \phi_N + a_s \phi_S]$$

$$= (\sum a_{nb} + \sum F_f) \phi_P - \sum a_{nb} \phi_{nb}$$



$F_e = \rho u_e A_e$ $D_e = \frac{\Gamma_e A_e}{(\Delta x)_e}$
 $F_w = \rho u_w A_w$ $D_w = \frac{\Gamma_w A_w}{(\Delta x)_w}$
 $F_n = \rho u_n A_n$ $D_n = \frac{\Gamma_n A_n}{(\Delta y)_n}$
 $F_s = \rho u_s A_s$ $D_s = \frac{\Gamma_s A_s}{(\Delta y)_s}$
 $u_e = \frac{u_E + u_P}{2}$
 $a_E = D_e - \frac{F_e}{2}$ $a_w = D_w + \frac{F_w}{2}$
 $a_n = D_n - \frac{F_n}{2}$ $a_s = D_s + \frac{F_s}{2}$
 $D_e + \frac{F_e}{2} = D_e - \frac{F_e}{2} + \frac{F_e}{2} + \frac{F_e}{2}$
 $D_w - \frac{F_w}{2} = D_w + \frac{F_w}{2} - \frac{F_w}{2} - \frac{F_w}{2} = a_w - F_w$

Finite Volume Discretization of Convection-Diffusion Equation

$$\int_{t_n}^{t_{n+1}} \int_V \frac{\partial(\rho\phi)}{\partial t} dV dt + \int_{t_n}^{t_{n+1}} \int_V \nabla \cdot \vec{J} dV dt = \int_{t_n}^{t_{n+1}} \int_V S dV dt$$

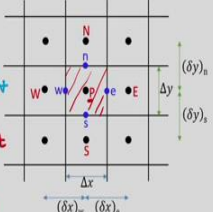
$$\int_{t_n}^{t_{n+1}} \frac{d(\rho\phi \Delta V)}{dt} dt + \int_{t_n}^{t_{n+1}} (\sum_f \vec{J}_f \cdot \vec{S}_f) dt = \int_{t_n}^{t_{n+1}} (\bar{S} \Delta V) dt$$

$$\rho \Delta V (\phi_P^{n+1} - \phi_P^n) + (\sum_f \vec{J}_f \cdot \vec{S}_f) \Delta t = (\bar{S} \Delta V) \Delta t$$

$$\frac{\rho \Delta V}{\Delta t} (\phi_P^{n+1} - \phi_P^n) + (\sum_f \vec{J}_f \cdot \vec{S}_f) = (\bar{S} \Delta V)$$

$$(\sum_f \vec{J}_f \cdot \vec{S}_f) = f_t (\sum_f \vec{J}_f \cdot \vec{S}_f)^{n+1} + (1-f_t) (\sum_f \vec{J}_f \cdot \vec{S}_f)^n$$

$$(\bar{S} \Delta V) = f_t \bar{S}^{n+1} \Delta V + (1-f_t) \bar{S}^n \Delta V$$



$\Delta t = t_{n+1} - t_n$
 $0 \leq f_t \leq 1$

So use central difference convection scheme. So, for that now what is J? J is rho u phi minus gamma grad phi, and now you find Jf dot Sf, Jf dot Sf if you find then you can write Je dot Se plus Jw dot Sw plus Jn dot Sn plus Js dot Ss. So, all the faces, four faces, so these now you find. So, what is Se? Se is your Aei, so what you will get?

So, Je dot Se so, Je dot Se now you are going to get using central difference so, phi e will take, phi e you will take average of phi E plus phi P divided by 2, because at this East place you are finding, so using uniform grid phi E plus phi P divided by 2 so, that if you

put so, then you will get rho so, you will get F_e by $2 \phi E$ plus ϕP this is your convection term and the diffusion term, you will get minus $D_e \phi E$ minus ϕP . What is F_e ? So, F_e we have written as $\rho u_e A_e$ and D_e is your $\gamma_e A_e$ by Δx and A is your Δy in this two dimensional case.

So now, similarly, you write $J_w \cdot S_w$ and other terms. So, if you see $J_w \cdot S_w$, so we will write minus F_w by 2. So, it will be average between this w so, if you take the central difference convection scheme it will be ϕP plus ϕw divided by 2. So, it will be ϕP plus ϕw divided by 2 plus the diffusion term will be $D_w \phi P$ minus ϕw , where F_w we are defining as $\rho u_w A_w$ and D_w is $\gamma_w A_w$ Δx_w .

Similarly, now you write for other two terms, so you will get $J_n \cdot S_n$ if you write so, it will be a F_n by 2 so, ϕN plus ϕP minus $D_n \phi N$ minus ϕP and for the last term you will get F_s by 2 ϕP plus ϕS minus, sorry it is plus $D_s \phi P$ minus ϕS . So now, rearrange it and write as. So, here F_n is $\rho u_n A_n$ and D_n is $\gamma_n A_n$ by Δy_n and A_n is in this case Δx and F_s is $\rho u_s A_s$ and D_s $\gamma_s A_s$ by Δy_s , and what is u_s ? u_n , u_w and u_e you will take central difference. So, u_e you calculate as u_e plus u_P divided by 2, here also very using central difference.

So now, if you rearrange it you are going to get, so all the P terms you take. So, for P terms you see, what are the P terms? So, this is your P , P , P , P , P , P and P . So, if you take all the P terms together, then you see it is will be minus minus plus, so it will be $D_e \phi P$ so, first we will write D_e plus F_e by 2, so for these ϕP this is F_e by 2 so, this is F_e by 2. So, this is one coefficient then for this you will get D_w minus F_w by 2, D_w minus F_w by 2, then for this north you will get D_n plus F_n by 2, then you will get plus D_s minus F_s by 2. So, this is the coefficient of ϕP .

Now, for the neighbor terms you can see, so this is your East, this is your East so, you can write it. So, you will write taking that minus, so we will get D_e so this is your minus so, we have taken outside, so D_e and here minus F_e by 2 because it is plus so minus I have taken outside, so it is minus F_e by 2 ϕE plus D_w plus F_w by 2 ϕW . Similarly, the other terms D_n minus F_n by 2 ϕN all our neighbor terms and D_s plus F_s by 2 ϕS .

So now, if you define the coefficients as a_E is equal to D_e minus F_e by 2 a_W as D_w plus F_w by 2 a_N as D_n minus F_n by 2 and a_S as D_s plus F_s by 2. So, you can see, this term you can write as a_E , this is your a_W , this is your a_N and this is your a_S . And this term, the first term, this you simplify so if you write first term this one, so you can write as you see D_e plus F_e by 2, I will write as a_E , so if you write a_E , so what is a_E ? So, D_e minus F_e by 2 so, I have subtracted minus F_e by 2, so again you add plus F_e by 2. So, let me write here, so I have D_e plus F_e by 2.

Now, I want to write in terms of coefficient a_E but a_E is D_e minus F_e by 2. So, if you write D_e minus F_e by 2, I am subtracting one F_e by 2, you add another F_e by 2. So, then there will be no change and one F_e by 2 is already there. So now, these D_e by F_e by 2 is a_E and F_e by 2 plus F_e by 2 it will be F_e .

And similarly, D_w minus F_w by 2 what you can do the second term, you can write D_w plus F_w by 2 and you subtract one F_w by 2 there will be no change then minus F_w by 2. So, this you can write as a_W minus F_w because these minus F_w by 2 and minus F_w by 2 will be minus F_w .

Similarly, for other coefficients you can show, and finally we can write the first term as a_E plus a_W plus a_N plus a_S minus, sorry plus F_e minus F_w plus F_n minus F_s . So, this is the coefficient of ϕ_P , and we have a_E and we have minus $a_E \phi_E$ plus $a_W \phi_W$ plus $a_N \phi_N$ plus $a_S \phi_S$.

So, now this we can write as, you can see this is your so, this a_E plus a_W plus a_N plus a_S we can write summation of $a_n b_n$ because all neighbor coefficient summation and this F_e minus F_w plus F_n minus F_s . So, these we can write as summation of F_f , what does it mean by summation F_f ? So, summation of F_f is net out flux.

So, this is your mass flux, so, obviously it satisfies the continuity. So, in the beginning of the solution it may be nonzero but at the convergence summation of F_f will be tend to 0 because it will satisfy the continuity equation. So this we have written then this is your ϕ_P and this you can write as minus summation of $a_n b_n \phi_n$.

So now, these fluxes now you substituted in this expression, so we will substitute in this expression this $\vec{J} \cdot \vec{S}_f$ with the term n plus 1 and n.

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Finite Volume Discretization of Convection-Diffusion Equation

$$\frac{\rho \Delta V}{\Delta t} (\phi_p^{n+1} - \phi_p^n) + \left(\sum \vec{J}_f \cdot \vec{S}_f \right)^{n+1} + (1-f_c) \left(\sum \vec{J}_f \cdot \vec{S}_f \right)^n$$

$$+ \sum_f \vec{S}_f^{n+1} \Delta V + (1-f_c) \sum_f \vec{S}_f^n \Delta V$$

Finite Volume Discretization of Convection-Diffusion Equation

Use central difference convection scheme

$$\vec{J} = \rho \vec{u} \phi - \tau \nabla \phi$$

$$\vec{S}_e = A_e \hat{i} \quad \phi_e = \frac{\phi_E + \phi_P}{2}$$

$$\sum \vec{J}_f \cdot \vec{S}_f = \vec{J}_e \cdot \vec{S}_e + \vec{J}_w \cdot \vec{S}_w + \vec{J}_n \cdot \vec{S}_n + \vec{J}_s \cdot \vec{S}_s$$

$$= \frac{F_e}{2} (\phi_E + \phi_P) - D_e (\phi_E - \phi_P)$$

$$- \frac{F_w}{2} (\phi_P + \phi_W) + D_w (\phi_P - \phi_W)$$

$$+ \frac{F_n}{2} (\phi_N + \phi_P) - D_n (\phi_N - \phi_P)$$

$$- \frac{F_s}{2} (\phi_P + \phi_S) + D_s (\phi_P - \phi_S)$$

$$= \left[\left(D_e + \frac{F_e}{2} \right) + \left(D_w - \frac{F_w}{2} \right) + \left(D_n + \frac{F_n}{2} \right) + \left(D_s - \frac{F_s}{2} \right) \right] \phi_P$$

$$- \left[\left(D_e - \frac{F_e}{2} \right) \right] \phi_E$$

$$+ \left(D_w + \frac{F_w}{2} \right) \phi_W + \left(D_n - \frac{F_n}{2} \right) \phi_N$$

$$+ \left(D_s + \frac{F_s}{2} \right) \phi_S$$

$$= (a_E + a_W + a_N + a_S + F_e - F_w + F_n - F_s) \phi_P$$

$$- [a_E \phi_E + a_W \phi_W + a_N \phi_N + a_S \phi_S]$$

$$= (\sum a_{nb} + \sum F_f) \phi_P - \sum a_{nb} \phi_{nb}$$

$$F_e = \rho u_e A_e \quad D_e = \frac{\Gamma_e A_e}{(\Delta x)_e}$$

$$F_w = \rho u_w A_w \quad D_w = \frac{\Gamma_w A_w}{(\Delta x)_w}$$

$$F_n = \rho u_n A_n \quad D_n = \frac{\Gamma_n A_n}{(\Delta y)_n}$$

$$F_s = \rho u_s A_s \quad D_s = \frac{\Gamma_s A_s}{(\Delta y)_s}$$

$$u_e = \frac{u_E + u_P}{2}$$

$$a_E = D_e - \frac{F_e}{2} \quad a_W = D_w + \frac{F_w}{2}$$

$$a_N = D_n - \frac{F_n}{2} \quad a_S = D_s + \frac{F_s}{2}$$

$$D_e + \frac{F_e}{2} = D_e - \frac{F_e}{2} + \frac{F_e}{2} + \frac{F_e}{2}$$

$$D_w - \frac{F_w}{2} = D_w + \frac{F_w}{2} - \frac{F_w}{2} - \frac{F_w}{2} = a_W - F_w$$

Finite Volume Discretization of Convection-Diffusion Equation

$$\frac{\rho \Delta V}{\Delta t} (\phi_p^{n+1} - \phi_p^n) + f_t (\sum J_f \cdot \vec{S}_f)^{n+1} + (1-f_t) (\sum \bar{J}_f \cdot \vec{S}_f)^n = \bar{J}_f \bar{S}^{n+1} \Delta V + (1-f_t) \bar{S}^n \Delta V + b$$

$$\sum \bar{J}_f \cdot \vec{S}_f = (\sum a_{nb} + \sum F_f) \phi_p - \sum a_{nb} \phi_{nb}$$

$$\frac{\rho \Delta V}{\Delta t} (\phi_p^{n+1} - \phi_p^n) + f_t [(\sum a_{nb} + \sum F_f) \phi_p^{n+1} - \sum a_{nb} \phi_{nb}^{n+1}] + (1-f_t) [(\sum a_{nb} + \sum F_f) \phi_p^n - \sum a_{nb} \phi_{nb}^n] = b$$

$$a_p^0 = \frac{\rho \Delta V}{\Delta t}$$

$$[a_p^0 + f_t (\sum a_{nb} + \sum F_f)] \phi_p^{n+1} = \sum a_{nb} [f_t \phi_{nb}^{n+1} + (1-f_t) \phi_{nb}^n] + [a_p^0 - (1-f_t) (\sum a_{nb} + \sum F_f)] \phi_p^n + b$$

$$a_p = a_p^0 + f_t (\sum a_{nb} + \sum F_f) \quad \phi_p = \phi_{i,j}$$

$$a_p^0 = \frac{\rho \Delta V}{\Delta t}$$

$$b = \bar{J}_f \bar{S}^{n+1} \Delta V + (1-f_t) \bar{S}^n \Delta V$$

So we have $\rho \Delta V$ by Δt $\phi_p^{n+1} - \phi_p^n$. Then we have summation of $J_f \cdot S_f$ with a factor F_t and at level $n+1$ and we have $1 - F_t$ summation of $J_f \cdot S_f$ at level n and we have the source term $f_t \bar{S}^{n+1} \Delta V + (1 - f_t) \bar{S}^n \Delta V$. So now, we have already written this term. So, the summation of a $J_f \cdot S_f$ we have written in this form, so, this let us substituted there.

So, summation of $J_f \cdot S_f$ we have written as summation of a_{nb} plus summation of $F_f \phi_p$ plus summation of $a_{nb} \phi_{nb}$, here it is minus. So now, if you put it there, so you are going to get $\rho \Delta V$ by Δt $\phi_p^{n+1} - \phi_p^n$ plus F_t now these old term at $n+1$ level. So, you have summation of a_{nb} plus summation of $F_f \phi_p$ minus summation of $a_{nb} \phi_{nb}$ at level $n+1$ and we have $1 - f_t$.

Again, we have summation of a_{nb} plus summation of $F_f \phi_p$ minus summation of $a_{nb} \phi_{nb}$ at n . So, n let us write in this way $n+1$, this is your $n+1$ and this is your n and this is your n and let us write this source term plus b , where b contains the source term. So now, you take the ϕ_p^{n+1} this terms in the left hand side and all other terms you take in the right hand side.

So, let us rearrange this equation. So, ϕ_p^{n+1} you see ϕ_p^{n+1} here we have and ϕ_p^{n+1} here we have. So, this you take in the left hand side so, you can see if you put a_p^0 is equal $\rho \Delta V$ by Δt , a_p^0 , if you write $\rho \Delta V$ by Δt ,

so now you see a_p naught this $\phi P_n + 1$. So this will take here so it will be Δt so, Δt into this, so plus Δt summation of a_b plus summation F_f into $\phi P_n + 1$. So, this is the whole bracket, so this is the coefficient of $\phi P_n + 1$.

Now you take the neighbor term in the right hand side. So, this is your neighbor term $n + 1$ and this is at $n + 1$. So, this you take in the right hand side, so equal to so, you can write summation of a_b so, you see this is minus and minus Δt and here also minus $1 - \Delta t$, so let us take in the right hand side so it will become plus.

So, it will be a Δt summation of, summation already we have taken. So, ϕP_{n+1} and this term if you take, so it will be $1 - \Delta t$. So, this we are taking in the right hand side so it will be ϕP_n and what is left so, ϕP at term level n . So, this is your ϕP at time level n and this is your ϕP term level n and so that you take in the right hand side, sorry here it will be equal to b .

So, it is equal to b please correct it. So, it is equal to and this is equal to b . So, this whole term we are considering as the source term b and that will be in the right hand side, equal to will be here and here also will be equal to b , so please correct it. So now, you write all this ϕP term at time level n . So, you can see this will come negative, so right hand side plus so, it will be a_p naught so, plus a_p naught.

Then, we have this if you take in the right hand side so it will be $1 - \Delta t$, then we have summation a_b plus summation of F_f , summation of F_f . So, ϕP_n and plus b . So the coefficient of ϕP is the diagonal coefficient that we can represent as a_p so you can write a_p as a_p naught plus Δt summation of a_b plus summation of F_f .

So, this is your diagonal coefficient and a_p naught is your $\rho \Delta b$ by Δt and we have b . So b is $\bar{S}_n + 1 \Delta b + 1 - \Delta t \bar{S}_n \Delta v$. So, this is the final discretize equation for this unsteady convection diffusion equation and we have written in terms of the coefficient. So, diagonal coefficient is a_p and $\phi P_n + 1$, obviously it represents the value at the cell center and that is your ij . So P , ϕP obviously, it is ϕ_{ij} and similarly, you can see that East will be $i + 1, j$, $i - 1, j$ and this North will be $ij + 1$ and South will be $ij - 1$.

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Finite Volume Discretization of Convection-Diffusion Equation

$f_t = 0$ **Explicit Scheme**

$$a_p \phi_p^{n+1} = \sum_f a_{nb} \phi_{nb}^n + [a_p^o - (\sum a_{nb} + \sum F_f)] \phi_p^n + b$$

$$a_p = a_p^o \quad b = \bar{S}^n \Delta V \Delta t$$

$f_t = 1$ **Implicit Scheme**

$$a_p = a_p^o + \sum a_{nb} + \sum F_f \quad b = \bar{S}^{n+1} \Delta V \Delta t$$

$$a_p \phi_p^{n+1} = \sum a_{nb} \phi_{nb}^{n+1} + a_p^o \phi_p^n + b$$

$f_t = \frac{1}{2}$ **Crank-Nicolson Scheme**

$$a_p = a_p^o + \frac{1}{2} (\sum a_{nb} + \sum F_f)$$

$$b = \frac{1}{2} (\bar{S}^n + \bar{S}^{n+1}) \Delta V \Delta t$$

$$a_p \phi_p^{n+1} = \frac{1}{2} \sum a_{nb} (\phi_{nb}^{n+1} + \phi_{nb}^n) + [a_p^o - \frac{1}{2} (\sum a_{nb} + \sum F_f)] \phi_p^n + b$$

Now, if you consider the special cases, so if f_t is 0, then it will become explicit. So, you can see if f_t is 0 special case if you consider f_t is 0, then it will become explicit, explicit scheme and obviously, you will get f_t is 0 so, you will get $a_p \phi_p^{n+1}$ and in the right hand side you will get summation of $a_{nb} \phi_{nb}^n$ at n all the faces f_t is 0.

So, you will get $a_p \phi_p^{n+1}$ minus summation of a_{nb} plus summation of F_f at ϕ_p^{n+1} plus b . So, you can see here right hand side all are at n th time level. So, all are known from the previous time level and left hand side you have ϕ_p^{n+1} . So, you are actually calculating at present time level ϕ_p^{n+1} .

So, one unknown and right hand side are known, so obviously it is explicit scheme and its order of accuracy is first order and this you can see if f_t is equal to 1, so you will get implicit scheme. So, in this special case it will become implicit. So, a_p will be a_p plus summation of a_{nb} plus summation of F_f .

So, you will get $a_p \phi_p^{n+1}$ is equal to and in this case if you represent a_p , so it will be a_p is equal to a_p plus summation of a_{nb} plus summation of F_f and b will be only $\bar{S}^{n+1} \Delta V \Delta t$ and here it will be $\bar{S}^{n+1} \Delta V \Delta t$. So, $a_p \phi_p^{n+1}$ and right hand side you will get summation of a_{nb} , all the neighbors ϕ_{nb}^n at n plus 1.

Then, the next term you will get plus $\Delta t \phi^n$ and plus b . So, b only $S^n + \Delta t v$ and you can see that right hand side all the neighbor terms at $n + 1$, left hand side it is $n + 1$ then you will get more than one unknown, so it is implicit scheme and it is also first order accurate and if Δt is equal to half Δt if you put, then you will get Crank Nicolson. So, in that case you will get Δt is equal to Δt half summation of Δt plus summation of F^n and b you will get half of $S^n + S^{n+1}$ into $\Delta t v$ and the equation you will be getting $\Delta t \phi^{n+1}$, then you will get half summation of $\Delta t \phi^{n+1} + \Delta t \phi^n$.

So, then you will get Δt minus half summation of Δt plus summation of $F^n \phi^n$ at n and plus the source term. So, today we consider two dimensional unsteady convection diffusion equation and we integrated this equation over the main cell P and over the time, then we have written the discretize form in general way and we have considered few special cases where if Δt is equal to 0 it is explicit, Δt is equal to 1 it is implicit and Δt is equal to half it is Crank Nicolson and obviously Crank Nicolson is second order time accurate.

While discretizing these equations we considered central difference convection scheme, but you can extend it to faster upwind scheme, so this may be your homework so you just consider first upwind scheme and discretize this equation and find the coefficients. Thank you.