

# Computational Fluid Dynamics for Incompressible Flows

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Module 10: Finite Volume Method- 1

Lecture 3

## Finite Volume Discretization of Unsteady Diffusion Equation

Hello everyone, in today's lecture, we will consider two Dimensional Unsteady Diffusion Equation and we will discretize using finite volume method. In last class, you learnt the finite volume discretization of steady state diffusion equation and you have also learned the assumptions made in different places.

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**Finite Volume Discretization of Diffusion Equation**

2-D unsteady diffusion equation

$$\frac{\partial(\rho\phi)}{\partial t} = \nabla \cdot (\Gamma \nabla \phi) + S$$

Flux vector

$$\vec{J} = -\Gamma \nabla \phi$$
$$\frac{\partial(\rho\phi)}{\partial t} + \nabla \cdot \vec{J} = S$$

Integrating the above equation over main cell P and time interval  $(t_{n+1} - t_n)$ ,

$$\int_{t_n}^{t_{n+1}} \int_V \frac{\partial(\rho\phi)}{\partial t} dV dt + \int_{t_n}^{t_{n+1}} \int_V \nabla \cdot \vec{J} dV dt = \int_{t_n}^{t_{n+1}} \int_V S dV dt \Rightarrow \bar{S} \Delta V$$
$$\sum_f \vec{J}_f \cdot \vec{S}_f = \vec{J}_w \cdot \vec{S}_w + \vec{J}_e \cdot \vec{S}_e + \vec{J}_n \cdot \vec{S}_n + \vec{J}_s \cdot \vec{S}_s$$

CV is stationary and non-deforming

$$\frac{\partial}{\partial t} \int_V (\rho\phi) dV = \frac{\partial}{\partial t} (\rho\phi \Delta V) = \frac{d}{dt} (\rho\phi) \Delta V$$

So, today let us consider, two dimensional unsteady diffusion equation. So, the equation is  $\rho \phi$  by  $\Delta t$ . So, this is your temporal term now, this term is added, and in last class already we have discretize these divergence of  $\Gamma \nabla \phi$  plus  $S$ . So, in last class this first term, temporal term was 0. So, you know how to discretize these using finite volume method, only now we have to see that how we will integrate in the time level.

Obviously, new time levels will be extrapolated from the previous time level values, so you can see that now, we will define the flux vector  $J$  as minus gamma grad phi. So, the modified equation you can write, as divergence of  $J$  is equal to  $S$ . So, now this equation you integrate in this main cell  $P$ , this is your main control volume. So, in this main cell  $P$  we will integrate both over in volume and over time. So, integrating the above equation over main cell  $P$  and time interval  $t_n$  plus 1 minus  $t_n$ , what you will get now? So you will integrate over time, so time from  $t_n$  to  $t_n$  plus 1.

So again over the control volume, so this is your volume integral  $\frac{d}{dt} \int_V \rho \phi \, dv$ , so this is the temporal term, the second term now you integrate over the time  $t_n$  to  $t_n$  plus 1 then volume integral then you have divergence of  $J \, dv$ , then  $dt$  and then you have source term so that also, you have to integrate from  $t_n$  to  $t_n$  plus 1 volume integral  $S \, dv \, dt$ .

So now, first let us integrate over control volume, each term. So, if you see the second term, this already we have learned how to discretize using finite volume method, what you do? You convert this volume integral to surface integral using Gauss divergence theorem and write the, in terms of face fluxes and these also we have integrated using the finite volume method and we have used average value at the cell center  $P$ . So, these two terms already we have learned in last class, how to discretize using finite volume method.

So, now let us see this term. So, now, let us integrate over the control volume. So, if you see this term, particularly, then this term you can write as say, in our case control volume is stationary, stationary and  $cv$  is stationary and non-deforming. So what we can do? We can write  $\frac{d}{dt} \int_V \rho \phi \, dv$ . So now we can see that now, this value  $\rho \phi$ , you have to take the average value at the main cell  $P$ .

So that will take the mean value theorem and it will take the average value at cell  $P$ , so we can write  $\frac{d}{dt} \int_V \rho \phi \, dv$ , so we will take average value, so it will be  $\rho \phi_P$ , so  $\phi_P$  is the average value at the cell center  $P$  and  $\rho$  is constants only  $\rho$  so into  $\Delta v$ , so that we can write and you can see now these value, we have to integrate over time so that you can write  $\frac{d}{dt} \rho \phi_P \Delta v$ , so this is the first term, you can write this way.

Now, you consider the second term. So, now we are not going to write in detail. So, you use the Gauss divergence theorem and convert these volume integral to surface integral and write in terms of face fluxes. So, this, you can write, summation of  $f \mathbf{J}_f \cdot \mathbf{S}_f$ , summation of all the faces, so in this case we have four faces, East, West, sorry, we have East, North, West and South. So, this way, this will be just  $\mathbf{J}_e \cdot \mathbf{S}_e$  plus  $\mathbf{J}_w \cdot \mathbf{S}_w$  plus  $\mathbf{J}_n \cdot \mathbf{S}_n$  plus  $\mathbf{J}_s \cdot \mathbf{S}_s$ .

Now, let us see the last term, which is your source term, and also we will, use the average value at the cell center. So all these, you can see that we have used second order accurate scheme. So, this term only in this bracket we will write  $\bar{S} \Delta v$ . So all this time we have discretize using volume integral, now we have to integrate over the time.

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### Finite Volume Discretization of Diffusion Equation

$$\int_{t_n}^{t_{n+1}} \frac{d}{dt} (\rho \phi_P) \Delta v dt + \int_{t_n}^{t_{n+1}} \left( \sum_f \vec{J}_f \cdot \vec{S}_f \right) dt = \int_{t_n}^{t_{n+1}} \bar{S} \Delta v dt$$

*The source term and fluxes are assumed as average value taken at a time level  $\tau$  with  $t_n \leq \tau \leq t_{n+1}$*

$$\rho (\phi_P^{n+1} - \phi_P^n) \Delta v + \left( \sum_f \vec{J}_f \cdot \vec{S}_f \right) (t_{n+1} - t_n) = (\bar{S} \Delta v) (t_{n+1} - t_n)$$

*time step  $\Delta t = t_{n+1} - t_n$*

*the average values may be calculated as follows*

$$\left( \sum_f \vec{J}_f \cdot \vec{S}_f \right) = (1 - f_b) \left( \sum_f \vec{J}_f \cdot \vec{S}_f \right)^n + f_b \left( \sum_f \vec{J}_f \cdot \vec{S}_f \right)^{n+1}$$

$$(\bar{S} \Delta v) = \left\{ (1 - f_b) \bar{S}^n + f_b \bar{S}^{n+1} \right\} \Delta v$$

$$0 \leq f_b \leq 1$$

$\Delta v = \Delta x_w \Delta x_e = \Delta x \Delta y$

So, let us write, so, it will be now,  $t_n$  to  $t_n + 1$ ,  $d$  of  $dt$   $\rho \phi_P$ ,  $\Delta v$  this is actually volume of the main cell  $P$ ,  $\Delta v$  is nothing but  $\Delta v$  at  $p$  and that is nothing but for this particular case  $\Delta x$  into  $\Delta y$ .

The second term  $t_n$  to  $t_n + 1$ , summation of  $f \mathbf{J}_f \cdot \mathbf{S}_f$ ,  $dt$ , here  $dt$  is missing so, you have to write  $dt$  here is equal to  $t_n$  to  $t_n + 1$   $\bar{S} \Delta v dt$ . So, now see we have to integrate each term from the time level  $t_n$  to  $t_n + 1$ . So, obviously if you want to

integrate, so we have to use some average value at tau which tau maybe in between  $t_n$  and  $t_n + 1$ .

So, the source term and fluxes are assumed as average value taken at a time level tau with  $t_n$  less than or equal to tau less or equal to  $t_n + 1$ . So, you can see that we are going from  $t_n$  to  $t_n + 1$  and we have to take the average value of this source term and fluxes at a time level tau, so tau maybe in between  $t_n$  and  $t_n + 1$ .

So we can write now the first term. So, if you integrate, so you can see, so these if you integrate over the time, so  $dt dt$  will cancel then it will be  $d$  of  $\rho \phi P$  and  $\Delta v$  does not depend on time. So obviously you can write  $\rho \phi P_{n+1} - \rho \phi P_n \Delta v$ . So, we are now marching in time from time level  $n$  to  $n + 1$ , so  $n$  is the previous time level and this is the present time level.

Now, second term, you see, so we will take the average value,  $S_f$  at time tau and as we integral over  $\Delta t$ , so it will be  $t_n + 1$ , minus  $t_n$ , is equal to  $\bar{S} \Delta v$ . So, this also, we have to take at a particular time tau,  $t_n + 1$ , minus  $t_n$ . So, now let us say that,  $\Delta t$ , the time step,  $\Delta t$  is equal to obviously  $t_n + 1$ , minus  $t_n$  and so now, how will you calculate this value of this fluxes and source term at time level tau.

So, obviously you have to take either from  $t_n$  or  $t_n + 1$ , because you are having the values at time level  $t_n$  or at  $t_n + 1$  which is your current time level. So, all these fluxes and the source term you have to take the average between these two points.

So, the average value, maybe calculated as follows. So, summation of  $J_f \cdot S_f$  at time level tau, will calculate as  $(1 - f_t)$ , summation of  $J_f \cdot S_f$  at time level  $n$  plus  $f_t$  summation  $f J_f \cdot S_f$  at time level  $n + 1$ . So, you can see, we are considering these average value which is at time level tau, because we have the values available either at  $n$  time level or  $n + 1$  time level, because we are moving from  $n$  to  $n + 1$ .

So, obviously, we are taking a factor  $(1 - f_t)$ , of  $n$ th time level and  $f_t$  times with  $n + 1$  time level. So, obviously, you can see that if  $f_t$  is 0, then obviously, you will get only a time level  $n$  and if  $f_t$  is 1, then you will take only from  $n + 1$  time level. So similarly,  $S \Delta v$ , you can take as,  $(1 - f_t) S$  at  $n$  time level plus  $f_t$  at  $S_{n+1}$ , into  $\Delta v$ .

So, you can see  $f_t$  lies between 0 and 1. So  $f_t$  is a factor which you are taking the average between  $n$  and  $n + 1$  time level values.

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**Finite Volume Discretization of Diffusion Equation**

$$\rho(\phi_P^{n+1} - \phi_P^n) \Delta V + \left\{ (1-f_t) \left( \sum_f J_f \bar{S}_f \right)^n + f_t \left( \sum_f J_f \bar{S}_f \right)^{n+1} \right\} \Delta t$$

$$= \left\{ (1-f_t) \bar{S}^n + f_t \bar{S}^{n+1} \right\} \Delta V \Delta t$$

*Divide both sides by Δt*

$$\frac{\rho \Delta V}{\Delta t} (\phi_P^{n+1} - \phi_P^n) + \left[ (1-f_t) \left\{ -D_e (\phi_E^n - \phi_P^n) + D_w (\phi_P^n - \phi_W^n) \right. \right.$$

$$\left. - D_n (\phi_N^n - \phi_P^n) + D_s (\phi_P^n - \phi_S^n) \right\} + f_t \left\{ -D_e (\phi_E^{n+1} - \phi_P^{n+1}) \right.$$

$$\left. + D_w (\phi_P^{n+1} - \phi_W^{n+1}) - D_n (\phi_N^{n+1} - \phi_P^{n+1}) + D_s (\phi_P^{n+1} - \phi_S^{n+1}) \right\} \Delta V$$

$$= \left\{ (1-f_t) \bar{S}^n + f_t \bar{S}^{n+1} \right\} \Delta V$$

$$\left\{ \frac{\rho \Delta V}{\Delta t} + f_t (D_e + D_w + D_n + D_s) \right\} \phi_P^{n+1} = f_t [D_e \phi_E^{n+1} + D_w \phi_W^{n+1}$$

$$+ D_n \phi_N^{n+1} + D_s \phi_S^{n+1}] + (1-f_t) (D_e \phi_E^n + D_w \phi_W^n + D_n \phi_N^n + D_s \phi_S^n) \Delta V$$

$$+ \left[ \frac{\rho \Delta V}{\Delta t} - (1-f_t) (D_e + D_w + D_n + D_s) \right] \phi_P^n + f_t \bar{S}^{n+1} \Delta V + (1-f_t) \bar{S}^n \Delta V$$

$$a_P \phi_P^{n+1} = \sum_n a_{nb} \left\{ f_t \phi_{nb}^{n+1} + (1-f_t) \phi_{nb}^n \right\} + \left[ a_P - (1-f_t) \sum_n a_{nb} \right] \phi_P^n + b$$

*where*

$$a_P = \frac{\rho \Delta V}{\Delta t} \quad a_E = D_e = \frac{\Gamma_e A_e}{(\delta x)_e} \quad a_W = D_w = \frac{\Gamma_w A_w}{(\delta x)_w}$$

$$a_N = D_n = \frac{\Gamma_n A_n}{(\delta y)_n} \quad a_S = D_s = \frac{\Gamma_s A_s}{(\delta y)_s}$$

$$a_P = a_P^* + f_t \sum_n a_{nb}$$

$$b = f_t \bar{S}^{n+1} \Delta V + (1-f_t) \bar{S}^n \Delta V$$

So, let us put all the values now,  $\rho \phi_P^{n+1} - \rho \phi_P^n \Delta V$  plus, so now you write  $1 - f_t$ , summation of  $J_f \cdot S_f$ , so which time level we will take? It is at  $n$ th time level, plus  $f_t$  summation of  $J_f \cdot S_f$  at time level  $n + 1$ , and we have  $\Delta t$ , is equal to now source term, we have  $1 - f_t$ ,  $\bar{S}^n$  plus,  $f_t \bar{S}^{n+1}$ . We have  $\Delta V$  and  $\Delta t$ . So now you divide both sides by  $\Delta t$ . So what you will get, divide both sides by  $\Delta t$ , so what you are going to get,  $\rho \Delta V$  by  $\Delta t$   $\phi_P^{n+1} - \rho \phi_P^n$ .

Now,  $\Delta t$  we have divided, so only we will have  $1 - f_t$  and  $f_t$ , and what is  $J_f \cdot S_f$ ? So we know already  $J_f \cdot S_f$  all the faces, so  $J_e \cdot S_e$  plus  $J_w \cdot S_w$ ,  $J_n \cdot S_n$  plus  $J_s \cdot S_s$ . So, that if you write, so you will get minus  $D$ . So, this we have derived in last class,  $\phi_P^{n+1} - \phi_P^n$  plus  $D_w \phi_P^n$ , sorry this time level, do not write.

So,  $\phi_P^{n+1} - \phi_P^n$  plus  $D_n \phi_P^n$  this is minus  $D_n$  and  $\phi_N^{n+1} - \phi_P^{n+1}$  plus  $D_s \phi_P^{n+1} - \phi_S^{n+1}$ , so, this is the fluxes at the faces,  $J_f \cdot S_f$ , so, that already we have discretize in last class so that we have just written here. Now you put these values at  $n$  and  $n + 1$  time level wherever applicable.

So, now you can see this is a time level  $n$  so you can write minus  $D_e \phi_E$ , so we are taking a time level  $n$  minus  $\phi_P^n$  plus  $D_w \phi_P^n$  minus  $\phi_W^n$ , then we have minus  $D_n \phi^n$  minus  $\phi_P^n$  plus  $D_s$ . So, these at time level  $n$   $\phi_P^n$  minus  $\phi_S^n$ . So, now we have this term so you can write plus  $f_t$ , so we can write minus  $D_e$ , so at which time level at  $n+1$  time level so  $\phi_E^{n+1}$  minus  $\phi_P^{n+1}$  plus  $D_w \phi_P^{n+1}$  minus  $\phi_W^{n+1}$  minus  $D_n \phi^n$  at  $n+1$  minus  $\phi_P$  at  $n+1$  plus  $D_s \phi_P^{n+1}$  minus  $\phi_S$  at  $n+1$ .

So, these are the flux terms and we have now this source term, so is equal to  $1 - f_t S$  bar  $n$  plus  $f_t S$  bar  $n$  plus  $f_t S$  bar  $1 \Delta v$ . So, now what we will do, we will take the values of  $\phi_P$  at time level  $n+1$  in the left hand side and all other terms will take in the right hand side so that we can write this equation in a  $\phi_P$  form.

So, you can see where we have  $\phi_P$  at  $n+1$ . So, this 1, this, this, this and this, so this will keep in the left hand side and rest all other terms will keep in the right hand side. So, we will write  $\rho \Delta v$  by  $\Delta t$  so, this is the first coefficient, then if you consider this one, so you will have minus minus plus so it will be plus  $D_e$ , then, so  $D_e$  but we will have  $f_t$  here, so you can write it as  $f_t D_e$ , then we have  $D_w$ , this is your  $D_n$  plus this is your  $D_s$ .

So, these all are at time level  $n+1$   $\phi_P^{n+1}$ . And now rest all other terms you will take in the right hand side. So, what we will write equal to, so equal to, now, first you see the  $n+1$  time level terms. So, we have, so  $f_t$  and you see minus  $D_e \phi$  so, if you take in the right hand side it will become plus.

So, you can write  $f_t$ , so it will be  $D_e \phi_E^{n+1}$ . Then plus this one  $\phi_W$ , so it is minus so this side it will be plus  $D_w \phi_W^{n+1}$ . Then we have plus this one. So,  $D_n \phi_N^{n+1}$  plus this one,  $D_s \phi_S$  at  $n+1$ . So, all the terms at  $n+1$ , so, these you can see all our neighbor times  $\phi_E$ ,  $\phi_W$ ,  $\phi_N$  and  $\phi_S$ .

Now, you take the neighbor terms up at  $n$  time level. So, you can consider this term, this term, this term and this term. So, if you write it, so, you will get  $1 - f_t$  is there, so,

one minus  $f_t$  and we can write  $D_e \phi_E$  at  $n$  plus  $D_w \phi_W$  at  $n$  plus  $D_n \phi_N$  at  $n$  and  $D_n \phi_N$  at sorry,  $\phi_{South}$  at  $n$ .

So, this is all neighbor terms now, we have some terms  $\phi_P$  at time level  $n$ , so, this one, this one, this one, and this one so, this we take in the right hand side. So, if we take in the right hand side you can see we have another term also  $\phi_P$  here. So, all these you take in the right hand side so, you can see, so, this is your minus which will become plus so, if you take right hand side it will be  $\rho \Delta v$  by  $\Delta t$ .

Then you see this is your,  $1 - f_t$ , so, it will become  $1 - f_t$  and again you can see this is your minus minus plus, so obviously right hand side, it is minus and we have put minus here. So you can write  $D_e$  plus  $D_w$  plus  $D_n$  plus  $D_o$ . So, this is your  $\phi_P$  at  $n$  and now the source terms. So you can write  $f_t S_{n+1} \Delta v$  plus  $1 - f_t \bar{S}_n \Delta v$ , so all other terms we have taken in the right hand side.

So, this is the final algebraic equation. So, now let us write in terms of the coefficients, so obviously, the coefficient associated with the  $\phi_P$  at  $n+1$  that is your diagonal coefficient and that will represent that  $a_p$ , so we will represent as  $a_p$ , so, this is your  $a_p \phi_P$  at  $n+1$ .

So, this is your  $a_p$  then you have right hand side you see you can write it as summation of all the face neighbor you have all  $a_{nb}$ . Then we have  $f_t \phi_{nb}$  at  $n+1$ , and also you have  $1 - f_t \phi_{nb}$  at time level  $n$ . So, up to this we have written. So now, you write plus so, this is your, we are denoting as  $a_p$ ,  $\rho \Delta v$  by  $\Delta t$ , then we have  $1 - f_t$  and this is your summation of  $a_{nb}$ , all the neighbors. So, we have now  $\phi_P$  at  $n$  plus the source term  $b$ . So, this is the final algebraic equation which you need to solve and you can see we have the terms associated with  $n$  and  $n+1$  time level.

So, now what are the coefficients? So, we told that we have  $a_p$  is  $\rho \Delta v$  by  $\Delta t$ , what is  $a_E$ ?  $a_E$  is your  $D_e$  which is, we have already calculated  $\gamma_E$ ,  $a_E$  by  $\Delta x_e$  and  $a_W$  is  $D_w \gamma_w$   $a_W$  by  $\Delta x_w$ . Similarly,  $a_N$  is  $D_n$  is equal to  $\gamma_N$   $A_n$  divided by  $\Delta y_n$ . So, this is the difference between  $\phi$  and minus  $\phi_P$ .

So, and a South is  $D_s$  is equal to  $\gamma_s A_s$  by  $\Delta y_s$ , and we have  $a_p$  so  $a_p$  you can write as  $a_p^0$  plus  $\sum a_{nb}$ , and the source term  $b$  you have  $\bar{S}_n \Delta v$  plus  $1 - f_t$   $\bar{S}_n \Delta v$ . So, now let us see some special cases where  $f_t$  is 0,  $f_t$  is 1, and  $f_t$  is half.

So, as you can see that if you put  $f_t$  is equal to 0, obviously, you are going to get only the terms associated with  $n$  in the right hand side. So, obviously, it will be explicit and if it is  $f_t$  equal to 1, then you have both  $n$  and  $n+1$  so, obviously, it will be implicit and if you put  $f_t$  is equal to half, then you are taking the average of at time level  $n$  and time level  $n+1$  values, so obviously it will become Crank Nicolson and you can see that obviously explicit and implicit are these are first order Euler time integration we have done so, it will be order of accuracy is  $\Delta t$ , but Crank Nicolson is  $\Delta t^2$ .

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### Finite Volume Discretization of Diffusion Equation

**Special Cases:**

$f_t = 0$  Explicit Scheme  $O(\Delta t)$

$a_p = a_p^0$

$b = \bar{S}_n \Delta v$

$a_p \phi_p^{n+1} = \sum_{nb} a_{nb} \phi_{nb}^n + (a_p^0 - \sum a_{nb}) \phi_p^n + b$

$f_t = 1$  Implicit Scheme  $O[(\Delta t)^2]$

$a_p = a_p^0 + \sum_{nb} a_{nb}$

$b = \bar{S}_n \Delta v$

$a_p \phi_p^{n+1} = \sum_{nb} a_{nb} \phi_{nb}^{n+1} + a_p^0 \phi_p^n + b$

$f_t = \frac{1}{2}$  Crank-Nicolson Scheme  $O[(\Delta t)^2]$

$a_p = a_p^0 + \frac{1}{2} \sum_{nb} a_{nb}$

$b = \frac{1}{2} (\bar{S}_n^{n+1} + \bar{S}_n^n) \Delta v$

$a_p \phi_p^{n+1} = \sum_{nb} \frac{a_{nb}}{2} (\phi_{nb}^{n+1} + \phi_{nb}^n) + (a_p^0 - \frac{1}{2} \sum_{nb} a_{nb}) \phi_p^n + b$

So now, explicit scheme, special cases. So,  $f_t$  is equal to 0 which is your explicit schemes so if you put  $f_t$  equal to 0. So, obviously, you can see  $a_p$  will be your only  $a_p$  naught, the source term will be  $\bar{S}_n \Delta v$  and the algebraic equation will become  $a_p \phi_P^{n+1}$  is equal to summation of all the neighbors  $a_{nb} \phi_{nb}$  at time level  $n$  plus  $a_p$  naught minus summation of  $a_{nb} \phi_P^n$  plus  $b$ .



So, you can see only  $\phi_P^{n+1}$  is unknown and all other terms are available at time level  $n$ . So, obviously this is explicit schemes and it is very easy to code, easy and convenient to code and but order of accuracy is  $\Delta t$ , so order of accuracy is  $\Delta t$ .

Now, you see  $f_t$  is equal to 1, so it will become implicit methods on implicit schemes. So, in this particular case your diagonal coefficient  $a_p$  will become  $a_p$  naught plus summation of all the neighbor coefficients. And the source term  $b$  will become  $\bar{S}^{n+1} \Delta v$  and the algebraic equation will become  $a_p \phi_P^{n+1}$  is equal to summation of all the neighbors  $a_{nb} \phi_{nb}^{n+1}$  and we have  $a_p$  naught  $\phi_P^{n+1}$  plus  $b$ .

So, you can see here  $n+1$  all the neighbor coefficient diagonal coefficients, sorry, all the neighbor values at  $n+1$ , the diagonal value  $\phi_P^{n+1}$ . So, obviously, it is implicit schemes and order of accuracy is  $\Delta t$ . Now,  $f_t$  is equal to half, if you consider then it will become Crank Nicolson method, Crank Nicolson schemes, so it is better to write scheme.

Now, in this particular case now  $a_p$  will become half of summation of  $a_{nb}$  neighbor  $a_p$  is equal to  $a_p$  naught plus half of summation of all neighbors and the  $b$  will become  $\bar{S}^{n+1}$  by  $S$  plus  $\bar{S}^n \Delta v$ . And the algebraic equation will become  $a_p \phi_P^{n+1}$  is equal to summation of all the neighbor  $a_{nb} \phi_{nb}^{n+1}$  plus  $\phi_{nb}^n$  plus  $a_p$  naught minus half summation of  $a_{nb} \phi_{nb}^n$  plus  $b$ . So, you can see here now half of  $\phi_{nb}^{n+1}$  by  $\phi_{nb}^n$ . So, this will become order of  $\Delta t$  square.

So, in today's class, we considered unsteady diffusion equation in two dimensional situation and we discretize the using finite volume method. So, we integrated the governing equation over control volume as well as over time and we wrote the fluxes, when we integrate over the time as the average value at time  $\tau$ , but the values at  $\tau$  you can take the average from either  $t_n$  time level or  $t_{n+1}$  time level.

So, obviously, because at these two points, we know the values of the variables. So, we took a factor  $f_t$  where we have taken  $1 - f_t$  times the fluxes at  $n$  time level and  $f_t$  times the fluxes at  $n+1$  time level and similarly for the source terms. So, then we have

discretize that equation and wrote the final algebraic equation and also we have written the coefficients.

As a special case, we have put  $f_t$  is equal to 0 and we got the explicit scheme where only one unknown is there  $\phi^P_{n+1}$  and when  $f_t$  is equal to 1 it will become implicit scheme and there are more than one unknowns at a time  $n+1$  level and when we put  $f_t$  is equal to half so it become Crank Nicolson method, we take half from  $n$  time level and  $n+1$  time level of this fluxes and the source terms and the order of accuracy becomes second order. Thank you.