

Computational Fluid Dynamics of Incompressible Flow

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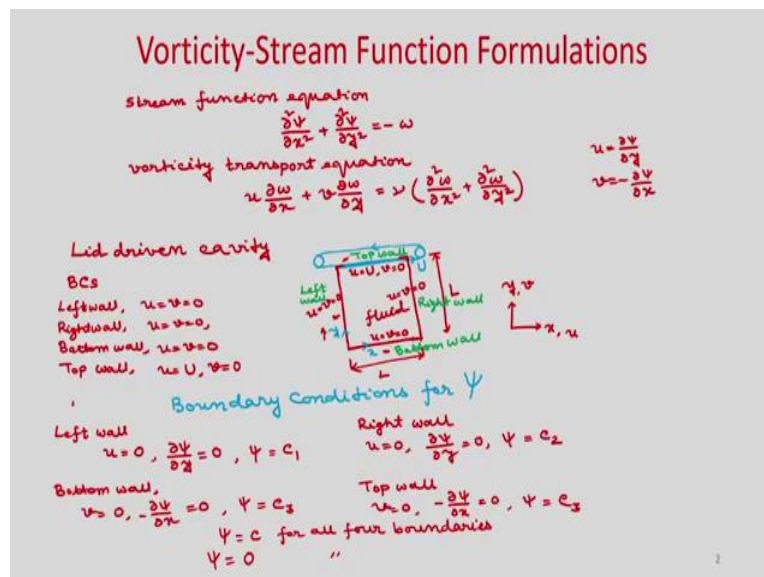
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Lecture 2

Boundary conditions for flow problems

Hello everyone. So, in last lecture, we discretize stream function equation and vorticity transport equation using finite difference method. So, both the equations to discretize both the equations we used central difference method. So, the order of accuracy was Δx^2 , Δy^2 . Today we will consider two problems, first we will consider lid driven cavity problem. Then we will consider flow between parallel plates, plane Poiseuille flow and we will discuss about the Boundary Conditions.

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The slide, titled "Vorticity-Stream Function Formulations", contains the following content:

- stream function equation**
$$\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} = -\omega$$
- vorticity transport equation**
$$u \frac{\partial \omega}{\partial x} + v \frac{\partial \omega}{\partial y} = \nu \left(\frac{\partial^2 \omega}{\partial x^2} + \frac{\partial^2 \omega}{\partial y^2} \right)$$
- Velocity components:
$$u = \frac{\partial \psi}{\partial y}$$

$$v = -\frac{\partial \psi}{\partial x}$$
- Lid driven cavity**
BCs:
Left wall, $u=v=0$
Right wall, $u=v=0$
Bottom wall, $u=v=0$
Top wall, $u=U, v=0$
- Boundary conditions for ψ**
Left wall: $u=0, \frac{\partial \psi}{\partial y} = 0, \psi = c_1$
Right wall: $u=0, \frac{\partial \psi}{\partial y} = 0, \psi = c_2$
Bottom wall: $v=0, -\frac{\partial \psi}{\partial x} = 0, \psi = c_3$
Top wall: $v=0, -\frac{\partial \psi}{\partial x} = 0, \psi = c_4$
 $\psi = c$ for all four boundaries
 $\psi = 0$ "

So, our governing equation stream function equation, stream function equation $\Delta^2 \psi = -\omega$, where ω is your vorticity and vorticity transport equation is, vorticity transport equation $u \frac{\partial \omega}{\partial x} + v \frac{\partial \omega}{\partial y} = \nu \Delta^2 \omega$.

So, now we will derive the boundary conditions for lid driven cavity problem. So, first let us discuss what is lid driven cavity problem? Lid driven cavity problem. So, we have enclosure, square enclosure. So, it is a square enclosure filled with some liquid, some fluid.

So, inside some fluid is there, some fluid is there and this is your square enclosure, say let us say this is your L and this is also L , square enclosure and this upper surface it is connected with some belt let us say and this lid, this is upper plate. So, this is moving with some velocity. So, this belt is moving like this, so we can imagine that these boundary is moving in x direction. So, if it is x and this is your y , then along x direction this lid is moving with some velocity u .

So, you have enclosure, square enclosure left bottom and right surfaces are having no slip boundary conditions, those are stationary walls and the top wall is moving with velocity u . So, in terms of velocity boundary conditions, what are the boundary conditions. So, in terms of velocity boundary conditions.

So, here u is equal to v is equal to 0. The left boundary also u is equal to v is equal to 0 and right also u is equal to v is equal to 0. So, all are no slip condition. Here also it is no slip condition on the top wall. So, it is u is equal to capital U . Because it is moving in the x direction with a velocity U and v is 0.

So, we have applied no slip boundary conditions in all four walls and this is your left wall. So, this is your left wall, this is your left wall, so this is your left wall, this is your right wall. This is your bottom wall and this is your top wall. So, if you write the boundary conditions mathematically for velocity only, then boundary conditions for the problem lid driven cavity.

So, you have left wall, you have u is equal to v is equal to 0, right wall u is equal to v is equal to 0. Bottom wall is u is equal to v is equal to 0 and top wall, u is equal to capital U and v is equal to 0. So, these are the boundary conditions in terms of velocities. But we are solving the governing equations for ψ which is your stream function and vorticity ω .

So, we need to know the boundary conditions in terms of ψ and ω to solve the discretized equation. So, first let us find the boundary conditions for ψ . Stream function boundary conditions. Boundary conditions for ψ , ψ is stream function. So, we need to find what are the boundary conditions for ψ ? So, to find the boundary conditions for ψ what we will do, first we will impose the normal velocity at each boundary and you will find the boundary conditions for ψ .

So, if you will consider left wall, if you consider left wall, what is your normal velocity? Normal velocity will be u , v is your tangential velocity because if this is your x , x direction u have velocity u and in y direction you have velocity v . So, if you consider the left wall then left wall the normal velocity is u .

So, you put normal velocity as 0 that means u is equal to 0. So, now you write u in terms of stream function. So that you know, so the gradient of stream function with respect to y will be 0. So, that means $\frac{\partial \psi}{\partial y}$ will be 0. So, you can see so this is your y .

In this direction it is y , so along this your, the gradient of ψ will be 0, that what does it mean, it means ψ is equal to constant, ψ is equal to constant. Let us say, some constant C_1 . So, from these $\frac{\partial \psi}{\partial y}$ is equal to 0, if you integrate it, then you will get ψ is equal to C_1 , C_1 is some constant.

Similarly, for right wall what is the normal velocity u , u is equal to 0. Similarly, $\frac{\partial \psi}{\partial y}$ is equal to 0 and you can write ψ is equal to C_1 or ψ is equal to some another constant. Let us say ψ is equal to C_2 and if you write the, so on the right wall ψ is equal to C_2 another constant let us say.

Now bottom wall, so for the bottom wall you see here. So, normal velocity will be in y direction that means v is equal to 0. So, you can write v is equal to 0 and v is in terms of ψ it is minus $\frac{\partial \psi}{\partial x}$ is equal to 0 and you can see so along x it is 0. So, ψ is also on the bottom wall will be constant and let us say the constant is C_3 .

Similarly, on the top wall, if you consider top wall, what is the normal velocity v is equal to 0 and you can write minus $\frac{\partial \psi}{\partial x}$ is equal to 0, why? Because u is $\frac{\partial \psi}{\partial y}$ and v is equal to minus $\frac{\partial \psi}{\partial x}$ you know right.

So, minus $\frac{\partial \psi}{\partial x}$ is equal to 0, so ψ is equal to C_3 . So, from the normal velocity boundary conditions we have found that in all boundaries you have some constant stream lines. So, constant stream functions, constant stream functions. So, left wall is C_1 , right wall is C_2 , bottom wall is C_3 and top wall is C_4 . Now you see, so these walls are connected.

So, now the constant stream function denotes to the stream line, we know that constant stream function is a stream line and two stream lines cannot cross each other that means you have. So, if you see right wall, left wall, so in the left wall, so this is continuous and to the bottom walls and bottom wall to right wall and right wall to top wall and top wall to left wall.

So, these are all connected, what does it mean? That means it is a single stream line and the stream function should be same. So, along all these four boundaries psi will be constant. So, you can write psi is equal to C for all four boundaries. So, that means psi is equal to constant and it is a same constant. Because two stream lines cannot cross each other. So, obviously it should be continuous.

So, that means psi is constant single value in all four boundaries and for simplicity let us take that psi is equal to 0 this constant value is 0 on these four walls. So, we are considering psi is equal to 0 for all four boundaries. Because we are considering that constant as 0. Because it does not matter, because you can take any constant. So, it will not effect your solution. So, we are considering the stream function value on the walls as 0 and accordingly your stream function value will be there at the interior points.

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Vorticity-Stream Function Formulations

Boundary Conditions for ω

$$\omega = -\left(\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2}\right)$$

Left wall

$$u = 0 \quad \frac{\partial \psi}{\partial y} = 0 \quad \frac{\partial \psi}{\partial x} = 0$$

$$\omega = -\frac{\partial^2 \psi}{\partial x^2} = -\frac{\psi_{2,j} - 2\psi_{1,j} + \psi_{0,j}}{(\Delta x)^2}$$

$$\omega_{1,j} = -\frac{\psi_{2,j} - 2\psi_{1,j} + \psi_{0,j}}{(\Delta x)^2}$$

Right wall

$$v = 0, \quad -\frac{\partial \psi}{\partial x} = 0, \quad \frac{\partial \psi}{\partial y} = 0$$

$$\frac{\psi_{2,j} - \psi_{0,j}}{2\Delta x} = 0$$

$$\psi_{0,j} = \psi_{2,j}$$

$$\omega_{1,j} = -\frac{\psi_{2,j} - 2\psi_{1,j} + \psi_{0,j}}{(\Delta x)^2} = -\frac{2}{(\Delta x)^2} (\psi_{2,j} - \psi_{1,j})$$

Diagram: A square domain with vertices labeled 0 (top-left), 1 (top-right), 2 (bottom-right), and 0 (bottom-left). The top boundary is labeled $\psi=0$ and the bottom boundary is labeled $\psi=0$. The horizontal axis is labeled x and the vertical axis is labeled y . The grid indices are $i=1, M$ and $j=1, N$.

So, now let us find the boundary conditions for vorticity. Boundary conditions for omega. So, we know that omega in terms of stream function, so minus del 2 psi by del x square plus del 2 psi by del y square and you have this square enclosure. Top wall is

moving, top wall is moving and ψ is equal to 0, we have considered on four walls and let us say that you have the values, i is equal to 1 to M and 1 to N .

So, i is equal to 1 to M it is varying and j is equal to 1 to N . So, i is equal to 1 that means it is j is varying. So, it is a lid boundary and i is equal to M and j is varying, so it is your right boundary. So, now let us consider left wall. So, in the left wall, if you consider, so what is normal velocity? Normal velocity is u is equal to 0. That means $\frac{\partial \psi}{\partial y}$ will be 0.

So, on the left wall $\frac{\partial \psi}{\partial y}$ is 0 and you can see that ψ is constant along this line. Along this line, so its gradient is 0 and its second derivative also will be 0. So, if you take the gradient of $\frac{\partial \psi}{\partial y}$ that will be also 0 along this line. So, $\frac{\partial^2 \psi}{\partial y^2}$ will be 0.

Because it is constant there is no variation and its gradient is 0 everywhere. So, obviously if you would take the derivative with respect to y of this gradient $\frac{\partial \psi}{\partial y}$ that will be also 0 along this left wall. So, $\frac{\partial^2 \psi}{\partial y^2}$ is 0. So, now you can see this equation, vorticity equation. So, if you put $\frac{\partial^2 \psi}{\partial y^2}$ is equal to 0 for the left wall.

So, ω will be minus $\frac{\partial^2 \psi}{\partial x^2}$. So, now you use central difference method to discretize this second derivative. So, if you use second derivative, so it will be, so now you can write ω for which wall, left wall, where i is equal to 1 and j will vary. So, ω $1 j$ will vary from 1 to N . So, for this it will be now minus second derivative, so you discretize using central difference method that means if you have, what we will do. So, if this is the left boundary, so it is your, i is equal to 1.

We will take a fictitious point outside this boundary. So, this is your one grid, and this is your interior grid. So, this is your fictitious point, so here you have i is equal to 1, here i is equal to 2 and this is i is equal to 0. So, and this way it is, if it is j , so it is j plus 1 and this is your j minus 1.

But it is $\frac{\partial^2 \psi}{\partial x^2}$, so it is be for point j only. So, you can write with this fictitious point. So, i is equal to 0 is fictitious point outside the boundary and with the same distant Δx . So, this is also Δx and this is also Δx .

So, if you considered this then you can write discretize form of the second derivative as ψ_j . So, $\psi_{j+2} - 2\psi_{j+1} + \psi_j$ divided by Δx^2 . But we do not know the value of ψ_0 because it is a fictitious point. So, to find the value of ψ_0 we will use another velocity boundary condition which is v is equal to 0.

So, you know that on the left wall v is equal to 0, what does it mean? $-\frac{\partial \psi}{\partial x}$ is 0. That means $\frac{\partial \psi}{\partial x}$ is 0. So, if $\frac{\partial \psi}{\partial x}$ is 0, so you can write the central difference. So, $\psi_{j+2} - \psi_j$ divided by distance between these two points $2\Delta x$. Because here i is equal to 2 and i is equal 0. So, $2\Delta x$ is equal to 0 .

So, we have used central difference. So, now ψ_0 you can find as ψ_0 is equal to ψ_{j+2} . So, now ψ_0 value you know, so you put it in this expression. So, you will get $\psi_{j+1} = \frac{1}{2}(\psi_{j+2} - \psi_j)$ and this is also ψ_{j+2} divided by Δx^2 .

So, that means you can write as $-\frac{2}{\Delta x^2}(\psi_{j+2} - \psi_{j+1})$. So, ψ_{j+1} , where it is ψ_{j+1} on the boundary. So, if you assume ψ is 0, then you can put it 0. But if you can take, if you consider some other values constant values on the wall of ψ then ψ_{j+1} will be having some constant value.

So, we are not going to type 0 here. Because how you are considering the constant depending on that ψ_{j+1} value will depend. So, that means if you considered stream function on the walls, then it will be 0 and if you consider some constant value then you can put it that constant value here. So, we are not going to put it 0 here. So, let us keep in terms of ψ_{j+1} . So, ψ_{j+1} we have found. Similarly, now you do for the right wall.

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Vorticity-Stream Function Formulations

Right wall

$$u=0, \frac{\partial \psi}{\partial y}=0, \frac{\partial^2 \psi}{\partial y^2}=0$$

$$\omega = -\frac{\partial^2 \psi}{\partial x^2}$$


$$\omega_{M,j} = -\frac{\psi_{M+1,j} - 2\psi_{M,j} + \psi_{M-1,j}}{(\Delta x)^2}$$

$$v=0, \frac{\partial \psi}{\partial x}=0$$

$$\frac{\psi_{M+1,j} - \psi_{M-1,j}}{2\Delta x} = 0$$

$$\psi_{M+1,j} = \psi_{M-1,j}$$

$$\omega_{M,j} = -\frac{\psi_{M-1,j} - 2\psi_{M,j} + \psi_{M-1,j}}{(\Delta x)^2}$$

$$\omega_{M,j} = -\frac{2}{(\Delta x)^2} [\psi_{M-1,j} - \psi_{M,j}]$$


So, right wall also, so for the right wall, so if it is right wall, i is equal to M in this case and you have, this is your original interior grid and this is your fictitious grid outside the boundary. So, you can see this is your i, j this is your $i+1, j$ and this is your $i-1, j$. So, $i-1, j$ means $M-1, j$ and this is your, $i+1, j$ means it is your $M+1, j$ and this is your M, j , this is your M, j .

So, similarly we will use the normal velocity. So, normal velocity is u is equal to 0. So, $\frac{\partial \psi}{\partial y}$ is 0 and $\frac{\partial^2 \psi}{\partial y^2}$ is equal to 0 and from the vorticity equation you can write ω is equal to minus $\frac{\partial^2 \psi}{\partial x^2}$. Now if we discretize it, so you will get $\omega_{M,j}$ on the right wall is equal to minus now central difference we are using. So, you can write $\psi_{M+1,j}$ minus twice $\psi_{M,j}$ plus $\psi_{M-1,j}$ divided by Δx^2 .

Now $\psi_{M+1,j}$ is unknown, because it is a fictitious point value. So, we will use again the tangential velocity v is equal to 0 $\frac{\partial \psi}{\partial x}$ is 0 and if you discretize it. So, you will get $\psi_{M+1,j}$ minus $\psi_{M-1,j}$ divided by $2\Delta x$. Because we have considered at the same distance this fictitious point. So, $2\Delta x$ is equal to 0 and $\psi_{M+1,j}$ which is outside the boundary is equal to $\psi_{M-1,j}$ and this value if you put it in this ω expression, then you will get $\omega_{M,j}$ is equal to minus.

So, it will be $\psi_{M+1,j}$ minus twice $\psi_{M,j}$ and plus $\psi_{M-1,j}$ now, this $M-1, j$ you just write because now we are putting the value of $\psi_{M+1,j}$ as $\psi_{M-1,j}$

$1 - j$ here $M - 1 - j$ divided by Δx^2 . So, you are going to get $\omega_{i,j}$ is equal to $-\frac{2}{\Delta x^2} \psi_{i,j} - \frac{1}{\Delta x^2} \psi_{i,j-1}$.

So, in this case, again $\psi_{i,j}$ actually the values of ψ on the right wall and if you choose 0. Then it will be 0, otherwise it will have some other value. Now let us consider the bottom boundary. Similarly, we will consider the fictitious point outside this bottom wall and first we will make the normal velocity 0.

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Vorticity-Stream Function Formulations

Bottom wall

$$u = 0$$

$$\frac{\partial \psi}{\partial x} = 0$$

$$\frac{\partial^2 \psi}{\partial x^2} = 0$$

$$\omega = -\frac{\partial^2 \psi}{\partial y^2}$$


$$\omega_{i,1} = -\frac{\psi_{i,2} - 2\psi_{i,1} + \psi_{i,0}}{(\Delta y)^2}$$

$$v = 0$$

$$\frac{\partial \psi}{\partial y} = 0$$

$$\frac{\psi_{i,2} - \psi_{i,0}}{2\Delta y} = 0 \Rightarrow \psi_{i,0} = \psi_{i,2}$$

$$\omega_{i,2} = -\frac{\psi_{i,2} - 2\psi_{i,1} + \psi_{i,2}}{(\Delta y)^2}$$

$$\omega_{i,1} = -\frac{2}{(\Delta y)^2} (\psi_{i,2} - \psi_{i,1})$$


So, for the bottom wall. So, normal velocity is, so this is your bottom wall and you have the fictitious point outside. So, this is your grid, so this is your j is equal to 1. So, it will be $i, 1$, this is your $i, 0$ and this is your $i, 2$. $i, 2$ is the interior point, $i, 1$ is the, on the wall and $i, 0$ is the fictitious point outside the boundary. So, similarly first we will consider the normal velocity as 0 that means v is the normal velocity. So, it is 0 and from here you will get $\frac{\partial \psi}{\partial x}$ is 0.

So, you can see along this bottom wall ψ is constant, ψ is constant $\frac{\partial \psi}{\partial x}$ is 0 and obviously $\frac{\partial \psi}{\partial x}$ is 0 everywhere. So, if you take the gradient of this $\frac{\partial \psi}{\partial x}$ with respect to x that will be also 0 that means you will get $\frac{\partial^2 \psi}{\partial x^2}$ is equal to 0.

So, your ω equation will be $-\frac{\partial^2 \psi}{\partial y^2}$. So, now if you discretize using central difference method. Then for $\omega_{i,1}$ you can write $\psi_{i,2}$

minus twice psi i 1 and psi i 0 and divided by delta y square. So, this delta y square, delta y is the constant step size.

So, similarly now psi i 0 is the value of psi at the fictitious point. Let us find using the tangential velocity condition. So, tangential velocity is u is equal to 0 that means del psi by del y is equal to 0 and if you discretize del psi by del y using second approximation then you will get psi i 2 minus psi i 0 divided by 2 delta y is equal to 0 and that will give psi i 0 is equal to psi i 2.

So, you can see that fictitious value is in, as the interior points psi i 2. So, now this value you put it in the omega expression. So, you will get omega i 1 is equal to psi i 2 minus twice psi i 1 and psi i 0 is psi i 2 divided by delta y square. So, you will get omega i 1 is equal to minus 2 by delta y square. So, you will get psi i 2 minus psi i 1 and here you can see psi i 1 is the boundary value. So, for our case we have chosen as 0. But you keep it as it is, so that if you have a, some constant value you can use it here so psi i 1.

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Vorticity-Stream Function Formulations

Top boundary

$$\psi = 0$$

$$\frac{\partial \psi}{\partial x} = 0$$

$$\frac{\partial^2 \psi}{\partial x^2} = 0$$

$$\omega = -\frac{\partial^2 \psi}{\partial x^2}$$

$$\omega_{i,N} = -\frac{\psi_{i,N+1} - 2\psi_{i,N} + \psi_{i,N-1}}{(\Delta y)^2}$$

$$u = U$$

$$\frac{\partial \psi}{\partial y} = U \quad \frac{\psi_{i,N+1} - \psi_{i,N-1}}{2\Delta y} = U$$

$$\Rightarrow \psi_{i,N+1} = \psi_{i,N-1} + 2U\Delta y$$

$$\omega_{i,N} = -\frac{\psi_{i,N-1} + 2U\Delta y - 2\psi_{i,N} + \psi_{i,N-1}}{(\Delta y)^2}$$

$$\omega_{i,N} = -\frac{2}{(\Delta y)^2} (\psi_{i,N-1} - \psi_{i,N} + U\Delta y)$$

Similarly, now considered the top boundary, so top boundary, so on the top boundary so you can see this is our boundary, and it is moving in the x direction, this is moving in this x direction with u is equal to capital U. So, now if you take the grid, so this is your fictitious point, this is on the wall and you have interior point.

So, this is your j is equal to N . So, it will be $i, N, i, N - 1$ and $i, N + 1$. So, similarly now first we will use the normal velocity condition. So, what is the velocity condition? Normal velocity is v is equal to 0 , so normal boundary condition is v is equal to 0 that means $\frac{\partial \psi}{\partial x}$ is equal to 0 .

So, along this boundary obviously $\frac{\partial \psi}{\partial x}$ is 0 . So, you can write $\frac{\partial^2 \psi}{\partial x^2}$ is also 0 . So, your omega condition is $-\frac{\partial^2 \psi}{\partial y^2}$ and if you use central difference method, then you can write ω_i, N is equal to $\frac{\psi_{i, N+1} - 2\psi_{i, N} + \psi_{i, N-1}}{\Delta y^2}$.

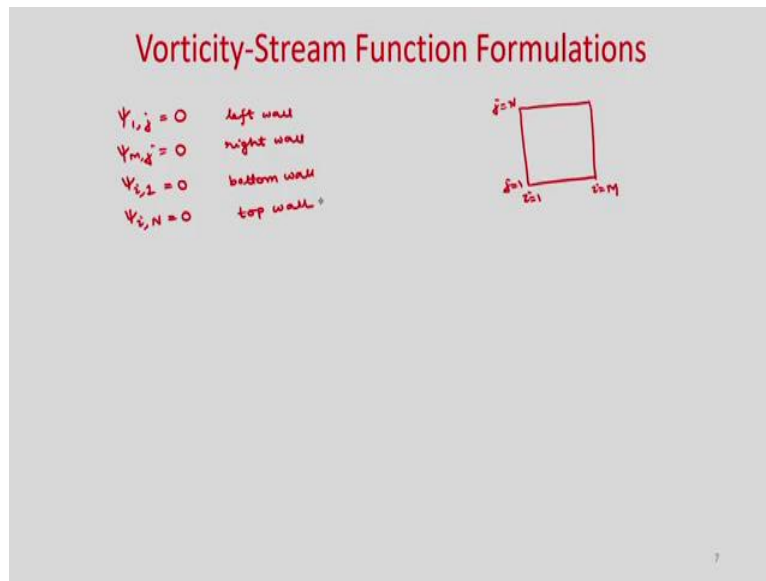
So, now we have to find the value of $\psi_{i, N+1}$ which is the value of ψ at the fictitious point. So, for that now we will use the tangential velocity, so what is tangential velocity u is equal to U . Because it is moving with a velocity u in the positive x direction.

So, that means $\frac{\partial \psi}{\partial y}$ is equal to U . So, that you can now discretize using central difference. So, it will be $\frac{\psi_{i, N+1} - \psi_{i, N-1}}{2\Delta y}$. So, this is your Δy and this is also Δy . So, $2\Delta y$ is equal to U . So, you can write $\psi_{i, N+1}$ is equal to $\psi_{i, N-1} + U \cdot 2\Delta y$, So $2U\Delta y \psi_{i, N+1}$ is equal to $\psi_{i, N-1} + 2U\Delta y$

Now this value you substitute in the omega expression and find the omega i, N . So, omega i, N will be, it will be $-\frac{\psi_{i, N+1} - 2\psi_{i, N} + \psi_{i, N-1}}{\Delta y^2}$. So, now you can see $\psi_{i, N-1} - \psi_{i, N}$, $\psi_{i, N-1}$ so it will be 2 . So, you can write omega i, N is equal to $-\frac{2\Delta y}{\Delta y^2}$. So, it will be $-\frac{2}{\Delta y}$ minus $\psi_{i, N} + 2U$, $U \cdot \Delta y$, $U \cdot \Delta y$, this term $U \cdot \Delta y$, 2 we have taken outside so it will be $2 \cdot \Delta y$.

So, along this top boundary, so the value of omega will be with this expression. Where U is the lead velocity. Now you have seen that to solve the vorticity transport equation we could derive all the boundary conditions for omega as well as ψ we have already found.

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So, psi for all the boundaries this is psi, we have found, so it is i is equal to 1 to i is equal to M and j is equal to 1 to j is equal to N. So, you can see that now for psi, psi is 0 in everywhere. So, it will be psi 1, psi 1 j is equal to 0, psi. So, it is for the left boundary, left wall, psi, so this is your M, so psi M j is equal to 0 for right boundary, right wall, psi i 1 is equal to 0 for bottom wall and psi i N is equal to 0 for top wall, for all i.

So, these are the discretize boundary conditions for psi and also we have shown the discretize boundary conditions for omega. So, you have the discretized equation for the interior points with these boundary conditions, you can solve using some linear solver and you will get the values of psi and omega for this lid driven cavity problem.

But you have to remember that psi and omega are coupled. Because when you are solving the stream function equation you have omega and when you are solving the vorticity transport equation you have psi. So, together this you have to converge and get the solution.

So, in today's lecture we will show later the computer code for this problem lid driven cavity and later we will show the results for this lid driven cavity and plane poiseuille flow. So, now let us consider plane poiseuille flow where the flow is taking place between two infinite parallel plates.

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Vorticity-Stream Function Formulations

Flow between two parallel plates

$\psi = c_1$ on bottom wall
 $\psi = 0$ on " " "

$\psi_2 = c_2$ on top wall

Left wall $u = U$
 $\frac{\partial \psi}{\partial y} = U$
 $\partial \psi = U \partial y \leftarrow$
 $\psi_2 - \psi_1 = U(H - 0)$
 $\Rightarrow \psi_2 = \psi_1 + UH$
 $\Rightarrow \psi_2 = UH$ considering $\psi_1 = 0$

Bottom wall $\psi_{i,1} = 0$ Top wall $\psi_{i,N} = \psi_{i,1} + UH$

Left wall $\psi_{2,j} = \psi_{2,j-1} + U \Delta y$
 $\frac{\partial \psi}{\partial y} = U \Rightarrow \frac{\psi_{2,j} - \psi_{2,j-1}}{\Delta y} = U \Rightarrow \psi_{2,j} = \psi_{2,j-1} + U \Delta y \leftarrow$

Outlet BC $\Rightarrow \psi_{m,i} = \psi_{m,i-1} + U_{m,i} \Delta x$
 $\frac{\partial \psi}{\partial x} = 0 \Rightarrow \psi_{m,i} - \psi_{m,i-1} = 0$ $\psi_{m,i} = \psi_{m,i-1}$

So, let us consider another problem, flow between two parallel plates. So, this is a flow between two parallel plates. So, we are considering two plates, these are two plates, parallel to each other, parallel to each other. Let us say that length of the plate is L and height is H.

This is your domain, this is your computational domain. Now obviously on the wall velocity will be 0. So, that means in terms of velocity if you give, so on the bottom wall u is equal to v is equal to 0. On the top wall also u is equal to v is equal to 0 and let us say that you have a uniform velocity inlet at the inlet.

So, this is your inlet, so you have a uniform velocity. So, this is your uniform velocity. So, this is your capital U with this velocity it is entering. So, now you can see that boundary layers will be developed over the plates and you will get a fully developed flow after a certain distance.

So, and obviously you will get a parabolic profile at the, when it will become fully developed. So, now numerical will solve this problem. So, we are not considering plane poiseuille flow, where we consider two infinite parallel plates and the assumption is that you have fully developed flow.

But here it is a initially it will be a developing flow, after that it will become fully developed. So, at the outlet now we will consider the gradients of velocity will be 0.

Because the flow, whatever it is going we are assuming that it is a fully developed and whatever way it is coming it will go out.

So, that means at the outlet we will consider $\frac{\partial u}{\partial x}$ is equal to $\frac{\partial v}{\partial x}$ is 0. So, it is a fully developed condition and this is your at the outlet. So, it is outflow boundary conditions, outflow boundary conditions. So, inlet U and let us say that you have i is equal to 1, i is equal to M and here j is equal to 1, j is equal to N and obviously your, let us say coordinate is here. So, it is your x and this is your y .

So, the out flow condition is fully developed. So, we have written $\frac{\partial u}{\partial x}$ is equal to, $\frac{\partial u}{\partial x}$ is equal to 0 because x is along the flow direction. So, now you can see that ψ value again we can consider some constant value on the bottom wall.

So, ψ is $C1$ on bottom wall. So, that you can find from the normal velocity. So, normal velocity is v , v is equal to 0 that means $\frac{\partial \psi}{\partial x}$ is 0 and so $\frac{\partial \psi}{\partial x}$ is 0 that means ψ will be constant along x . So, ψ is equal to $C1$ and for simplicity, let us say ψ is equal to 0 on bottom wall.

So, you can take any value. Because it does not matter, so whatever value you will consider on the bottom wall accordingly we will get the ψ on top wall. So, now we have to find what is the value of ψ on the top wall. But anyway from the normal velocity condition you can find that ψ will be some another constant $C2$ on top wall. Because that will be also constant, because from normal velocity v is equal to 0. So, obviously $\frac{\partial \psi}{\partial x}$ is 0, so along this x direction. So, ψ will be constant, so let us say ψ is equal to $C2$ on top wall.

Now how we will find this constant. So, now you know that inlet velocity U from there you can consider that how you can find the $C2$. So, now you can see that what is this condition, so you have the, on the left wall, you have u is equal to U which is your normal velocity that means $\frac{\partial \psi}{\partial y}$ is equal to U. So, that means you can write that $\frac{\partial \psi}{\partial y}$ is equal to U $\frac{\partial \psi}{\partial y}$.

So, if you will integrate it from 0 to H, then you will get. So, if it is ψ_1 and this is your ψ_2 , if it is ψ_1 and this is your ψ_2 then you can get that $\psi_2 - \psi_1$ is equal to U into H. So, y is varying 0 to H, so it will be H minus 0 let us say. So, that

means ψ_2 will be $\psi_1 + UH$. So, ψ_2 value we have found, so ψ_2 is equal to $\psi_1 + UH$.

So, if you consider ψ_1 as 0 then ψ_2 will be UH , for considering ψ_1 is equal to 0, considering ψ_1 is equal to 0. So, in terms of discretize boundary condition you can write ψ , so on the left wall, bottom wall. So, bottom wall is ψ_1 is equal to 0 we have assumed. So, ψ_1 is nothing but, so i, j is equal to 1. So, ψ_i it will vary from 1 to M this is equal to 0 and top wall ψ will be ψ_i, j is equal to N . So, it will be U into H , or you can write ψ_{i+1} plus U into H .

So, for a given flow inlet velocity U and the distance between two parallel plates each you can find the value of ψ_{iN} . Now we need to find what is the value of ψ on the left wall? So, left wall you can find from the integration of this equation. So, that means you can write ψ , so i is equal to 1, so ψ_{1j} is nothing but $\psi_{1j-1} + U \Delta y$.

So, Δy is the grid point, so you can see. So, it is your left wall, left boundary, inlet boundary. So, you have, if this is your $1j$ and this is your $1j-1$. So, you can write, so from this integration you can write from here. You can write ψ_{1j} is equal to $\psi_{1j-1} + U \Delta y$.

So, we have used first order scheme. So, you can see $\frac{\Delta \psi}{\Delta y}$. So, $\frac{\Delta \psi}{\Delta y}$ is you have written ψ . So, you can write here also $\frac{\Delta \psi}{\Delta y}$ is equal to U that means $\psi_{1j} - \psi_{1j-1}$. So, backward difference we have used, we have used backward difference and distance between these two points is Δy .

So, is equal to U , so you can write ψ_{1j} is equal to $\psi_{1j-1} + U \Delta y$ and on the outlet boundary. So, at outlet what you can write for the ψ value. Outlet boundary conditions. So, these you can actually find from this equation only so you can write ψ_i is equal to M, Mj is equal to $\psi_{Mj-1} +$, if you know the value of U , so then you can find u , small u .

So, you have to find at which point, at Mj into Δy . So, this way also you can use. But better condition is that it is fully developed, so $\frac{\Delta \psi}{\Delta x}$ also will be 0 so that you can write. $\frac{\Delta \psi}{\Delta x}$ is equal to 0, so $\frac{\Delta \psi}{\Delta x}$ is equal to 0, so $\frac{\Delta \psi}{\Delta x}$ is equal to 0, so you can write ψ , so j will vary. So, it will ψ_j will

vary so it will be M_j is equal to or minus $\psi_{M-1,j}$. So, if you use first order, so it will be Δx is equal to 0 that means $\psi_{M,j}$ will be $\psi_{M-1,j}$.

So, it will be easy to use this condition, so this you can $\psi_{M,j}$ is equal to $\psi_{M-1,j}$. So, we have found all the stream function boundary conditions in all four boundaries. So, on the bottom wall we have considered some constant value of ψ that is 0.

So, we can see here and on the top wall you can calculate from the given inlet boundary condition and the distance between two parallel plates. Left wall you can find from this relation and outlet it is easier to use this one. So, you can use $\psi_{M,j}$ is equal to $\psi_{M-1,j}$.

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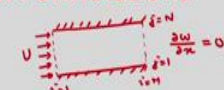
Vorticity-Stream Function Formulations

Bottom wall
 $\omega_{i,2} = -\frac{2}{(\Delta y)^2} (\psi_{i,2} - \psi_{i,1})$

Top wall
 $\omega_{i,N} = -\frac{2}{(\Delta y)^2} (\psi_{i,N} - \psi_{i,N-1})$

Outlet boundary
 $\frac{\partial \omega}{\partial x} = 0 \quad \frac{\omega_{M,j} - \omega_{M-1,j}}{\Delta x} = 0$
 $\Rightarrow \omega_{M,j} = \omega_{M-1,j}$

Inlet boundary
 $u = U$
 $\frac{\partial \psi}{\partial x} = U$
 $\frac{\partial^2 \psi}{\partial x^2} = 0$
 $v = 0$
 $\frac{\partial \psi}{\partial y} = 0$
 $\frac{\psi_{2,j} - \psi_{0,j}}{2\Delta x} = 0 \Rightarrow \psi_{0,j} = \psi_{2,j}$



$\omega = -\left(\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2}\right)$

$\omega = -\frac{\partial^2 \psi}{\partial x^2} - 2\psi_{1,j} + \psi_{0,j} + \psi_{2,j}$

$\omega_{2,j} = -\frac{2}{(\Delta x)^2} (\psi_{2,j} - \psi_{1,j})$

$\omega_{1,j} = -\frac{2}{(\Delta x)^2} (\psi_{2,j} - \psi_{1,j})$

Vorticity-Stream Function Formulations

Bottom wall

$$v = 0$$

$$\frac{\partial \psi}{\partial x} = 0$$

$$\frac{\partial^2 \psi}{\partial x^2} = 0$$

$$\omega = -\frac{\partial^2 \psi}{\partial y^2}$$

$$\omega_{i,1} = -\frac{\psi_{i,2} - 2\psi_{i,1} + \psi_{i,0}}{(\Delta y)^2}$$

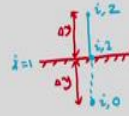
$$u = 0$$

$$\frac{\partial \psi}{\partial y} = 0$$

$$\frac{\psi_{i,2} - \psi_{i,0}}{2\Delta y} = 0 \Rightarrow \psi_{i,0} = \psi_{i,2}$$

$$\omega_{i,1} = -\frac{\psi_{i,2} - 2\psi_{i,1} + \psi_{i,2}}{(\Delta y)^2}$$

$$\omega_{i,1} = -\frac{2}{(\Delta y)^2} (\psi_{i,2} - \psi_{i,1})$$



5

Vorticity-Stream Function Formulations

Top boundary

$$v = 0$$

$$\frac{\partial \psi}{\partial x} = 0$$

$$\frac{\partial^2 \psi}{\partial x^2} = 0$$

$$\frac{\partial^2 \psi}{\partial x^2} = 0$$

$$\omega = -\frac{\partial^2 \psi}{\partial y^2}$$

$$\omega_{i,N} = -\frac{\psi_{i,N+1} - 2\psi_{i,N} + \psi_{i,N-1}}{(\Delta y)^2}$$

$$u = U$$

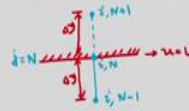
$$\frac{\partial \psi}{\partial y} = U$$

$$\frac{\psi_{i,N+1} - \psi_{i,N-1}}{2\Delta y} = U$$

$$\Rightarrow \psi_{i,N+1} = \psi_{i,N-1} + 2U\Delta y$$

$$\omega_{i,N} = -\frac{\psi_{i,N-1} + 2U\Delta y - 2\psi_{i,N} + \psi_{i,N-1}}{(\Delta y)^2}$$

$$\omega_{i,N} = -\frac{2}{(\Delta y)^2} (\psi_{i,N-1} - \psi_{i,N} + U\Delta y)$$



6

Now similarly, you can find the boundary condition for omega. For omega so if you use, say bottom wall and top wall it will be almost similar whatever we have derived for the lid driven cavity. So, this is your parallel plates and you have inlet velocity U and this is your, i is equal to 1. So, this is your, i is equal to M and j is equal to 1 and j is equal to N. So, for the bottom wall you can write the similar way, so you can see which way you have written for the bottom.

So, bottom wall obviously the no slip condition so v will be 0. So, this equation you can write, omega i 1, so you can write omega i 1, it is bottom, so it will be, so bottom one will be omega, so j is equal to 1. So, i 1 it will be minus 2 by delta y square, so it will be psi i 2 minus psi i 1. This is your omega boundary condition same as the earlier and for the top boundary. So, top boundary also you can write similar way only.

So, top wall will be ω . So, j is equal to N , so $i = N$ is equal to -2 by Δy square and you will get $\psi_{i=N} - \psi_{i=N-1}$.

So, $\omega_{i=N}$ you can find like this. Now outlet boundary, so outlet boundary we will consider, at outlet boundary we will consider fully developed flow, so if it is a fully developed flow you can write $\frac{\partial \omega}{\partial x} = 0$. So, $\frac{\partial \omega}{\partial x} = 0$, so it is a fully developed condition. So, $\frac{\partial \omega}{\partial x}$ will be 0 and you can use first order accurate scheme. So, it will $\omega_{M,j} - \omega_{M-1,j}$ divided by Δx will be 0.

So, that means $\psi_{i=N}$, $\omega_{M,j}$ will be just $\omega_{M-1,j}$. You can use also second order, one sided difference also. So, but we have used for simplicity the first order equate method. So, you can write $\omega_{M,j}$ is equal to $\omega_{M-1,j}$. What about the inlet boundary?

So, similarly for inlet boundary, so your u is U normal velocity. So, you can write $\frac{\partial \psi}{\partial y} = U$. Now you can take the derivative, so it is U is equal to constant. So, you can write $\frac{\partial^2 \psi}{\partial y^2} = 0$. So, as you have a constant velocity inlet. So, $\frac{\partial^2 \psi}{\partial y^2}$ will be 0.

So, from the vorticity equation ω is equal to $-\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2}$. So, you can put $\frac{\partial^2 \psi}{\partial y^2} = 0$. So, ω is equal to $-\frac{\partial^2 \psi}{\partial x^2}$. So, now you can discretize ω , so it will be i is equal to 1.

So, 1_j is equal to $-\omega$, so it will be ψ . So, if you take a fictitious point, so this is your fictitious and this is your interior grid separated with a uniform step size, Δx , this is also Δx . So, this is your 1_j , this is your 2_j and this is your 0_j . So, you can write $\psi_{2,j} - 2\psi_{1,j} + \psi_{0,j}$.

So, $\psi_{0,j}$ we do not know. So, you can find from tangential velocity. So, that means v is equal to 0. So, that means $\frac{\partial \psi}{\partial x} = 0$, so it means $\frac{\partial \psi}{\partial x}$ will be 0. Then you can write $\psi_{2,j} - \psi_{1,j}$ divided by $2\Delta x = 0$ and that implies, $\psi_{2,j} = \psi_{0,j}$. So, distance between these points $2\Delta x$, so it is 0, so $\psi_{0,j}$ is nothing, but $\psi_{2,j}$.

So, that means ω_{1j} will be $-\frac{2}{\Delta x^2}$, $\frac{2}{\Delta x^2}$. So, it will be ψ_{0j} you put and ψ_{2j} here. So, you will get $\psi_{2j} - \psi_{1j}$. So, now we have found the boundary conditions for ω in all boundaries and in discretize form we have written. Thank you.