

Computational Fluid Dynamics for Incompressible Flows
Professor. Amaresh Dalal
Department of Mechanical Engineering
Indian Institute of Technology, Guwahati
Lecture 1

Discretization vorticity-stream function equations using FDM

Hello everyone. So, in last few lectures we have learned the discretization of hyperbolic equations, parabolic equations and elliptic equations. So, we have learned some explicit and implicit schemes and you know that the finite difference approximation of a certain order to discretize any derivative.

So, now, today we will discretize the Stream Function Vorticity Equation, just to solve some fluid flow problem. So, we will consider steady incompressible flow and we will consider momentum equation first. Then we will convert it into stream function vorticity equations. Then we will discretize that equations using finite difference method. So, you know what is vorticity, right from fluid mechanics knowledge. So, vorticity is twice of the angular velocity or curl of velocity.

(Refer Slide Time: 01:38)

Vorticity-Stream Function Formulations

Vorticity, $\omega = \omega_z = \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}$ 2-D
 Steady
 Newtonian
 Incompressible flow

Continuity eqn
 $\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$

Momentum equations
 x-mom : $u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = -\frac{1}{\rho} \frac{\partial p}{\partial x} + \nu \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right)$ (1) $-\frac{1}{\rho} \frac{\partial p}{\partial x}$
 y-mom : $u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} = -\frac{1}{\rho} \frac{\partial p}{\partial y} + \nu \left(\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} \right)$ (2) $-\frac{1}{\rho} \frac{\partial p}{\partial y}$

$\frac{\partial}{\partial y} [Eq.(1)] - \frac{\partial}{\partial x} [Eq.(2)]$

Vorticity equation
 $u \frac{\partial \omega}{\partial x} + v \frac{\partial \omega}{\partial y} = \nu \left(\frac{\partial^2 \omega}{\partial x^2} + \frac{\partial^2 \omega}{\partial y^2} \right)$
 convection terms Diffusion terms ν - Kinematic viscosity
 ω - vorticity

So, in two dimensional situation you can write, the vorticity, vorticity will denote at omega. So, we are considering two dimensional. So omega z, so this is your del v del x minus del u del y. So, this omegas, we are just representing as vorticity and we are considering 2D, steady, Newtonian and incompressible flow.

So, with these assumptions let us write the governing equations of free flow. So, that you know that one is continuity equation. Then we have two momentum equations. So, we have continuity equation $\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$. So, this is your continuity equation and you have momentum equations. So, momentum equations now x momentum, x momentum equations is steady we are writing. So, $u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = -\frac{\partial p}{\partial x} + \nu \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2}$.

So, this 1 by row you can write minus 1 by row and this is your ν by row, so kinematic viscosity $\nu \frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2}$. So x momentum equation, so we considered u velocity and y momentum equation you can write $u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} = -\frac{\partial p}{\partial y} + \nu \frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2}$. So, in this equation there are 3 equations and 3 variables u, v and p, we have to solve. So, and fluid properties ρ and dynamic viscosity μ are known.

So, to eliminate the pressure from this term because we want to solve by simplified equations which conductive diffusive equation first. So, we will eliminate pressure from these momentum equations. To eliminate pressure, so what you do? You just take that derivative with respect to y of this equation 1 and you take the derivative with respect to x of this equation 2.

So, what you do? You take the derivative with respect to y, with respect to y of equation 1 equation 1 and subtract the second equation with respect to, so you take the derivative with respect to x of equation 2. Then you if you subtract then you can see that the pressure term will get eliminated, so you can see this will get $\frac{\partial p}{\partial x} \frac{\partial}{\partial y} - \frac{\partial}{\partial x} \frac{\partial p}{\partial y}$ and here you will get $-\frac{\partial}{\partial x} \frac{\partial p}{\partial y}$.

So, you can see that if you subtract then this will be eliminated and if you rearrange then you will get the equations vorticity equations. So, you will get $u \frac{\partial \omega}{\partial x} + v \frac{\partial \omega}{\partial y} = \nu \frac{\partial^2 \omega}{\partial x^2} + \frac{\partial^2 \omega}{\partial y^2}$.

So, from these two momentum equations, we could derive the vorticity equation just eliminating the pressure. So, you can see it is u is the x direction velocity and v is the y direction velocity and ω is the vorticity and ν is the kinematic viscosity. So, here ν is kinematic viscosity and ω is in this case vorticity. So, we could derive this vorticity equation and which is

conductive diffusive equation. Because in the left hand side convection term and this is your convection term and this is your diffusion term and this is your vorticity equation.

But how to find u and v now? Because vorticity equation you can solve for vorticity omega. But you need to know values of u and v. So, for that we will use stream function. So, you know what is stream function, stream function is a function which satisfies the continuity equation such that u is equal to del psi by del y and v is equal to minus del psi del x.

(Refer Slide Time: 08:16)

Vorticity-Stream Function Formulations

Stream function, ψ

$$u = \frac{\partial \psi}{\partial y} \quad v = -\frac{\partial \psi}{\partial x}$$

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = \frac{\partial}{\partial x} \left(\frac{\partial \psi}{\partial y} \right) - \frac{\partial}{\partial y} \left(\frac{\partial \psi}{\partial x} \right) = 0$$

$$\omega = \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} = \frac{\partial}{\partial x} \left(-\frac{\partial \psi}{\partial x} \right) - \frac{\partial}{\partial y} \left(\frac{\partial \psi}{\partial y} \right) = -\left(\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} \right)$$

Equation for stream function,

$$\Rightarrow \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} = -\omega$$

$$u \frac{\partial \omega}{\partial x} + v \frac{\partial \omega}{\partial y} = \nu \left(\frac{\partial^2 \omega}{\partial x^2} + \frac{\partial^2 \omega}{\partial y^2} \right) \quad u = \frac{\partial \psi}{\partial y} \quad v = -\frac{\partial \psi}{\partial x}$$

Vorticity-Stream Function Formulations

Vorticity, $\omega = \omega_z = \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}$ 2-D
Steady
Newtonian
Incompressible flow

Continuity eqn

$$\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$$

Momentum equations

x-mom : $u \frac{\partial u}{\partial x} + v \frac{\partial u}{\partial y} = -\frac{1}{\rho} \frac{\partial p}{\partial x} + \nu \left(\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} \right)$ (1) $-\frac{\partial p}{\partial x}$

y-mom : $u \frac{\partial v}{\partial x} + v \frac{\partial v}{\partial y} = -\frac{1}{\rho} \frac{\partial p}{\partial y} + \nu \left(\frac{\partial^2 v}{\partial x^2} + \frac{\partial^2 v}{\partial y^2} \right)$ (2) $-\frac{\partial p}{\partial y}$

$$\frac{\partial}{\partial y} [Eq.(1)] - \frac{\partial}{\partial x} [Eq.(2)]$$

Vorticity equation

$$u \frac{\partial \omega}{\partial x} + v \frac{\partial \omega}{\partial y} = \nu \left(\frac{\partial^2 \omega}{\partial x^2} + \frac{\partial^2 \omega}{\partial y^2} \right)$$

convection terms Diffusion terms

ν - Kinematic viscosity
 ω - vorticity

So, we define stream function, so stream function is psi. So, it satisfied the continuity equation such that u is equal to del psi by del y and v is equal to minus del psi by del x. So, you see

whether it is satisfied continuity equation or not. So, you have $\frac{\partial u}{\partial x} + \frac{\partial v}{\partial y}$ is the plus $\frac{\partial v}{\partial y}$, if it is 0 then it satisfies the continuity equation. So, if it is, so $\frac{\partial}{\partial x} \left(\frac{\partial \psi}{\partial y} \right) - \frac{\partial}{\partial y} \left(\frac{\partial \psi}{\partial x} \right)$. So, obviously this is equal to 0.

So, that means it satisfies the continuity equation. So, now we know the definition of vorticity. So, vorticity already we have defined ω is equal to $\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}$. So, now you substitute v here. So, $\frac{\partial}{\partial x}$ what is v ? v is nothing but $-\frac{\partial \psi}{\partial x}$ and $-\frac{\partial}{\partial y}$ what is u ? u is $\frac{\partial \psi}{\partial y}$.

So, you can see that it is $-\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2}$. That means you can write as $\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} = -\omega$. So, this is the equation for stream function, equation for stream function. So, now this equation this is your obviously elliptic type equation. So, it is Poisson equation, so for a known value of vorticity ω you can solve this equation can find the stream function ψ .

Once you know the ψ , you can find the velocities. Because u, v you can represent as a gradient of ψ . So, from these expressions u is equal to $\frac{\partial \psi}{\partial y}$ and v is equal to $-\frac{\partial \psi}{\partial x}$ you can find the velocities once you solve this equation ψ is known.

So, now if you go back and see the vorticity equation. Now, in vorticity equation u, v you can get from the stream function. So, that means you need to solve two equations, one is your stream function equation, $\nabla^2 \psi = -\omega$ and the vorticity equation which is your convective diffusive equation.

So, that means now your this set of equations. So, you have $u \frac{\partial \omega}{\partial x} + v \frac{\partial \omega}{\partial y} = \nu \nabla^2 \omega$. So, this is your vorticity equation and this is your stream function equation and you can find u, v from $\frac{\partial \psi}{\partial y}$ and v is equal to $-\frac{\partial \psi}{\partial x}$.

So obviously, we have written the governing equations in dimensional form and these equations if you solve. Then you will be able to get the velocity distribution in a domain. So, we will discuss two problems and we will discuss about the boundary conditions and we will show some results after solving those, these equations for a particular problem.

So, now we need to discretize this equation using finite difference method. So, you have two different equations. One is stream function equation and one is vorticity equations and now these two equations used as discretize, using finite difference method you know different scheme. So, depending on your choice you can use some scheme and discretize using finite difference method. So, first, let us discretize the stream function equation using finite difference method.

(Refer Slide Time: 13:10)

Vorticity-Stream Function Formulations

stream function equation

$$\frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} = -\omega$$

$$\frac{\psi_{i+1,j} - 2\psi_{i,j} + \psi_{i-1,j}}{(\Delta x)^2} + \frac{\psi_{i,j+1} - 2\psi_{i,j} + \psi_{i,j-1}}{(\Delta y)^2} = -\omega_{i,j}$$

$$\psi_{i+1,j} - 2\psi_{i,j} + \psi_{i-1,j} + \beta^2 (\psi_{i,j+1} - 2\psi_{i,j} + \psi_{i,j-1}) = -\omega_{i,j} (\Delta x)^2$$

$$\beta^2 \psi_{i,j+1} + \psi_{i-1,j} - 2(1+\beta^2)\psi_{i,j} + \psi_{i+1,j} + \beta^2 \psi_{i,j-1} = -\omega_{i,j} (\Delta x)^2$$

$$\psi_{i,j} = \frac{\psi_{i+1,j} + \psi_{i-1,j} + \beta^2 (\psi_{i,j+1} + \psi_{i,j-1}) + (\Delta x)^2 \omega_{i,j}}{2(1+\beta^2)}$$

So, you have stream function equation, del 2 psi by del x square plus del 2 psi by del y square is equal to minus omega. So, we will take a grid 2 dimensional grid. So, you have ij i plus 1 j and this is your, i minus 1 j this is your i j plus 1 and this is your i j minus 1 and we will consider uniform grid.

So, that means here the delta x is constant, this is your delta x, this is your delta x and this is your delta y and this is also delta y. So, delta x may not be equal to delta y. So, we will define the beta is equal to delta x by delta y, the ratio of the step size in x and y direction delta x by delta y we are denoting as beta.

So, this is second derivative of psi. So, second derivative, so what we will use? We will use central difference approximation to discretize this second derivative. So, if you use second order method or central difference, so which is second order. So, you can write psi i plus 1 j minus 2 psi i j plus psi i minus 1 j divided by delta x square.

So, this is second order accurate this is you know, this approximation. Similarly in $\Delta y^2 \psi$ by Δy^2 . So, $\psi_{i,j+1} - 2\psi_{i,j} + \psi_{i,j-1}$ divided by Δy^2 is equal to $-\omega_{i,j}$. So, this is the discretized equation, now you just write $\psi_{i,j+1} - 2\psi_{i,j} + \psi_{i,j-1}$ and we have defined $\beta = \Delta x / \Delta y$. So, we can write $\beta^2 \psi_{i,j+1} - 2\psi_{i,j} + \psi_{i,j-1} = -\omega_{i,j} \Delta x^2$.

Because Δx^2 I have multiplied both side. So, Δx^2 by Δy^2 I have denoted as β^2 and Δx^2 is coming here $-\omega_{i,j}$ into Δx^2 . So, after simplification you can write it as $\psi_{i,j+1} - 2\psi_{i,j} + \psi_{i,j-1} = -\beta^2 \omega_{i,j} \Delta x^2$. So, this is your β^2 and $\psi_{i,j+1} - 2\psi_{i,j} + \psi_{i,j-1} = -\beta^2 \omega_{i,j} \Delta x^2$.

So, this is the discretize equation. So, now you have to choose whether you will use implicit method or explicit method to solve it. So, you can see that $\psi_{i,j}$ you need to find. So, if you explicit type method if you use then you can write $\psi_{i,j}$, which is your unknown. So, $\psi_{i,j}$ is equal to you can write $\psi_{i,j+1} - 2\psi_{i,j} + \psi_{i,j-1} = -\beta^2 \omega_{i,j} \Delta x^2$. So, these I am taking this side was $\beta^2 \omega_{i,j} \Delta x^2$ and this minus it will become plus. So, it will be $\Delta x^2 \omega_{i,j} / (2 + \beta^2)$.

So, if you add using some solver which is explicit type solver. Then you can solve for $\psi_{i,j}$ at time at iteration level $k+1$ and all you can take at best available value. So, here you can see that although you are solving for a ψ . But you have to know the value of ω and this ω you have to get it from the vorticity equation and you can see that this way, this two equation are couple. Because when you are solving, the stream function equation you need to know the value of vorticity ω .

Similarly, when you are solving for Vorticity equation you need to know the velocity is u, v which in turn actually you need to know the stream function. So, these are coupled equations, so you have to solve simultaneously. So, now we have discretized this equation, stream function equation and $\psi_{i,j}$ you can solve using this expression for the known value of $\omega_{i,j}$. Now let us discretize the vorticity equation.

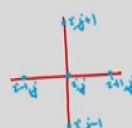
(Refer Slide Time: 19:30)

Vorticity-Stream Function Formulations

Vorticity equation

$$u \frac{\partial \omega}{\partial x} + v \frac{\partial \omega}{\partial y} = \nu \left(\frac{\partial^2 \omega}{\partial x^2} + \frac{\partial^2 \omega}{\partial y^2} \right)$$

$$u = \frac{\partial \psi}{\partial y} \quad u_{i,j} = \frac{\psi_{i,j+1} - \psi_{i,j-1}}{2\Delta y}$$

$$v = -\frac{\partial \psi}{\partial x} \quad v_{i,j} = -\frac{\psi_{i+1,j} - \psi_{i-1,j}}{2\Delta x}$$


$$u_{i,j} \frac{\omega_{i+1,j} - \omega_{i-1,j}}{2\Delta x} + v_{i,j} \frac{\omega_{i,j+1} - \omega_{i,j-1}}{2\Delta y} = \nu \left(\frac{\omega_{i+1,j} - 2\omega_{i,j} + \omega_{i-1,j}}{(\Delta x)^2} + \frac{\omega_{i,j+1} - 2\omega_{i,j} + \omega_{i,j-1}}{(\Delta y)^2} \right)$$

$\beta = \frac{\Delta x}{\Delta y}$
Multiply both side by $(\Delta x)^2$, divide both side by ν

$$-\left(1 - \frac{u_{i,j} \Delta x}{2\nu}\right) \omega_{i+1,j} - \left(1 + \frac{u_{i,j} \Delta x}{2\nu}\right) \omega_{i-1,j} - \left(1 - \frac{v_{i,j} \Delta y}{2\nu}\right) \beta^2 \omega_{i,j+1}$$

$$- \left(1 + \frac{v_{i,j} \Delta y}{2\nu}\right) \beta^2 \omega_{i,j-1} = -2(1 + \beta^2) \omega_{i,j}$$

$$\omega_{i,j} = \frac{\left(1 - \frac{u_{i,j} \Delta x}{2\nu}\right) \omega_{i+1,j} + \left(1 + \frac{u_{i,j} \Delta x}{2\nu}\right) \omega_{i-1,j} + \left(1 - \frac{v_{i,j} \Delta y}{2\nu}\right) \beta^2 \omega_{i,j+1} + \left(1 + \frac{v_{i,j} \Delta y}{2\nu}\right) \beta^2 \omega_{i,j-1}}{2(1 + \beta^2)}$$

So, now we will discretize Vorticity equation, so what is the Vorticity equation? $u \frac{\partial \omega}{\partial x} + v \frac{\partial \omega}{\partial y}$ is equal to $\nu \frac{\partial^2 \omega}{\partial x^2} + \nu \frac{\partial^2 \omega}{\partial y^2}$. So, u is $\frac{\partial \psi}{\partial y}$. So, you can discretize this, in 2 dimensional grid if you use.

This is your, $i + 1, j$, this is your $i - 1, j$, this is $i, j + 1$ and this is $i, j - 1$. So, now if you want to find the value u at point i, j , then you can discretize using $u_{i,j}$, $u_{i,j} = \frac{\partial \psi}{\partial y}$ it is the first derivative, so you can use some second order approximation, so that you can use central difference, so you can write $u_{i,j}$.

So, it is a first derivative $\frac{\partial \psi}{\partial y}$. So, you can use some second order approximation. So, we will use central difference. So, you can write $\psi_{i,j+1} - \psi_{i,j-1}$ and the distance between this 2 points is $2\Delta y$. So, this is your $u_{i,j}$. Similarly, v is equal to $-\frac{\partial \psi}{\partial x}$. So, $v_{i,j}$ and if you use the central difference approximation then you can write $\psi_{i+1,j} - \psi_{i-1,j}$ divided by $2\Delta x$.

Because we have use central difference approximation for this first derivative and the distance between that 2 points is $2\Delta x$. So, $\psi_{i+1,j} - \psi_{i-1,j}$ by $2\Delta x$ and ψ you already you can find it, so ψ you can find it from the stream function equation. Already we have discretized that equation. So, from the known values of ψ you can find the values of u and v . Now you discretize this equation.

So, here also let us use the central difference for the first derivative and central difference for the second derivative. So, that overall order of accuracy will be Δx^2 and Δy^2 . Because the stream function equation already we have discretize using second order approximation. So, the order of accuracy is Δx^2 and Δy^2 and when we are going to discretize this vorticity equation. Let us take central difference approximation, so the order of accuracy will be $\Delta x^2 \Delta y^2$.

So, we will take the central difference approximation for both first derivative and second derivative. So, you can write, $u_{i,j} = \frac{\omega_{i+1,j} - \omega_{i-1,j}}{2\Delta x} + v_{i,j}$ then $\omega_{i,j} = \frac{\omega_{i,j+1} - \omega_{i,j-1}}{2\Delta y}$. Now in the right hand side we have the second derivative. So, you will use central difference, so you will use $\omega_{i+1,j} - 2\omega_{i,j} + \omega_{i-1,j}$ divided by Δx^2 .

So, this is the first term and the second term you can write, $\omega_{i,j+1} - 2\omega_{i,j} + \omega_{i,j-1}$ divided by Δy^2 . So, you can see overall the approximation is second order $\Delta x^2 \Delta y^2$. So now, you rearrange it and you define β is equal to $\Delta x / \Delta y$. So, if you rearrange it you are going to get, so you can see what are the terms are there. So, $\omega_{i+1,j}$ here, $\omega_{i-1,j}$ here. This we will take common, then $\omega_{i,j+1}$ here, $\omega_{i,j-1}$ here.

Then we have $\omega_{i,j+1}$ here $\omega_{i,j+1}$ and now we have $\omega_{i,j-1}$ and we have $\omega_{i,j-1}$ and we have $\omega_{i,j}$ here and $\omega_{i,j}$ here. So now, you rearrange it and write. So, multiply both side by Δx^2 . So, that we can write in terms of β , so Δx^2 . So you can see this $\omega_{i+1,j}$. So $\omega_{i+1,j}$, so if you write $\omega_{i+1,j}$. So, it will be multiply both side by Δx^2 and divide both side by ν . So, that here ν will come in the denominator.

So, now you can write, so what is the coefficient of $\omega_{i+1,j}$? So $\omega_{i+1,j}$ so it will be $-\frac{1}{2} - \frac{u_{i,j}}{\nu}$. So, we have multiplied by Δx^2 . So, it will be Δx^2 divided by 2ν it will be $\omega_{i+1,j}$. Similarly, $\omega_{i-1,j}$. So, $\omega_{i-1,j}$, so it will be negative, so you can, it will be $\omega_{i-1,j} + \frac{u_{i,j}}{\nu} \Delta x^2$ by 2ν .

Similarly now you write $\omega_{ij} + 1$. So, $\omega_{ij} + 1$, if you write then it will be minus. So, it will be β^2 , so it will be β^2 , so it will be $1 - u_{ij} v_{ij}$. So we have multiply with Δx^2 . So, it will be Δx^2 . So Δx by Δy will be β , so it will be Δx^2 . So, β^2 and this is Δy if you get it. So, it will be v_{ij} , so Δy divided by 2 into ν and β^2 will be there. So, it will be $\beta^2 \omega_{ij} + 1$.

Because here Δy is there. So, if you multiply both side Δy . So, denominator and numerator, so it will be Δy by Δy^2 and Δx^2 you have a. So, Δx^2 by Δy^2 will be β^2 and that we are taking common. Similarly you write for $\omega_{ij} - 1$. So, it will be $-1 + v_{ij} \Delta y$ by twice ν . Then $\beta^2 \omega_{ij} - 1$ and now you have in right hand side. So, you have ω_{ij} , so ν we have already divided. So, it will be ω_{ij} . So, it will be -2 into $1 + \beta^2$ into ω_{ij} .

So, now you can write for ω_{ij} . So, ω_{ij} is equal to so minus in both side it is there. So, 2 into $1 + \beta^2$ and so it will be, so it will be $1 - u_{ij} \Delta x$ divided by twice $\nu \omega_{ij} + 1 + u_{ij} \Delta x$ by twice $\nu \omega_{ij} - 1 + v_{ij} \Delta y$ by twice $\nu \beta^2 \omega_{ij} + 1$ and you have $1 + v_{ij} \Delta y$ by twice ν and multiplied by $\beta^2 \omega_{ij} - 1$. So, now you can substitute the values of u_{ij} and v_{ij} . Because u_{ij} , v_{ij} you have written in terms of ψ . So, if you substitute it then you are going to get

(Refer Slide Time: 30:29)

Vorticity-Stream Function Formulations

$$u_{i,j} = \frac{\psi_{i,j+1} - \psi_{i,j-1}}{2\Delta y}$$

$$v_{i,j} = -\frac{\psi_{i+1,j} - \psi_{i-1,j}}{2\Delta x}$$

$$\omega_{i,j} = \frac{1}{2(\nu\Delta^2)} \left[\left(1 - \frac{\psi_{i,j+1} - \psi_{i,j-1}}{2\Delta y} \frac{\Delta x}{2\nu}\right) \omega_{i,j+1} + \left(1 + \frac{\psi_{i,j+1} - \psi_{i,j-1}}{2\Delta y} \frac{\Delta x}{2\nu}\right) \omega_{i,j-1} \right. \\ \left. + \left(1 + \frac{\psi_{i+1,j} - \psi_{i-1,j}}{2\Delta x} \frac{\Delta y}{2\nu}\right) \beta^2 \omega_{i,j+1} \right. \\ \left. + \left(1 - \frac{\psi_{i+1,j} - \psi_{i-1,j}}{2\Delta x} \frac{\Delta y}{2\nu}\right) \beta^2 \omega_{i,j-1} \right]$$

$$\beta = \frac{\Delta x}{\Delta y}$$

$$\omega_{i,j} = \frac{1}{2(\nu\Delta^2)} \left[\left\{1 - \left(\frac{\psi_{i,j+1} - \psi_{i,j-1}}{2\Delta y}\right) \frac{\beta}{\Delta y}\right\} \omega_{i,j+1} \right. \\ \left. + \left\{1 + \left(\frac{\psi_{i,j+1} - \psi_{i,j-1}}{2\Delta y}\right) \frac{\beta}{\Delta y}\right\} \omega_{i,j-1} \right. \\ \left. + \left\{1 + \left(\frac{\psi_{i+1,j} - \psi_{i-1,j}}{2\Delta x}\right) \frac{1}{\Delta y\beta}\right\} \omega_{i,j+1} \right. \\ \left. + \left\{1 - \left(\frac{\psi_{i+1,j} - \psi_{i-1,j}}{2\Delta x}\right) \frac{1}{\Delta y\beta}\right\} \beta^2 \omega_{i,j-1} \right]$$

$$\psi_{i,j} = \frac{1}{2(\nu\Delta^2)} \left[\psi_{i+1,j} + \psi_{i-1,j} + \beta^2 (\psi_{i,j+1} + \psi_{i,j-1}) + (\Delta x)^2 \omega_{i,j} \right]$$

Omega i j we have written, but now we can substitute u i j and v i j. So, u i j is psi i j plus 1 minus psi i j minus 1 divided by twice delta y and v i j as minus psi i plus 1 j minus psi i minus 1 j divided by 2 delta x. So, now we will substitute this u v in the expression of omega i j. So, omega i j is 1 by 2 into 1 plus beta square 1 minus.

Now you put the value of u. So, it will be psi i j plus 1 minus psi i j minus 1 divided by 2 delta y delta x by twice nu. So, this is omega i plus 1 j, then plus now 1. So, minus so it will be plus. So, psi i j plus 1 minus psi i j minus 1 divided by 2 delta y delta x by 2 nu omega i minus 1 j. Then next term now 1, so it was minus but velocity v i j is having negative sign so this minus minus will become plus.

So, psi i plus 1 j minus psi i minus 1 j divided by 2 delta x delta y by twice nu. Then beta square omega i j plus 1 and 1 minus. So, here plus was there but now v i j you have negative. So, this negative here it will appear negative. So, psi i plus 1 j minus psi i minus 1 j divided by 2 delta x delta y by twice nu beta square omega i j minus 1.

So, now you substitute beta is equal to delta x by delta Y, here delta x delta y and here delta y delta x, so you can put beta. So, if you write omega ij, so you will get 1 by 2 into 1 plus beta square 1 minus psi i j plus 1 minus psi i j minus 1 beta by 4 nu, delta x by delta y we are putting beta omega i plus 1 j plus 1 plus psi i j plus 1 minus psi i j minus 1 beta by 4 nu omega i minus 1 j plus 1 plus psi i plus 1 j minus psi i minus 1 j by 1 by 4 nu beta.

So, $\beta \frac{\Delta y}{\Delta x}$ is equal to $\frac{1}{2} \beta \omega_{i,j+1} + \frac{1}{2} \beta \omega_{i,j-1} - \psi_{i,j+1} - \psi_{i,j-1} + \frac{1}{4} \nu \beta^2 \omega_{i,j}^2$. So, this is the discretized equation for $\omega_{i,j}$ in dimensional form and whatever ψ we already derived, $\psi_{i,j}$ is $\frac{1}{2} \nu \beta^2 \psi_{i,j+1} + \frac{1}{2} \nu \beta^2 \psi_{i,j-1} + \beta^2 \psi_{i,j} + \frac{1}{2} \Delta x^2 \omega_{i,j}$. So, this is the discretized equation for ψ . So, you can see in the ψ discretized equation $\omega_{i,j}$ is coming and in ω ψ values are coming. So, it is coupled and here you need to solve this $\psi_{i,j}$ and $\omega_{i,j}$ for all interior points.

So, whatever these equations we have derived for the stream function and the vorticity. So, these are applicable for the interior points. But we need to solve some free flow problem. So, in the next class we define, some free flow problem and for those, problem we will define the boundary conditions and we will discuss the boundary conditions for ω and ψ .

Because when you are solving these equations these vorticity equation and stream function equation, you need the boundary as well in terms of ψ when you are solving for stream function and you need the values of vorticity at the boundary when you are solving the vorticity equation.

But, you may know the boundary condition in terms of velocities for a particular fluid flow problem. So, from these velocity boundary condition you need to derive the stream function boundary condition as well as the vorticity boundary condition. Then only these two equations whatever we have derived today for the interior points you will be able to solve for a particular problem provided the boundary conditions are given in terms of stream function and vorticity. So, in next class we will discuss two problems and we will describe the, we will discuss the boundary conditions for those two problems. Thank you.