

# Aircraft Propulsion

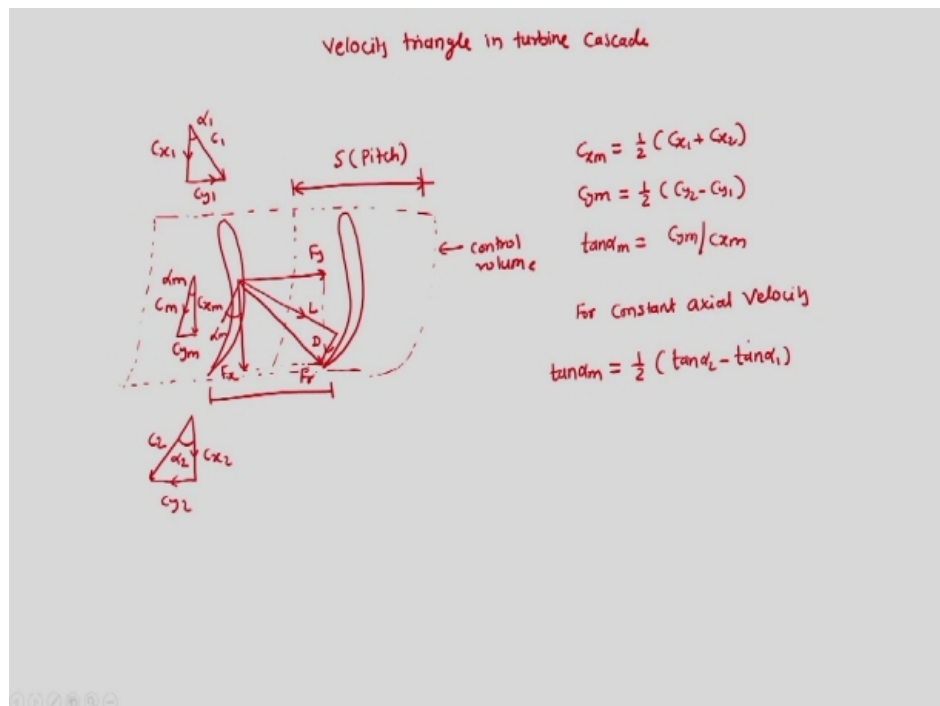
Prof. Vinayak N. Kulkarni

Department of Mechanical Engineering  
Indian Institute of Technology-Guwahati

## Lecture-40

### Velocity diagrams of Turbine Cascade, Compressor cascade.

Welcome to the class, we will continue our discussion about blade design and cascade theory till time what we had seen was the cascade tunnel parts of cascade tunnel and then we had seen about the turbine cascade, different angles of the turbine cascade.



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control volume now which includes 1 pitch at least so this is turbine cascade and now this is control volume now this is spacing between the blades and then this what we call as the pitch S.

And here, we have velocity absolute approaching like this, which is  $C_1$ . So, this is  $C_{x1}$  and then this is  $C_{y1}$  approaching at an angle  $\alpha_1$ , we have seen these are the  $\alpha_1$  and  $\alpha_2$  are the air angles at inlet and outlet. So, these are the angles made by the absolute velocity with axial direction. So, similarly this is absolute velocity  $C_2$  leaving the blade, then this is  $\alpha_2$ , since this is  $C_{x2}$  and this is  $C_{y2}$  here from these 2 velocity triangles or mean velocity triangle is constructed.

So, this is  $C_{xm}$  and then this is  $C_m$  and this is  $C_{ym}$  and the angle included is  $\alpha_m$ . In such case the forces acting are this is  $F_y$ , this is  $F_x$  and hence the resultant is  $F_r$  but this resultant can further be decomposed into 2 more components, which are lift and drag. So, for that these components will be made, which are normal and perpendicular to the velocity mean velocity. So, this is lift and this is drag.

So, this total force  $F_r$  actually gets decomposed into two ways one is  $F_x$ ,  $F_y$  which has the coordinate system associated with laboratory coordinate system and then it gets further decomposed into 2 more components which are lift and drag, which are normal to the velocity which are based upon the velocity direction. So, this angle so, this included angle is  $\alpha_m$ . Now, having said this we will write down the expressions.

So,  $C_{xm}$  which is axial mean velocity is  $\frac{1}{2}(C_{x1} + C_{x2})$   $C_{ym}$ . Since  $Y_m$ s are in opposite direction, we will have  $C_{y2} - C_{y1}$  and  $\tan \alpha_m = \frac{C_{ym}}{C_{xm}}$  and then we have for constant axial velocity we have  $\tan \alpha_m = \frac{1}{2} (\tan \alpha_2 - \tan \alpha_1)$ . So, this is for the case of constant axial velocity. So, now the having these control volume with us, we will go ahead and find out what are the forces acting on the cascade.

(refer time: 05:49). So, let us apply mass conservation or continuity equation and in the continuity equation we have  $\dot{m} = \rho_1 C_{x1} \times S$  is the pitch  $\times 1$  is the blade height and we are consulting it to be unit. So, similarly we have  $\rho_2 C_{x2} \times S \times 1$  for incompressible flow we have  $\rho_1 = \rho_2 = \rho$ . So, we have  $\dot{m} = \rho C_{x1} S = \rho C_{x2} S$ . So, we have  $C_{x1} = C_{x2} = C_{xm}$ . So,  $\dot{m} = \rho C_{xm} S$ .

Now, let us apply Y moment Y momentum equation or let us apply momentum equation in Y direction. Here by this we can write the force  $F_y$  will have its magnitude as  $\dot{m}(C_{y2} - (-C_{y1}))$ . So, this is  $F_y$  so, what we will have is  $F_y = \dot{m}(C_{y2} + C_{y1})$ . We can put the expression for  $\dot{m}$ ,  $\rho C_{xm} S [C_{y2} + C_{y1}]$  from the triangle we can write down  $C_{y2} = C_{x1} \tan \alpha_1$ .

So, it is equal to  $C_{xm} \tan \alpha_1$  similarly,  $C_{y1}$  sorry, similarly,  $C_{y2} = C_{x2} \tan \alpha_2 = C_{xm} \tan \alpha_2$ . So, putting this in the expression for momentum, we can get  $F_y = \rho C_{xm} S [C_{xm} \tan \alpha_1 + C_{xm} \tan$

$\hookrightarrow$  continuity equation  
 $\dot{m} = \rho_1 C_{x1} S \times 1 = \rho_2 C_{x2} S \times 1$   
 For incompressible flow  
 $\rho_1 = \rho_2 = \rho$   
 $\therefore \dot{m} = \rho C_{x1} S = \rho C_{x2} S$   
 $C_{x1} = C_{x2} = C_{xm}$   
 $\therefore \dot{m} = \rho C_{xm} S$

$\hookrightarrow$  Momentum equation in y direction  
 $F_y = \dot{m} [C_{y2} - (-C_{y1})]$   
 $\therefore F_y = \dot{m} [C_{y2} + C_{y1}] = \rho C_{xm} S [C_{y2} + C_{y1}]$   
 $C_{y2} = C_{x1} \tan \alpha_1 = C_{xm} \tan \alpha_1$   
 $C_{y1} = C_{x2} \tan \alpha_2 = C_{xm} \tan \alpha_2$

$F_y = \rho C_{xm} S [C_{xm} \tan \alpha_1 + C_{xm} \tan \alpha_2]$   
 $\therefore F_y = \rho C_{xm}^2 S [\tan \alpha_1 + \tan \alpha_2]$   
 $F_y = \left(\frac{1}{2} \rho S C_{xm}^2\right) \cdot 2 \left(\frac{S}{l}\right) [\tan \alpha_1 + \tan \alpha_2]$   
 $\therefore C_{Fy} = \frac{F_y}{\frac{1}{2} \rho l C_{xm}^2} = \underline{\underline{2 \left(\frac{S}{l}\right) [\tan \alpha_1 + \tan \alpha_2]}}$

$\hookrightarrow$  Momentum equation in x direction  
 $F_x = (P_1 - P_2) S \times 1 + \rho (S \times 1) C_{xm} (C_{x1} - C_{x2})$   
 $\therefore$  For incompressible flow  
 $P_1 = P_01 - \frac{1}{2} \rho C_1^2 \rightarrow P_2 = P_02 - \frac{1}{2} \rho C_2^2$   
 $\therefore P_1 - P_2 = (P_01 - P_02) + \frac{1}{2} \rho (C_2^2 - C_1^2)$   
 $F_x = (P_1 - P_2) S + \rho S C_{xm} (C_{x1} - C_{x2}) \rightarrow$  if  $C_{x1} = C_{x2}$   
 $\therefore F_x = (P_1 - P_2) S$   
 $\therefore F_x = (P_01 - P_02) S + \frac{1}{2} \rho S (C_2^2 - C_1^2)$

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$\alpha_2$ ]. So,  $F_y = \rho C_{xm}^2 S [\tan \alpha_1 + \tan \alpha_2]$ , we can rearrange the term and write down as  $\left(\frac{1}{2} \rho l C_{xm}^2\right) \times 2 \times \left(\frac{S}{l}\right) \times [\tan \alpha_1 + \tan \alpha_2]$ .

We are purposely doing since we can write down these coefficient of force  $C_{Fy} = \frac{F_y}{\frac{1}{2}(\rho l C_{xm}^2)}$  and then this gives us answer as  $2 \times \left(\frac{S}{l}\right) \times [\tan \alpha_1 + \tan \alpha_2]$ . so, if we know  $\alpha_1$  and  $\alpha_2$ , for given geometry and arrangement, we can find out what is the coefficient of the normal force, which is acting under blades. Then we can write down the momentum equation in x direction.

So, we will have  $F_x$  is equal to, here we just would not have force momentum change equal to force we will have pressure change as well. So,  $P_1 - P_2$  with the change in pressure in the direction of flow which is  $S \times 1$  is the area +  $\rho \times (S \times 1)$  for unit blade height  $C_{xm} \times (C_{x1} - C_{x2})$ . So, we have for incompressible flow we can write down  $P_1 = P_01 - \frac{1}{2} \rho C_1^2$ .

Basically, this expression says that  $P_01 = P_1 + \frac{1}{2} \rho V^2$ . This is strictly for incompressible flows, where total pressure is equal to static plus dynamic pressures. So, we are finding our static pressure from the total pressure and the dynamic pressure after subtracting them. Similarly we can write down  $P_2 = P_02 - \frac{1}{2} \rho C_2^2$ . So, we can write  $P_1 - P_2$  is equal to  $P_01 - P_02 + \frac{1}{2}(C_2^2 - C_1^2)\rho$ .

So, we can write down expression for  $F_x$  again  $(P_1 - P_2)S + \rho S C_{xm}(C_{x1} - C_{x2})$ . Take a partic-

ular case where we have if  $C_{x1} = C_{x2}$  we can write down  $F_x$  only as  $(P_1 - P_2)S$ . So,  $F_x = (P_{01} - P_{02})S + \frac{1}{2}\rho S(C_2^2 - C_1^2)$ .

The image shows handwritten mathematical derivations for the axial force coefficient  $F_x$  and the lift coefficient  $L$ .

**Left side (Axial Force  $F_x$ ):**

- $F_x = S\Delta P_0 + \frac{1}{2}\rho S(C_2^2 - C_1^2)$
- $C_2^2 = C_{x2}^2 + C_{y2}^2, C_1^2 = C_{x1}^2 + C_{y1}^2$
- $F_x = S\Delta P_0 + \frac{1}{2}\rho S(C_{x2}^2 + C_{y2}^2 - C_{x1}^2 - C_{y1}^2)$
- $F_x = S\Delta P_0 + \frac{1}{2}\rho S(C_{y2}^2 - C_{y1}^2)$
- $C_{y2} = C_{x2} \tan \alpha_2 = C_{xm} \tan \alpha_2$
- $C_{y1} = C_{x1} \tan \alpha_1 = C_{xm} \tan \alpha_1$
- $F_x = S\Delta P_0 + \frac{1}{2}\rho S(C_{xm}^2 \tan^2 \alpha_2 - C_{xm}^2 \tan^2 \alpha_1)$
- $F_x = S\Delta P_0 + \frac{1}{2}\rho S C_{xm}^2 [\tan^2 \alpha_2 - \tan^2 \alpha_1]$  — (1)
- $F_x = S\Delta P_0 + \frac{1}{2}\rho S C_{xm}^2 (\tan \alpha_2 - \tan \alpha_1) (\tan \alpha_2 + \tan \alpha_1)$
- $F_x = S\Delta P_0 + \rho S C_{xm}^2 \tan \alpha_m (\tan \alpha_2 + \tan \alpha_1)$  — (2)
- $C_{F_x} = \frac{F_x}{\frac{1}{2}\rho l C_{xm}^2} = \frac{S\Delta P_0}{\frac{1}{2}\rho l C_{xm}^2} + \left(\frac{S}{l}\right) (\tan^2 \alpha_2 - \tan^2 \alpha_1)$

**Right side (Lift  $L$ ):**

- $L = F_y \cos \alpha_m + F_x \sin \alpha_m$
- $L = \rho S C_{xm}^2 (\tan \alpha_2 + \tan \alpha_1) \cos \alpha_m + S\Delta P_0 \sin \alpha_m$
- $L = \rho S C_{xm}^2 (\tan \alpha_2 + \tan \alpha_1) \left[ \frac{\cos \alpha_m + \sin^2 \alpha_m}{\cos \alpha_m} \right] + S\Delta P_0 \sin \alpha_m$
- $L = \rho S C_{xm}^2 \cos \alpha_m (\tan \alpha_2 + \tan \alpha_1) [1 + \tan^2 \alpha_m] + S\Delta P_0 \sin \alpha_m$
- $L = \rho S C_{xm}^2 \sec \alpha_m (\tan \alpha_2 + \tan \alpha_1) + S\Delta P_0 \sin \alpha_m$
- $C_L = \frac{L}{\frac{1}{2}\rho l C_{xm}^2} = 2 \left(\frac{S}{l}\right) \sec \alpha_m (\tan \alpha_2 + \tan \alpha_1) + \frac{S\Delta P_0}{\frac{1}{2}\rho l C_{xm}^2} \sin \alpha_m$

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(refer time: 13:24). We can write down the same expression  $F_x = S\Delta P_0 + \frac{1}{2}\rho S(C_2^2 - C_1^2)$  here, we know that  $C_2^2 = C_{x2}^2 + C_{y2}^2$  and  $C_1^2 = C_{x1}^2 + C_{y1}^2$ . So,  $F_x = S\Delta P_0 + \frac{1}{2}\rho S(C_{x2}^2 + C_{y2}^2 - C_{x1}^2 - C_{y1}^2)$  and both being equal, it would get cancelled. So,  $F_x = S\Delta P_0 + \frac{1}{2}\rho S(C_{y2}^2 - C_{y1}^2)$ , but from velocity triangle. We can write down  $C_{y2} = C_{x2} \tan \alpha_2 = C_{xm} \tan \alpha_2$ .

Similarly,  $C_{y1} = C_{x1} \tan \alpha_1 = C_{xm} \tan \alpha_1$ . So,  $F_x = S\Delta P_0 + \frac{1}{2}\rho S(C_{xm}^2 \tan^2 \alpha_2 - C_{xm}^2 \tan^2 \alpha_1)$ . So, we will have  $F_x = S\Delta P_0 + \frac{1}{2}\rho S C_{xm}^2 (\tan^2 \alpha_2 - \tan^2 \alpha_1)$ . So,  $F_x = S\Delta P_0 + \frac{1}{2}\rho S C_{xm}^2 (\tan \alpha_2 - \tan \alpha_1)(\tan \alpha_2 + \tan \alpha_1)$ . These are the 2 expressions for us further, we know  $\frac{1}{2}$  into this term can lead to  $\tan \alpha_m$ .

So, we can write down it as  $S\Delta P_0 + \frac{1}{2}\rho S C_{xm}^2 \tan \alpha_m (\tan \alpha_2 + \tan \alpha_1)$ . We will remember this expression and we will name it as number 1 and we will name it as number 2. So, we can write down  $C_{F_x} = \frac{F_x}{\frac{1}{2}(\rho l C_{xm}^2)}$  and this gives us if we use equation number 1. We can get as  $\frac{S_0}{\frac{1}{2}(\rho l C_{xm}^2)}$  will cancel and we will have  $\frac{S}{l}(\tan^2 \alpha_2 - \tan^2 \alpha_1)$ .

So, this is our non-dimensional form of axial force coefficient having said this, we can move ahead and find out what is lift and drag from the coefficients whatever we have got or from

the forces whatever we have arrived at so, L lift from the force resultant force and its decomposition we can clearly mentioned a lift is equal to  $F_y \cos \alpha_m + F_x \sin \alpha_m$ . So, L is equal to  $F_y$  we have found out  $F_x$ .

We have found out we can keep those expressions written over here  $\rho S C_{x_m}^2 \cdot (\tan \alpha_1 + \tan \alpha_2) \cdot \cos \alpha_m + \rho S C_{x_m}^2 \cdot \tan \alpha_m \cdot (\tan \alpha_1 + \tan \alpha_2) \cdot \sin \alpha_m + S \Delta P_0 \cdot \sin \alpha_m$ . Now here, we can rearrange and then we can take  $\rho S C_{x_m}^2 \cdot (\tan \alpha_1 + \tan \alpha_2)$  common from first 2 terms of this.

And then we can have  $\rho S C_{x_m}^2 \cdot (\tan \alpha_1 + \tan \alpha_2) \cdot [\cos \alpha_m + \frac{\sin^2 \alpha_m}{\cos \alpha_m} + S \Delta P_0 \cdot \sin \alpha_m]$ . Now, we can take  $\cos \alpha_m$  also common from this bracket and then we have  $L = \rho S C_{x_m}^2 \cdot \cos \alpha_m (\tan \alpha_1 + \tan \alpha_2) \cdot [1 + \tan^2 \alpha_m] + S \Delta P_0 \cdot \sin \alpha_m$ .

So, 1 but we know  $1 + \tan^2 \alpha$  is equal to  $1 + \sec^2 \alpha$ . So, here we will have  $\sec^2 \alpha_m$ , but we have one more  $\cos \alpha_m$  here. So, that is a reciprocal so, we can write it as  $\rho S C_{x_m}^2 \cdot \sec \alpha_m (\tan \alpha_1 + \tan \alpha_2) + S \Delta P_0 \cdot \sin \alpha_m$ . So, this is our expression for the lift, OK.

So, we can write down  $C_L$ , which is lift coefficient as  $\frac{L}{\frac{1}{2} \rho l C_{x_m}^2}$  and then we can get it as  $2 \left( \frac{S}{l} \right) \sec \alpha_m (\tan \alpha_1 + \tan \alpha_2) + \frac{S_0}{\frac{1}{2} \rho l C_{x_m}^2} \cdot \sin \alpha_m$ . So, knowing this, we can find out what is the lift coefficient which is getting acted on the cascade which is a turbine cascade.

Drag force

$$D = \text{drag} = F_x \cos \alpha_m - F_y \sin \alpha_m$$

$$D = \rho S C_{x_m}^2 \tan \alpha_m (\tan \alpha_1 + \tan \alpha_2) \cos \alpha_m + S \Delta P_0 \cos \alpha_m - \rho S C_{x_m}^2 (\tan \alpha_1 + \tan \alpha_2) \sin \alpha_m$$

$$D = S \Delta P_0 \cos \alpha_m$$

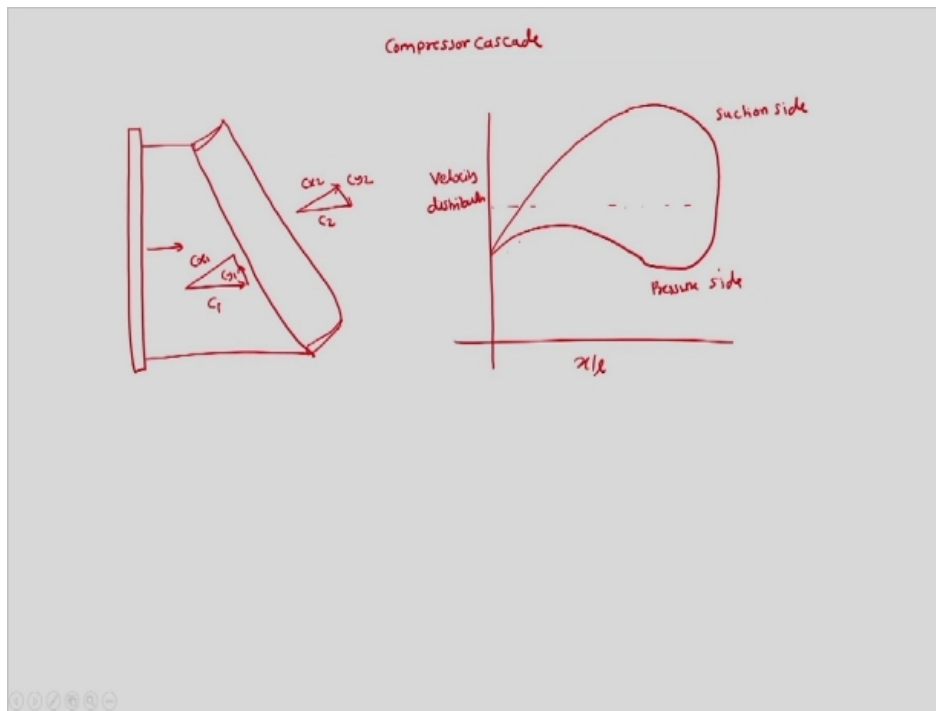
$$C_D = \frac{D}{\frac{1}{2} \rho l C_{x_m}^2} = \frac{S \Delta P_0 \cos \alpha_m}{\frac{1}{2} \rho l C_{x_m}^2}$$

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(refer time: 22:40). Now, we will move ahead and then we can find out drag force for the turbine cascade. So, far drag force again from the force resultant force and its decomposition, we can write down D which is a drag which is equal to  $F_x \cos \alpha_m - F_y \sin \alpha_m$ . So, we have D we can put the expression for  $F_x$  as  $\frac{\rho S C_{x_m}^2 \tan \alpha_m (\tan \alpha_1 + \tan \alpha_2) \cos \alpha_m}{S_0 \cos \alpha_m - \rho S C_{x_m}^2 (\tan \alpha_1 + \tan \alpha_2) \sin \alpha_m}$ .

Here, it is good point to note that we have  $\tan \alpha_m$  and  $\cos \alpha_m$  so, this would ultimately be  $\sin \alpha_m$ . So, we can see that the first term and third term they are same. So,  $D = S_0 \cos \alpha_m$ . So, we can get drag coefficient,  $C_D$  as  $\frac{D}{\frac{1}{2} \rho l C_{x_m}^2} = \frac{S_0 \cos \alpha_m}{\frac{1}{2} \rho l C_{x_m}^2}$  and this is how we would have found out different coefficients and the forces which are acting on the cascade.

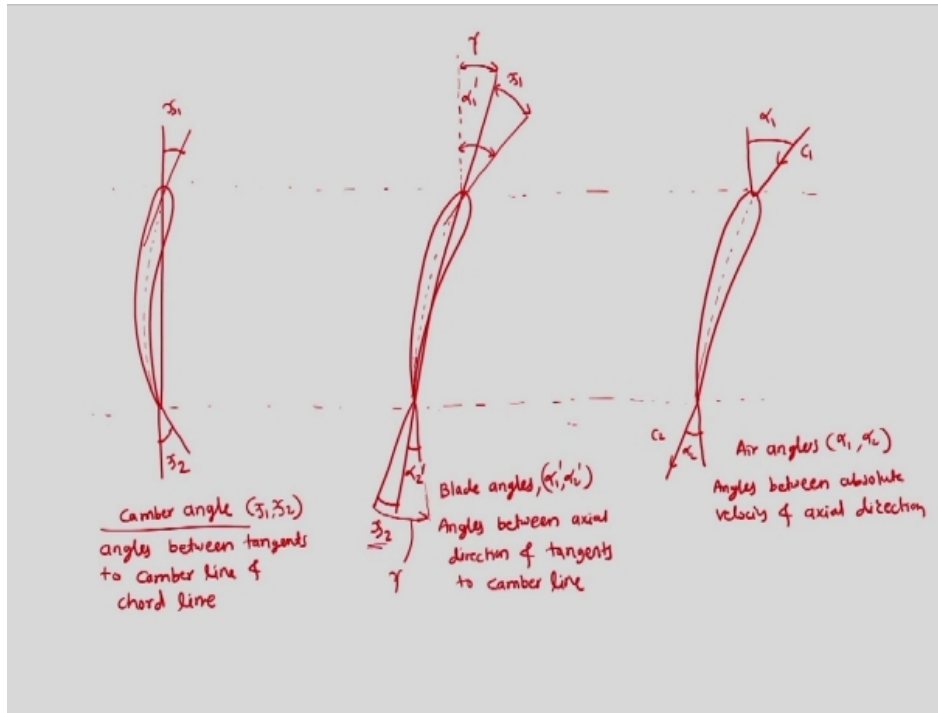
And we can see or hear that everything depends upon the setting of the angle and a blade angles and the angles whatever we would have used for the blade or the air and also for the setting.



**(Slide Time: 25:08)**

(refer time: 25:08). So, we will move ahead for the compressor cascade and we can note that for the compressor cascade first it will be the same wind tunnel which is the cascade tunnel getting applied for the application of the compressor. So, we will change our turbine cascade to the compressor cascade in case of compressors. So, this is the  $C_1$  so, this will be  $C_{x1}$  this will be  $C_{y1}$  this will be  $C_{x2}$  this will be  $C_{y2}$  and then this will be  $C_2$  on a particular Aerofoil.

If we plot the velocity distribution then this velocity distribution for  $\frac{x}{l}$  can be seen like this. This is a typical velocity distribution for suction side on top and for pressure side on bottom, pressure side now again.



(Slide Time: 27:11)

(refer time: 27:11). We will draw the angles for the cascade which is a compressor cascade as what we did earlier. In case of turbine we do 3 types of angles first angle was camber angle. So, for camber angle so, this is camber line, so, this is the vertical line. So, this is the tangent to the camber at the inlet, so, this is  $\chi_1$  and this is  $\chi_2$ . Basically this vertical line is the chord line in this specific case. So, this is  $\chi_2$ .

So, this is camber angle as what we had seen earlier this is  $\chi_1$  and  $\chi_2$ , they are the angles between tangents to camber line and chord line then, we will show the blade angles. So, for the blade angle we will make the chord little bit tilted. So, now this is chord and now this is camber line and now this is tangent to the camber. So, as what we know this angle is  $\chi_1$  and this is we will put it here  $\chi_2$ .

Now, I will plot the axial direction here and then we will have the tangent and the tangent and the axial direction we will have angle  $\alpha'_1$ . So, similarly here also it is  $\alpha'_2$  so, this is blade angle which are  $\alpha'_1$  and  $\alpha'_2$  angles between axial direction and tangents to camber line and then we can have third which is for air angle. So, same way we will replicate.

This is camber line and this is axial direction and this is supposed the direction of velocity, which is  $C_1$  so, this angle is  $\alpha_1$ . Similarly, if velocity is leaving, absolute velocity in this direction and this is  $\alpha_2$  this is the axial direction. So, we have air angles which are  $\alpha_1, \alpha_2$ . So they are angles between absolute velocity and axial direction. So, having said this, we can continue here showing the other angles among that first was the stagger angle.

And, stagger angle can be shown over here this  $\gamma$  will be the stagger angle which is between the chord line and the axial direction.

Handwritten mathematical derivations:

$$\theta = \chi_1 + \chi_2 = \alpha_1' - \alpha_2'$$

$$\alpha_1' = \chi_1 + \gamma \quad \alpha_2' = \gamma - \chi_2$$

$$\alpha_1' - \alpha_2' = \chi_1 + \chi_2 = \theta$$

$$i = \text{incidence angle} = \alpha_1 - \alpha_1'$$

$$\delta = \text{deviation angle} = \alpha_2 - \alpha_2'$$

$$\epsilon = \text{deflection angle} = \alpha_1 - \alpha_2$$

$$\epsilon = (\alpha_1' + i) - (\alpha_2' + \delta)$$

$$\epsilon = (\alpha_1' - \alpha_2') + (i - \delta)$$

$$\underline{\underline{\epsilon = \theta + i - \delta}}$$

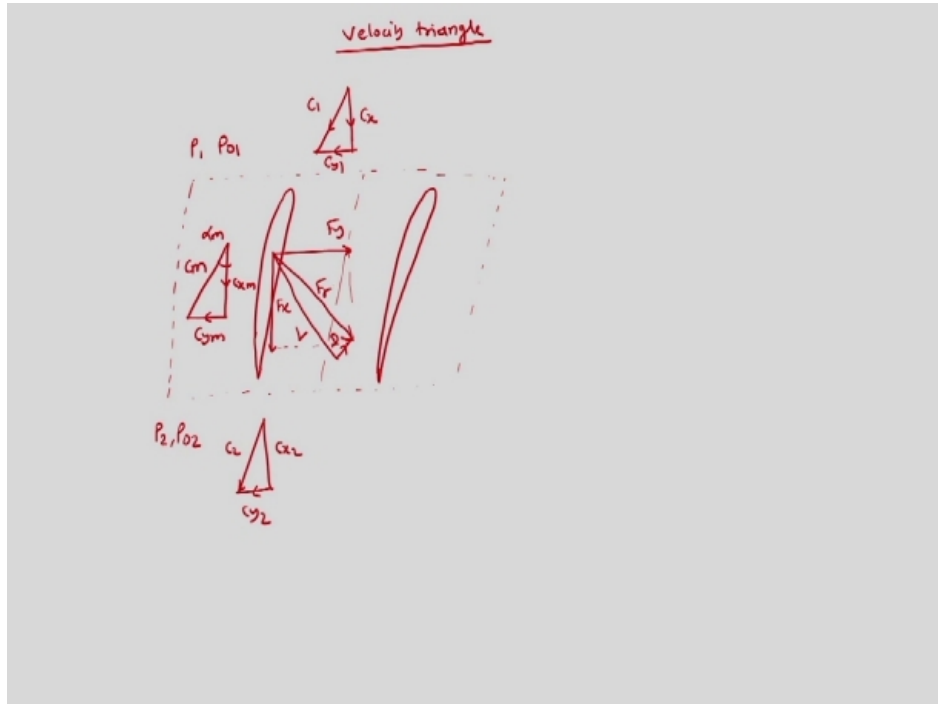
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(refer time: 32:58). So, other settings we can easily make out as what we did earlier camber angle  $\theta = \chi_1 + \chi_2$ , which is  $\alpha_1' - \alpha_2'$  further  $\alpha_1' = \chi_1 + \gamma$  and  $\alpha_2' = \gamma - \chi_2$ , it should be noted here that, this angle which is between the tangent which is between the so, this big angle is  $\gamma$  in the outside in the outlet. So, this  $\gamma$  is equal to sorry this  $\gamma - \chi_2$  is  $\alpha_2'$ .

So, we can write down  $\alpha_1' - \alpha_2' = \chi_1 + \chi_2$ , which is  $\theta$ . So, i we have defined as incidence angle so, it is  $\alpha_1 - \alpha_1'$ , we have defined  $\delta$  which is deviation angle and then that is  $\alpha_2 - \alpha_2'$ . So, we have  $\epsilon$ , which is deflection angle and then that is  $\alpha_1 - \alpha_2$ . So,  $\epsilon = \alpha_1' + i$ , from the expression of this  $(\alpha_1' + i) - (\alpha_2' + \delta)$ .



So, deflection is  $(\alpha'_1 - \alpha'_2) + (i - \delta)$ , but  $\alpha'_1 - \alpha'_2$  is equal to  $\theta$ . So, it is  $\theta + i - \delta$ . This is the deflection angle so, now for the compressor cascade.



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We can again plot the velocity triangle. So, velocity triangle for compressor cascade here as well, we will draw a control volume so, far that let us first plot the Aerofoil and let us complete the control volume. So, this is the control volume for the compressor cascade. So, here we will say this is  $P_{01}$  and  $P_1$  and this is  $P_2$  and  $P_{02}$  say this is inlet, so, for the inlet we know that this is the direction of absolute velocity.

This is direction of axial velocities, this is  $C_{y1}$ . Similarly, if we take outlet, then for outlet, this is  $C_2$ , this is  $C_{x2}$  and this is  $C_{y2}$ . So, we can draw the mean velocity triangle as  $C_{xm}$  this is  $C_m$  and this is  $C_{ym}$  and this angle mean is  $\alpha_m$ . So, this is the typical velocity triangle for compressor cascade. Now, we can show the forces same way as what we did for turbine, this is  $F_y$  this is  $F_x$ .

So, this becomes the resultant  $F_r$  now here we are again needs to decompose them into lift and drag. So, this is lift and this is drag this resultant is  $F_r$ .

(refer time: 38:25). So, we can continue by applying momentum equations for the compressor cascade. So, let us apply mass first mass conservation here we get  $\dot{m} = \rho_1 C_{x1} S \times 1$ ,

↳ mass conservation

$$\dot{m} = \rho_1 C_{x1} (S x_1) = \rho_2 C_{x2} (S x_2)$$

↳ incompressible

$$\rho_1 = \rho_2 = \rho$$

$$C_{x1} = C_{x2} = C_{xm}$$

↳ Tangential or y momentum

$$F_y = \dot{m} [C_{y1} - C_{y2}]$$

$$F_y = \rho C_{xm} S [C_{y1} - C_{y2}]$$

$$F_y = \rho C_{xm} S [C_{xm} \tan \alpha_1 - C_{xm} \tan \alpha_2]$$

$$F_y = \rho C_{xm}^2 S [\tan \alpha_1 - \tan \alpha_2]$$

$$C_{Fy} = \frac{F_y}{\frac{1}{2} \rho l C_{xm}^2} = 2 \left( \frac{S}{l} \right) (\tan \alpha_1 - \tan \alpha_2)$$

↳ axial momentum equation or x-momentum equation

$$F_x = (P_2 - P_1) S + \rho (S) C_{xm} (C_{x2} - C_{x1})$$

$$F_x = (P_2 - P_1) S$$

$$F_x = \left\{ (P_2 - \frac{1}{2} \rho C_{x2}^2) - (P_1 - \frac{1}{2} \rho C_{x1}^2) \right\} S$$

$$F_x = \frac{1}{2} \rho S (C_1^2 - C_2^2) - \Delta P_0 S$$

$$F_x = \frac{1}{2} \rho S [C_1^2 - C_2^2] - \Delta P_0 S$$

$$F_x = \frac{1}{2} \rho S [C_{xm}^2 \tan^2 \alpha_1 - C_{xm}^2 \tan^2 \alpha_2] - \Delta P_0 S$$

$$F_x = \frac{1}{2} \rho S C_{xm}^2 [\tan^2 \alpha_1 - \tan^2 \alpha_2] - \Delta P_0 S \quad \text{--- (1)}$$

$$F_x = \frac{1}{2} \rho S C_{xm}^2 (\tan \alpha_1 + \tan \alpha_2) (\tan \alpha_1 - \tan \alpha_2) - \Delta P_0 S$$

$$F_x = \frac{1}{2} \rho S C_{xm}^2 \tan \alpha_m (\tan \alpha_1 - \tan \alpha_2) - S \Delta P_0 \quad \text{--- (2)}$$

$$C_{Fx} = \frac{F_x}{\frac{1}{2} \rho l C_{xm}^2} = \left( \frac{S}{l} \right) (\tan^2 \alpha_1 - \tan^2 \alpha_2) - \frac{S \Delta P_0}{\frac{1}{2} \rho l C_{xm}^2}$$

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we know at the outlet  $\rho_2 C_{x2} S \times 1$  for unit blade height but again we will say that  $C_{x1} = C_{x2} = C_{xm}$  for the constant axial velocity, we can get it as  $\rho C_{xm} S$  where we are saying it as incompressible which has led to  $\rho_1 = \rho_2 = \rho$  and then we have said  $C_{x1} = C_{x2} = C_{xm}$ .

So, this is how we would have the mass conservation equation for compressor cascade which is same as what for the turbine cascade. Now, let us apply tangential or y momentum equation if we apply we can get  $F_y = \dot{m}(C_{y1} - C_{y2})$  here it should be remembered that both y components are in the same direction, which was in the case of turbine both were in opposing, so, we had  $C_{y1} + C_{y2}$ .

So, we can write down put m dot expression for  $\dot{m}$ ,  $\rho C_{xm} S (C_{y1} - C_{y2})$ . So, we will have  $F_y = \rho C_{xm} S$ , we can put  $C_y$  as  $[C_{xm} \tan \alpha_1 + C_{xm} \tan \alpha_2]$ . So,  $F_y = \rho C_{xm}^2 S [\tan \alpha_1 - \tan \alpha_2]$ . So, we can write down coefficient of force  $F_y$ ,  $C_{Fy} = \frac{F_y}{\frac{1}{2} (\rho l C_{xm}^2)} = 2 \times \left( \frac{S}{l} \right) \times [\tan \alpha_1 - \tan \alpha_2]$ . Similarly, we can apply the axial momentum equation or x momentum equation.

There, we can write down  $F_x$  is equal to  $(P_2 - P_1) S + \rho S C_{xm} (C_{x2} - C_{x1})$  but again, we would have same axial coefficient velocity so,  $F_x = P_2 - P_1$  we already know that  $P_2 = [P_{02} - \frac{1}{2} \rho C_2^2 - (P_{01} - \frac{1}{2} \rho C_1^2)] \times S$ . Here, we can mention  $F_x$  as  $\frac{1}{2} \rho S (C_1^2 - C_2^2) - \Delta P_0 S$ . So,  $\Delta P_0$  is  $P_{01} - P_{02}$  are here, since we have minus sign here. So,  $F_x$  can be written as the  $\frac{1}{2} \rho S$ .

So,  $C_1^2$  can be split into  $C_{x1}^2 + C_{y1}^2$ . Similarly  $C_2^2$  will also be split into  $C_{x2}^2 + C_{y2}^2$  and  $C_{x1}$  and  $C_{x2}$  both are equal. So, we will have only  $y_2$  remaining so,  $C_{y1}^2 - C_{y2}^2 = \Delta P_0 S$ , we can put the value of  $C_{y1}$  using the velocity triangle as  $[C_{xm}^2 \tan^2 \alpha_1 - C_{xm}^2 \tan^2 \alpha_2] - \Delta P_0 S$ . So,  $F_x = \frac{1}{2} \rho S C_{xm}^2 [\tan^2 \alpha_1 - \tan^2 \alpha_2] - \Delta P_0 S$ .

Similarly, we can decompose we can keep this equation 1 remembered, we can decompose the bracket and then write  $\frac{1}{2} \rho S C_{xm}^2 [\tan \alpha_1 + \tan \alpha_2] [\tan \alpha_1 - \tan \alpha_2] - \Delta P_0 S$ . So,  $F_x$  we can make this  $\frac{1}{2}$  and this bracket and then we can have  $\rho S C_{xm}^2 \cdot \tan \alpha_m [\tan \alpha_1 - \tan \alpha_2] - \Delta P_0 S$ .

So, we can write down force coefficient in x direction as  $\frac{F_x}{\frac{1}{2}(\rho l C_{xm}^2)}$ , and then we will get it as if we use this expression. This will be remember as 2 and then we can get it as,  $\frac{S}{l} [\tan^2 \alpha_1 - \tan^2 \alpha_2] - \frac{S_0}{\frac{1}{2}(\rho l C_{xm}^2)}$ . Now, having this force coefficient in axial direction known and similarly force coefficient in normal direction known.

↳ Lift

$$L = F_y \cos \alpha_m + F_x \sin \alpha_m$$

$$L = \frac{1}{2} \rho S C_{xm}^2 \cos \alpha_m (\tan \alpha_1 - \tan \alpha_2) + \frac{1}{2} \rho S C_{xm}^2 (\tan \alpha_1 - \tan \alpha_2) \cdot \tan \alpha_m \cdot \sin \alpha_m - \Delta P_0 \cdot \sin \alpha_m$$

$$L = \frac{1}{2} \rho S C_{xm}^2 (\tan \alpha_1 - \tan \alpha_2) \left[ \cos \alpha_m + \frac{\sin^2 \alpha_m}{\cos \alpha_m} \right] - \Delta P_0 \sin \alpha_m$$

$$L = \frac{1}{2} \rho S C_{xm}^2 (\tan \alpha_1 - \tan \alpha_2) \cos \alpha_m [1 + \tan^2 \alpha_m] - \Delta P_0 \sin \alpha_m$$

$$L = \frac{1}{2} \rho S C_{xm}^2 (\tan \alpha_1 - \tan \alpha_2) \sec \alpha_m - \Delta P_0 \sin \alpha_m$$

$$C_L = \frac{L}{\frac{1}{2} \rho l C_{xm}^2} = 2 \left( \frac{S}{l} \right) \cdot (\tan \alpha_1 - \tan \alpha_2) \cdot \sec \alpha_m - \frac{\Delta P_0 \cdot \sin \alpha_m}{\frac{1}{2} \rho l C_{xm}^2}$$

↳ drag

$$D = F_y \sin \alpha_m - F_x \cos \alpha_m$$

$$D = \frac{1}{2} \rho S C_{xm}^2 (\tan \alpha_1 - \tan \alpha_2) \cdot \sin \alpha_m - \frac{1}{2} \rho S C_{xm}^2 (\tan \alpha_1 - \tan \alpha_2) \cdot \tan \alpha_m \cdot \cos \alpha_m + \Delta P_0 \cos \alpha_m$$

$$D = \frac{1}{2} \rho S C_{xm}^2 (\tan \alpha_1 - \tan \alpha_2) \cdot \sin \alpha_m - \frac{1}{2} \rho S C_{xm}^2 (\tan \alpha_1 - \tan \alpha_2) \sin \alpha_m + \Delta P_0 \cos \alpha_m$$

$$\therefore D = \Delta P_0 \cos \alpha_m \rightarrow C_D = \frac{D}{\frac{1}{2} \rho l C_{xm}^2} = \frac{\Delta P_0 \cos \alpha_m}{\frac{1}{2} \rho l C_{xm}^2}$$

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(refer time: 46:27). We can find out lift and drag so, lift L is equal to from the decomposition of forces,  $F_y \cos \alpha_m + F_x \sin \alpha_m$ . So, lift from the  $F_y$  expression  $\rho S C_{xm}^2 \cdot \cos \alpha_m \cdot (\tan \alpha_1 - \tan \alpha_2) + \rho S C_{xm}^2 \cdot (\tan \alpha_1 - \tan \alpha_2) \cdot \tan \alpha_m \cdot \sin \alpha_m - \Delta P_0 \cdot \sin \alpha_m$ . So, L is equal to we can take this common from both the terms and then we can take.

We can write down  $\rho SC_{xm}^2 \cdot (\tan \alpha_1 - \tan \alpha_2) \cdot [\cos \alpha_m + \frac{\sin^2 \alpha_m}{\cos \alpha_m}]$ . Here, we have 2 terms, these 2 terms mean this single term which is  $\sin \alpha_m$  divided by  $\cos \alpha_m$ . So,  $S\Delta P_0 \cdot \sin \alpha_m$ . so, we have  $\rho SC_{xm}^2 (\tan \alpha_1 + \tan \alpha_2)$ , we can still take  $\cos \alpha_m$  common once, we take  $\cos \alpha_m$  common.

We know it will decompose into  $\tan^2$ ,  $1 + \tan^2 \alpha_m - S\Delta P_0 \cdot \sin \alpha_m$ . So, we can say that  $\rho SC_{xm}^2 (\tan \alpha_1 - \tan \alpha_2)$  so, this  $\cos \alpha_m$  and then this term is  $\sec^2 \alpha_m$ . So, we will have  $\sec \alpha_m - S\Delta P_0 \cdot \sin \alpha_m$ . So, we can write down L as this and hence, we can further say that we can further find out the lift coefficient as  $\frac{L}{\frac{1}{2} \rho l C_{xm}^2}$ , and then we can have it as  $2 \times \frac{S}{l} \rho C_{xm}$  will get cancelled.

Then, we have  $(\tan \alpha_1 - \tan \alpha_2) \sec \alpha_m - \frac{S_0}{\frac{1}{2} \rho l C_{xm}^2} \cdot \sin \alpha_m$ . So, this is the expression for lift coefficient now, we can find out drag so, drag from the decomposition of force is same  $F_y \sin \alpha_m - F_x \cos \alpha_m$ . So, we can put the expression for  $F_y$ ,  $\rho SC_{xm}^2 \cdot (\tan \alpha_1 - \tan \alpha_2) \cdot \sin \alpha_m - \rho SC_{xm}^2 \cdot (\tan \alpha_1 - \tan \alpha_2) \cdot \tan \alpha_m \cdot \cos \alpha_m + S\Delta P_0 \cdot \cos \alpha_m$ .

So, we have  $D = \rho SC_{xm}^2 \cdot (\tan \alpha_1 - \tan \alpha_2) \cdot \sin \alpha_m - \rho SC_{xm}^2 \cdot (\tan \alpha_1 - \tan \alpha_2) \cdot \sin \alpha_m + S\Delta P_0 \cdot \cos \alpha_m$ . So, we will have  $S\Delta P_0 \cdot \cos \alpha_m$ , we can see that the first 2 terms are same so, we have drag is equal to  $S\Delta P_0 \cdot \cos \alpha_m$ . So, we can write down drag coefficient  $C_D = \frac{D}{\frac{1}{2} \rho l C_{xm}^2}$ . So, it is  $\frac{S\Delta P_0 \cos \alpha_m}{\frac{1}{2} \rho l C_{xm}^2}$ .

So, this is the method by which we can find out the different forces acting on the cascade. So, here what we have seen that what are the different angles associated with the cascade testing or blade design or blade testing the angles like camber angle, air angle, blade angle. So, these are the angles and once we set these angles, we measure the direction of velocities and we can get the drag forces acting on the drag and lift or the  $F_x$  and  $F_y$  forces acting on the cascade.

So, having said this, we have talked about the complete blade design and cascade tunnel and the gas flow through a cascade. Here, we have completed the portion which is related with the blade design and cascade theory. Thank you.