

Two-Phase flow with phase change in conventional and miniature channels

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Lecture - 06

The Drift Flux Model

We meet again for the course on Two-phase flow with phase change in conventional and miniature channels. We were discussing the modeling of two phase flow in that we have discuss the homogeneous model, the separated flow model with two correlations Lockhart Martinelli correlation and Martinelli Nelson correlation and now we will discuss another model called the Drift Flux Model.

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The Drift Flux

Volumetric flux of gas phase

$$j_g = \frac{Q_g}{A} = \frac{U_g A_g}{A} = \alpha U_g = \alpha [j + (U_g - j)]$$
$$j_g = \alpha j + \alpha (U_g - j)$$

Gas drift flux and liquid drift flux

$$j_{gf} = \frac{A_g (U_g - j)}{A} = \alpha (U_g - j) = j_g - \alpha j$$
$$j_{fg} = \frac{A_f (U_f - j)}{A} = (1 - \alpha) (U_f - j) = j_f - (1 - \alpha) j$$
$$j_{gf} + j_{fg} = j_g + j_f - j = 0$$
$$j_{gf} = -j_{fg}$$

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The drift flux model combines the simplicity of homogeneous model, but it takes into account the relative motion between the phases like the separated flow model. So, it is simple as well as relatively accurate, it involves the concept of drift flux. So, let us first understand what is drift flux.

The Drift Flux before we define the drift flux let us consider the volumetric flux of the gas phase which is j_g it is Q_g upon A and Q_g is $U_g A_g$. So, it is $U_g A_g$ upon A and A_g by A is equal to α . So, this is equal to αU_g and U_g we can rewrite as j plus U_g minus j . So, it becomes α into j plus U_g minus j ; now, why I have written U_g in this form? Because this U_g minus j is the difference between the gas velocity and the total volumetric flux, this j is the total volumetric flux and it is also equal to the mean velocity of the two phase mixture. And therefore, U_g minus j is the relative velocity of the gas phase with respect to the plane which is moving with the mean velocity j ok.

Now j_g can be written as this αj plus α times U_g minus j , now consider the gas drift flux. The gas drift flux is denoted by j_{gf} and it is the volumetric flux of the gas phase with respect to the plane moving with the mean velocity j . The volumetric flow rate of the gas phase with respect to this plane which is moving with the mean velocity j a volumetric flow rate will be equal to the area A_g occupied by the gas phase multiplied by the relative velocity of the gas phase with respect to the plane moving with mean velocity j .

And this whole thing has to be divided by the total cross sectional area A . But A_g by A is nothing, but α ; so, this is equal to α times U_g minus j . And we recognized that αU_g is nothing, but j_g ; do, get j_g minus αj . In a similar manner j_{fg} the liquid drift flux is defined as the volumetric flux of the liquid phase with respect to the same plane, which is moving with mean velocity j . So, this will be equal to A_f multiplied by the relative velocity U_f minus j divided by the total cross sectional area A , but A_f by A is nothing, but 1 minus α .

So, we get 1 minus α times U_g minus j and 1 minus α U_f is nothing, but j_f . So, this is equal to j_f minus 1 minus αj . So, if we add j_{gf} plus j_{fg} then we get 0 ok. So,

therefore, j_g is the negative j_l . This means that the gas drift flux and liquid drift flux are equal and opposite that is the volumetric flux of the gas phase with respect to the plane moving with mean velocity j is equal and opposite of the volumetric flux of the liquid phase with respect to the same plane.

Now let us come back to this equation j_g is equal to αj plus α times U_g minus j .

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$$j_g = \alpha j + \alpha (U_g - j)$$

Area average $\langle \xi \rangle = \frac{1}{A} \int_A \xi dA$

$$\langle j_g \rangle = \langle \alpha j \rangle + \langle \alpha (U_g - j) \rangle$$

$$\langle j_g \rangle = \frac{\langle \alpha j \rangle}{\langle \alpha \rangle \langle j \rangle} \langle \alpha \rangle \langle j \rangle + \frac{\langle \alpha (U_g - j) \rangle}{\langle \alpha \rangle} \langle \alpha \rangle$$

$$\langle j_g \rangle = C_0 \langle \alpha \rangle \langle j \rangle + V_{gj} \langle \alpha \rangle$$

where C_0 = two-phase distribution coefficient or concentration coefficient (accounts for global or overall slip due to flow area averaging)

$$C_0 = \frac{\langle \alpha j \rangle}{\langle \alpha \rangle \langle j \rangle}$$

V_{gj} = gas drift velocity (represents local slip)

$$V_{gj} = \frac{\langle \alpha (U_g - j) \rangle}{\langle \alpha \rangle}$$

Handwritten notes: $\langle j+g \rangle = \langle j \rangle + \langle g \rangle$
 $\langle jg \rangle \neq \langle j \rangle \langle g \rangle$

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Let us define area average and denote it by this angular brackets suppose ψ is some variable, then the area average of this variable is equal to the integral of ψ over the area A divided by the area A . And from this we can clearly see that the sum of the area averages of two quantities will be equal to the area average of the sum of the two quantities. Suppose f and g are two functions, then area average of f plus g will be equal to area average of f plus area average of g .

But if we multiply two functions then the area average of the product of the two functions will not be equal to the product of the area averages. So, in an equation we can take area averages of each term and the equation will still be valid. So, now, in this equation we take area average of each term. So, we get area average $j g$ or angular bracket $j g$ is equal to angular bracket αj plus angular bracket α times U_g minus j .

Now, we divide and multiply all the terms by αj . A first term we multiply by angular bracket α times angular bracket j and also divide by angular bracket α and angular bracket j and the second term we multiply by angular bracket α and divide by angular bracket α . And now this quantity we define a C_0 and this quantity here we called it V_{gj} . So, now, we have this equation angular bracket $j g$ is equal to C_0 angular bracket α , angular bracket j plus V_{gj} angular bracket α now these two new parameters which we have defined they are as follows.

C_0 is called the two phase distribution coefficient or concentration coefficient and it accounts for global or overall slip due to flow area averaging. C_0 is equal to area average of αj upon area average α into area average j and as we have discussed area average of the product of two functions is not equal to the product of the area averages of the two functions, it can happen sometimes only if the two functions vary in a similar manner.

If the void fraction profile or the concentration profile and the velocity profile are similar, then the numerator and the denominator will be equal and in that case C_0 will be

equal to 1; otherwise if there void fraction profile or the concentration profile and the velocity profile are not similar then C_0 will be different from 1. So, therefore, C_0 accounts for the global effect of the slip between the two phases that is due to that this similarity of the concentration profile and velocity profile and this effect comes due to flow area averaging.

Now, the second parameter V_{gj} it is called the gas drift velocity and it represents the local slip. As we can see here V_{gj} involves the related local relative velocity U_g minus j and therefore, it represents the effect of local slip that is the relative velocity between the phases at every point on the cross section and then it is averaged over the cross section.

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$$\langle j_g \rangle = C_0 \langle \alpha \rangle \langle j \rangle + V_{gj} \langle \alpha \rangle = \langle \alpha \rangle [C_0 \langle j \rangle + V_{gj}]$$

$$\langle \alpha \rangle = \frac{\langle j_g \rangle}{C_0 \langle j \rangle + V_{gj}}$$

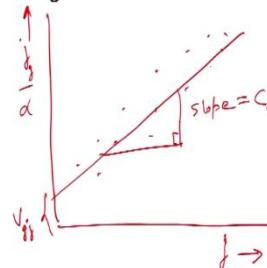
$$\frac{\langle j_g \rangle}{\langle \alpha \rangle} = C_0 \langle j \rangle + V_{gj} \quad \checkmark$$

C_0 and V_{gj} can be found experimentally by plotting the line

$$\frac{j_g}{\alpha} = C_0 j + V_{gj} \quad \checkmark \checkmark$$

The void fraction can be evaluated from

$$\alpha = \frac{j_g}{C_0 j + V_{gj}} \quad \checkmark$$



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So, now j_g is equal to $C_0 \alpha j + V_{gj} \alpha$. So, now we take this angular bracket α common and we have $C_0 \alpha j + V_{gj} \alpha$ and then we can express angular bracket α as angular bracket j_g upon angular $C_0 \alpha j + V_{gj} \alpha$. This is the expression for the void fraction that we have obtained.

We can rewrite this expression in this form also j_g upon α is equal to $C_0 j + V_{gj}$. Now for simplicity of notation we can get rid of angular brackets and rewrite all the expressions without the angular brackets, but it is understood that we are dealing with one dimensional modeling and therefore, area averaging has been done.

So, j_g by α is equal to $C_0 j + V_{gj}$ this is the expression we are having; and α can be expressed as j_g upon $C_0 j + V_{gj}$. So, if we can somehow obtain C_0 and V_{gj} , we should be able to obtain α the void fraction from this expression. So, how to obtain C_0 and V_{gj} for that consider this equation. If we plot j on the x axis and j_g upon α from the y axis, then and we plot the data point obtained experimentally and then fit a line.

So, then this y intercept will be equal to V_{gj} and the slope of this line will be equal to C_0 . So, this is the way correlations for C_0 and V_{gj} have been obtained, researchers have done experiments collected data and then after plotting graph, they have obtained V_{gj} and C_0 . Some researchers have used data containing different flow regimes whereas, some others researchers have used data for some specific flow regimes. So, if the data for different flow regimes has been used, then the resulting correlation will be applicable regardless of the flow regime. If the data has been used for a particular flow regime, then that correlation obtained from that data will be applicable for that particular flow regime.

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DFM parameters without reference to any flow regime

$$C_0 = 1.13, \quad V_{gj} = 1.41 \left(\frac{\sigma g \Delta \rho}{\rho_f^2} \right)^{0.25}$$

$\Delta \rho = \rho_f - \rho_g$
if $\rho_g \gg \rho_f$ then $\Delta \rho \approx \rho_g$

For turbulent cocurrent slug flow in vertical pipes

$$C_0 = 1.2, \quad V_{gj} = 0.35 \left(\frac{\sigma g \Delta \rho}{\rho_f} \right)^{0.5}$$

For turbulent cocurrent slug flow in vertical pipes, with $\rho_f \gg \rho_g$

$$C_0 = 1.2, \quad V_{gj} = 0.35 \sqrt{g D}$$

For turbulent slug flow in horizontal pipes $C_0 = 1.2, \quad V_{gj} = 0$

$$\Rightarrow \alpha = \frac{j_g}{1.2 j} = 0.833 \beta$$

For bubbly and slug flow in minichannels ($D < 1 \text{ mm}$)

$$C_0 = 1.2 + 0.510 e^{-0.692 D}, \quad V_{gj} = 0$$

For subcooled and saturated flow boiling

$$C_0 = 1 + 0.12 (1 - x), \quad V_{gj} = 1.18 (1 - x) \left(\frac{\sigma g \Delta \rho}{\rho_f^2} \right)^{0.25}$$

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So, different correlations have been given for C_{naught} and v_{gj} here I am listing only a few of them. There is a correlation which is without reference to any particular flow regime and this gives C_{naught} as 1.13 and V_{gj} is given by $1.41 \frac{\sigma g \Delta \rho}{\rho_f^2}$ raised to 0.25. Here as we know $\Delta \rho$ is the density difference $\rho_f - \rho_g$ and in many cases the density of the liquid phase is much greater than the density of the gas phase.

So, then $\Delta \rho$ is approximately equal to ρ_f only and if this approximation is used, then the numerator becomes the ρ_f and $1 - \rho_f$ will get canceled with the ρ_f in the denominator. Then there are some other correlations which are for some particular flow regime. For turbulent concurrent slug flow in vertical pipes we have this correlation C_{naught} is equal to 1.2 V_{gj} is equal to $0.35 \frac{\sigma g \Delta \rho}{\rho_f}$ raised to 0.5. And if it is the same type of flow turbulent slug flow in vertical pipes, but the density of the liquid phase is much larger than the density of the gas phase then as we have discussed $\Delta \rho$ will be equal to ρ_f and $\rho_f - \rho_f$ will get canceled.

So, in that case we get V_{gj} is equal to $0.35 \sqrt{g D}$. For turbulent flow slug flow in horizontal pipes we have C_{naught} equal to 1.2 and V_{gj} is equal to 0. So, if we go back to this expression and we put here V_{gj} equal to 0 then we get C_{naught} upon this will be equal to $j g$ upon $C_{naught} j$, but $j g$ by j is nothing, but β . So, we get β upon C_{naught} , ok.

So, in this case we get α is equal to 0.833β . So, bubbly and slug flow in many channels which is less than 1 and diameter, there is a correlation C_{naught} equal to $1.2 + 0.510 e^{-0.692 D}$ and V_{gj} is equal to 0.

For subcooled and saturated flow boiling there is another correlation which involves the quality x C_{naught} is equal to $1 + 0.12 (1 - x)$ and V_{gj} is equal to $1.18 (1 - x) \frac{\sigma g \Delta \rho}{\rho_f^2}$ raised to 0.25. These are a few correlations for drift flux model there are many more and you can use whichever is most suitable for your problem.

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Calculation of pressure gradients due to acceleration and gravity

$$\alpha = \frac{j_g}{C_0 j + V_{gj}}$$

$$\alpha = \frac{Gxv_g}{C_0 \{Gxv_g + G(1-x)v_f\} + V_{gj}}$$

$$\left(\frac{\partial \alpha}{\partial x}\right)_p = \frac{\alpha}{x} - \frac{\alpha^2 C_0 v_{fg}}{xv_g}$$

$$v^* = \left\{ \frac{2xv_g}{\alpha} - \frac{2(1-x)v_f}{1-\alpha} \right\} + \left(\frac{\partial \alpha}{\partial x}\right)_p \left\{ \frac{(1-x)^2 v_f}{(1-\alpha)^2} - \frac{x^2 v_g}{\alpha^2} \right\}$$

$$-\left(\frac{dP}{dz}\right)_a = \frac{1}{(1-M^2)} \left\{ G^2 \frac{dx}{dz} v^* \right\}$$

$$-\left(\frac{dP}{dz}\right)_z = [\rho_g \alpha + \rho_f (1-\alpha)] g \sin \theta$$

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Now, how to calculate pressure gradients due to acceleration and gravity? For pressure gradient due to friction we use the same expression as in the homogeneous model, because the frictional pressure gradient does not involve with the void fraction alpha. But the acceleration and gravity pressure gradients will be different if we use the drift flux model and these will involve the effect of relative motion between the phases. The alpha is

calculated from this expression and if we substitute for j then we get this j is equal to Gx v g plus G 1 minus x v f; this one this is j g and this is j f ok.

This is also j g if we differentiate in this expression with respect to quality x at constant pressure, that constant pressure v g and v f will be constant and C naught is constant and V gj let us look at the expressions here, the expressions for V gj involve rho f rho g and sigma. So, these properties if the pressure is constant and we are considering thermal equilibrium. So, it is saturated liquid vapor mixture. So, if the pressure is constant the temperature will also be constant that it will be equal to the saturation temperature and therefore, the surface tension also will be constant and rho f and rho j will be constant.

So, therefore, V gj will be constant. So, considering V gj v g and v fs constant we differentiate alpha with respect to x and after some algebra we get this expression, alpha upon x minus alpha square C naught v fg upon x vg. Now this expression is used to calculate v star and then this value of v star is used to calculate the pressure gradient due to acceleration and the value of alpha obtained from the drift flux model is used here to calculate the pressure gradient due to gravity.

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Example-1: Water+steam @100 kPa, horizontal flow, D = 2 mm, L = 10 cm
 $G = 100 \text{ kg/m}^2\text{s}$, $x(0)=0$, $q'' = 20 \text{ kW/m}^2$
 To find the pressure gradient at $z = 5 \text{ cm}$

Solution:

Properties of water+ steam @100 kPa

$$\mu_f = 282.9 \times 10^{-6} \text{ Pa.s}, \mu_g = 12.26 \times 10^{-6} \text{ Pa.s}, h_{fg} = 2257.45 \text{ kJ/kg}$$

$$v_f = 1.043 \times 10^{-3} \text{ m}^3/\text{kg}, v_g = 1.6939 \text{ m}^3/\text{kg}, v_{fg} = 1.693 \text{ m}^3/\text{kg} \quad \checkmark$$

$$\frac{dx}{dz} = \frac{4q''}{GDh_{fg}} = 0.443 \text{ m}^{-1}, x(5\text{cm}) = 0.0221 \quad \checkmark$$

$$\frac{dv_g}{dP} \approx \frac{\Delta v_g}{\Delta P} = \frac{1.6782 - 1.6939}{1000} = -1.57 \times 10^{-5} \text{ m}^3 \text{ kg}^{-1} \text{ Pa}^{-1} \quad \checkmark$$

$$M^2 = G^2 x \left| \frac{dv_g}{dP} \right| = 3.47 \times 10^{-3} \ll 1, 1 - M^2 \approx 1, (1 - M^2)^{-1} \approx 1$$

$$\frac{1}{\bar{\mu}} = \frac{x}{\mu_g} + \frac{1-x}{\mu_f} \Rightarrow \bar{\mu} = 190.1 \times 10^{-6} \text{ Pa.s} \quad Re_{TP} = \frac{GD}{\bar{\mu}} = 1052 \Rightarrow \text{Laminar flow}$$

$$f_{TP} = 16/Re = 0.01521 \quad \checkmark$$

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Now, we will consider some numerical examples based on the drift flux model. Again we consider the first example with the same data as we have considered before for the other models, water steam mixture at a 100 kilo Pascal, horizontal flow 2 mm diameter channel and length 10 centimeter, mass flux 100 kg per meter square second. In let quality 0, drift flux 20 kilo Watt per meter square and we have to find the pressure gradient at the midpoint z is equal to 5 centimeter the properties of water steam at 100 kilo Pascal. At the same as before, dx by dz is already calculated before 0.0221, dx by dz is 0.443 per meter and x is equal to 0.221.

And dv_g by dp is calculated using the numerical differentiation as minus 1.57 into 10 raise to minus 5 s i units and then we calculate m square which is 3.47 into 10 raise to minus 3 and this is much less than 1. And therefore, $1 - M$ square is a approximately equal to 1 and 1 upon $1 - M$ square is also approximately equal to 1. Then we calculate the mixture viscosity and then the two phase Reynolds number, and it is laminar flow therefore, we use the laminar flow correlation 16 upon Re and got this friction factor and then using this friction factor we get the frictional pressure gradient.

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$$-\left(\frac{dP}{dz}\right)_F = \frac{2f_{TP}}{D} G^2 \bar{v} = 2.74 \text{ kPa/m}$$

$$\alpha = \frac{j_g}{C_0 j + V_{gj}} = 0.82, \quad \beta = 0.973$$

$$C_0 = 1.13, \quad V_{gj} = 1.41 \left(\frac{\sigma g \Delta \rho}{\rho_f^2} \right)^{0.25}$$

Now,

$$-\left(\frac{dP}{dz}\right)_a = \frac{1}{(1-M^2)} \left\{ G^2 \frac{dx}{dz} v^* \right\}$$

$$1 - M^2 \approx 1$$

$$v^* = \left\{ \frac{2xv_g}{\alpha} - \frac{2(1-x)v_f}{1-\alpha} \right\} + \left(\frac{\partial \alpha}{\partial x} \right)_p \left\{ \frac{(1-x)^2 v_f}{(1-\alpha)^2} - \frac{x^2 v_g}{\alpha^2} \right\}$$

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Now alpha the void fraction is obtained from this expression j_g upon $C_0 j + V_{gj}$, and C_0 here is equal to 1.13 and V_{gj} is calculated using this expression and after substituting we get alpha is equal to 0.82 and beta is equal to $x v_g$ upon \bar{v} . So, beta is calculated is 0.973 and we note that here alpha is less than beta irrespective. Now using this alpha we calculate the pressure gradient due to acceleration and for that we will have to calculate v^* first.

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$$\alpha = \frac{j_g}{C_0 j + V_{gj}} \quad \checkmark$$

$$\alpha = \frac{Gxv_g}{C_0 \{Gxv_g + G(1-x)v_f\} + V_{gj}} \quad \checkmark$$

$$\left(\frac{\partial \alpha}{\partial x}\right)_p = \frac{\alpha}{x} - \frac{\alpha^2 C_0 v_{fg}}{xv_g} = 2.75 \quad \checkmark$$

$$v^* = \left\{ \frac{2xv_g}{\alpha} - \frac{2(1-x)v_f}{1-\alpha} \right\} + \left(\frac{\partial \alpha}{\partial x}\right)_p \left\{ \frac{(1-x)^2 v_f}{(1-\alpha)^2} - \frac{x^2 v_g}{\alpha^2} \right\} = \underline{0.1612} \text{ m}^3/\text{kg}$$

$$-\left(\frac{dP}{dz}\right)_a = \frac{1}{(1-M^2)} \left\{ G^2 \frac{dx}{dz} v^* \right\} = \underline{0.714} \text{ kPa/m}$$

$$\alpha = \frac{j_g}{C_0 j + V_{gj}}$$

$$\alpha = \frac{Gxv_g}{C_0 \{Gxv_g + G(1-x)v_f\} + V_{gj}}$$

$$\frac{\partial \alpha}{\partial x} = \frac{\alpha}{x} - \frac{\alpha^2 C_0 v_{fg}}{xv_g} = 2.75$$

$$v^* = \left\{ \frac{2xv_g}{\alpha} - \frac{2(1-x)v_f}{1-\alpha} \right\} + \left(\frac{\partial \alpha}{\partial x}\right)_p \left\{ \frac{(1-x)^2 v_f}{(1-\alpha)^2} - \frac{x^2 v_g}{\alpha^2} \right\} = 0.1612 \text{ m}^3/\text{s}$$

$$-\left(\frac{dP}{dz}\right)_a = \frac{1}{(1-M^2)} \left\{ G^2 \frac{dx}{dz} v^* \right\} = 0.714 \text{ kPa/m}$$

So, alpha is given by this and after substituting we get this and then after differentiating and substituting the values we get 2.75, the slope of the alpha x curve at constant pressure is 2.75.

Now, using this value of the slope we get v star equal to 0.1612 meter cube per kg and then using this value of v star we get the acceleration pressure gradient as 0.714 kilo Pascal per meter.

(Refer Slide Time: 27:39)

$$-\left(\frac{dP}{dz}\right)_F = \underline{2.74} \text{ kPa/m}$$

$$-\left(\frac{dP}{dz}\right)_a = \underline{0.714} \text{ kPa/m}$$

$$-\left(\frac{dP}{dz}\right)_z = \textcircled{0}$$

$$-\left(\frac{dP}{dz}\right) = 5.85 + 0.714 + 0 = \underline{\underline{3.45}} \text{ kPa/m}$$

$$-\left(\frac{dP}{dz}\right)_F = 2.74 \text{ kPa/m}$$

$$-\left(\frac{dP}{dz}\right)_a = 0.714 \text{ kPa/m}$$

$$-\left(\frac{dP}{dz}\right)_z = 0$$

$$-\left(\frac{dP}{dz}\right) = 2.74 + 0.714 + 0 = 3.45 \text{ kPa/m}$$

So, now here are the results frictional pressure gradient 2.74 kilo Pascal per meter acceleration pressure gradient 0.714 kilo Pascal per meter and the gravitational pressure gradient is 0 because it is a horizontal pipe, and we now add all 3 pressure gradients and get the total pressure gradient as 3.45 kilo Pascal per meter.

(Refer Slide Time: 28:09)

Example-2: Water+steam @100 kPa, vertical upward flow, D=2 cm, L=2m
 $G = 1000 \text{ kg/m}^2\text{s}$, $x(0)=0$, $x(L)=2\%$
 To find the pressure gradient at $z=1\text{m}$

Solution: Properties of water+ steam @100 kPa

$$\mu_f = 282.9 \times 10^{-6} \text{ Pa.s}, \mu_g = 12.26 \times 10^{-6} \text{ Pa.s}$$

$$v_f = 1.043 \times 10^{-3} \text{ m}^3/\text{kg}, v_g = 1.6939 \text{ m}^3/\text{kg}, v_{fg} = 1.693 \text{ m}^3/\text{kg}$$

$$x(1\text{m}) = 0.01, \quad \frac{dx}{dz} = \frac{0.02}{2} = 0.01 \text{ m}^{-1}$$

$$\frac{dv_g}{dP} \approx \frac{\Delta v_g}{\Delta P} = \frac{1.6782 - 1.6939}{1000} = -1.57 \times 10^{-5} \text{ m}^3 \text{ kg}^{-1} \text{ Pa}^{-1}$$

$$M^2 = G^2 x \left| \frac{dv_g}{dP} \right| = 0.157, 1 - M^2 = 0.843, (1 - M^2)^{-1} = 1.186$$

$$\frac{1}{\bar{\mu}} = \frac{x}{\mu_g} + \frac{1-x}{\mu_f} \Rightarrow \bar{\mu} = 232 \times 10^{-6} \text{ Pa.s}$$

$$Re_{TP} = \frac{GD}{\bar{\mu}} = 8.62 \times 10^4 \Rightarrow \text{Turbulent flow}$$

$$f_{TP} = 0.079 Re_{TP}^{-0.25} = 4.61 \times 10^{-3}$$

Example-2: Water+steam @ 100 kPa, vertical upward flow, D=2 cm, L=2m

$G = 1000 \text{ kg/m}^2\text{s}$, $x(0)=0$, $x(L)=2\%$. To find the pressure gradient at $z=1\text{m}$

Solution: Properties of water+ steam @100 kPa

$$\mu_f = 282.9 \times 10^{-6} \text{ Pa.s}, \mu_g = 12.26 \times 10^{-6} \text{ Pa.s},$$

$$v_f = 1.043 \times 10^{-3} \text{ m}^3/\text{kg}, v_g = 1.6939 \text{ m}^3/\text{kg}, v_{fg} = 1.693 \text{ m}^3/\text{kg}$$

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$$M^2 = G^2 x \left| \frac{dv_g}{dP} \right| = 0.157, 1 - M^2 = 0.843, (1 - M^2)^{-1} = 1.186$$

$$\frac{1}{\bar{\mu}} = \frac{x}{\mu_g} + \frac{1-x}{\mu_f} \Rightarrow \bar{\mu} = 232 \times 10^{-6} \text{ Pa.s}$$

$$Re_{TP} = \frac{GD}{\bar{\mu}} = 8.62 \times 10^4 \Rightarrow \text{Turbulent flow}$$

$$f_{TP} = 0.079 Re_{TP}^{-0.25} = 4.61 \times 10^{-3}$$

Now, consider the second example with the same data as before 100 kilo Pascal pressure vertical upward flow, diameter 2 centimeter, length 2 meter, mass flux 1000 units, inlet quality 0, outlet quality 2 percent and we have to find the pressure gradient at the midpoint. So, property is at 100 kilo Pascal are same as before their quality at the midpoint is 1 percent dx by dz is 0.01 per meter and dv_g by dP is minus 1.157 into 10 raise to minus 5 units and M square is equal to 0.157 which is not negligible $1 - M$ square is equal to 0.843 and 1 upon $1 - M$ square is 1.186.

The average viscosity of the mixture is calculated as this and then the two phase Reynolds number is 8.62 into 10 raise to 4. So, it is turbulent flow and then using the turbulent flow correlation the two phase friction factor is 4.61 into 10 raise to minus 3 and using this frictional pressure gradient is obtained is 9.85 kilo Pascal per meter.

(Refer Slide Time: 29:49)

$$-\left(\frac{dP}{dz}\right)_F = \frac{2f_{TP}}{D} G^2 \bar{v} (1 - M^2)^{-1} = \underline{9.85 \text{ kPa/m}}$$

$$\alpha = \frac{j_g}{C_0 j + V_{gj}} = \underline{0.82}, \quad \beta = \underline{0.943}$$

Now,

$$-\left(\frac{dP}{dz}\right)_a = \frac{1}{(1 - M^2)} \left\{ G^2 \frac{dx}{dz} v^* \right\}$$

$$1 - M^2 \approx 0.843$$

$$v^* = \left\{ \frac{2xv_g}{\alpha} - \frac{2(1-x)v_f}{1-\alpha} \right\} + \left(\frac{\partial \alpha}{\partial x} \right)_p \left\{ \frac{(1-x)^2 v_f}{(1-\alpha)^2} - \frac{x^2 v_g}{\alpha^2} \right\}$$

$$-\left(\frac{dP}{dz}\right)_F = \frac{2f_{TP}}{D} G^2 \bar{v} (1 - M^2)^{-1} = 9.85 \text{ kPa/m}$$

$$\alpha = \frac{j_g}{C_0 j + V_{gj}} = 0.82, \quad \beta = 0.943$$

Now,

$$-\left(\frac{dP}{dz}\right)_a = \frac{1}{(1 - M^2)} \left\{ G^2 \frac{dx}{dz} v^* \right\}$$

$$1 - M^2 \approx 0.843$$

$$v^* = \left\{ \frac{2xv_g}{\alpha} - \frac{2(1-x)v_f}{1-\alpha} \right\} + \left(\frac{\partial \alpha}{\partial x} \right)_P \left\{ \frac{(1-x)^2 v_f}{(1-\alpha)^2} - \frac{x^2 v_g}{\alpha^2} \right\}$$

Now alpha is obtained by using this expression and it is 0.82 and beta is equal to 0.943. The acceleration pressure gradient is obtained as using this expression and v star is given by this expression.

(Refer Slide Time: 30:31)

$$\alpha = \frac{j_g}{C_0 j + V_{gj}}$$

$$\alpha = \frac{Gxv_g}{C_0 \{Gxv_g + G(1-x)v_f\} + V_{gj}}$$

$$\left(\frac{\partial \alpha}{\partial x} \right)_P = \frac{\alpha}{x} - \frac{\alpha^2 C_0 v_{fg}}{xv_g} = 5.62$$

$$v^* = \left\{ \frac{2xv_g}{\alpha} - \frac{2(1-x)v_f}{1-\alpha} \right\} + \left(\frac{\partial \alpha}{\partial x} \right)_P \left\{ \frac{(1-x)^2 v_f}{(1-\alpha)^2} - \frac{x^2 v_g}{\alpha^2} \right\} = 0.216$$

$$-\left(\frac{dP}{dz} \right)_a = \frac{1}{(1-M^2)} \left\{ G^2 \frac{dx}{dz} v^* \right\} = 2.56 \text{ kPa/m}$$

$$\alpha = \frac{j_g}{C_0 j + V_{gj}}$$

$$\alpha = \frac{Gxv_g}{C_0 \{Gxv_g + G(1-x)v_f\} + V_{gj}}$$

$$\frac{\partial \alpha}{\partial x} = \frac{\alpha}{x} - \frac{\alpha^2 C_0 v_{fg}}{xv_g} = 5.62$$

$$v^* = \left\{ \frac{2xv_g}{\alpha} - \frac{2(1-x)v_f}{1-\alpha} \right\} + \left(\frac{\partial \alpha}{\partial x} \right)_P \left\{ \frac{(1-x)^2 v_f}{(1-\alpha)^2} - \frac{x^2 v_g}{\alpha^2} \right\} = 0.216$$

$$-\left(\frac{dP}{dz} \right)_a = \frac{1}{(1-M^2)} \left\{ G^2 \frac{dx}{dz} v^* \right\} = 2.56 \text{ kPa/m}$$

So, alpha we find the slope of the alpha x square and it is 5.62 and then substitute the slope here and get v star as 0.216 and then substitute v star here and get the acceleration pressure gradient as 2.56 kilo Pascal per meter.

(Refer Slide Time: 30:59)

$$-\left(\frac{dP}{dz}\right)_F = 9.85 \text{ kPa/m}$$

$$-\left(\frac{dP}{dz}\right)_a = 2.56 \text{ kPa/m}$$

$$-\left(\frac{dP}{dz}\right)_z = [\rho_g \alpha + \rho_f (1 - \alpha)] g \sin \theta = 1.648 \text{ kPa/m}$$

$$-\left(\frac{dP}{dz}\right) = 9.85 + 2.56 + 1.648 = 14.06 \text{ kPa/m}$$

$$-\left(\frac{dP}{dz}\right)_F = 9.85 \text{ kPa/m}$$

$$-\left(\frac{dP}{dz}\right)_a = 2.56 \text{ kPa/m}$$

$$-\left(\frac{dP}{dz}\right)_z = [\rho_g \alpha + \rho_f (1 - \alpha)] g \sin \theta = 1.648 \text{ kPa/m}$$

$$-\left(\frac{dP}{dz}\right) = 9.85 + 2.56 + 1.648 = 14.06 \text{ kPa/m}$$

So, now the results are the following, the fictional pressure gradient is 9.85 kilo Pascal per meter, acceleration pressure gradient is 2.56 kilo Pascal per meter and for the gravitational pressure gradient we use this expression and we get 1.648 kilo Pascal per meter and adding all this pressure gradients we get the total pressure gradient as 14.06 kilo Pascal per meter.

(Refer Slide Time: 31:35)

Example-3: Water+steam @10 MPa, vertical upward flow, D=2 cm, L=2m

$$G = 1000 \text{ kg/m}^2\text{s}, x(0)=0, x(L)=2\%$$

To find the pressure gradient at z=1m

Solution:

Properties of water+ steam @10 MPa

$$\mu_f = 81.80 \times 10^{-6} \text{ Pa.s}, \mu_g = 20.27 \times 10^{-6} \text{ Pa.s}$$

$$v_f = 1.453 \times 10^{-3} \frac{\text{m}^3}{\text{kg}}, v_g = 1.803 \times 10^{-2} \text{ m}^3/\text{kg}, v_{fg} = 0.01658 \text{ m}^3/\text{kg}$$

$$x(1\text{m}) = 0.01, \frac{dx}{dz} = \frac{0.02}{2} = 0.01 \text{ m}^{-1}$$

$$\frac{dv_g}{dP} \approx \frac{\Delta v_g}{\Delta P} = \frac{0.01781 - 0.01803}{1 \times 10^5} = -2.20 \times 10^{-9} \text{ m}^3 \text{ kg}^{-1} \text{ Pa}^{-1}$$

$$M^2 = G^2 x \left| \frac{dv_g}{dP} \right| = 2.20 \times 10^{-5} \ll 1, 1 - M^2 \approx 1, (1 - M^2)^{-1} \approx 1$$

$$\frac{1}{\bar{\mu}} = \frac{x}{\mu_g} + \frac{1-x}{\mu_f} \Rightarrow \bar{\mu} = 79.4 \times 10^{-6} \text{ Pa.s}$$

$$Re_{TP} = \frac{GD}{\bar{\mu}} = \frac{2.52 \times 10^5}{79.4 \times 10^{-6}} \Rightarrow \text{Turbulent flow}$$

$$f_{TP} = 0.079 Re_{TP}^{-0.25} = 3.5 \times 10^{-3}$$

Example-3: Water+steam @ 10 MPa, vertical upward flow, D=2 cm, L=2m

$G = 1000 \text{ kg/m}^2\text{s}$, $x(0)=0$, $x(L)=2\%$. To find the pressure gradient at $z=1\text{m}$

Solution: Properties of water+ steam @10 MPa

$$\mu_f = 81.80 \times 10^{-6} \text{ Pa.s}, \mu_g = 20.27 \times 10^{-6} \text{ Pa.s},$$

$$v_f = 1.453 \times 10^{-3} \frac{\text{m}^3}{\text{kg}}, v_g = 1.803 \times 10^{-2} \text{ m}^3/\text{kg}, v_{fg} = 0.01658 \text{ m}^3/\text{kg}$$

$$x(1\text{m}) = 0.01, \frac{dx}{dz} = \frac{0.02}{2} = 0.01 \text{ m}^{-1}$$

$$\frac{dv_g}{dP} \approx \frac{\Delta v_g}{\Delta P} = \frac{0.01781 - 0.01803}{1 \times 10^5} = -2.20 \times 10^{-9} \text{ m}^3 \text{ kg}^{-1} \text{ Pa}^{-1}$$

$$M^2 = G^2 x \left| \frac{dv_g}{dP} \right| = 2.20 \times 10^{-5} \ll 1, 1 - M^2 \approx 1, (1 - M^2)^{-1} \approx 1$$

$$\frac{1}{\bar{\mu}} = \frac{x}{\mu_g} + \frac{1-x}{\mu_f} \Rightarrow \bar{\mu} = 79.4 \times 10^{-6} \text{ Pa.s}$$

$$Re_{TP} = \frac{GD}{\bar{\mu}} = 2.52 \times 10^5 \Rightarrow \text{Turbulent flow}$$

$$f_{TP} = 0.079 Re_{TP}^{-0.25} = 3.5 \times 10^{-3}$$

Now the third example data same as before 10 mega Pascal pressure vertical upward flow, diameter 2 centimeter length 2 meter, mass flux 1000 kg per meter square second, inlet quality 0, outlet quality 2 percent pressure gradient has to be found at the midpoint. At 10 mega Pascal these are the properties at the midpoint the quality is 1 percent and dx by dz is equal to 0.01 per meter dv g by dP is of the order of the 10 raise to minus 9 m square is of the order of 10 raise to minus 5 negligible.

So, $1 - M^2$ is approximately equal to 1, the mean viscosity of the mixture is 79.4×10^{-6} Pascal second, the Reynolds numbers is of the order of 10^5 . So, it is turbulent flow.

And using the turbulent flow correlation we get the two phase friction factor 3.5×10^{-3} .

(Refer Slide Time: 32:51)

$$-\left(\frac{dP}{dz}\right)_F = \frac{2f_{TP}}{D} G^2 \bar{v} (1 - M^2)^{-1} = \underline{0.572 \text{ kPa/m}}$$

$$\alpha = \frac{j_g}{C_0 j + V_{gj}} = \underline{0.091}, \quad \beta = \underline{0.1114} \quad \alpha < \beta$$

Now,

$$-\left(\frac{dP}{dz}\right)_a = \frac{1}{(1 - M^2)} \left\{ G^2 \frac{dx}{dz} v^* \right\}$$

$$1 - M^2 \approx 1$$

$$v^* = \left\{ \frac{2xv_g}{\alpha} - \frac{2(1-x)v_f}{1-\alpha} \right\} + \left(\frac{\partial \alpha}{\partial x} \right)_p \left\{ \frac{(1-x)^2 v_f}{(1-\alpha)^2} - \frac{x^2 v_g}{\alpha^2} \right\}$$

$$-\left(\frac{dP}{dz}\right)_F = \frac{2f_{TP}}{D} G^2 \bar{v} (1 - M^2)^{-1} = 0.572 \text{ kPa/m}$$

$$\alpha = \frac{j_g}{C_0 j + V_{gj}} = 0.091, \quad \beta = 0.1114$$

Now,

$$-\left(\frac{dP}{dz}\right)_a = \frac{1}{(1 - M^2)} \left\{ G^2 \frac{dx}{dz} v^* \right\}$$

$$1 - M^2 \approx 1$$

$$v^* = \left\{ \frac{2xv_g}{\alpha} - \frac{2(1-x)v_f}{1-\alpha} \right\} + \left(\frac{\partial \alpha}{\partial x} \right)_P \left\{ \frac{(1-x)^2 v_f}{(1-\alpha)^2} - \frac{x^2 v_g}{\alpha^2} \right\}$$

Now, using this, we get the frictional pressure gradient as 0.572 kilo Pascal per meter, alpha using this expression we get 0.091 and beta is equal to 0.1114. So, here also alpha is less than beta as expected. And for the frictional pressure gradient this is the expression it involves v star.

(Refer Slide Time: 33:37)

$$\alpha = \frac{j_g}{C_0 j + V_{gj}}$$

$$\alpha = \frac{Gxv_g}{C_0 \{Gxv_g + G(1-x)v_f\} + V_{gj}}$$

$$\left(\frac{\partial \alpha}{\partial x} \right)_P = \frac{\alpha}{x} - \frac{\alpha^2 C_0 v_{fg}}{xv_g} = 8.22$$

$$v^* = \left\{ \frac{2xv_g}{\alpha} - \frac{2(1-x)v_f}{1-\alpha} \right\} + \left(\frac{\partial \alpha}{\partial x} \right)_P \left\{ \frac{(1-x)^2 v_f}{(1-\alpha)^2} - \frac{x^2 v_g}{\alpha^2} \right\} = 0.0132 \text{ m}^3/\text{kg}$$

$$-\left(\frac{dP}{dz} \right)_a = \frac{1}{(1-M^2)} \left\{ G^2 \frac{dx}{dz} v^* \right\} = 0.1317 \text{ kPa/m}$$

$$\alpha = \frac{j_g}{C_0 j + V_{gj}}$$

$$\alpha = \frac{Gxv_g}{C_0 \{Gxv_g + G(1-x)v_f\} + V_{gj}}$$

$$\frac{\partial \alpha}{\partial x} = \frac{\alpha}{x} - \frac{\alpha^2 C_0 v_{fg}}{xv_g} = 8.22$$

$$v^* = \left\{ \frac{2xv_g}{\alpha} - \frac{2(1-x)v_f}{1-\alpha} \right\} + \left(\frac{\partial \alpha}{\partial x} \right)_P \left\{ \frac{(1-x)^2 v_f}{(1-\alpha)^2} - \frac{x^2 v_g}{\alpha^2} \right\} = 0.0132 \text{ m}^3/\text{s}$$

$$-\left(\frac{dP}{dz} \right)_a = \frac{1}{(1-M^2)} \left\{ G^2 \frac{dx}{dz} v^* \right\} = 0.1317 \text{ kPa/m}$$

And now we get this the slope of the alpha x square 8.22 and use the slope to calculate v star which is 0.0132 meter cube per kg and the acceleration pressure gradient is calculated as 0.1317 kilo Pascal per meter.

(Refer Slide Time: 34:11)

$$\begin{aligned}
 -\left(\frac{dP}{dz}\right)_F &= \underline{0.572 \text{ kPa/m}} \\
 -\left(\frac{dP}{dz}\right)_a &= \underline{0.1317 \text{ kPa/m}} \\
 -\left(\frac{dP}{dz}\right)_z &= [\underline{\rho_g \alpha + \rho_f(1 - \alpha)}]g \sin \theta = \underline{6.18 \text{ kPa/m}} \\
 -\left(\frac{dP}{dz}\right) &= 0.572 + 0.1317 + 6.18 = \underline{\underline{6.88 \text{ kPa/m}}}
 \end{aligned}$$

$$-\left(\frac{dP}{dz}\right)_F = 0.572 \text{ kPa/m}$$

$$-\left(\frac{dP}{dz}\right)_a = 0.1317 \text{ kPa/m}$$

$$-\left(\frac{dP}{dz}\right)_z = [\rho_g \alpha + \rho_f(1 - \alpha)]g \sin \theta = 6.18 \text{ kPa/m}$$

$$-\left(\frac{dP}{dz}\right) = 0.572 + 0.1317 + 6.18 = 6.88 \text{ kPa/m}$$

So, the frictional pressure gradient is 0.572 kilo Pascal per meter, acceleration pressure gradient is 0.1317 kilo Pascal per meter and the gravitational pressure gradient is calculated using this expression and it is 6.18 kilo Pascal per meter. Now we add all these pressure gradients and get the total pressure gradient as 6.88 kilopascal per meter.

So, now let us summarize the results of all the numerical examples.

(Refer Slide Time: 34:57)

Pressure Gradients (kPa) obtained for Example 1 using Different Models

	Homogeneous Model ✓	Lockhart-Martinelli Correlation ✓	Martinelli-Nelson Correlation ✓	Drift Flux Model ✓
$-\left(\frac{dP}{dz}\right)_F$ ✓	2.74 ✓	2.05 ✓	3.06 ✓	2.74 ✓
$-\left(\frac{dP}{dz}\right)_a$ ✓	7.50 ✓	1.045 ✓	0.743 ✓	0.714 ✓
$-\left(\frac{dP}{dz}\right)_z$ ✓	0 ✓	0 ✓	0 ✓	0 ✓
$-\left(\frac{dP}{dz}\right)$ ✓	10.24 ✓	3.10 ✓	3.80 ✓	3.45 ✓

Pressure Gradients (kPa) obtained for Example 1 using Different Models

	Homogeneous Model	Lockhart-Martinelli Correlation	Martinelli - Nelson Correlation	Drift Flux Model
$-\left(\frac{dP}{dz}\right)_F$	2.74	2.05	3.06	2.74
$-\left(\frac{dP}{dz}\right)_a$	7.5	1.045	0.743	0.714
$-\left(\frac{dP}{dz}\right)_z$	0	0	0	0
$-\left(\frac{dP}{dz}\right)$	10.24	3.10	3.80	3.45

Here the results of example 1 are shown in tabular form. The frictional pressure gradient, acceleration pressure gradient, gravitational pressure gradient and the total pressure

gradient using homogeneous model, Lockhart-Martinelli correlation, Martinelli-Nelson correlation and the Drift Flux model. The frictional pressure the gradient from homogeneous model is 2.74 kilo Pascal per meter, ok. It is 2.74 kilo Pascal per meter and you from Lockhart-Martinelli it is 2.05 kilo Pascal per meter, Martinelli Nelson give 3.06 kilo Pascal per meter, the Drift Flux model and that we use the same expression as that for homogeneous model.

So, it is the same 2.74 kilo Pascal per meter. The acceleration pressure gradient from homogeneous model is 7.50 kilo Pascal per meter and Lockhart Martinelli give 1.045 Martinelli Nelson give 0.743 and Drift Flux model give 0.714. Here the homogeneous model give very high value of the acceleration pressure gradient that is because alpha was highly overestimated using homogeneous model.

The gravitational pressure gradient is 0 for all the model because it is the horizontal pipe and the total pressure gradients are 10.24 3.10 3.80 3.45. The total pressure gradient using homogeneous model is high mainly because of the pressure gradient due to acceleration.

(Refer Slide Time: 37:25)

Pressure Gradients (kPa) obtained for Example 2 using Different Models

	Homogeneous Model	Lockhart-Martinelli Correlation	Martinelli-Nelson Correlation	Drift Flux Model
$-\left(\frac{dP}{dz}\right)_f$ ✓	9.85 ✓	5.46 ✓	0.1322 ✓	9.85 ✓
$-\left(\frac{dP}{dz}\right)_a$	20.1 ✓	2.92 ✓	1.167 ✓	2.56 ✓
$-\left(\frac{dP}{dz}\right)_z$ ✓	0.646 ✓	2.99 ✓	3.105 ✓	1.648 ✓
$-\left(\frac{dP}{dz}\right)$	30.6 ✓	11.37 ✓	4.40 ✓	14.06 ✓

Pressure Gradients (kPa) obtained for Example 2 using Different Models

	Homogeneous Model	Lockhart-Martinelli Correlation	Martinelli - Nelson Correlation	Drift Flux Model
$-\left(\frac{dP}{dz}\right)_F$	9.85	5.46	0.1322	9.85
$-\left(\frac{dP}{dz}\right)_a$	20.1	2.92	1.167	2.56
$-\left(\frac{dP}{dz}\right)_z$	0.646	2.99	3.105	1.648
$-\left(\frac{dP}{dz}\right)$	30.6	11.37	4.40	14.06

In the second example here it is kilo Pascal per meter and homogeneous model give the frictional pressure gradient as 9.85 kilo Pascal per meter Lockhart-Martinelli give 5.46, Martinelli-Nelson give a very low value 0.1322, Drift Flux same as homogenous 9.85. The acceleration pressure gradient from homogeneous model of 20.1 kilo Pascal per meter, Lockhart-Martinelli 2.92, Martinelli-Nelson 1.167 and Drift Flux same as Drift Flux 2.56.

The gravitational pressure gradient for homogeneous it is 0.646 kilo Pascal per meter Lockhart-Martinelli to 2.99, Martinelli=Nelson 3.105 and Drift Flux 1.648. The total pressure gradient from homogeneous model is 30.6 kilo Pascal per meter, Lockhart-Martinelli 11.37, Martinelli-Nelson 4.40 and Drift Flux 14.06. Here also we see that the acceleration pressure gradient from homogeneous model is very high compared to the

other models and the frictional pressure gradient given by the Martinelli-Nelson correlation as very low compared to the other models.

(Refer Slide Time: 39:29)

Pressure Gradients (kPa) obtained for Example 3 using Different Models

	Homogeneous Model	Lockhart-Martinelli Correlation	Martinelli-Nelson Correlation	Drift Flux Model
$-\left(\frac{dP}{dz}\right)_F$	0.572	1.047	0.697	0.572
$-\left(\frac{dP}{dz}\right)_a$	0.166	7.56	0.035	0.1317
$-\left(\frac{dP}{dz}\right)_z$	6.05	4.83	6.12	6.18
$-\left(\frac{dP}{dz}\right)$	6.79	13.44	6.85	6.88

Pressure Gradients (kPa) obtained for Example 3 using Different Models

	Homogeneous Model	Lockhart-Martinelli Correlation	Martinelli - Nelson Correlation	Drift Flux Model
$-\left(\frac{dP}{dz}\right)_F$	0.572	1.047	0.697	0.572
$-\left(\frac{dP}{dz}\right)_a$	0.166	7.56	0.035	0.1317
$-\left(\frac{dP}{dz}\right)_z$	6.05	4.83	6.12	6.18
$-\left(\frac{dP}{dz}\right)$	6.79	13.44	6.85	6.88

The results of the example 3, frictional pressure gradient, homogenous model give 0.572 kilo Pascal per meter, Lockhart-Martinelli 1.047, Martinelli-Nelson 0.697 and Drift Flux 0.572. Pressure gradient due to acceleration, homogenous 0.166, Lockhart-Martinelli 7.56, Martinelli-Nelson 0.035 and Drift Flux 0.3117; pressure gradient due to gravity homogenous 6.05, Lockhart-Martinelli 4.83, Martinelli-Nelson 6.12 and Drift Flux 6.18.

The total pressure gradient from homogenous model is 6.79 kilo Pascal per meter, Lockhart-Martinelli 13.44, Martinelli-Nelson 6.85 and Drift Flux 6.88. In this example we see that Lockhart-Martinelli model gives very high value of the acceleration pressure gradient, where is Martinelli-Nelson correlation gives a very low value of acceleration pressure gradient.

As I mention before these numerical examples are only to illustrate the methods how to use the models, how to calculate the pressure gradients using different models and these are not meant to be used for design purposes and there may be errors in the calculations, because calculations are very complicated. So, you may do your own calculations and check.

Thank you.