

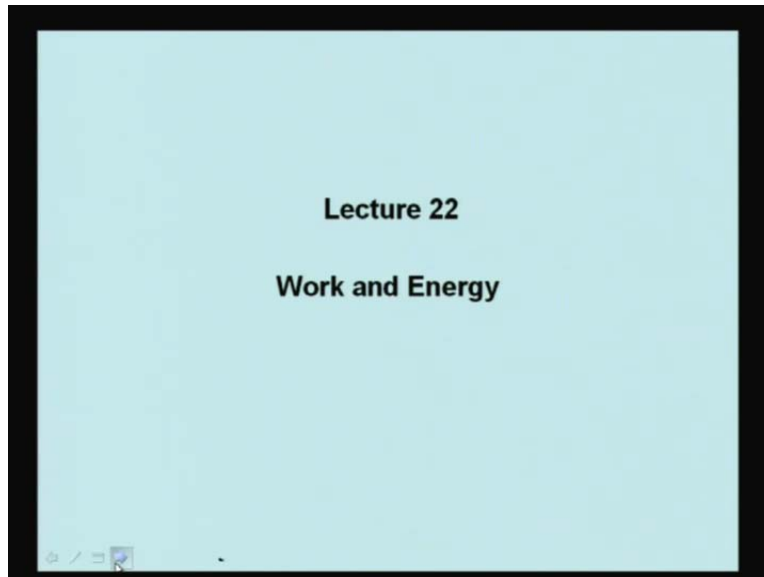
**Engineering mechanics**  
**Prof. U S Dixit**  
**Department of Mechanical Engineering**  
**Indian Institute of Technology, Guwahati**  
**Kinetics-1**

**Module 11 Lecture - 28**  
**Work and Energy**

We have studied dynamics of a particle. For solving the dynamics problem we used Newton's law. Particularly in solving numerical problems, Newton's second law was used. In place of Newton's second law, we can also use D'Alembert's principle. D'Alembert's principle is not much different from Newton's law. It basically considers minus  $m$  into  $a$ , equal to a force that is called inertia force. Thus, a problem of dynamics can be converted into a problem of a statics, but essentially, both the approaches are making use of the equation  $F$  is equal to mass times acceleration, where  $F$  is the net resultant force acting on the body and  $m$  is the mass,  $a$  is the acceleration.

Here, one point has to be noted that the acceleration has to be measured from an inertial frame of reference. An inertial frame of reference will be moving with a constant velocity including zero. It may be moving with a zero velocity with respect to a primary inertial system. A primary inertial frame of reference is that frame of reference which is at rest. Thus, we see the D'Alembert's principle and Newton's second law essentially are the same. Today, we discuss a different approach of solving the dynamics problem that is work and energy.

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You will see that this formulation is quite different although it can be derived from Newton's second law. In one kind of energy, we will not talk about the same type of things as we talk in Newton's second law. In Newton's second law, one requires acceleration. Here, another approach maybe there. When you require only the measurement of velocities and force and distance, then also the same type of solution can come.

In case, if you have to just find out the velocity after the application of force at some distance then instead of going by Newton's second law, we can directly use work and energy. The Newton's second law provides a differential equation; that is  $F$  is equal to  $m \frac{d^2 r}{dt^2}$ . Here, you will get a type of integral equation which can be solved in a different manner. Although, Newton's second law and work and energy approach follow different procedures, they are basically derivable from each other. I will explain the basic concepts.

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Newton's law for a particle moving relative to an inertial reference is given by,

$$F = m \frac{d^2 r}{dt^2} = m \frac{dv}{dt}$$

Multiplying each side of the equation by  $dr$  as a dot product and integrating from  $r_1$  to  $r_2$  along the path of motion:

$$\int_{r_1}^{r_2} F \cdot dr = m \int_{r_1}^{r_2} \frac{dv}{dt} \cdot dr = m \int_{r_1}^{r_2} \frac{dv}{dt} \frac{dr}{dt} dt = m \int_{t_1}^{t_2} \left( \frac{dv}{dt} \cdot v \right) dt$$

$$\frac{d}{dt} (\vec{v} \cdot \vec{v}) = \vec{v} \cdot \frac{d\vec{v}}{dt} + \vec{v} \cdot \frac{d\vec{v}}{dt} = 2 \vec{v} \cdot \frac{d\vec{v}}{dt} = 2 \frac{d}{dt} (v \cdot v)$$

Newton's law for a particle moving relative to an inertial reference frame is given by  $F$  is equal to  $m \frac{d^2 r}{dt^2}$ , which is equal to  $m \frac{dv}{dt}$ , where  $r$  is a displacement vector and  $v$  is the velocity vector. Here, I am considering the mass of the particle to be constant. If we multiply each side of the equation by a vector  $dr$  as a dot product and integrate from  $r_1$  to  $r_2$  along the path of motion, the particle may be moving from  $r_1$  to  $r_2$ .

This is 1/2. Position vector of 1 is  $r_1$ ; position vector of 2 is  $r_2$ . We can integrate. If dot product of  $F$  and  $dr$  from  $r_1$  to  $r_2$ , which gives us  $\int_{r_1}^{r_2} F \cdot dr$  is equal to  $m \int_{r_1}^{r_2} \frac{dv}{dt} \cdot dr$ . I have taken  $m$  outside the integral sign, considering mass to be constant. Therefore, this has become  $m \int_{r_1}^{r_2} \frac{dv}{dt} \cdot \frac{dr}{dt} dt$  is equal to  $m \int_{t_1}^{t_2} \frac{dv}{dt} \cdot v dt$ , where  $dr/dt$  has been written as  $v$ .

This is equal to half  $m \int_{t_1}^{t_2} \frac{d}{dt} (v \cdot v) dt$ , where the product rule of vector calculus has been used; that is  $\frac{d}{dt} (v \cdot v)$  is equal to  $v \cdot \frac{dv}{dt} + v \cdot \frac{dv}{dt}$ , which gives you half  $m \int_{t_1}^{t_2} \frac{d}{dt} (v \cdot v) dt$ .

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$$= \frac{1}{2} m \int_{t_1}^{t_2} \frac{d}{dt} v^2 dt = \frac{1}{2} m \int_{v_1}^{v_2} d(v^2) = \frac{1}{2} m (v_2^2 - v_1^2)$$

**Thus, the work done on the particle is equal to change in its kinetic energy**

**If we write Newton's law in component form, then,**

$$F_x \hat{i} + F_y \hat{j} + F_z \hat{k} = m \frac{dv_x}{dt} \hat{i} + m \frac{dv_y}{dt} \hat{j} + m \frac{dv_z}{dt} \hat{k}$$

or

$$F_x \hat{i} = m \frac{dv_x}{dt} \hat{i}$$

This expression becomes half  $m \int_{t_1}^{t_2} \frac{d}{dt} v^2 dt$ . See the dot product  $\mathbf{v} \cdot \mathbf{v}$  is a scalar that is  $v^2$  which is the magnitude squared of velocity. This becomes half  $m \int_{t_1}^{t_2} \frac{d}{dt} v^2 dt$  is equal to half  $m \int_{v_1}^{v_2} d(v^2)$ . That gives, here I have changed the limits of the integration, because the variable  $dt$  has vanished from here. This becomes half  $m (v_2^2 - v_1^2)$ . Thus, the work done on the particle is equal to change in its kinetic energy.

Therefore, this essentially has come from Newton's second law. Therefore, instead of Newton's second law, one can use this theorem that work done on the particle is equal to change in its kinetic energy. If you know the force acting on the particle and you know the displacement vector, then you can find out the work done. You can equate it to change in the kinetic energy. Thus, you can find out the final velocity, if you know the initial velocity.

We have to see that is it possible for us to find out, if instead of magnitude, can we also find out the components, well just components of the velocity. If we write Newton's law in component form, then  $F_x \hat{i} + F_y \hat{j} + F_z \hat{k}$  is equal to  $m \frac{dv_x}{dt} \hat{i} + m \frac{dv_y}{dt} \hat{j} + m \frac{dv_z}{dt} \hat{k}$ . This Newton's law has been written in the component form.  $F_x \hat{i}$  is therefore,  $F_x \hat{i}$  is equal to  $m \frac{dv_x}{dt} \hat{i}$ .

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$$\hat{F}_x \hat{i} = m \frac{d\hat{v}_x \hat{i}}{dt}$$

**Taking the dot product of this equation with**

$$dx \hat{i} + dy \hat{j} + dz \hat{k} (=d\mathbf{r})$$

$$\int_{x_1}^{x_2} F_x dx = \frac{m}{2} [(V_x)_2^2 - (V_x)_1^2]$$

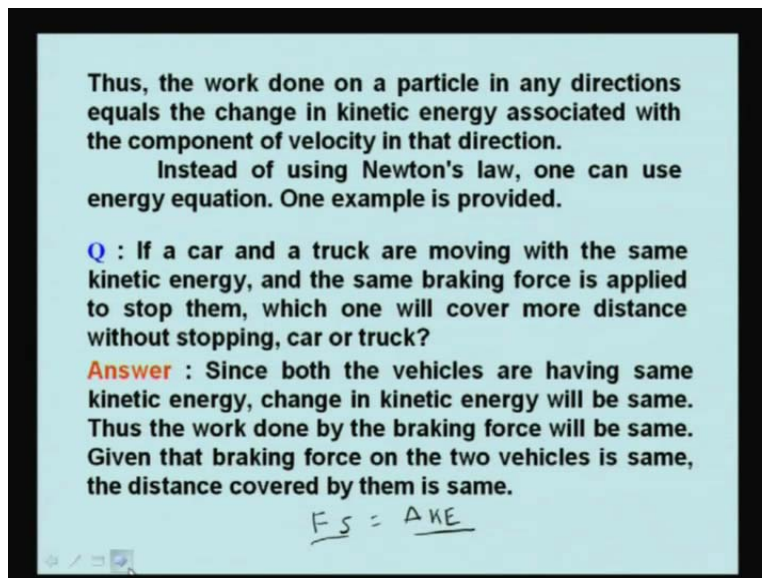
**Similarly**

$$\int_{y_1}^{y_2} F_y dy = \frac{m}{2} [(V_y)_2^2 - (V_y)_1^2]$$

$$\int_{z_1}^{z_2} F_z dz = \frac{m}{2} [(V_z)_2^2 - (V_z)_1^2]$$

Thus, if we take the dot product of this equation with  $dx \hat{i} + dy \hat{j} + dz \hat{k}$  is equal to  $d\mathbf{r}$ , if we take the dot product of equation not the dot product of  $\mathbf{F}$  with  $d\mathbf{r}$ , but dot product of  $F_x \hat{i}$  is equal to  $m dv_x$  by  $dt$ . I with this  $d\mathbf{r}$ , then we will be getting that  $x_1$  to  $x_2$ ,  $F_x dx$  will be equal to  $m dv_x$  by  $dt$  into  $dx$ , which, when it will be integrated, it will provide us  $m$  by  $\frac{1}{2} V_{x2}^2$  square minus  $V_{x1}^2$  square. Similarly, we can also get  $y_1$  to  $y_2$ ,  $F_y dy$  is equal to  $m$  by  $\frac{1}{2} V_{y2}^2$  square minus  $V_{y1}^2$  square. Similarly,  $z_1$  to  $z_2$ ,  $F_z dz$  is equal to  $m$  by  $\frac{1}{2} V_{z2}^2$  square minus  $V_{z1}^2$  square. Thus, the Newton's law is also valid in the component form.

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Thus, the work done on a particle in any directions equals the change in kinetic energy associated with the component of velocity in that direction.

Instead of using Newton's law, one can use energy equation. One example is provided.

**Q :** If a car and a truck are moving with the same kinetic energy, and the same braking force is applied to stop them, which one will cover more distance without stopping, car or truck?

**Answer :** Since both the vehicles are having same kinetic energy, change in kinetic energy will be same. Thus the work done by the braking force will be same. Given that braking force on the two vehicles is same, the distance covered by them is same.

$F s = \Delta KE$

Thus, the work done on a particle in any directions equals the change in kinetic energy associated with the component of velocity in that direction. This is very important because we can find out the velocity components by using the work and kinetic energy equation. One example is provided. If a car and a truck are moving with the same kinetic energy and the same breaking force is applied to stop them, which one will cover more distance without stopping, car or truck?

In this case, if we want to apply Newton's law, we have to do many arithmetic computations. We have to find out the acceleration of the truck and cars. Even after that, we have to apply equation of motion in order to find out the result, but in this case, we can quickly find out the answer by applying the work kinetic energy theorem. Since, both the vehicles are having same kinetic energy, change in kinetic energy will be same. Thus, the work done by the breaking force will be same. Given that breaking force on the two vehicles is same, the distance covered by them is same.

Here, you get a simple relation that  $F$  is the breaking force.  $F$  into  $s$  is equal to change in kinetic energy  $\Delta KE$ . Since change in the kinetic energy is same, breaking force comes out to be same. This one, this shows that how efficiently we apply to work energy theorem in many cases.

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**Power :** Power is the time rate of doing work.

Accordingly, the power  $P$  developed by a force  $F$  which does an amount of work  $W$  is

$$P = \frac{dW}{dt} = \frac{d(F \cdot dr)}{dt} = F \cdot \frac{dr}{dt} + \frac{dF}{dt} \cdot dr$$

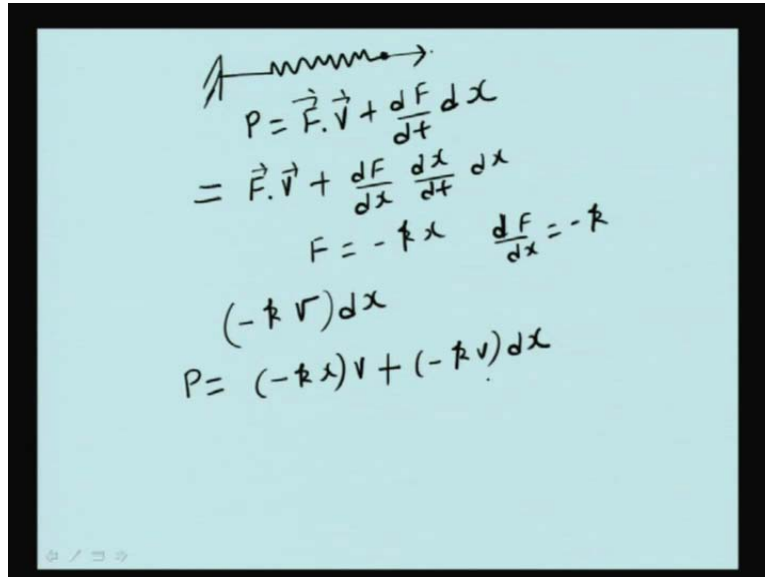
$$P = \frac{dW}{dt} = F \cdot \frac{dr}{dt} = F \cdot v = F \cdot v + \frac{dF}{dt} \cdot dr$$

Suppose a man rises walking on the top of a hill and another man goes riding on a motor bike. Both have done the same amount of work, but the second man has used more power, because he would have done the same amount of work faster.

We will discuss something more about these things, but let us go to another definition. We define power is the time rate of doing work. Accordingly, the power  $P$  developed by a force  $F$  which does an amount of work  $W$ , is  $P$  is equal to  $dW$  by  $dt$ ; that means, this is equal to  $F \cdot dr$  by  $dt$  that is equal to  $F \cdot v$ . Here, in this case, we have obtained this expression in the following manner;  $P$  is equal to  $dW$  by  $dt$ .  $dW$  by  $dt$  is  $dF \cdot dr$  by  $dt$ , we take  $F$  as a constant, which does not depend on time. Take it out and then you get  $F \cdot dr$  by  $dt$ .

In case  $F$  is dependent on time, then this can be written as  $F \cdot dr$  by  $dt$  plus  $dF$  by  $dt$  dot  $dr$ . In this expression, we get additional term.  $F \cdot dr$  by  $dt$  surely can be written as  $F \cdot v$ . However, we will get extra term; that is  $dF$  by  $dt$  dot  $dr$ . If  $F$  is changing with time, then  $dF$  by  $dt$  must also be evaluated and multiplied with  $dr$ . However, if we consider the  $dr$  to be very small, then in that case, this term becomes insignificant.

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The image shows a handwritten derivation of power for a spring force. At the top, there is a diagram of a spring attached to a wall on the left and a mass on the right, with an arrow pointing to the right indicating displacement. Below the diagram, the equations are written as follows:

$$P = \vec{F} \cdot \vec{v} + \frac{dF}{dt} dx$$

$$= \vec{F} \cdot \vec{v} + \frac{dF}{dx} \frac{dx}{dt} dx$$

$$F = -kx \quad \frac{dF}{dx} = -k$$

$$(-k v) dx$$

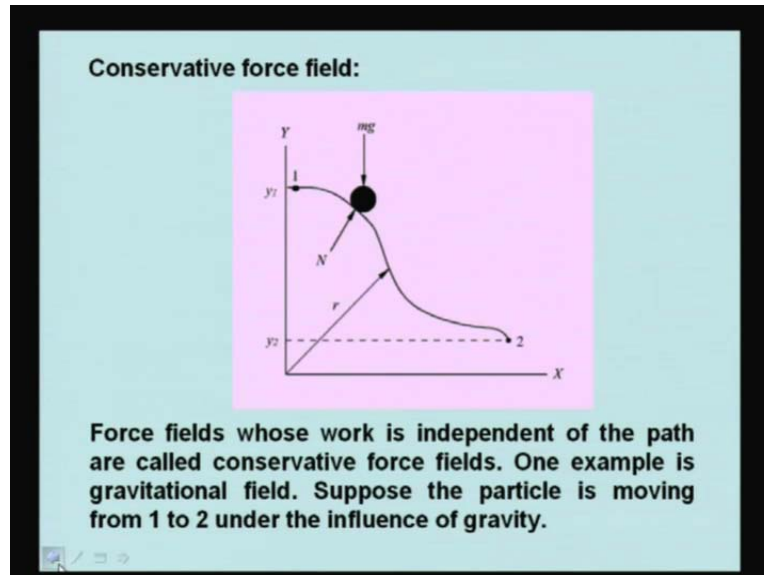
$$P = (-kx)v + (-kv) dx$$

We give one example, where the force is changing with time or distance. For example, consider a spring. If particle attached with spring is moved, then the spring force is dependent on the extension of the spring. In this case,  $P$  will be equal to  $F$  times dot  $V$  plus  $dF$  by  $dt$  dot  $dr$ , which for one dimensional case, we can write as  $dx$ . This can be written as  $F$  dot  $V$  plus  $dF$  by  $dt$ ,  $dF$  by  $dt$  can be written as  $dF$  by  $dx$ ; this can be written as  $dx$  by  $dt$   $dx$ , in which  $dF$ ,  $F$  is equal to minus  $kx$ . Therefore,  $dF$  by  $dx$  is equal to minus  $k$ . Therefore, the second term can be written as minus  $k$   $dx$  by  $dt$ ; that is minus  $kv$   $dx$ .

Therefore, the expression for power is, in this case minus force into velocity minus  $kx$  into  $V$  plus minus  $kV$  into  $dx$ ,  $dx$  is infinite small. In that case, the power will be minus  $kx$   $V$ . However, if force has some singularity, in that case, this term maybe significant. Concept of power is found very useful in certain context. Suppose a man rises walking on the top of a hill and another man goes riding on a motor bike, both have done the same amount of work but the second man used more power because he would have done the same amount of work faster.



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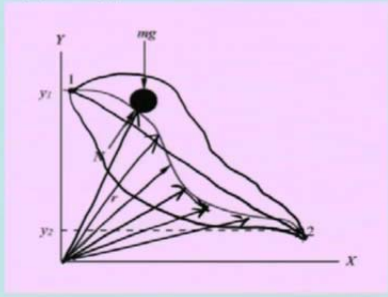
We discuss about the conservative force field. Force fields whose work is independent of the path are called conservative force fields. Force field, we understand the force field, but what is force field? Force can be expressed as a function of  $x$ ,  $y$  and  $z$  and time also. Therefore, when we express force as a function of  $x$ ,  $y$ ,  $z$ , then it becomes a vector field. Force is a vector; this is a vector field. Force field may, in this there can be different type of force fields. Fields, in which work is independent of the path are called conservative force fields. One example is a gravitational field.

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Work done by the force is  $W = \int_1^2 \mathbf{F} \cdot d\mathbf{r} = \int_1^2 (-mg\hat{j}) \cdot (dx\hat{i} + dy\hat{j} + dz\hat{k})$

$$= -mg \int_1^2 dy = -mg(y_2 - y_1) = mg(y_1 - y_2)$$

Thus, the work done is dependent on the end coordinates  $y_1$  and  $y_2$ .



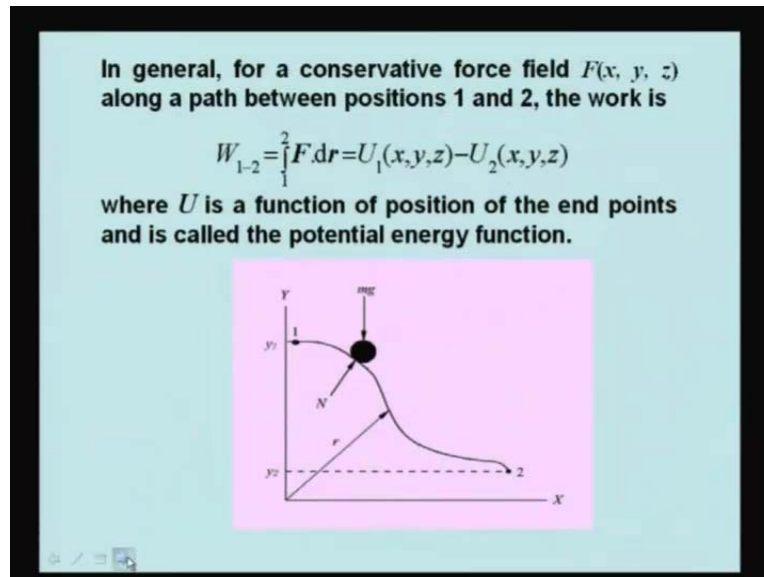
Suppose, the particle is moving from 1 to 2 under the influence of gravity like this, in that case work done by the force is  $W$  is equal to  $\int_1^2 \mathbf{F} \cdot d\mathbf{r}$ ,  $\mathbf{r}$  is the position vector of particle at any instant.  $\mathbf{r}$  keeps on changing as the particle is moving towards paths, like that.  $d\mathbf{r}$  is the distance between two position vectors like at this instance, this is  $d\mathbf{r}$ . Provided the difference is very small it is indicated by  $d\mathbf{r}$ . Work done by the force  $W$  is equal to  $\int_1^2 \mathbf{F} \cdot d\mathbf{r}$  which is  $\int_1^2 -mg\hat{j} \cdot d\mathbf{r}$ , because the force is always constant and that is directed downwards in the  $y$  direction.

Therefore, it is  $-mg\hat{j} \cdot d\mathbf{r}$ ,  $d\mathbf{r}$  can be written as  $dx\hat{i} + dy\hat{j} + dz\hat{k}$ . Dot product of  $\hat{j}$  and  $\hat{i}$  is 0, dot product of  $\hat{i}$  and  $\hat{k}$  is 0; therefore, we are left with  $-mg dy$  because  $\hat{j} \cdot \hat{j}$  is one which gives us  $-mg \int_{y_1}^{y_2} dy$  that is equal to  $-mg(y_2 - y_1)$ . Thus, the work done is dependent on the end coordinates  $y_1$  and  $y_2$  only. It is not dependent on the path of motion. The particle can adapt any path. It could have gone like this to the same point, it could have gone like this or it could have gone in a straight line.

The work done will be same. It is dependent on the end coordinates  $y_1$  and  $y_2$ . Therefore, gravitational field will be called a conservative force field. Similarly, if the particle goes from 1 to 2, the work done is  $-mg(y_2 - y_1)$ . Work done by the gravity on the particle is  $-mg(y_2 - y_1)$ . If the particle would have gone from 2 to 1, the work done by the gravity on the particle will be  $-mg(y_1 - y_2)$ ; add both the works, you get 0.

In the conservative force field, if a particle undergoes a cyclic motion like this, completes the cycle and comes back to the original position, the network is 0, which is not the case in non conservative force field.

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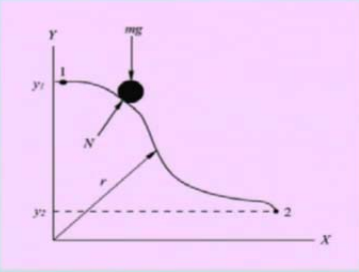


So, in general, for a conservative force field  $F(x, y, z)$  along a path between positions 1 and 2, the work is  $W_{1-2} = \int_1^2 F \cdot dr$ ; that is equal to  $U_1(x, y, z) - U_2(x, y, z)$ , where  $U$  is a function of position of the end points and is called the potential energy function. In the previous case, you got work done between 1 and 2. It is  $mg y_1 - mg y_2$ . So,  $mg y_1$  can be written as the potential energy  $U_1$ ;  $mg y_2$  is the potential energy  $U_2$ .  $U_1$  and  $U_2$  are functions of  $x, y, z$ .  $U_1$  is equal to  $mg y_1$  is a function of  $x, y, z$  in which other coefficients are 0.  $U$  is a function of position of the end points and is called the potential energy function.

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Noting that  $-\int_1^2 \mathbf{F} \cdot d\mathbf{r} = U_2(x, y, z) - U_1(x, y, z) = \Delta U$

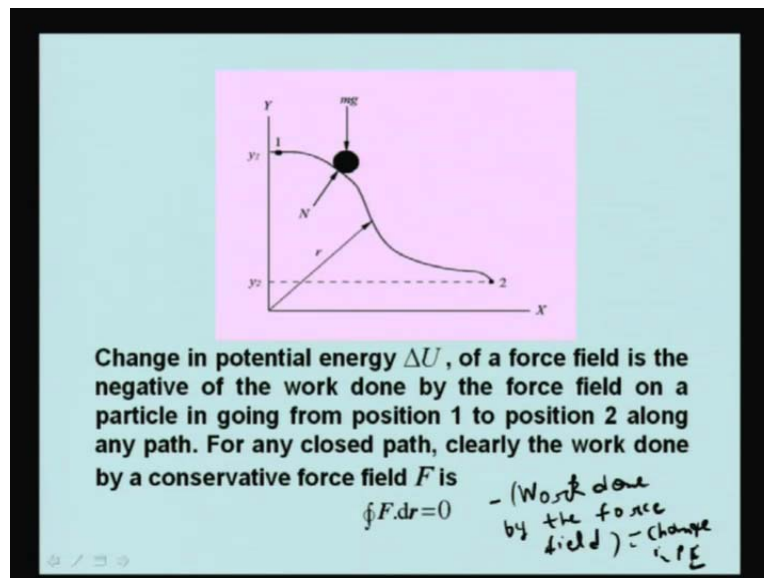
Note that the potential energy function  $U$  depends on the reference  $xyz$  used or the datum used. However, the change in potential energy  $\Delta U$  is independent of the datum used.



The diagram shows a particle of mass  $m$  moving along a curved path in a 2D coordinate system with axes  $X$  and  $Y$ . The particle is at point 1, which is at a height  $y_1$  on the  $Y$ -axis. A downward arrow labeled  $mg$  represents the gravitational force. A normal force vector  $N$  points from the particle perpendicular to the path. A position vector  $r$  points from the origin to the particle. The path ends at point 2, which is at a height  $y_2$ . A dashed horizontal line connects point 2 to the  $Y$ -axis at  $y_2$ .

We note that minus 1 to 2  $\mathbf{F} \cdot d\mathbf{r}$  is equal to  $U_2, x, y, z$  minus  $U_1, x, y, z$  which is  $\Delta U$ . Note that potential energy function depends on the reference  $xyz$  used or the datum used. However, we will be concerned only with the change in potential energy  $\Delta U$ , because work done between two points will always be the difference of the two potential energies. Therefore, it is independent of the datum used. Any datum can be used for defining the zero potential energy. In kinetic energy calculation, the zero velocity uses zero kinetic energy, but this is not the case with the potential energy. In potential energy, you can take any datum and you can find out  $\Delta U$  which will be independent of the datum used.

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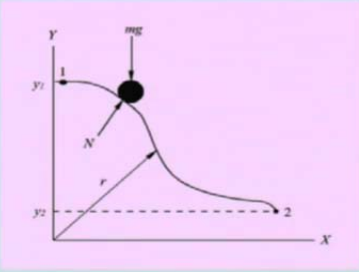


Change in potential energy  $\Delta U$  of a force field is the negative of the work done by the force field or a particle in going from position 1 to position 2 along any path. Therefore, force field, if when the force field does some positive work then the potential energy increases. If the force field does the negative work then the potential energy decreases. Change in potential energy  $\Delta U$ , negative of the work done by the force field is equal to change in potential energy. For any closed path theory, the work done by a conservative force field  $F$  is  $\oint F \cdot dr$  equal to 0.

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For an infinitesimal path difference  $dr$  starting from 1,

$$F \cdot dr = -dU$$

$$F_x dx + F_y dy + F_z dz = -\left(\frac{\partial U}{\partial x} dx + \frac{\partial U}{\partial y} dy + \frac{\partial U}{\partial z} dz\right)$$


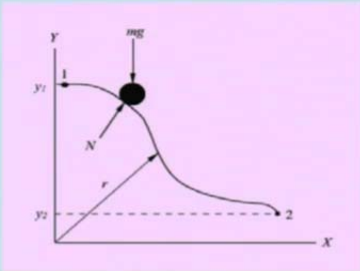
For an infinitesimal path difference  $dr$  starting from 1,  $F \cdot dr$  is equal to minus  $dU$ ; that is  $F_x dx$  plus  $F_y dy$  plus  $F_z dz$  is equal to minus  $\Delta U$  by  $\Delta x$  plus  $\Delta U$  by  $\Delta y$  plus  $\Delta U$  by  $\Delta z$ .

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Thus

$$F_x = -\frac{\partial U}{\partial x}$$

$$F_y = -\frac{\partial U}{\partial y}$$

$$F_z = -\frac{\partial U}{\partial z}$$


or  $F = -\text{grad } U = -\nabla U$  where  $\nabla$  the operator is called the gradient operator and is given as follows for rectangular coordinates.

$$\text{grad} = \nabla = \left(\frac{\partial}{\partial x} \hat{i} + \frac{\partial}{\partial y} \hat{j} + \frac{\partial}{\partial z} \hat{k}\right)$$

Thus  $F_x$  is equal to minus  $\Delta U$  by  $\Delta x$ ,  $F_y$  is equal to minus  $\Delta U$  by  $\Delta y$ ,  $F_z$  is equal to minus  $\Delta U$  by  $\Delta z$ , or  $F$  is equal to minus gradient of  $U$ , where  $U$  is a scalar field and this is indicated

by minus delta U. Gradient is defined as, this is a gradient operator which is defined by del by del x i plus del y by del y j plus del by del z k. This is in the Cartesian coordinate system. In the same way, it can be defined in the polar coordinate or spherical coordinate system. If you have got a conservative force field, it can be expressed as a gradient of some potential energy function.

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Thus, a conservative force field must be a function of position and expressed as the gradient of a scalar function.

If a force field is a function of position and the gradient of a scalar field, it must then be a conservative force field.

$\text{curl } \vec{F} = 0$

$\vec{\nabla} \times \vec{F}$

$$\begin{vmatrix} \hat{i} & \hat{j} & \hat{k} \\ \frac{\partial}{\partial x} & \frac{\partial}{\partial y} & \frac{\partial}{\partial z} \\ \frac{\partial U}{\partial x} & \frac{\partial U}{\partial y} & \frac{\partial U}{\partial z} \end{vmatrix} = \hat{i} \left( \frac{\partial^2 U}{\partial y \partial z} - \frac{\partial^2 U}{\partial z \partial y} \right) - \hat{j} \left( \frac{\partial^2 U}{\partial x \partial z} - \frac{\partial^2 U}{\partial z \partial x} \right) + \hat{k} \left( \frac{\partial^2 U}{\partial x \partial y} - \frac{\partial^2 U}{\partial y \partial x} \right) = 0$$

Thus, a conservative force field must be a function of position and expressed as the gradient of a scalar function. If a force field is a function of position and the gradient of a scalar field, it must then be a conservative force field. In some books, it will be shown that for a conservative force field curl F is equal to 0; that is curl of the force is equal to 0 for a conservative force field. Curl F of vector is basically del cross F, where the del operator has been defined here. This can be written as i j k del by del x del by del y del by del z minus del u by del x minus del u by del y minus del u by del z, because a conservative force has got these components.

Therefore, this becomes i minus del square U del y del z plus del square U del z by del y; but for continuous U this term will become 0. Similarly, the other terms will also come out to be 0. minus j minus del square U del x del z plus del square U del x, del z plus k minus del square U del x del y plus del square U del x del y is equal to 0.

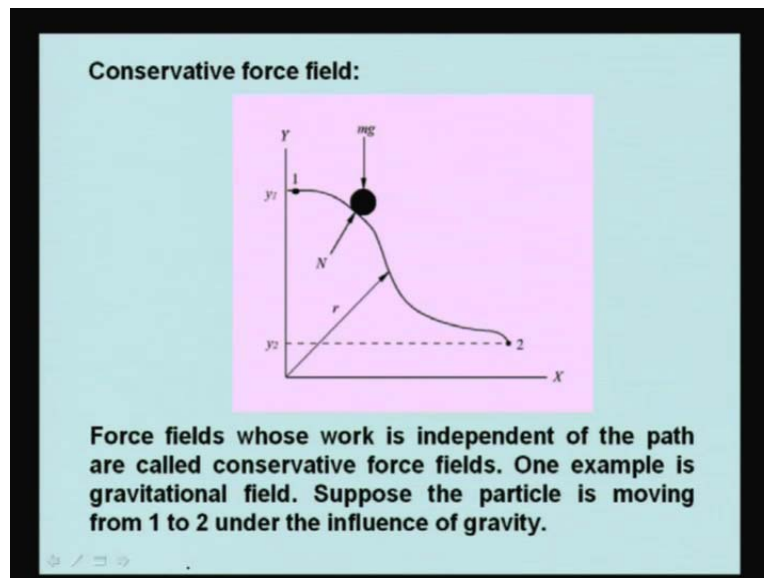
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$$\begin{aligned} \text{curl } \vec{F} \\ F_x = \frac{y}{x^2+y^2}, F_y = \frac{-x}{x^2+y^2}, F_z = 0 \\ \text{curl } \vec{F} = 0 \end{aligned}$$

Thus, you can judge whether a force field is conservative or not by finding out the curl of the force. Although if this is too far, if this function which is not having the singularity, in certain cases, like if you have  $F_x$  is equal to  $y$  divided by  $x$  square plus  $y$  square  $F_y$  is equal to minus  $x$  divided by  $x$  square plus  $y$  square  $F_z$  is equal to 0. This force field provides  $\text{curl } F$  equal to 0. However, this force field is not conservative. Thus,  $\text{curl } F$  is equal to 0 can be called a necessary condition but not a sufficient condition; sufficient condition is from the basic definition.



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A force field whose work is independent of the path is called conservative force field. What are the examples of the conservative force fields?

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**Linear force :** If the force is given by  $F = ax\hat{i} + by\hat{j} + cz\hat{k}$  then it can be expressed as the gradient of

$$U = \frac{ax^2}{2} + \frac{by^2}{2} + \frac{cz^2}{2}$$

where  $a$ ,  $b$  and  $c$  are constants. One example is spring force. If we put  $b = 0$ ,  $c = 0$  and  $a = -k$

$$F = -kx\hat{i}$$

The corresponding potential energy is

$$U = \frac{kx^2}{2}$$

Linear force: if the force is given by  $F$  is equal to  $ax\hat{i} + by\hat{j} + cz\hat{k}$ , then it can be expressed as the gradient of  $U$  is equal to  $\frac{ax^2}{2} + \frac{by^2}{2} + \frac{cz^2}{2}$ , where  $a$ ,  $b$  and  $c$  are constant. If you take the gradient of  $U$ , this becomes  $ax\hat{i} + by\hat{j} + cz\hat{k}$ .

One example is a spring force. If we put  $b$  is equal to 0,  $c$  is equal to 0 and  $a$  is equal to minus  $k$ , we get a force field that is  $F$  is equal to minus  $ax$ . The corresponding potential energy from this is  $U$  is equal to  $kx^2$  by 2. So, the spring force is a conservative force.

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**Conservation of mechanical Energy:**

We know that  $\int_{r_1}^{r_2} F \cdot dr = \frac{1}{2}mV_2^2 - \frac{1}{2}mV_1^2 \rightarrow -(U_2 - U_1)$

Using the definition of potential energy,

$$U_1 - U_2 = \frac{1}{2}mV_2^2 - \frac{1}{2}mV_1^2$$

$$U_1 + \frac{1}{2}mV_1^2 = U_2 + \frac{1}{2}mV_2^2$$

Thus, the sum of the potential energy and the kinetic energy for a particle remains constant at all times during the motion of the particle, provided the force field is conservative.

Now, we discuss the conservation of mechanical energy which is valid for a conservative force field. We know, from the basic concepts developed in the beginning that  $\int_{r_1}^{r_2} F \cdot dr$  is equal to half  $m V_2^2$  minus half  $m V_1^2$  square. This has directly come from the Newton's second law. Using the definition of potential energy that is  $\int_{r_1}^{r_2} F \cdot dr$  is equal to, actually negative of the change in the  $U$  is the work done on the particle by the field and it has the tendency to decrease the potential energy. The force field is trying to decrease the potential energy; therefore,  $\int_{r_1}^{r_2} F \cdot dr$  minus  $U_2$  minus  $U_1$ .

If the particle works against the applied force field, if you could raise that particle in the gravitational field, then the potential energy of the particle will increase. But if the particle is allowed to freely fall then the potential energy will decrease. Force field is doing some work on the particle and therefore, the particle's potential energy is decreasing. It will try to make the potential energy minimum. Therefore, if you do not constraint the motion, the particle will keep on moving.

If you drop one object from the sixth floor of a building, it will reach the ground floor. If the ground, was a pit at the ground then the particle would have gone in the pit. If the pit was even deeper, it would how gone even more deep, it is like that. Force field tries to minimize the potential energy. Therefore, this point is important that  $\int_{r_1}^{r_2} \mathbf{F} \cdot d\mathbf{r}$  is basically  $U_1$  minus  $U_2$ , not  $U_2$  minus  $U_1$ . Using the definition of potential energy,  $U_1$  minus  $U_2$  is equal to  $\frac{1}{2} m V_2^2$  square minus  $\frac{1}{2} m V_1^2$  square.

We can take this to this side and we say,  $U_1$  plus  $\frac{1}{2} m V_1^2$  square is equal to  $U_2$  plus  $\frac{1}{2} m V_2^2$  square. Thus, the sum of the potential energy and the kinetic energy for a particle remains constant at all times during the motion of the particle provided the force field is conservative. Thus, in a gravitational field that kinetic energy plus potential energy of a particle will remain constant. If you throw a projectile then during the motion of the projectile, this kinetic energy plus potential energy will remain conjunct. It will remain same throughout. Therefore, this conservation principle can be applied at many places.

However, if there are non conservative forces such as friction force, which of course is non-conservative, because it depends on the path of motion. In that case, conservation of energy cannot be implied. However, this basic thing that  $\int_{r_1}^{r_2} \mathbf{F} \cdot d\mathbf{r}$  is equal to  $\frac{1}{2} m V_2^2$  square minus  $\frac{1}{2} m V_1^2$  square holds good, because it directly comes from Newton's second law. Condition is that you must identify proper  $\mathbf{F}$ . If you have, it is not only what you have applied, but also friction forces or the external forces have to be included in this.

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$$\text{curl } \vec{F}$$

$$F_x = \frac{y}{x^2+y^2}, F_y = \frac{-x}{x^2+y^2}, F_z = 0$$

$$\text{curl } \vec{F} = 0$$


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$$\int \vec{F} \cdot d\vec{r} = \frac{1}{2} m (v_2^2 - v_1^2)$$
$$\begin{aligned} \int \vec{F} \cdot d\vec{r} &= \int m \frac{d\vec{V}}{dt} \cdot d\vec{r} \\ &= \int m \frac{d(\vec{V}_t + \vec{V}_{rel})}{dt} \cdot d\vec{r} = \int m \frac{d\vec{V}_{rel}}{dt} \cdot d\vec{r} \\ &= \int m v_{rel} \cdot \frac{d\vec{r}}{dt} dt = \frac{1}{2} m (v_{rel,2}^2 - v_{rel,1}^2) \end{aligned}$$

We have learnt that in this case,  $\int \vec{F} \cdot d\vec{r}$  is equal to this half  $m V_2$  square minus  $V_1$  square is a very basic equation. Let us see whether what velocity should be taken, if I am doing the work on a train. The train is moving with a constant velocity. You have a particle and displaced by some amount  $dr$  or some amount  $r$ , some work has been done there on this thing. Particle has displaced by an amount of  $r$  in a relative frame of reference; that means in the frame of reference of the train. In the absolute frame, for an observer who is standing outside the particle has displaced even a greater distance.

Similarly, the velocities, the question is that in this case, which velocity should be taken? Whether the velocities relative to the frame of reference or the absolute velocity of the particle in calculating the kinetic energy. Let us go to basics for understanding this.  $\vec{F} \cdot d\vec{r}$  is equal to  $m$ . Here, that force  $m \frac{dV}{dt}$ , this  $V$  is the absolute velocity and this  $\frac{dV}{dt}$  is the absolute acceleration; that means as observed by standing in the inertial frame of reference, which is ground in this case and this is  $dr$ .

This can be written as  $m \frac{dV_t}{dt}$ , that velocity of the train plus  $V$  is relative  $V_{rel}$  by  $dt$  dot  $dr$ . Let us consider that we are taking this  $r$  is a relative  $r$ . So,  $\vec{F} \cdot d\vec{r}$  is the work done in the relative frame of reference. Work done will depend on the reference frame chosen. If you have applied a force

F on the particle, the work seen by the outside observer maybe more but in this case, it is  $\vec{F} \cdot d\vec{r}$  in the relative frame. So,  $m \frac{dV_t}{dt} + V_{rel} \frac{dr}{dt}$ .

If  $V_t$  is constant, then in that case this equation just simplifies to  $m \frac{dV_{rel}}{dt} dr$ , which can be written as  $m V_{rel} \frac{dr}{dt} \frac{dr}{dt}$ .  $\frac{dr}{dt}$  is again  $V_{rel}$ , which will give you half  $m V_{rel2}^2$  minus  $V_{rel1}^2$  square. Therefore, this relation that work kinetic energy theorem can be applied in the inertial frame of reference; that means, which is moving at a constant velocity; provided, you take the work done in the relative reference frame. That means, force multiplied by the displacement in the relative frame and is equal to change in kinetic energy in the relative frame.

In this case,  $V_2$  and  $V_1$  should be the relative velocity in the frame not the absolute velocity. What happens if the frame is non-inertial?

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The image shows a handwritten derivation on a light blue background. At the top, it states  $\frac{d\vec{v}_t}{dt} = \vec{a}_t$ . Below this, the work done  $\int \vec{F} \cdot d\vec{r}$  is equated to  $\int m \frac{d(\vec{v}_t + \vec{v}_{rel})}{dt} \cdot d\vec{r}$ . This is then expanded to  $\int (m \vec{a}_t + m \frac{d\vec{v}_{rel}}{dt}) \cdot d\vec{r}$ . The next line shows  $\int (\vec{F} - m \vec{a}_t) \cdot d\vec{r} = \int m \frac{d\vec{v}_{rel}}{dt} \cdot d\vec{r}$ . Finally, an arrow points to the result:  $= \frac{1}{2} m (v_{rel2}^2 - v_{rel1}^2)$ .

That means  $\frac{dV_t}{dt}$  is not 0. Train is having some acceleration and that is indicated by  $a_t$ . Acceleration of the train, in that case it becomes  $m \vec{F} \cdot d\vec{r}$  is equal to  $m \frac{dV_a}{dt} V_t + V_{rel} \frac{dr}{dt}$ , which will be equal to  $m a_t + m \frac{dV_{rel}}{dt} \frac{dr}{dt}$ . I can take this  $m a_t$  to this side and can write that  $\vec{F} - m \vec{a}_t \cdot d\vec{r}$  is equal to  $m \frac{dV_{rel}}{dt} \frac{dr}{dt}$ , which gives you this half  $m V_{rel2}^2$  square minus  $V_{rel1}^2$  square. Thus, the work energy theorem in the case of non inertial

frame can be implied, provided you replace the force by effective force; that is  $F$  minus  $m a_t$ , where  $a_t$  is the acceleration of the frame. Then, in this case it will be  $V_{rel2}$  square minus  $V_{rel1}$  square. These are the relative velocities. Let us discuss some interesting things about the energy approach. We have discussed the work energy theorem that work done, if you do a work on this one that is change in the energy.

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Minimum Potential Energy

$\Pi = \frac{1}{2} k x^2 - \frac{F x}{\text{Potential}}$

$\frac{d\Pi}{dx} = 0 = kx - F = 0$   
 $kx = F \quad x = \frac{F}{k}$

$\frac{d^2\Pi}{dx^2} = k$

$\frac{1}{2} F x = \frac{1}{2} k x^2$   
 $x = \frac{F}{k}$

We also have a principle for conservative system that is principle of minimum potential energy in statics. When a conservative force field is applied, the potential energy of the system tries to be minimum. If you take a spring and apply a force  $F$ , in this case, if you give a displacement  $x$ , potential energy as a function of  $x$  can be written as this;  $U$  is equal to total potential energy. If I indicate it by big pi, big pi is equal to half  $k x$  square strain energy of the spring minus  $F_x$ . Remember, the potential associated with the force has been written as minus  $F_x$  because if you take the negative gradient of minus  $F_x$  you get equal to plus  $F$ .

Here, many times, students may get confused. Then, they put half  $kx$  square minus plus  $F_x$ . Similarly, sometimes the students may put half  $F_x$  considering that, the force is gradually applied and therefore, the work done is equal to half  $F_x$ . Remember that this is not a work done, rather it can be called a potential; this is potential.

This  $x$  is the value of  $x$  which minimizes the potential energy will be the solution of the problem. For that, the necessary condition is that  $d\tau$  by  $dx$  is equal to 0. Thus,  $kx$  minus  $F$  is equal to 0 and we get  $kx$  is equal to  $F$ . If the potential energy gets minimized then it is a stable equilibrium situation. In this case, we can see that  $d^2\tau$  by  $dx^2$  is equal to  $k$ . As,  $k$  is positive, this is minimization problem. This is a minima and therefore,  $x$  is equal to  $F$  by  $k$  provides the stable equilibrium of the spring force system.

If the same problem is done by the work energy concept then it can be done in this way. If the force is equal to, in this case, force is  $F$ ; however, this  $F$  is applied very gradually. Spring force is linear; therefore, the force is also linear. At each and every stage, applied force and the stretching force is, it is a Cauchy's static situation. No dynamic effects are there and therefore, it is linear. In this case, the work done will be equal to half  $F_x$ . This work done on the spring is stored as the potential energy of the spring.


The potential energy of the spring increases because external force is doing a work against the conservative force field; that is in the opposite direction, minus  $kx$ . This half  $F_x$  will be half  $kx^2$  and which gives us the solution as  $x$  is equal to  $F$  by  $k$ . We get the same thing by the work energy theorem. Similar to the case of statics, we have the principle in the dynamics too.

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Principle of Least action

$$\text{Action} = (KE - PE)$$

$$S = \int_{t_1}^{t_2} (KE - PE) dt$$

$$= (KE_{av} - PE_{av})(t_2 - t_1)$$


That is called principle of least action. Action along the particle trajectory is defined as kinetic energy minus potential energy;  $KE$  minus  $PE$  and this is instantaneous action, the total value of action. A class the trajectory from  $t_1$  to  $t_2$  time,  $t_1$  to  $t_2$  is  $KE$  minus  $PE$   $dt$ . One can write the same as,  $KE$  average minus  $PE$  average, where the average is appropriately defined. Average is not just the arithmetic average, it is by the integration. If now, the particle follows that path which minimizes its action, this is the thing.

This theorem can be used. Therefore, if a particle is moving from here to here, there can be various paths; but out of all paths, the particle will follow that path which provides the minimum action. We can show from the principle of least action that a free particle that starts from one position at time  $t_1$  and arrives at a different position moves along a straight line here. So, this is the particle's position, this is time and this is particles  $x$ . How to find out the minimum path? This requires some advanced study.

It is not possible to find out this by means of, by taking infinite number of paths and evaluating that. Even the fastest computer will not be able to study these types of things. One has to find, study calculus of variation, however that we will not be discussing at this point.