

Non-Linear Vibration
Prof. S. K. Dwivedy
Department of Mechanical Engineering
Indian Institute of Technology, Guwahati

Module - 3
Solution of Nonlinear Equation of Motion
Lecture - 6
Method of Averaging

So, welcome to today's class of non-linear vibration. So, we are continuing with the solution of non-linear equation of motion, and already we have covered different methods; and today class we are particularly study about this method of averaging. Before studying this method of averaging, let me briefly review the other methods what I told you in the previous classes.

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Review of Previous Lectures

Method of Multiple Scales

$$T_n = \varepsilon^n t \quad \underline{T_0, T_1, T_2}$$

$$\frac{d}{dt} = \frac{dT_0}{dt} \frac{\partial}{\partial T_0} + \frac{dT_1}{dt} \frac{\partial}{\partial T_1} + \dots = D_0 + \varepsilon D_1 + \dots$$

$$\frac{d^2}{dt^2} = D_0^2 + 2\varepsilon D_0 D_1 + \varepsilon^2 (D_1^2 + 2D_0 D_2) + \dots$$

So particularly, this method of multiple scales and harmonic balance method, so in case of method of multiple scales, we are taking different time scales which can be represented by using the time terms T_n equal to $\varepsilon^n t$ to the power t , where t is the time, ε is the book keeping parameter and these time scales are T_0 , T_1 , T_2 like this. So, where this T_0 equal to t and T_1 equal to εt and T_2 equal to $\varepsilon^2 t$. So, by substituting these different time scales in the original differential equation of motion and separating the terms with different order of ε then, we can

get a set of equations. So, from those equations we can eliminate the secular terms to get the frequency equations.

So, in this case initially taking this T_n equal to epsilon to the power n t, so one can write this d by d t equal to D_0 plus epsilon D_1 , and d square by d t square equal to D_0 square plus 2 epsilon $D_0 D_1$ plus epsilon square D_1 square plus 2 $D_0 D_2$ plus the higher order terms.

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$$x(t; \epsilon) = \epsilon x_1(T_0, T_1, T_2, \dots) + \epsilon^2 x_2(T_0, T_1, T_2, \dots) + \epsilon^3 x_3(T_0, T_1, T_2, \dots) + \dots$$

$$D_0^2 x_1 + \omega_0^2 x_1 = 0$$

$$D_0^2 x_2 + \omega_0^2 x_2 = -2D_0 D_1 x_1 - \alpha_2 x_1^2$$

$$D_0^2 x_3 + \omega_0^2 x_3 = -2D_0 D_1 x_2 - D_1^2 x_1 - 2D_0 D_2 x_1 - 2\alpha_2 x_1 x_2 - \alpha_3 x_1^3$$

$$\ddot{x} + \alpha_1 x + \alpha_2 x^2 + \alpha_3 x^3 = 0$$

$$\alpha_1 = \omega_0^2$$

By taking these terms and assuming that the equation in this form x equal to epsilon x_1 plus epsilon square x_2 plus epsilon cube x_3 , and by substituting it in the original equation. So, the equation what we have studied in this case is x double dot plus alpha 1 x plus alpha 2 x square plus alpha 3 x cube. So, we have taken three terms and we have equated that to 0. So, in this case this alpha 1 equal to omega 0 square by substituting this equation, in this equation we and we got these three equations, so by taking the different order of epsilon. So, we know the solution of this equation D_0 square x_1 plus omega 0 square x_1 .

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$$x_1 = A(T_1, T_2) \exp(i\omega_0 T_0) + \bar{A} \exp(-i\omega_0 T_0)$$

$$D_0^2 x_2 + \omega_0^2 x_2 = -2i\omega_0 D_1 A \exp(i\omega_0 T_0) - \alpha_2 [A^2 \exp(2i\omega_0 T_0) + A\bar{A}] + cc$$

Eliminating Secular term

$$D_1 A = 0$$

$$x_2 = \frac{\alpha_2 A^2}{3\omega_0^2} \exp(2i\omega_0 T_0) - \frac{\alpha_2}{\omega_0^2} A\bar{A} + cc$$

So, this solution by substituting it in this equation second equation and getting the solution of x_1 and x_2 , and substituting those two in the third equation, we can get different frequency response equations. So, for example, in this case by the solution of this equation equal to a $1 e$ to the power $i\omega_0 t_0$ plus a $1 \bar{e}$ to the power minus $i\omega_0 t_0$. So, the solution by substituting in the second equation, so you got the equation to be written in this form, and in this case one can see that this is the secular term; that means, presence of this term will leads to a infinite will leads to infinite response. So, one has to eliminate this term. So, to eliminate this term one can find that this $D_1 A$ equal to 0. So, by taking $D_1 A$ equal to 0, so that means, making this term equal to 0.

So, one can find this expression for x_2 . So, x_2 can be written as $\alpha_2 A^2$ by $3\omega_0^2$ e to the power $2i\omega_0 T_0$ minus α_2 by ω_0^2 $AA \bar{A}$ plus it is complex conjugate. So, now, we know the solution of solution in terms of x_1 and x_2 , but we do not know what is the exact value of a and b .

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$$D_0^2 x_3 + \omega_0^2 x_3 = - \left[\frac{2i\omega_0 D_2 A - \frac{10\alpha_2^2 - 9\alpha_3 \omega_0^2}{3\omega_0^2} A^2 \bar{A}}{\frac{3\alpha_3 \omega_0^2 + 2\alpha_2^2}{3\omega_0^2} A^3 \exp(3i\omega_0 T_0)} + cc \right] \exp(i\omega_0 T_0)$$

To eliminating secular term

$$2i\omega_0 D_2 A - \frac{10\alpha_2^2 - 9\alpha_3 \omega_0^2}{3\omega_0^2} A^2 \bar{A} = 0$$

$$A = \frac{1}{2} a \exp(i\beta)$$

So, by substituting these two x_1 and x_2 in the third equation, so one can find this expression, so here also one can note that this term is a secular term, because the coefficient of e to the power $i\omega_0 T_0$. So, one has to eliminate this term. So, to eliminate this term one can write this $2i\omega_0 D_2 A - \frac{10\alpha_2^2 - 9\alpha_3 \omega_0^2}{3\omega_0^2} A^2 \bar{A} = 0$, now substituting this A equal to half a e to the power $i\beta$ in this equation.

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$$\omega_0 a \beta' = 0$$

$$\omega_0 a \beta' + \frac{10\alpha_2^2 - 9\alpha_3 \omega_0^2}{24\omega_0^3} a^3 = 0$$

$$\beta = \frac{9\alpha_3 \omega_0^2 - 10\alpha_2^2}{24\omega_0^3} a^2 T_2 + \beta_0$$

$$A = \frac{1}{2} a \exp \left[i \frac{9\alpha_3 \omega_0^2 - 10\alpha_2^2}{24\omega_0^3} \epsilon^2 a^2 t + i\beta_0 \right]$$

So, one can separating the real and imaginary parts, so one can get these two equations that is $\omega_1 \dot{a} = 0$ and $\omega_0 \dot{\beta} + 10\alpha_2 a^2 \sin 2\omega_0 t - 9\alpha_3 \omega_0^2 a \cos 2\omega_0 t = 0$. So, for steady state, so one can substitute this \dot{a} and $\dot{\beta}$ equal to 0. So, from this equation one can known thus that this \dot{a} equal to 0. So, a is not a function of t_0 and t_1 . So, now, from this equation one can find, so in this case as \dot{a} is da/dt , so it is equal to 0, so \dot{a} is not a function of t_0 , t_1 and t_2 .

But beta one can find this $\dot{\beta}$ equal to $-\frac{10\alpha_2 a^2 \sin 2\omega_0 t - 9\alpha_3 \omega_0^2 a \cos 2\omega_0 t}{\omega_0}$ and integrating that thing so one can get this expression. So, from this expression of a and β , so one can write a equal to $\frac{1}{2} \epsilon a_1 \cos(\omega_0 t + \beta_0) + O(\epsilon^3)$ to the power i $9\alpha_3 \omega_0^2 a^2 \sin 2\omega_0 t - 10\alpha_2 a^2 \cos 2\omega_0 t$ by $24\alpha_1^2 \epsilon^2 a^2$ plus i beta 0.

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$$x = \epsilon a \cos(\omega t + \beta_0) - \frac{\epsilon a^2 \alpha_2}{2\alpha_1} \left[1 - \frac{1}{3} \cos(2\omega t + 2\beta_0) \right] + O(\epsilon^3)$$

$$\omega = \omega_0 \left[1 + \frac{9\alpha_3 \alpha_1 - 10\alpha_2^2}{24\alpha_1^2} \epsilon^2 a^2 \right] + O(\epsilon^3)$$

So, one can find finally, the solution in this form. So, x equal to $\epsilon a_1 \cos \omega_0 t + \beta_0 - \frac{\epsilon a_1^2 \alpha_2}{2\alpha_1} \left[1 - \frac{1}{3} \cos 2\omega_0 t + 2\beta_0 \right] + O(\epsilon^3)$ and these are the higher order terms, so neglecting these higher order terms the solution can be written in this form where this ω equal to ω_0 into $1 + \frac{9\alpha_3 \alpha_1 - 10\alpha_2^2}{24\alpha_1^2} \epsilon^2 a_1^2$ square a square. So, in this case of method of multiple scale by taking different time scales one can find the solution of the equation motion and the frequency response curve.

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THE METHOD OF HARMONIC BALANCE

$$x = \sum_{m=0}^M A_m \cos(m \omega t + m \beta_0)$$

Example $\ddot{x} + \alpha_1 x + \alpha_2 x^2 + \alpha_3 x^3 = 0$

$\alpha_1 = \omega_0^2$

Single term $x = A_1 \cos(\omega t + \beta_0) = A_1 \cos \phi$

Two term $x = A_1 \cos \phi + A_0$

Similarly in case of the harmonic balance method, one can take the solution in this form, so x equal to summation of different harmonics, so either one can write in terms of cos and sin or one can use this cos and this phase difference beta, so one can write x equal to m equal to 0 to M $A_m \cos m \omega t + m \beta_0$. So, last class we have studied this method by taking this example of x double dot plus $\alpha_1 x$ plus $\alpha_2 x^2$ plus $\alpha_3 x^3$ equal to 0, so one can take a single term or two terms or multiple terms. So depending on the complexity of the problem, so one can consider different terms or different order of terms by taking single term and multiple terms.

Last time we have shown that depending on the number of terms the accuracy of the solution can be achieved. So, by taking a single term that is x equal to $A_1 \cos \phi$ or by taking a two terms $A_1 \cos \phi$ plus A_0 or taking three terms that is A_0 plus $A_1 \cos \phi$ plus $A_2 \cos 2 \phi$. So, one can find the solution of this equation, the disadvantage of this harmonic balance method is that one has to know the solution or up to what order one has to take one should have the idea or but in case of method of multiple scale, one can one enough have to know the solution of theory and one more advantage of method of multiple scale above this method of harmonic balance is that one can study the stability of the systems by taking the reduced equations.

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$$\begin{aligned} \omega \alpha' &= 0 \\ \omega_0 a \beta' + \frac{10\alpha_2^2 - 9\alpha_1 \omega_0^2}{24\omega_0^3} a^3 &= 0 \end{aligned} \quad \left. \vphantom{\begin{aligned} \omega \alpha' &= 0 \\ \omega_0 a \beta' + \frac{10\alpha_2^2 - 9\alpha_1 \omega_0^2}{24\omega_0^3} a^3 &= 0 \end{aligned}} \right\} \text{Reduced eq}$$
$$\beta = \frac{9\alpha_1 \omega_0^2 - 10\alpha_2^2}{24\omega_0^3} a^2 T_2 + \beta_0$$
$$A = \frac{1}{2} a \exp \left[i \frac{9\alpha_1 \omega_0^2 - 10\alpha_2^2}{24\omega_0^3} \varepsilon^2 a^2 t + i \beta_0 \right]$$

So, in this case taking these two as the reduced equation, so one can study the stability of this reduced equation and one can study the stability of the whole system, but in case of this by using this harmonic balanced method after getting the solution one has to perturb the original equation of motion to study the stability of the system.

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METHOD OF AVERAGING

- Van der Pol's technique ✓
- Krylov-Bogoliubov technique ✓
- the generalized method of averaging
- the Krylov-Bogoliubov-Mitropolsky technique
- averaging using canonical variables
- averaging using the Lie series
- transforms and averaging using Lagrangians

So, today class we will study about the method of averaging, so in case of method of averaging. So, which is a variational technique, so there are several techniques available to study or to find the non-linear solution of the non-linear differential equation, so in

this case one can use this Van der Pol technique or Krylov-Bogoliubov technique, the generalized method of averaging the Krylov-Bogoliubov-Mitropolsky technique, averaging using canonical variables, averaging using the Lie series, transform and averaging using Lagrangians. So, there are many methods available and we are going to study today about this Van der Pol's technique and the Krylov-Bogoliubov technique and next class we will study about this generalized method of averaging and Krylov-Bogoliubov-Mitropolsky technique.

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• Van der Pol's Technique (1926)

$$\frac{d^2 u}{dt^2} + \omega_0^2 u + \varepsilon (u^2 - 1) \frac{du}{dt} = \varepsilon f \Omega \cos \Omega t$$

$$\Omega = \omega_0 + \varepsilon \sigma$$

$$u(t) = a_1(t) \cos \Omega t + a_2(t) \sin \Omega t$$

$a_1(t)$ $a_2(t)$ slowly varying function of time

$$\frac{da_i}{dt} = o(\varepsilon) \quad \frac{d^2 a_i}{dt^2} = o(\varepsilon^2)$$

*A/H Nayfeh
Perturbation
technique*

$\Omega = \omega_0$

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So, in 1926 Van der pol propose this method to solve this type of equation. So, this equation is generally known as the Van der Pol equation in which one can write the equation in this form that is $d^2 u / dt^2 + \omega_0^2 u + \varepsilon (u^2 - 1) du / dt = \varepsilon f \Omega \cos \Omega t$, here f is the amplitude of the excitation external excitation, ω_0 is the frequency of external excitation and this ω_0 is the natural frequency of the linear system, and then this ε is a book keeping parameter, so depending on the value of this u . So, if u^2 is greater than 1, so this term will be positive, so this may be this term will be positive and, but when u^2 is less than this term then this may be negative; that means, this bracketed term if it is negative.

So, this equation is similar to that of a equation in which we have negative damping and if u^2 is greater than 1 then we have a simple equation similar to that of a single

degree of freedom system, but one can use or one can add a non-linear terms also to this equation. So, in this equation was originally solved in 1926 by Van der pol using the following technique, this method is adapted from the book by A H Nayfeh Perturbational Perturbation technique, so it is adapted from this book perturbation technique by A H Nayfeh technique, so in this method, so one can assume the solution u to be a $1 \cos \omega t$ plus a $2 \sin \omega t$, so in case of method of variation technique, so one by putting this epsilon equal to 0.

So, here we are assuming this epsilon to be small and if one can take this epsilon equal to 0, the equation will reduce to $d^2 u / dt^2 + \omega_0^2 u = 0$. The solution of that linear equation equal to $u = a_1 \cos \omega t + a_2 \sin \omega t$, so where this a_1 and a_2 are constant; that means, if we are considering epsilon equal to 0 then one can write the solution equal to $u(t) = a_1 \cos \omega t + a_2 \sin \omega t$, where this ω is nearly equal to this ω_0 or in that case if you are taking only the linear term that is $d^2 u / dt^2 + \omega_0^2 u = 0$ then the solution will be $u(t) = a_1 \cos \omega_0 t + a_2 \sin \omega_0 t$, where this ω equal to ω_0 .

That is the natural frequency of the system, but if we are taking a variation; that means, if we are taking non-zero epsilon, but a very small value of epsilon still this equation will be valid, but in this case, this a_1 and a_2 can be considered as function of time so this. So, one has to consider this a_1 and a_2 in such way that they are slowly varying function of time. So, that this da_1/dt ; that means, da_1/dt and da_2/dt will be of the order of epsilon and this $d^2 a_1/dt^2$ and $d^2 a_2/dt^2$ or $d^2 a_1/dt^2$ by dt^2 will be further slowly varying term that is it should be of the order of epsilon square. So, by taking this $a_1(t)$ and $a_2(t)$ as slowly varying function of time and using this $u(t) = a_1(t) \cos \omega t + a_2(t) \sin \omega t$ and substituting this equation in the original equation.

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$$\begin{aligned}\dot{u} &= (\dot{a}_1 + a_2 \Omega) \cos \Omega t + (\dot{a}_2 - a_1 \Omega) \sin \Omega t \\ \ddot{u} &= (-\Omega^2 a_1 + 2\Omega \dot{a}_2 + \ddot{a}_1) \cos \Omega t + (-2\Omega \dot{a}_1 + \ddot{a}_2 - \Omega^2 a_2) \sin \Omega t \\ &= (-\Omega^2 a_1 + 2\Omega \dot{a}_2 + \ddot{a}_1) \cos \Omega t + (-2\Omega \dot{a}_1 + \ddot{a}_2 - \Omega^2 a_2) \sin \Omega t + \\ &\quad \omega_0^2 (a_1(t) \cos \Omega t + a_2(t) \sin \Omega t) + \\ &\quad \varepsilon ((a_1(t) \cos \Omega t + a_2(t) \sin \Omega t)^2 - 1) \\ &\quad ((\dot{a}_1 + a_2 \Omega) \cos \Omega t + (\dot{a}_2 - a_1 \Omega) \sin \Omega t) = \varepsilon f \Omega \cos \Omega t\end{aligned}$$

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• Van der Pol's Technique (1926)

$$\frac{d^2 u}{dt^2} + \omega_0^2 u + \varepsilon (u^2 - 1) \frac{du}{dt} = \varepsilon f \Omega \cos \Omega t$$

$$\Omega = \omega_0 + \varepsilon \sigma$$

$$u(t) = a_1(t) \cos \Omega t + a_2(t) \sin \Omega t$$

$a_1(t)$ $a_2(t)$ slowly varying function of time

$$\frac{da_i}{dt} = o(\varepsilon) \quad \frac{d^2 a_i}{dt^2} = o(\varepsilon^2)$$

Handwritten note: A-Hilbert's Perturbation technique

Handwritten note: $\Omega = \omega_0$

So, one can write this equation in this following form. So, this u as we have taking this u equal to a 1 t cos omega t plus a 2 t sin omega t. So, this u dot by differentiating this thing, so one can get so differentiation of a 1 will be a 1 dot. So, it will be a 1 dot cos omega t and then minus a 1 omega sin omega t, similarly differentiating this term this a 2 t sin omega t 1 can write. So, this will be a 2 dot sin omega t plus a 2 omega cos omega t. So, we have four terms for this u dot t, so u dot can be written as by combining this coefficient of cos and sin one can write this u dot equal to a dot a 2 omega into cos omega t plus a 2 dot minus a 1 omega sin omega t.

So, again differentiating this equation one can write this u double dot equal to, so by differentiating this inside term one can write this will be a 1, so this is a 1. So, a 1 double dot plus a 2 dot omega and into cos omega t, now differentiating this keeping this constant and differentiating this cos omega t 1 can write this as minus a 1 dot plus a 2 omega into sin omega t, similarly differentiating this term one can write this is a 2 double dot minus a 1 dot omega into sin omega t plus a 2 dot minus a 1 omega into omega cos omega t. So, by collecting the coefficient of cos omega t and sin omega t 1 can write this u double dot in this form, so this is equal to minus omega square a 1 plus 2 omega a 2 dot plus a 1 double dot into cos omega t plus minus 2 omega a 1 dot plus a 2 dot minus omega square a 2 into sin omega t.

Now, substituting this equation of u double dot and u t in this original equation that is u d square u by d t square plus omega 0 square u plus epsilon into u square minus 1 into d u by d t equal to epsilon f omega cos omega t, so one can write this equation in this form. So, this will be this for u double dot initially it is written then the term containing this u square, so this is u square minus 1 and this term is for d u by d t, so this is equal to epsilon f omega cos omega t.

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$$\begin{aligned} & \left((-\Omega^2 + \omega_0^2)a_1 + 2\Omega\dot{a}_2 + \ddot{a}_1 \right) \cos \Omega t + \\ & \left(-2\Omega\dot{a}_1 + \ddot{a}_2 + (-\Omega^2 + \omega_0^2)a_2 \right) \sin \Omega t + \\ & \varepsilon \Omega \left(\begin{aligned} & a_1^2 a_2 \cos^3 \Omega t - a_2^2 a_1 \sin^3 \Omega t + \\ & (a_2^3 - 2a_1^2 a_2) \sin^2 \Omega t \cos \Omega t - \\ & (a_1^3 - 2a_1 a_2^2) \cos^2 \Omega t \sin \Omega t - \\ & a_2 \Omega \cos \Omega t + a_1 \Omega \sin \Omega t \end{aligned} \right) + h.o.t = \varepsilon f \Omega \cos \Omega t \end{aligned}$$

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$$\begin{aligned}\dot{u} &= (\dot{a}_1 + a_2 \Omega) \cos \Omega t + (\dot{a}_2 - a_1 \Omega) \sin \Omega t \\ \ddot{u} &= (-\Omega^2 a_1 + 2\Omega \dot{a}_2 + \ddot{a}_1) \cos \Omega t + (-2\Omega \dot{a}_1 + \ddot{a}_2 - \Omega^2 a_2) \sin \Omega t \\ &= \underbrace{(-\Omega^2 a_1 + 2\Omega \dot{a}_2 + \ddot{a}_1)}_{\omega_0^2 (a_1(t) \cos \Omega t + a_2(t) \sin \Omega t) + \varepsilon ((a_1(t) \cos \Omega t + a_2(t) \sin \Omega t)^2 - 1)} \cos \Omega t + \underbrace{(-2\Omega \dot{a}_1 + \ddot{a}_2 - \Omega^2 a_2)}_{\varepsilon f \Omega \cos \Omega t} \sin \Omega t + \\ &= \varepsilon f \Omega \cos \Omega t\end{aligned}$$

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$$\begin{aligned}& ((-\Omega^2 + \omega_0^2) a_1 + 2\Omega \dot{a}_2 + \ddot{a}_1) \cos \Omega t + \\ & (-2\Omega \dot{a}_1 + \ddot{a}_2 + (-\Omega^2 + \omega_0^2) a_2) \sin \Omega t + \\ & \varepsilon \Omega \left(\begin{aligned} & a_1^2 a_2 \cos^3 \Omega t - a_2^2 a_1 \sin^3 \Omega t + \\ & (a_2^3 - 2a_1^2 a_2) \sin^2 \Omega t \cos \Omega t - \\ & (a_1^3 - 2a_1 a_2^2) \cos^2 \Omega t \sin \Omega t - \\ & a_2 \Omega \cos \Omega t + a_1 \Omega \sin \Omega t \end{aligned} \right) + \text{hot} = \varepsilon f \Omega \cos \Omega t\end{aligned}$$

$\cos 2t = \cos(2t + t)$
 $= \cos 2t \cdot \cos t - \sin 2t \cdot \sin t$
 $= (2\cos^2 t - 1)\cos t - 2(1 - \cos^2 t)\sin t$

So, now, one can arrange these terms and one can write the equation in this form. So, arranging the terms with $\cos \omega t$, so this is the term with $\cos \omega t$, and there are some other terms also. So, this term also contain $\cos \omega t$, and one can find as minus 1 into this term is there, so this is also $\cos \omega t$, so similarly by collecting the terms with $\sin \omega t$, so with double line. So, this is $\sin \omega t$ this term is also for $\sin \omega t$ and. So, with minus 1 if it is multiplied then this term also contain $\sin \omega t$, so one can write these terms in that way.

So, this becomes minus ω^2 plus ω_0^2 into $a_1 + 2\omega$ into $a_2 \dot{+}$ plus $a_1 \ddot{+}$, so $a_1 \ddot{+}$ into $\cos \omega t$ plus minus $2\omega a_1 \dot{+}$ plus $a_2 \dot{+}$ plus minus ω^2 plus ω_0^2 into a_2 into $\sin \omega t$ plus $\epsilon \omega$ into $a_1^2 \cos^3 \omega t$ minus $a_2^2 a_1 \sin^3 \omega t$ plus $\sin^2 \omega t a_2^2 a_1 \sin \omega t$ then a_2^3 minus $2 a_1^2 a_2$ into $\sin^2 \omega t$ into $\sin \omega t \cos \omega t$ then a_1^3 minus $2 a_1 a_2^2$ square into $\cos^2 \omega t$ into $\sin \omega t$ minus $a_2 \omega \cos \omega t$ plus $a_1 \omega \sin \omega t$ and plus these higher order terms.

So, this will be equal to $\epsilon f \cos \omega t$, now we have to convert this $\cos^3 \omega t$ term $\sin^3 \omega t$ terms and the $\sin^2 \omega t \cos \omega t$ and this $\cos^2 \omega t \sin \omega t$ into the respective cosine and sine terms. So, for example, this $\cos^3 \omega t$ can be broken down by using this $\cos 3t$. So, we know this $\cos 3t$ equal to $\cos^2 t$ plus t , so we can expand this thing in this form. So, $\cos a + b$, so it will be equal to $\cos a \cos b$ minus $\sin a \sin b$. So, in that way one can write or this will be equal to $\cos^2 t$ into $\cos t$ minus $\sin^2 t$ into $\sin t$.

So, now, again breaking down this $\cos^2 t$, so $\cos^2 t$ equal to $\cos^2 t$ minus $\sin^2 t$, so or one can write this equal to $2 \cos^2 t$ minus 1 into $\cos t$, so minus this becomes $2 \sin^2 t$ $2 \sin^2 t$. So, the $\sin^2 t$ 1 can write equal to 1 minus $\cos^2 t$ 1 minus $\cos^2 t$ into $\cos t$, so this becomes $2 \sin t$ into $\cos t$ into $\sin t$, so that means, this becomes $2 \sin^2 t$ into $\cos t$, so this $2 \sin^2 t$ one can replace it by 2 into 1 minus $\cos^2 t$ into $\cos t$. So, by expanding this thing, so one can get this becomes $2 \cos^2 t$ and plus again $2 \cos^2 t$ this is $2 \cos^3 t$ and again plus $2 \cos^3 t$ that is this becomes $4 \cos^3 t$ minus.

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$$\begin{aligned}\cos^3 \Omega t &= (\cos 3\Omega t + 3\cos \Omega t) / 4 \\ \sin^3 \Omega t &= (3\sin \Omega t - \sin 3\Omega t) / 4 \quad \checkmark \\ \cos^2 \Omega t \sin \Omega t &= (\sin \Omega t - \sin 3\Omega t) / 4 \\ \sin^2 \Omega t \cos \Omega t &= (\cos \Omega t - \cos 3\Omega t) / 4\end{aligned}$$

$$\begin{aligned}\cos^2 \Omega t \sin \Omega t &= (1 - \sin^2 \Omega t) \sin \Omega t \\ &= \sin \Omega t - \sin^3 \Omega t\end{aligned}$$

So, this becomes 4 cos cube t, so 4 cos cube t minus 3 cos omega t. So, one can write this cos cube sin omega t equal to cos 3 omega t plus 3 cos omega t by 4. Similarly one can find the sin cube omega t, so by expanding this sin 3 omega t, so one can find this expression. So, the sin omega cube sin cube omega t will be equal to 3 sin omega t minus sin 3 omega t by 4, similarly this cos square omega t sin omega t can be written as sin 3 t minus sin t minus sin 3 t by 4, similarly cos square omega t into cos omega t equal to cos omega t minus cos 3 omega t by 4. For example, this cos square omega t cos square omega t into sin omega t can be written as, so for this cos square omega t one can write this is equal to 1 minus sin square omega t into sin omega t.

So, this becomes sin omega t minus sin cube omega t and already we have found the expression for the sin omega cube t equal to 3 sin omega t minus sin 3 omega t whole divided by 4. So, by substituting this thing in this equation, so one can get the expression for this cos square omega t into sin omega t equal to sin omega t minus sin 3 omega t by 4

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$$\begin{aligned}
 & \left((-\Omega^2 + \omega_0^2) a_1 + 2\Omega \dot{a}_2 + \ddot{a}_1 \right) \cos \Omega t + \\
 & \left(-2\Omega \dot{a}_1 + \ddot{a}_2 + (-\Omega^2 + \omega_0^2) a_2 \right) \sin \Omega t + \\
 & \varepsilon \Omega \left(\begin{aligned} & a_1^2 a_2 \cos^3 \Omega t - a_2^2 a_1 \sin^3 \Omega t + \\ & (a_2^3 - 2a_1^2 a_2) \sin^2 \Omega t \cos \Omega t - \\ & (a_1^3 - 2a_1 a_2^2) \cos^2 \Omega t \sin \Omega t - \\ & a_2 \Omega \cos \Omega t + a_1 \Omega \sin \Omega t \end{aligned} \right) + h.o.t = \varepsilon f \Omega \cos \Omega t
 \end{aligned}$$

$\cos 2t = \cos(2t + t)$
 $= \cos 2t \cdot \cos t - \sin 2t \cdot \sin t$
 $= (2\cos^2 t - 1)\cos t - 2(1 - \cos^2 t)\sin t$

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Similarly, one can get the sin square omega t into cos omega t equal to cos omega t minus cos 3 omega t by 4. So, substituting these expressions for cos cube omega t sin cube omega t cos square omega t sin omega t and sin square omega t into cos omega t in the previous equation, so in this equation if you substitute this thing and separate the terms with cos and sin.

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$$\begin{aligned}
 & 2\dot{a}_1 + \left(\frac{\Omega^2 - \omega_0^2}{\Omega} \right) a_2 - \varepsilon a_1 \left(1 - \frac{a_1^2 + a_2^2}{4} \right) = 0 \\
 & 2\dot{a}_2 - \left(\frac{\Omega^2 - \omega_0^2}{\Omega} \right) a_1 - \varepsilon a_2 \left(1 - \frac{a_1^2 + a_2^2}{4} \right) = \varepsilon f
 \end{aligned}
 \quad \left. \begin{aligned} & a_1 \rightarrow a_1 e^{i\Omega t} \\ & a_2 \rightarrow a_2 e^{i\Omega t} \\ & f \rightarrow f_0 e^{i\Omega t} \end{aligned} \right\}$$

$$\frac{\Omega^2 - \omega_0^2}{\Omega} = \frac{(\omega_0 + \varepsilon \sigma)^2 - \omega_0^2}{\Omega} = \frac{\cancel{\omega_0^2} + 2\varepsilon \omega_0 \sigma + \varepsilon^2 \sigma^2 - \cancel{\omega_0^2}}{\Omega} \approx \frac{2\varepsilon \omega_0 \sigma}{\Omega}$$

$$\rho = \frac{a_1^2 + a_2^2}{4}$$

$\Omega = \omega_0 + \varepsilon \sigma$
 $\sigma \rightarrow \text{detuning parameter}$

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One can obtain two equations, so by collecting the terms with coefficient cos omega t and by collecting the terms with coefficient sin omega t, so we can get these two

equations. So, these two equations can be written or these are two $\frac{1}{2} + \omega^2$ minus ω_0^2 square by ω into $\frac{1}{2} - \epsilon$ into $1 - \frac{1}{2}$ square plus $\frac{1}{2}$ square by 4 equal to 0, and the second equation becomes $\frac{1}{2} + \omega^2$ minus ω_0^2 square by ω into $\frac{1}{2} - \epsilon$ into $1 - \frac{1}{2}$ square plus $\frac{1}{2}$ square by 4 equal to ϵf . So, in this case one can note that we have consider the equation or we have finding the solution near ω equal ω_0 ; that means, this external frequency nearly equal to the natural frequency of the system that is ω_0 and in this case we know that the solution to be periodic.

So, this $u(t)$ whatever we what we have considered, so it is equal to $\frac{1}{2} t \cos \omega t$ plus $\frac{1}{2} t \sin \omega t$. So, this amplitude of vibration will be equal to $\frac{1}{2}$ square plus $\frac{1}{2}$ square root over, so or one can write this equation $u(t)$ equal to $\frac{1}{2} t \cos \omega t$ plus $\frac{1}{2} t \cos \omega t$ or in terms of a single term $\frac{1}{2} t \cos \omega t$ plus ϕ . So, in that case the amplitude one can find from these two expressions, so instead of writing $\frac{1}{2} + \frac{1}{2}$. So, one can write using a single term $\frac{1}{2}$ also or one can use this term ρ to represent this $\frac{1}{2}$ square plus $\frac{1}{2}$ square by 4, so this ω^2 that is external frequency square minus ω_0^2 square by ω can be approximately can be written like this. So, we are using a detuning parameter to express the nearness of this external excitation to that of the natural frequency.

So, if it is very close to this natural frequency then this σ will tends to 0, so we are writing this ω that is external frequency equal to natural frequency plus this $\epsilon \sigma$, where σ is known as the detuning parameter, so detuning parameter, so using this detuning parameter, so one can write this ω equal to ω_0 plus $\epsilon \sigma$, so this detuning parameter shows the nearness of this external frequency towards the natural frequency of the system, so this ω^2 minus ω_0^2 square by ω can be written as by substituting this expression here, so ω_0 plus $\epsilon \sigma$ whole square minus ω_0^2 square by ω , so which is equal to, so this becomes ω_0^2 square plus $2 \epsilon \omega_0 \sigma$ plus $\epsilon^2 \sigma^2$ minus ω_0^2 square by ω , so this ω_0^2 square ω_0^2 square get cancelled and one can write this term is nearly equal to $2 \epsilon \sigma$, so by putting this term nearly equal to $2 \epsilon \sigma$.

So these two equations can be written as $2\ddot{a}_1 + 2\epsilon\sigma\dot{a}_2 - \epsilon a_2 = 0$ and $\ddot{a}_1 = 0$ similarly the second equation can be written as $2\ddot{a}_2 - 2\epsilon\sigma\dot{a}_1 - \epsilon a_1 = 0$ and $\ddot{a}_2 = 0$ so from these two equations. So, for steady state one knows that \dot{a}_1 and \dot{a}_2 will be equal to 0 as the response or equilibrium position. Equilibrium point a_1 and a_2 can be for steady state can be written as a_1 can be written as a_1^0 and a_2 can be written as a_2^0 , so this ρ can be written as ρ_0 , so for steady state when they are not function of time this \dot{a}_1 will be equal to 0 and \dot{a}_2 also will be equal to 0.

So, by substituting this expression here, so one can write these two equations in this form, so one can write $2\epsilon\sigma\dot{a}_2$, so this equation becomes $2\epsilon\sigma\dot{a}_2^0 - \epsilon a_2^0 = 0$ and $\ddot{a}_1 = 0$, similarly the second equation becomes this term 0, so this becomes $-2\epsilon\sigma\dot{a}_1^0 - \epsilon a_1^0 = 0$ and $\ddot{a}_2 = 0$ into $1 - \rho_0 = 0$ equal to ϵf the purpose of taking this f . So, instead of taking a forcing amplitude f , in this case the forcing amplitude is taken as ϵf , so that one can cancel from ω from both sides and one can get a simpler expression, so one can write this equation in this form.

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$$\begin{aligned}
 &2\epsilon\sigma\dot{a}_2 - \epsilon a_2 = 0 \\
 &-2\epsilon\sigma\dot{a}_1 - \epsilon a_1 = \epsilon f \\
 &4\sigma^2(a_1^2 + a_2^2) + (1 - \rho_0)^2(a_1^2 + a_2^2) = f^2 \\
 &\boxed{\rho_0(4\sigma^2 + (1 - \rho_0)^2) = \frac{f^2}{4}} \\
 &\boxed{4\sigma^2 + (1 - \rho_0)^2 = 0}
 \end{aligned}$$

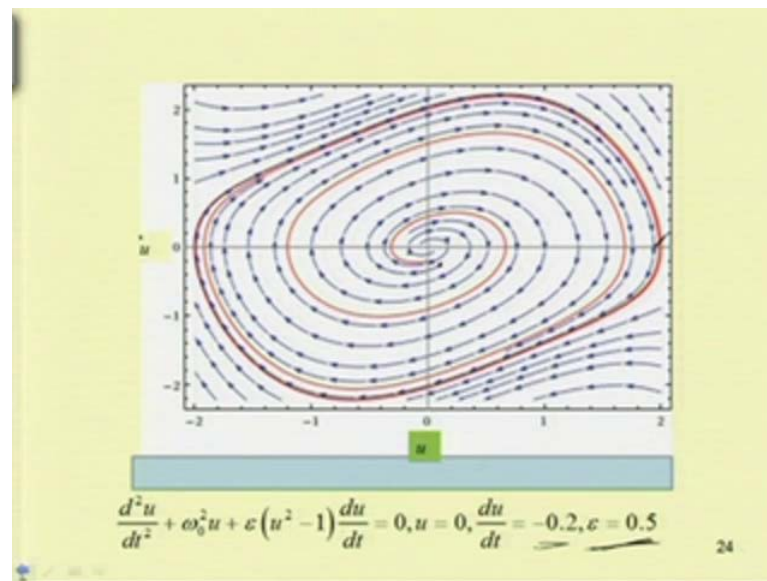
So $2\epsilon\sigma\dot{a}_2^0 - \epsilon a_2^0 = 0$ and second equation becomes $-2\epsilon\sigma\dot{a}_1^0 - \epsilon a_1^0 = \epsilon f$ and $\ddot{a}_2 = 0$ into $1 - \rho_0 = 0$

equal to ϵf . Now by squaring these two equations and adding, so one can get, so by squaring this. So, it becomes $4\epsilon^2 \sigma^2 a^2 \omega^2 \sin^2 \theta - 2 \times 2\epsilon \sigma a^2 \omega^2 \sin \theta \cos \theta + \epsilon^2 a^2 \omega^2 \cos^2 \theta = (1 - \rho)^2$, so by adding to the second equation, it becomes $4\epsilon^2 \sigma^2 a^2 \omega^2 \sin^2 \theta + \epsilon^2 a^2 \omega^2 \cos^2 \theta = (1 - \rho)^2 + 2 \times 2\epsilon \sigma a^2 \omega^2 \sin \theta \cos \theta$ that is $4\epsilon^2 \sigma^2 a^2 \omega^2 \sin^2 \theta + \epsilon^2 a^2 \omega^2 \cos^2 \theta = (1 - \rho)^2 + 4\epsilon \sigma a^2 \omega^2 \sin \theta \cos \theta$.

So, right hand side it will be equal to $\epsilon^2 f^2$, so this equation, so from this equation, so by cancelling the term that is 2×2 this into this with 2 into this term into this term, so one can finally, get this equation. So, this equation becomes $4\sigma^2 a^2 \omega^2 \sin^2 \theta + a^2 \omega^2 \cos^2 \theta = f^2$, now taking this term common that is $a^2 \omega^2 \sin^2 \theta + a^2 \omega^2 \cos^2 \theta$ equal to f^2 , so one can write this is $a^2 \omega^2 \sin^2 \theta + a^2 \omega^2 \cos^2 \theta = 4\sigma^2 a^2 \omega^2 \sin^2 \theta + (1 - \rho)^2$, but we have we can use this equation $\rho = 0$ equal to $\rho = 0$ equal to $a^2 \omega^2 \sin^2 \theta + a^2 \omega^2 \cos^2 \theta = 4$, so by using this expression, so one can write this equation in this form. So, this becomes $\rho = 0$ into $4\sigma^2 a^2 \omega^2 \sin^2 \theta + (1 - \rho)^2 = f^2$, so this equation shows.

So, for with the external amplitude equal to 0, so this equation again reduce to $4\sigma^2 a^2 \omega^2 \sin^2 \theta + (1 - \rho)^2 = 0$, so this is the equation of a circle with center at 1 and 0 so by. So, for the case of external forcing equal to 0, so one can find one can find this frequency equation, so for $f \neq 0$. So, this represent the frequency response curve, so from these one can plot the phase portrait of the system. So, this shows the response to be periodic as it is represent the equation of a circle with this one equation of a circle with center 1 and 0, so this is circular or in this case one can get some harmonics, so one can solve the equation directly by using Runge-Kutta method or use some numerical techniques to find the solution also and one can compare the resulting solution with the solution.

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So, for example, using the Runge-Kutta method with taking different initial conditions, so one can see the response to be periodic or one can obtain a limit cycle, if one consider this equation that is that is $d^2 u$ by $d t^2$ plus $\omega_0^2 u$ plus $\epsilon(u^2 - 1) \frac{du}{dt} = 0$ by substituting u equal to 0 and $\frac{du}{dt}$ equal minus 0.2 and ϵ equal to 0.5, so one can get this plot this phase portrait, so this phase portrait clearly show the existence of the limit cycle limit cycle means. So, if you start from interior point, let this is the interior point, so one can see that finally, it will trace this curve and finally, it will come to this limit cycle, so by taking different initial points, so finally, one come to this limit cycle similarly by taking a point outside of this limit cycle also one can land of with this limit cycle, so by using this Van der pol equation. So, in this case by using this equation one can find the solution of the system, so another method that is this k b method.

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The Krylov-Bogoliubov Technique

$$\frac{d^2 u}{dt^2} + \omega_0^2 u = \varepsilon f(u, \dot{u}) \quad \text{--- (1)}$$

$$u = a(t) \cos[\omega_0 t + \beta_0(t)] \quad \text{--- (2)}$$

$$\phi = \omega_0 t + \beta_0(t) \quad \text{--- (3)}$$

Subject to condition

$$\frac{du}{dt} = -\omega_0 a(t) \sin \phi \quad \text{--- (4)}$$

$$\frac{du}{dt} = -a\omega_0 \sin \phi + \frac{da}{dt} \cos \phi - a \frac{d\beta}{dt} \sin \phi \quad \text{--- (5)}$$

That is Krylov and Bogoliubov technique also will study now to find the equation find the solutions of the non-linear systems, so in case of this method or in this technique. So, let us consider a general equation. So, where $\frac{d^2 u}{dt^2} + \omega_0^2 u = \varepsilon f(u, \dot{u})$. So, in this case will assume the solution similar to the previous case, in the previous case we have taken two terms, in this case also we are taking two terms, but we are writing the solution in a different way; that means, we can write the solution equal to $u = a(t) \cos[\omega_0 t + \beta_0(t)]$ in the previous case that is in the Van der pol equation.

We have assume the solution of the form $u = a(t) \cos[\omega_0 t + \beta_0(t)]$ or $u = a(t) \sin[\omega_0 t + \beta_0(t)]$, but one can use this phase difference to show or write this equation in this form also, so by writing the equation or assuming the equation in this form that is $u = a(t) \cos[\omega_0 t + \beta_0(t)]$. So, here we are assuming that when this epsilon equal to small or when the epsilon equal to 0 the solution is $u = a \cos[\omega_0 t + \beta_0]$, where this a and beta are constant; that means, for small value for 0 value of epsilon the solution equal to $a \cos[\omega_0 t + \beta_0]$. So, here this a and beta are constant, but when you are considering the terms other than 0 for this epsilon; that means, epsilon is assume to be small, in this case one can assume the solution to be this same as that of when epsilon equal to 0, but here a and beta are the function of time.

So, here we are assuming a and β are slowly varying function of time. So, one can take this $\omega_0 t + \beta_0 t$ as ϕ , so one can write this u equal to $a \cos \phi$ so, where a and β are function of time. So, in this case Krylov and Bogoliubov propose that, so this will be the solution; that means, u will be the solution $a t \cos \omega_0 t + \beta_0 t$ subjected to the condition that is this $\frac{d u}{d t}$ or \dot{u} equal to, so \dot{u} will be equal to $-\omega_0 a t \sin \phi$ so; that means, this is differentiation of this thing is assume in this form that is $\omega_0 a t \sin \phi$, because it is assume that this \dot{a} equal to 0. So, the differentiation of this thing will be.

So now, differentiating this equation one can write this $\frac{d u}{d t}$ equal to $-\omega_0 a \sin \phi + \dot{a} t \cos \phi$, so $-\omega_0 a \sin \phi$, so as we have $a t$, so u equal $a t$ which is a is a function of time. Similarly this \cos inside is a function of time and this β_0 is a function of time, so we have this three terms. But Krylov and Bogoliubov, so they have made this assumption that this equation will be valid or this equation with the solution of the original equation that is $\frac{d^2 u}{d t^2} + \omega_0^2 u = \epsilon f(u, \dot{u})$ subjected to the condition that this \dot{u} will be equal to $-\omega_0 a \sin \phi$, so comparing this two equation.

So, let this is equation 2, so this is equation 3 and this equation 4 and 5, so comparing equation 4 and 5, so we can write we can write these two terms should be equal to 0, so here we have putting this condition that is $\frac{d u}{d t}$ equal to $-\omega_0 a t \sin \phi$, so to satisfy this equation that is u equal to $a t \cos \omega_0 t + \beta_0 t$, so the differentiation which comes to be $-\omega_0 a \sin \phi + \dot{a} t \cos \phi - a \dot{\beta} \sin \phi$, so comparing this 4 and 5. So, we can take this term to be equal to 0 so; that means, we can take this $\dot{a} t \cos \phi - a \dot{\beta} \sin \phi$ equal to 0, so now, differentiating this \dot{u} equal to $-\omega_0 a t \sin \phi$.

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$$\frac{da}{dt} \cos \phi - a \frac{d\beta}{dt} \sin \phi = 0$$

Differentiating $\dot{u} = -\omega_0 a(t) \sin \phi$

$$\frac{d^2 u}{dt^2} = -a\omega_0^2 \cos \phi - \omega_0 \frac{da}{dt} \sin \phi - a\omega_0 \frac{d\beta}{dt} \cos \phi \quad \checkmark$$

$$\frac{d^2 u}{dt^2} + \omega_0^2 u = \varepsilon f(u, \dot{u}) \quad \checkmark$$

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So, one can write, so this becomes, so a is a function of time, so if you differentiate it, so this $d^2 u$ by dt^2 will be equal to $a\omega_0^2 \cos \phi$. So, $a\omega_0^2 \cos \phi$ minus $a\omega_0^2 \cos \phi$ into minus ω_0 into $\frac{da}{dt} \sin \phi$ minus $a\omega_0$ into $\frac{d\beta}{dt} \cos \phi$, so substituting this expression for this $d^2 u$ by dt^2 and this \dot{u} in this equation that is $d^2 u$ by dt^2 plus $\omega_0^2 u$ equal to $\varepsilon f(u, \dot{u})$.

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$$\frac{d^2 u}{dt^2} + \omega_0^2 u = \varepsilon f(u, \dot{u})$$

$$\cancel{-a\omega_0^2 \cos \phi} - \omega_0 \frac{da}{dt} \sin \phi - a\omega_0 \frac{d\beta}{dt} \cos \phi + \omega_0^2 u \cancel{a \cos \phi} = -\varepsilon f(\underline{a \cos \phi}, \underline{-\omega_0 a \sin \phi}) - \delta$$

$$-\omega_0 \frac{da}{dt} \sin \phi - a\omega_0 \frac{d\beta}{dt} \cos \phi = -\varepsilon f(a \cos \phi, -\omega_0 a \sin \phi) - \delta \quad (8)$$

$$\frac{da}{dt} \cos \phi - a \frac{d\beta}{dt} \sin \phi = 0 \quad \text{--- (9)}$$

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So, one can write this equation; that means, $d^2 u / dt^2$ equal to $-\omega_0^2 u \cos \phi - \omega_0 d a / dt \sin \phi - \omega_0 d b / dt \cos \phi + \omega_0^2 u a \cos \phi$ equal to $-\epsilon f$. So, for u we have substituted $a \cos \phi$ and for \dot{u} we have substituted $-\omega_0 a \sin \phi$. So, this is the point to be noted that this u is substituted by $a \cos \phi$ and \dot{u} equal to $-\omega_0 a \sin \phi$, so now taking this equation, so we can write this by separating the coefficient of \cos and \sin , so we can write, so from this equation, so we have two equation. So, this is one equation and previously we have this equation $d^2 u / dt^2 + \omega_0^2 u$ equal to ϵf , so in this equation we have substituted the expression for u square, but this is one equation $d a / dt \cos \phi$. So, let us put the number, so this becomes equation 6 and substituting equation 7 in equation 1.

So, we can obtain this equation, so that is this is equation number 8, so in this equation number 8. We can write or we can rearrange this thing and we can write this equation in this form that is $-\omega_0 d a / dt \sin \phi$, so this the term $-\omega_0 a$. So, in this case $-\omega_0^2 u \cos \phi$, so this is plus and this is minus they cancel. So, this equation reduce to this form that is $-\omega_0 d a / dt \sin \phi - \omega_0 d b / dt \cos \phi$, so this will be equal to $-\epsilon f a \cos \phi - \omega_0 a \sin \phi$, so this is equation 9 and again we have written this equation 6 here, so this is equation 6, so from this equation 9 and 6 we can find the equation with $d a / dt$ and $d b / dt$ to obtain that thing.

Now, by multiplying this first equation by $\cos \phi$ and the second equation by $-\omega_0 \sin \phi$ and adding these two equation will get the expression for $d b / dt$. Similarly by eliminating this two terms; that means, by multiplying this equation by $\sin \phi$ and the second equation by $\omega_0 \cos \phi$, so we can eliminate these two terms and we can get the expression in terms of $d a / dt$.

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$$\begin{aligned}\frac{da}{dt} &= -\frac{\varepsilon}{\omega_0} \sin \phi f(a \cos \phi, -\omega_0 a \sin \phi) \\ \frac{d\beta}{dt} &= -\frac{\varepsilon}{a\omega_0} \cos \phi f(a \cos \phi, -\omega_0 a \sin \phi) \\ \frac{da}{dt} &= -\frac{\varepsilon}{2\omega_0} \left[\frac{2}{T} \int_0^T \sin \phi f(a \cos \phi, -\omega_0 a \sin \phi) dt \right] \\ \frac{d\beta}{dt} &= -\frac{\varepsilon}{2a\omega_0} \left[\frac{2}{T} \int_0^T \cos \phi f(a \cos \phi, -\omega_0 a \sin \phi) dt \right] \\ T &= \frac{2\pi}{\omega_0} \quad \omega_0 = \frac{2\pi}{T}\end{aligned}$$

So, we can write this equation now $\frac{da}{dt}$ equal to minus epsilon by omega 0 sin phi, so function of $a \cos \phi$ and minus omega 0 $a \sin \phi$. Similarly this $\frac{d\beta}{dt}$ equal to minus epsilon by $a \omega_0$ into cos phi into function of $a \cos \phi$ minus omega 0 $a \sin \phi$ as already we have pointed out this a and β are slowly varying function of time and this right side one can observe these are periodic as we have the sin phi and cos phi terms, so both these are periodic terms, so these are periodic with a period of 2π , so one can write this T or this T equal to 2π by omega 0, so 2π by omega 0 one know this omega 0 equal to 2π by T , so as these are slowly varying function of slowly varying function this a and β , so one can obtain or one can see or one can find that this value will not vary mass in a cycle.

So, averaging this value over a cycle, so one can write this equation $\frac{da}{dt}$ equal to minus epsilon by $2\omega_0$ $\frac{2}{T}$ by T integration 0 to T sin phi into $f(a \cos \phi, -\omega_0 a \sin \phi) dt$ or the second equation by averaging over the cycle one can write this $\frac{d\beta}{dt}$ equal to minus epsilon by $a \omega_0$, so this becomes $\frac{2}{T}$ by T integration 0 to T cos phi $f(a \cos \phi, -\omega_0 a \sin \phi)$. So, it can be noted that it is integrated over 0 to T and just divided by the time t , so this 2 2 cancels, so this is just written by integrating this equation and dividing this by the term t , similarly also this equation is written by integrating the right hand side and dividing that thing by t . So, now, one can express this time in terms of this frequency omega 0 and one can write this equation.

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$$T = \frac{2\pi}{\omega_0}$$

$$\frac{da}{dt} = -\frac{\varepsilon}{2\omega_0} \left[\frac{1}{\pi} \int_0^{2\pi} \sin \phi f(a \cos \phi, -\omega_0 a \sin \phi) d\phi \right]$$

$$\frac{d\beta}{dt} = -\frac{\varepsilon}{2a\omega_0} \left[\frac{1}{\pi} \int_0^{2\pi} \cos \phi f(a \cos \phi, -\omega_0 a \sin \phi) d\phi \right]$$

This $\frac{da}{dt}$ equal to minus epsilon by 2 omega 0, so by substituting this T equal 2 pi by omega 0, so one can write this equation, so this $\frac{da}{dt}$ will be equal to minus epsilon by 2 omega 0 into 1 by pi 0 to 2 pi sin phi f a cos phi minus omega 0 a sin phi d phi. Similarly this $\frac{d\beta}{dt}$ will be equal to minus epsilon by 2 a omega 0 into 1 by pi integration 0 to 2 pi cos phi into f a cos phi minus omega 0 a sin phi d phi.

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$$\frac{da}{dt} = -\frac{\varepsilon}{2\omega_0} f_1$$

$$\frac{d\beta}{dt} = -\frac{\varepsilon}{2a\omega_0} f_2$$

$$f_1 = \left[\frac{1}{\pi} \int_0^{2\pi} \sin \phi f(a \cos \phi, -\omega_0 a \sin \phi) d\phi \right]$$

$$f_2 = \left[\frac{1}{\pi} \int_0^{2\pi} \cos \phi f(a \cos \phi, -\omega_0 a \sin \phi) d\phi \right]$$

So, one can write this equation again in this form that is $\frac{da}{dt}$ equal to minus epsilon by 2 omega 0 into f 1 and this $\frac{d\beta}{dt}$ equal to minus epsilon by 2 a omega 0 into f

2, where this f_1 equal to $\frac{1}{\pi} \int_0^{2\pi} \sin \phi f_a \cos \phi \sin \phi d\phi$, and this f_2 equal to $\frac{1}{\pi} \int_0^{2\pi} \cos \phi f_a \cos \phi \sin \phi d\phi$, so by taking these functions are by evaluating these function, one can find this $\frac{da}{dt}$ and $\frac{d\beta}{dt}$, and for steady state one can substitute this $\frac{da}{dt}$ equal to 0 and $\frac{d\beta}{dt}$ equal to 0 and one can obtain the frequency equation. So, one can take a simple example for example, let us take the simple linear system.

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Handwritten notes on a yellow background showing the derivation of the frequency equation for a damped harmonic oscillator. The notes include the equation of motion, the method of averaging, and the final steady-state solution.

$$\text{KB Ex } \ddot{u} + \omega_0^2 u + 2\epsilon \mu \dot{u} = 0$$

$$\ddot{u} + \omega_0^2 u = -2\epsilon \mu \dot{u} \quad \epsilon f$$

$$f = -2\mu \dot{u}$$

$$\dot{a} = -\frac{\epsilon \mu}{\pi} \int_0^{2\pi} \sin^2 \phi d\phi = -\epsilon \mu$$

$$\dot{\beta} = -\frac{\epsilon \mu}{\pi} \int_0^{2\pi} \sin \phi \cos \phi d\phi = 0$$

$$a = a_0 \exp(-\epsilon \mu t)$$

$$u = a_0 \exp(-\epsilon \mu t) \cos(\omega_0 t + \beta_0)$$

Diagram of a mass-spring-damper system: A mass m is connected to a fixed support by a spring with constant k and a damper with coefficient μ in parallel.

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So, in this linear system linear damping, let us take $\ddot{u} + \omega_0^2 u + 2\epsilon \mu \dot{u} = 0$, so this is the equation for simple harmonic for a simple spring mass damper system, so this is not a non-linear system. But let us see how we can use this method of averaging for this case, so one can write the equation of motion in this form that is $\ddot{u} + \omega_0^2 u + 2\epsilon \mu \dot{u} = 0$, where this is equal to $\epsilon \mu \dot{u} = 2\omega_0^2 u$ into $\zeta \dot{u}$. So, now one can write this equation in this form $\ddot{u} + \omega_0^2 u = -2\epsilon \mu \dot{u}$, so our f equal to $-2\epsilon \mu \dot{u}$, so this is ϵf , so f equal to this, so this by substituting this f equal in that form, so one can write this \dot{a} equal to $-\epsilon \mu a$ by π integration $\int_0^{2\pi} \sin^2 \phi d\phi$, so this becomes $-\epsilon \mu a$.

Similarly, $\dot{\beta}$ one can write equal to $-\epsilon \mu \beta$ by π integration $\int_0^{2\pi} \sin \phi \cos \phi d\phi$, so this becomes as $\sin \phi$ integration $\int_0^{2\pi} \sin \phi \cos \phi d\phi = 0$.

becomes 0, so one can find this thing, so as a dot equal to minus epsilon mu a, so one can write this a equal to a 0 e to the power, so e to the power minus epsilon mu t and this beta dot equal to 0, so beta will be a constant, so beta will be equal to beta 0, so beta will be beta 0 as beta dot is a constant. Similarly a dot one can obtain, so the original the solution becomes a 0 e to the power minus epsilon mu t into cos omega 0 t plus beta 0, so plus order of epsilon. So, this way one can use this method of averaging or this K B method. So, here this example shows this K B method how this Krylov-Bogoliubov method can be use to apply for this equation. So, next class we will study or we will take few more examples to study this K B method and also I will tell you about this generalized method of averaging and other methods of averaging also.

Thank you.