

# Ideal Fluid Flows Using Complex Analysis

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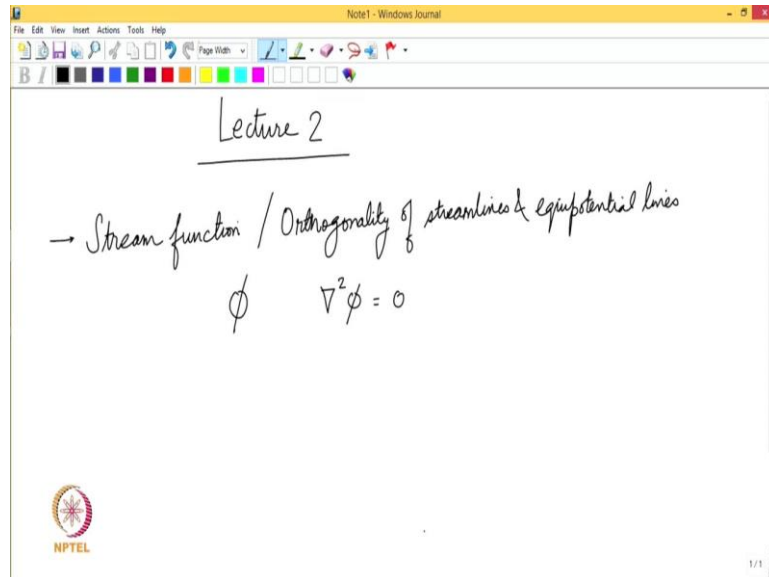
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## Lecture 2

### Stream Function, Orthogonality of streamlines and equipotential lines

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Welcome to this second lecture on ideal fluid flows using complex analysis. Now, recall that in the previous lecture, we looked at the concept of the velocity potential and how the velocity potential satisfies the Laplace equation. So, we have defined a variable or a function  $\phi$ , which is governed by the equation  $\nabla^2 \phi = 0$ . So, in this lecture, I will talk about a second complimentary function which is stream function, and then I will discuss the concept of orthogonality of streamlines and equipotential lines.

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→ Stream function / Orthogonality of streamlines

$$\phi \quad \nabla^2 \phi = 0$$

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Consider 2D flow, irrotational  $\vec{u} = u\hat{i} + v\hat{j} = (u, v)$

$$\nabla \cdot \vec{u} = 0 \quad \frac{\partial u}{\partial x} + \frac{\partial v}{\partial y} = 0$$

$\psi \equiv$  Stream Function

$$u = \frac{\partial \psi}{\partial y}, \quad v = -\frac{\partial \psi}{\partial x}$$

If  $\psi$  is twice differentiable,

So, to begin, let us consider a 2D flow and let us say this flow is irrotational. Now, in this 2D flow, we can write the velocity vector  $\mathbf{u}$  to be either say  $u \hat{i} + v \hat{j}$  or I would use a shorthand notation where I would write this as  $u, v$ . Now, for an ideal fluid, we know that this flow field must satisfy the continuity equation, which was  $\nabla \cdot \mathbf{u} = 0$ , which if you use in a vector sense will be  $\frac{du}{dx} + \frac{dv}{dy} = 0$ .

Now, what I will do is I will define a new function, which I am going to call as  $\psi$  and I am going to call it the stream function so, we will define this new function  $\psi$  which is a stream function, but it is given in a very specific form it is given in the following manner say I define  $u$  to be  $\frac{d\psi}{dy}$  and  $v$  to be minus  $\frac{d\psi}{dx}$ .

Now, if the function  $\psi$  is twice differentiable, so, we can say if  $\psi$  is twice differentiable which would mean we can take 2 derivatives especially the mixed type with this function then we can see that if I plug in these notations now, into the continuity equation.

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$$\frac{\partial}{\partial x} \left( \frac{\partial \psi}{\partial y} \right) + \frac{\partial}{\partial y} \left( -\frac{\partial \psi}{\partial x} \right) = \frac{\partial^2 \psi}{\partial x \partial y} - \frac{\partial^2 \psi}{\partial y \partial x} = 0$$

$$\text{Vorticity } \vec{\omega} = (\xi, \eta, \zeta) = \xi \hat{i} + \eta \hat{j} + \zeta \hat{k}$$

$$\text{In 2D flows: } \zeta = \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}$$

$$\text{If Flow is IRROTATIONAL, } \Rightarrow \zeta = 0$$

$$\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} = 0 \quad u = \frac{\partial \psi}{\partial y} \quad ; \quad v = -\frac{\partial \psi}{\partial x}$$

What we will get now will be  $dy$  by  $dx$  of  $d\psi$  plus  $dx$  by  $dy$  of minus  $d\psi$  which would be  $d^2\psi \, dx \, dy$  minus  $d^2\psi \, dy \, dx$  which is 0. So, this function  $\psi$  or this formulation of taking velocities components in terms of the function  $\psi$  which is a stream function is such that in the 2D flow it automatically satisfies the continuity equation.

Now, though this function  $\psi$  has nothing to do with flow rotationality or irrotationality so, this function  $\psi$  would also be valid for a rotational flow in 2D flows. But for the purpose of this course, because we are dealing with ideal fluids, we will be stressing again and again on irrotational flows. And so, our focus will be on deriving properties of this function  $\psi$  for a irrotational flows only. So, we will limit ourselves to the usage of this function  $\psi$  for the case of irrotational flows.

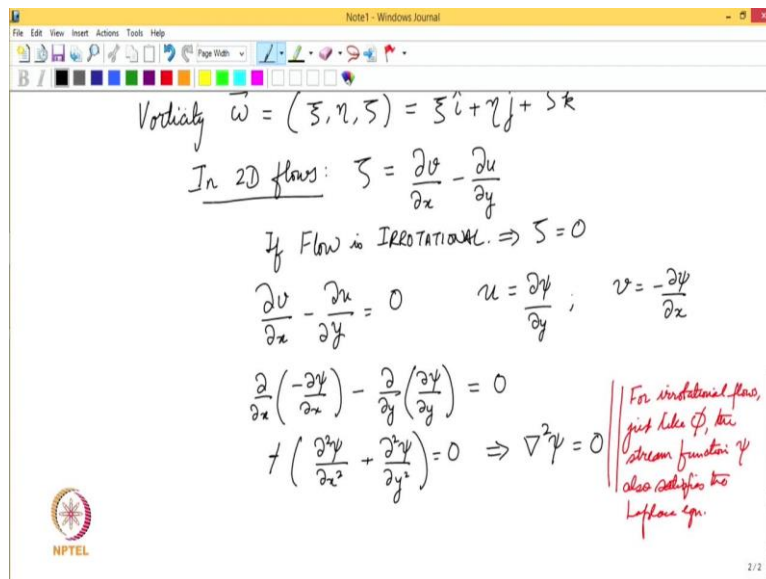
Now, recall that vorticity which I denoted by the symbol  $\omega$ ,  $\omega$  vector which would have three components. Let us say the 3 components are the following. We have  $\xi$ ,  $\eta$ , and  $\zeta$  as the 3 components of the vorticity vector, which I could write as for instance  $\xi \hat{i} + \eta \hat{j} + \zeta \hat{k}$ .

Now, the only non-zero component of this vorticity function or this vorticity vector in 2D coordinates or in 2D flows it is easy for you to work it out and I am sure you would have also covered this as part of your undergraduate fluid mechanics that in 2D flows, the only non-zero component or the only the single component that would be possible to derive would be  $\zeta$ , which would be given as  $\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}$ .

Further, if I was to say that the flow is irrotation, so, going beyond even the definition, if I say that the flow is irrotational what it means is that zeta should also be 0 that is the only way the flow can be irrotational what it means is  $dv dx - du dy$  is 0.

Now, what I am going to do is, I am going to substitute the definition of the velocity components  $u$  and  $v$  in terms of the stream function. So, we will say  $u$  is  $d\psi/dy$  and  $v$  was  $-d\psi/dx$ .

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Handwritten derivation in a Notepad window:

Vorticity  $\vec{\omega} = (\zeta, \eta, \zeta) = \zeta \hat{i} + \eta \hat{j} + \zeta \hat{k}$

In 2D flows:  $\zeta = \frac{\partial v}{\partial x} - \frac{\partial u}{\partial y}$

If Flow is IRROTATIONAL  $\Rightarrow \zeta = 0$

$\frac{\partial v}{\partial x} - \frac{\partial u}{\partial y} = 0$        $u = \frac{\partial \psi}{\partial y}$  ;       $v = -\frac{\partial \psi}{\partial x}$

$\frac{\partial}{\partial x} \left( -\frac{\partial \psi}{\partial x} \right) - \frac{\partial}{\partial y} \left( \frac{\partial \psi}{\partial y} \right) = 0$

$-\left( \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} \right) = 0 \Rightarrow \nabla^2 \psi = 0$

For irrotational flow, just like  $\phi$ , the stream function  $\psi$  also satisfies the Laplace eqn.

So, if you plug it in into the definition of zeta we get  $d$  by  $dx$  of minus  $d\psi/dx$  minus  $d$  by  $dy$  of  $d\psi/dy$  to be 0. And then if we take it one step further, we will get minus of  $d^2\psi/dx^2$  plus  $d^2\psi/dy^2$  is 0 when the negative sign is immaterial, so, let me just cancel it off. So, what we get is in 2D flows  $\nabla^2 \psi = 0$  if the flow is irrotational.

So, this very important result that we have derived now is that for irrotational flows just like the velocity potential the stream function also satisfies the Laplace equation just like  $\phi$  the stream function which is  $\psi$  also satisfies the Laplace equation and this is going to be again a very important tool for us in the following lectures.

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$$\nabla^2 \left( \frac{\partial^2 \psi}{\partial x^2} + \frac{\partial^2 \psi}{\partial y^2} \right) = 0 \Rightarrow \psi \text{ stream function } \psi \text{ also satisfies the Laplace eqn.}$$

Flow lines corresponding to  $\psi = \text{constant}$  are streamlines of the flow field.

Proof:  $d\psi = \frac{\partial \psi}{\partial x} dx + \frac{\partial \psi}{\partial y} dy = v dy - u dx$

If  $\psi = \text{const}$ ,  $d\psi = 0 \Rightarrow (v dy - u dx) \Big|_{\psi=c} = 0$

$\left( \frac{dy}{dx} \right) \Big|_{\psi=c} = \left( \frac{v}{u} \right) \quad \psi = \text{const} \rightarrow \text{Streamline}$

"STREAM FUNCTION"

Now, let me state one key fact here and then we will try and prove it using the definitions that we have come up with so far. The theory that I am going to put up is or the phi that I am going to put up are or is that that flow lines corresponding to a fixed value of psi. So, for instance a psi is some constant. So, flow lines corresponding to a given value of psi are streamlines of the flow field.

So, if there is a line over which psi is constant, then that line is going to be a streamline of the flow field. And we can easily prove this by the definitions that we have come up with, we could say d psi, the total change in psi could be written as d psi dx dx plus d psi dy dy. Now, note that d psi dx is minus v and d psi d y is u. So, we can write this as u dy minus v dx.

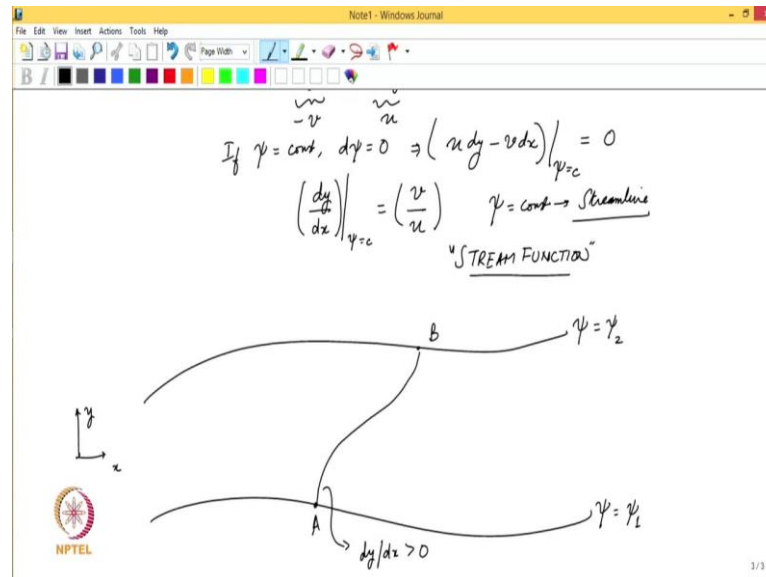
Now, let us consider the case that the stream function is a constant. So, if psi was a constant, then d psi must be 0 or the change of psi on that line must be 0, which gives us u dy minus v dx over this line psi equal to constant must be 0, which if we do a little bit of rearrangement, now, you can see that dy by dx for a constant psi is given as the ratio of the two velocity components. So, v by u.

Now, this is precisely how the streamline is defined, if you recall we said the streamline is an imaginary curve, the slope of which or the tangent to which gives us a direction of the velocity vector. So, it is this which is the slope of the streamline. And so, we can say that psi being a constant corresponds to a streamline.

Furthermore, each value of that we assign to a stream line each unique value means or pertains to a new stream line. So, hence, we get the name stream function. That is the whole

idea why this is called a stream function that this function can take a range of values. And for each value or each unique value, you get a unique streamline.

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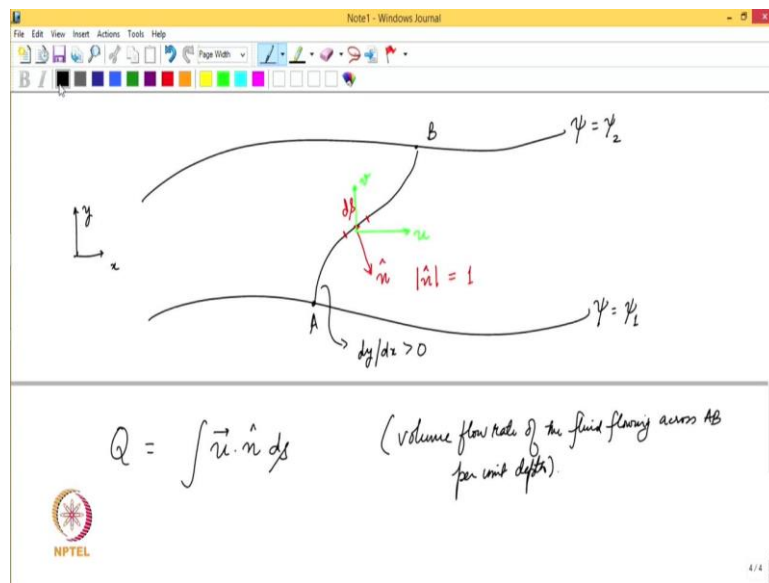
So, now, let us consider a scenario where we have 2 streamlines in a flow field, say  $\psi$  equal to  $\psi_1$  and  $\psi$  equal to  $\psi_2$ . And let us define a coordinate system as well. So, let us say we have  $x$  here  $y$  here. So this is a 2D flow field. Now, the question that I am going to pose is this is the following. I want to know the flow rate of the fluid flowing between these 2 streamlines.

Note that we had said earlier in the previous lecture, we said the dot product of the velocity field with the normal vector on a streamline is 0. So that means that there must be no mass flow rate across a streamline. So between these 2 streamlines, there must be some fluid that is flowing. And you want to quantify how much is that flow rate.

So what we will do is, say we take any 2 points A and B, A on streamline  $\psi_1$ , B on streamline  $\psi_2$ . And let me draw a convex surface connecting A and B. So say we take a surface which has the following shape. And the shape is chosen specifically so that the slope on this surface is always positive. That is just for convenience, we could have chosen some other surface as well. But just to make a point just to prove it quickly. Let me say this is what I am going to choose.

So it is a curve, this line AB is of arbitrary shape, but it is such that it has a positive slope. For all we care, it could even be a line connecting in a straight line. But for now, let us just take a curvilinear section.

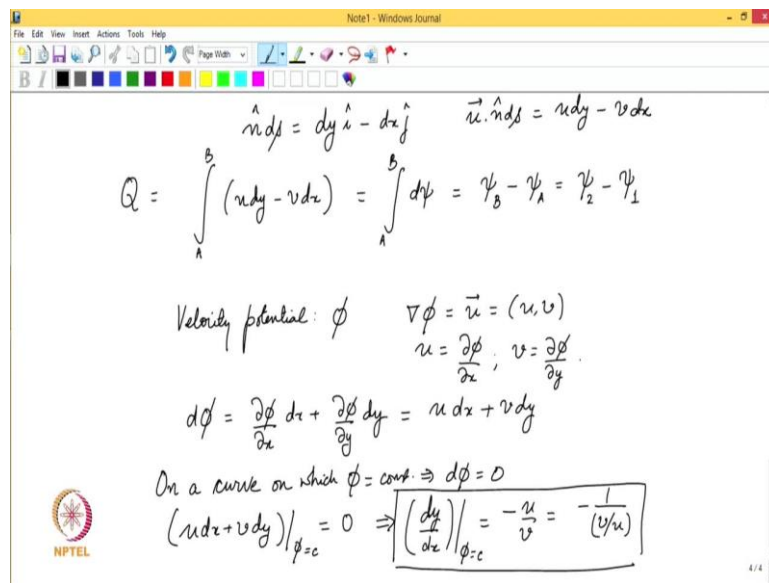
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What do you want to determine? We want to determine the following quantity we want to know what is the flow rate across this line AB and that would be given by this function Q and let me now show you what Q is. So basically, if we take a small element, and which I am going to represent by a large element here for only for the purpose of for visual clarity, if I take an element which has a length of  $ds$  along this curve, AB, say we are going from A to B, so anywhere the  $ds$  will be pointing from A to B, the normal vector to this point would be in this direction. And that is a unit normal, which means modulus of unit normal is 1.

If I can decompose the flow at this location, I could say, this is, for instance, the x component of velocity, this is the y component of velocity,  $u$  and  $v$ . I am interested in knowing what is the flow rate across this element, small element infinitesimal element  $ds$  and that too per unit depth. So, this is the volume flow rate of the fluid flowing across AB per unit depth and that is what we are trying to calculate.

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$$\hat{n} ds = dy \hat{i} - dx \hat{j} \quad \vec{u} \cdot \hat{n} ds = u dy - v dx$$

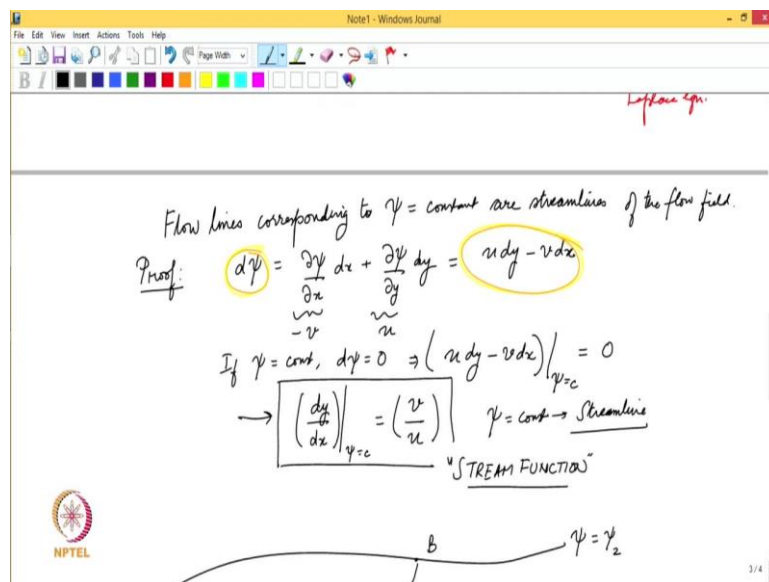
$$Q = \int_A^B (u dy - v dx) = \int_A^B d\psi = \psi_B - \psi_A = \psi_2 - \psi_1$$

Velocity potential:  $\phi \quad \nabla \phi = \vec{u} = (u, v)$   
 $u = \frac{\partial \phi}{\partial x}, \quad v = \frac{\partial \phi}{\partial y}$

$$d\phi = \frac{\partial \phi}{\partial x} dx + \frac{\partial \phi}{\partial y} dy = u dx + v dy$$

On a curve on which  $\phi = \text{const.} \Rightarrow d\phi = 0$

$$(u dx + v dy)|_{\phi=c} = 0 \Rightarrow \left( \frac{dy}{dx} \right) \Big|_{\phi=c} = -\frac{u}{v} = -\frac{1}{(v/u)}$$



Flow lines corresponding to  $\psi = \text{constant}$  are streamlines of the flow field.

Proof:  $d\psi = \frac{\partial \psi}{\partial x} dx + \frac{\partial \psi}{\partial y} dy = u dy - v dx$

If  $\psi = \text{const.}, d\psi = 0 \Rightarrow (u dy - v dx)|_{\psi=c} = 0$

$$\rightarrow \left( \frac{dy}{dx} \right) \Big|_{\psi=c} = \left( \frac{v}{u} \right) \quad \psi = \text{const.} \rightarrow \text{Streamline}$$

"STREAM FUNCTION"

$\psi = \psi_2$

Now, it can be easily proved and which I am going to not do here, but you can easily work it out. That the normal vector are in fact, the normal vector times ds can be easily written as dy i cap minus dx, j cap where dx and dy are the length when you project ds along x and y direction.

So, knowing this, let us calculate Q. So Q would be, now remember, it is u dot n ds. So clearly u dotted with n ds would now be u dy minus v dx. So, u dot n ds will be u dy minus v dx. And now it can integrate over the path A to B so, we will get u dy minus v dx as the integral from A to B and now I note that u dy minus v dx is the total change in the stream function.



This comes from the definition that we had written right here. So, if you see this, this is what we are using now, with using this definition of  $d\psi$ . So, we can plug this in and then we know that it is a total change in  $\psi$ . So, that must be  $\psi$  at B minus  $\psi$  at A which would be  $\psi_2$  at B minus  $\psi_1$  at A. And so, the total flow rate per unit depth between any 2 stream lines is the difference in the stream function values and that is the significance of the stream function. So, we can calculate the volume flowing between any 2 stream lines by the difference between the stream values or the stream function values.

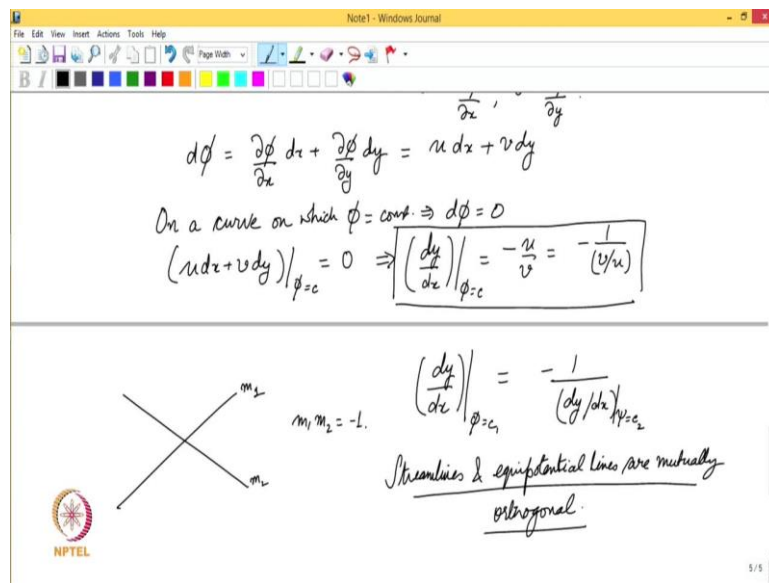
Now, let us consider the, we bring we go back to the velocity potential and let us look at a specific relationship between velocity potential and stream function. Recall it velocity potential was  $\phi$  and gradient of  $\phi$  was  $\mathbf{u}$  which is basically  $u, v$ . So, clearly  $u$  must be  $d\phi/dx$  and  $v$  must be  $d\phi/dy$  that is just from the definition.

So, let us consider the total change in the function  $\phi$ . So,  $d\phi$  would be  $d\phi/dx dx + d\phi/dy dy$  which we can now write as  $d\phi/dx = u$  I can write this as  $u dx + d\phi/dy = v$ , we can write this as  $v dy$ .

And we can now say that if I am on a curve on which  $\phi$  is a constant on a curve on which  $\phi$  is a constant would mean  $d\phi$  on that curve would be 0 what it means is? That  $u dx + v dy$  on this specific curve where  $\phi$  is a constant is 0 if there is a curve in space of that type, which means, if you rearrange again we can prove the  $dy/dx$  for this curve where  $\phi$  is a constant is minus  $u$  by  $v$ , which I could also write as minus  $1$  by  $v$  by  $u$ .

So, I want you to put this specific result in perspective from what we have done just a little while ago, which is to talk about the stream function. So, together with what we have done here, which is to prove that the slope of the stream function or slope of the stream line is given as a ratio of the  $y$  component of velocity to the  $x$  component of velocity, we now see that for a stream or for a potential function or for a potential flow problem or curve on which  $\phi$  is constant, the slope is given by the negative of the  $x$  component of velocity versus the  $y$  component of velocity.

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The image shows a handwritten derivation in a software window titled 'Note1 - Windows Journal'. At the top, the total differential of the velocity potential  $\phi$  is given as  $d\phi = \frac{\partial \phi}{\partial x} dx + \frac{\partial \phi}{\partial y} dy = u dx + v dy$ . Below this, it states that on a curve where  $\phi = \text{const.} \Rightarrow d\phi = 0$ . This leads to the equation  $(u dx + v dy)|_{\phi=c} = 0 \Rightarrow \left( \frac{dy}{dx} \right) \Big|_{\phi=c} = -\frac{u}{v} = -\frac{1}{(v/u)}$ . A horizontal line separates this from the next part. Below the line, a diagram shows two intersecting lines with slopes  $m_1$  and  $m_2$ , and the text  $m_1 m_2 = -1$ . To the right of the diagram, the slope  $\left( \frac{dy}{dx} \right) \Big|_{\phi=c_1}$  is equated to  $-\frac{1}{(dy/dx)|_{\psi=c_2}}$ . The final conclusion is underlined: 'Streamlines & equipotential lines are mutually orthogonal.' The NPTEL logo is visible in the bottom left corner of the journal window.

If you recall from coordinate geometry, if there are 2 lines across which let us say the 2 lines have slopes  $m_1$  and  $m_2$  the condition that these 2 lines are orthogonal to each other would be that  $m_1 m_2$  is minus 1. The same thing has happened here that if you look at the slope of the line at this specific location on a given value of velocity potential then that is negative 1 by negative inverse of the slope of the line which is given by the constant stream function.

What it means is that streamlines and equipotential lines by equipotential we mean lines where the potential is a constant are mutually orthogonal or are orthogonal to each other. So, this will also be very handy when we start drawing flow diagrams in the next few lectures. So, I will finish this lecture here. And I hope to see you in the next lecture. Thank you.