

Advanced Strength of Materials
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Lecture – 6

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Handwritten equations on a whiteboard:

$$\sigma_{ij}' \quad l_j \sigma_{ij}' - l_j \delta_{ij} \sigma' = 0 \quad \dots (1)$$

$l_j \rightarrow l, m, n$

$$l(\sigma_x' - \sigma') + m \tau_{xy} + n \tau_{xz} = 0$$

$$l \tau_{yx} + m(\sigma_y' - \sigma') + n \tau_{yz} = 0$$

$$l \tau_{zx} + m \tau_{zy} + n(\sigma_z' - \sigma') = 0$$

$$\det | \sigma_{ij}' - \delta_{ij} \sigma' | = 0 \quad \dots (2)$$

Last time we saw that the Deviatoric part of the stress tensor are σ_{ij}' . We will have their principal directions and also the associated stresses, which are going to be given by these set of equations. Where, the σ dash represents the principle stress and principle directions are given by l, m, n . And also we saw that the l, m, n has components like l, m and n . Then, we are going to have these equations in the expanded form. And if we have to have the non zero solutions for the directed cosines l, m and n , we must have determinant of the position matrix equal to 0.

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$$\det \left| \sigma_{ij}' - \delta_{ij} \sigma' \right| = 0 \quad \dots (3)^{(2)}$$

$$(3) \rightarrow \det \begin{vmatrix} \sigma_x' - \sigma' & \tau_{xy} & \tau_{xz} \\ \tau_{yx} & \sigma_y' - \sigma' & \tau_{yz} \\ \tau_{zx} & \tau_{zy} & \sigma_z' - \sigma' \end{vmatrix} = 0$$

This can be written in the expanded form like this, which is very similar to the earlier case, where we have considered that the total stress. And the principle stress is given by sigma. And these stress components were sigma x sigma y and sigma z, not if the dashes.
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$$\det \left| \sigma_{ij}' - \delta_{ij} \sigma' \right| = 0 \quad \dots (3)^{(2)}$$

$$(\sigma')^3 - I_1'(\sigma')^2 - I_2'(\sigma') - I_3' = 0 \quad \dots (4)$$

$$I_1' = 0 = \sigma_x' + \sigma_y' + \sigma_z'$$

$$= \sigma_x - \sigma_m + \sigma_y - \sigma_m + \sigma_z - \sigma_m$$

$$= 0$$

$$I_2' = \tau_{xy}^2 + \tau_{yz}^2 + \tau_{zx}^2 - (\sigma_x' \sigma_y' + \sigma_y' \sigma_z' + \sigma_z' \sigma_x')$$

Now, this can be expanded and if we expand it, comes out to be like this. Sigma dash cube minus I 1 dash into sigma dash square minus I 2 dash, sigma dash minus I 3 dash equal to 0. This is equation number 4, where in again I 1 dash, I 2 dash and I 3 dash are stress in variance. And these are given by I 1 dash is equal to 0, because this is equal to

σ_x dash plus σ_y dash plus σ_z dash. And this is nothing but σ_x minus σ_m plus σ_y minus σ_m plus σ_z minus σ_m .

Since, σ_m is equal to σ_x plus σ_y plus σ_z by 3. So, this is sum of these three stresses to cancel and therefore, this is going to be 0. Similarly, the other invariance as been formed as we have seen earlier is going to be τ_{xy} square plus τ_{yz} square, τ_{zx} square minus σ_x dash into σ_y dash plus σ_y dash into σ_z dash plus σ_z dash into σ_x dash.

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$$I_3' = \sigma_x' \sigma_y' \sigma_z' + 2\tau_{xy}\tau_{yz}\tau_{zx} - \sigma_x' \tau_{yz}^2 - \sigma_y' \tau_{zx}^2 - \sigma_z' \tau_{xy}^2$$

Roots of (4) $\rightarrow \sigma_1', \sigma_2', \sigma_3'$

$$\sigma_1 = \sigma_1' + \sigma_m \quad \sigma_2 = \sigma_2' + \sigma_m$$

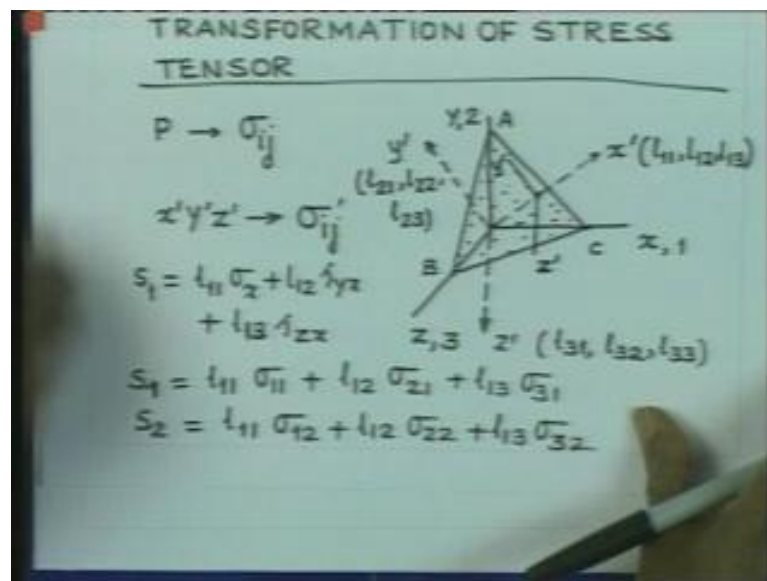
$$\sigma_3 = \sigma_3' + \sigma_m$$

That will be second stress invariance and we can now write the other component. I_3 dash σ_x dash, σ_y dash, σ_z dash plus twice τ_{xy} , τ_{yz} into τ_{zx} minus σ_x dash τ_{yz} square minus σ_y dash, τ_{zx} square minus σ_z dash into τ_{xy} square. So, we can solve the equation number 4, using these coefficients I_1 dash, I_2 dash and I_3 dash. We are since, I_1 dash, we will get some advantage in solving this equation 4.

And finally, suppose these roots are σ_1 dash, σ_2 dash and σ_3 dash. Then, we can write that principle stresses for the total stress stanza is nothing but σ_1 dash plus σ_m . σ_2 equal to σ_2 dash plus σ_m and σ_3 equal to σ_3 plus σ_m . You can also, after having obtained this 3 Deviatoric principle stresses, you can get back to the equation number 2.

You can get back to this equation, therein; you can substitute the value of σ_1 dash. In these three equations, you can extract the roots l, m, n from this equation to get back the corresponding directed cosines. And it can be repeated for σ_2 dash and σ_3 dash. This is the way, how you can solve for the principle directions starting from the Deviatoric part of the stress components. Next, we will consider the transformation of stresses.

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So, we will consider the transformation of stress tensors or the transformation of stresses, from one system of coordinates to another system of coordinates. Let us consider that, we have three directions like x, y and z , which you can also write as direction 1, 2 and 3. Suppose, the stress at a certain point is given by the tensor σ_{ij} , you have nine components of stresses in three dimensions.

Now, we would like to transform these stresses to another coordinates. Let us say that, we have another set of coordinates, which is given by x dash, y dash and z dash. So, if the stresses are given in this system x, y, z . We would like to get the system in x dash, y dash, z dash, which are arbitrarily oriented with respect to the original system. Now, in the system x dash, y dash and z dash, we would like to write the stresses by this symbol σ_{ij} dash.

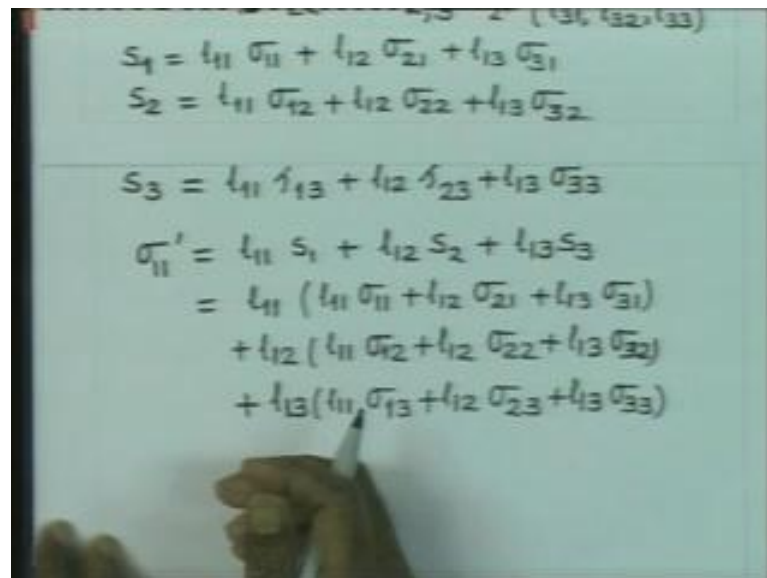
Now, let us see, how we can go about getting these stresses in the system x dash, y dash and z dash. Let us write, that the directed cosines of the new set of axis, say this x dash.

These have directed cosines like l_{11}, l_{12}, l_{13} . So, this is the directed cosines of the system x dash with respect to the x, y, z . So, also, the directed cosines of this axis y dash is l_{21}, l_{22}, l_{23} and the directed cosines of z dash is l_{31}, l_{32} and l_{33} .

Now, if you consider the plane perpendicular to x dash. Let us consider the plane, which is perpendicular to x dash. So, this is the plane A, B, C. So, this is the plane, which is perpendicular to x dash. Then, obviously, we can see that on this plane; will have the axis Y and Z line. So, therefore, this is the outer normal to the plane. Then, these two axis y dash is going to lie and on this plane, so also, z dash is also going to lie on this plane.

Now, we can write the fraction on this plane to be given by, let us say S_1 is given by l_{11} into σ_{11} plus l_{12} , τ_{yx} plus l_{13} , τ_{zx} . If we make use of the indicial notation, we can write now little in the simplest form $l_{11}, \sigma_{11}, l_{12}, \sigma_{21}$ plus l_{13}, σ_{31} . And the S_2 , the other traction s going to be nothing but $l_{11}, \sigma_{12}, l_{12}, \sigma_{22}$ plus l_{13}, σ_{32} .

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$$\begin{aligned}
 S_1 &= l_{11} \sigma_{11} + l_{12} \sigma_{21} + l_{13} \sigma_{31} \\
 S_2 &= l_{11} \sigma_{12} + l_{12} \sigma_{22} + l_{13} \sigma_{32} \\
 S_3 &= l_{11} \sigma_{13} + l_{12} \sigma_{23} + l_{13} \sigma_{33} \\
 \sigma'_{11} &= l_{11} S_1 + l_{12} S_2 + l_{13} S_3 \\
 &= l_{11} (l_{11} \sigma_{11} + l_{12} \sigma_{21} + l_{13} \sigma_{31}) \\
 &\quad + l_{12} (l_{11} \sigma_{12} + l_{12} \sigma_{22} + l_{13} \sigma_{32}) \\
 &\quad + l_{13} (l_{11} \sigma_{13} + l_{12} \sigma_{23} + l_{13} \sigma_{33})
 \end{aligned}$$

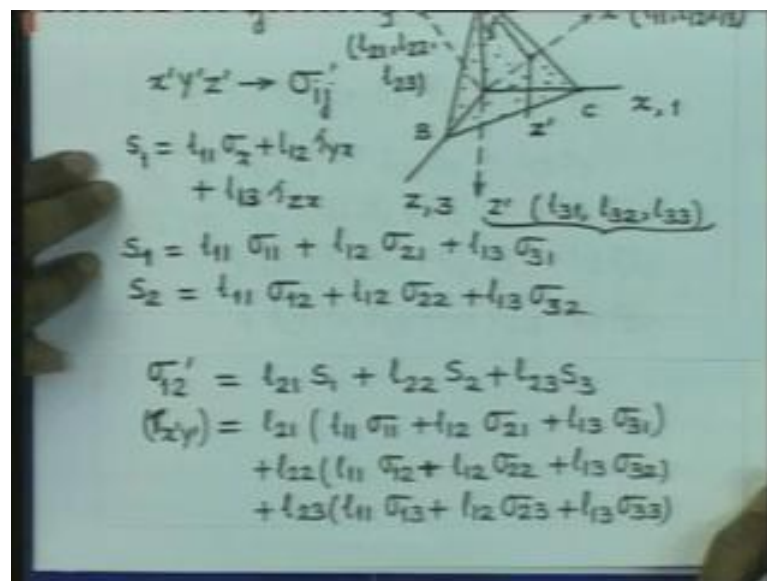
Similarly, we will have traction S_3 , l_{11} , τ_{13} , l_{12} , τ_{23} , plus l_{13} σ_{33} . So, these are the tractions on this plane. So, we are going to have the tractions on this plane to the x direction as X is S_1 , Y direction as S_2 and Z direction as S_3 . Now, we can get the components of the tractions, along the x dash. So, if we try to consider the

components of these tractions in x dash direction. That will give us the normal stress or that will give us the stress in the x dash direction.

So, therefore, we are going to get now, the stresses, sigma 11 dash is nothing but it is l 11, S 1 plus l 12 into S 2 plus l 13 into S 3. So, if you substitute the values of the three tractions. We are going to get now, l 11 sigma 11, l 12, sigma 21 plus l 13, sigma 31 plus l 12, l 11 sigma 12, l 12 sigma 22 plus l 13 sigma 32, plus, we are going to have l 13, l 11, sigma 13, l 12 sigma 23, l 13 sigma 33.

So, this is what is going to give us the expression for the normal stress on the plane, which is perpendicular to new axis, x dash. Similarly, we can get now, after having got the stress in this direction by taking the components of the three tractions. In this direction, we got sigma 11 dash. Similarly, if we take the components of those traction in the direction Y dash, that will give the traction sigma 12 dash.

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$$x'y'z' \rightarrow \sigma'_{ij}$$

$$S_1 = l_{11}\sigma_{11} + l_{12}\sigma_{21} + l_{13}\sigma_{31}$$

$$S_2 = l_{11}\sigma_{12} + l_{12}\sigma_{22} + l_{13}\sigma_{32}$$

$$S_3 = l_{11}\sigma_{13} + l_{12}\sigma_{23} + l_{13}\sigma_{33}$$

$$\sigma'_{11} = l_{11}S_1 + l_{12}S_2 + l_{13}S_3$$

$$(l_{21}) = l_{21}(l_{11}\sigma_{11} + l_{12}\sigma_{21} + l_{13}\sigma_{31}) + l_{22}(l_{11}\sigma_{12} + l_{12}\sigma_{22} + l_{13}\sigma_{32}) + l_{23}(l_{11}\sigma_{13} + l_{12}\sigma_{23} + l_{13}\sigma_{33})$$

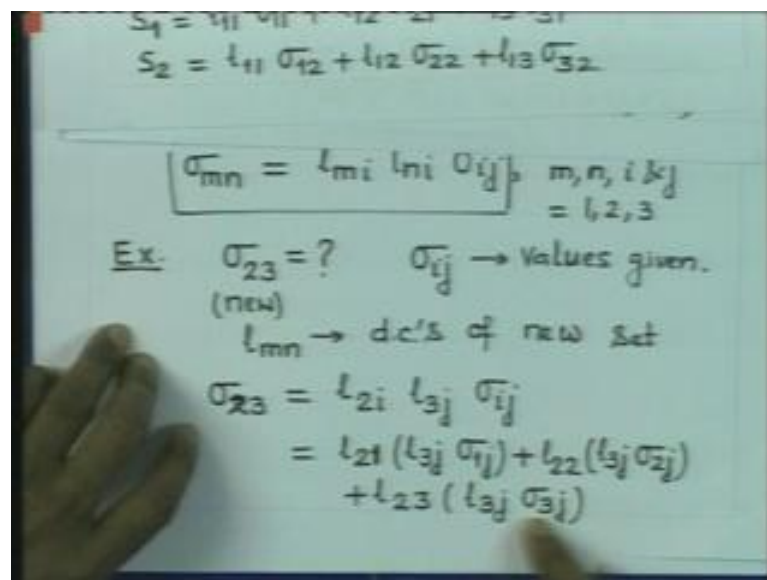
So, we can now write sigma 12 dash. So, for the convenience, I can probably write here, this is nothing but sigma x dash and this nothing but tau x dash y dash. That is the stress and this is going to be given by, will now make use of the direction cosine of this y dash which is nothing but these set.

So, we will now write, that the component sigma 12 dash is nothing but l 21, S 1 plus l 22, S 2 plus l 23, S 3. And again, if you substitute the value of the three tractions, that

you have, what we find is, nothing but the expression $l_{21}, l_{11}, \sigma_{11}$ plus l_{12}, σ_{12} plus l_{13}, σ_{13} plus l_{22}, σ_{22} plus l_{23}, σ_{23} plus l_{32}, σ_{32} plus l_{33}, σ_{33} .

So, we can get now the stress components σ_{12} dash. And you should have no difficulty in appreciating the fact that, if you want to find out the component of the stress in the z dash direction. We can make in the z dash direction, we can make use of the direction cosine given by this set. So, this will give us the component σ_{13} dash. Now, all these can be written very conveniently by making use of the tensor notation and it will give us finally, the total of all these components that we are going to write.

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$$S_1 = l_{11}\sigma_{11} + l_{12}\sigma_{12} + l_{13}\sigma_{13}$$

$$S_2 = l_{21}\sigma_{11} + l_{22}\sigma_{12} + l_{23}\sigma_{13}$$

$$\sigma_{mn} = l_{mi} l_{ni} \sigma_{ij} \quad m, n, i, j = 1, 2, 3$$

Ex: $\sigma_{23} = ?$ $\sigma_{ij} \rightarrow$ values given.

(new) $l_{mn} \rightarrow$ d.c's of new set

$$\sigma_{23} = l_{2i} l_{3j} \sigma_{ij}$$

$$= l_{21}(l_{3j} \sigma_{1j}) + l_{22}(l_{3j} \sigma_{2j}) + l_{23}(l_{3j} \sigma_{3j})$$

We can write in this form, which is nothing but σ_{mn} . Suppose, these are the stresses in the new system, X dash, Y dash and Z dash. And we have the original set given in the σ_{ij} , let us say and it is given in the system x, y, z. Then, we can write the connectivity between the two set, using the earlier expressions in the compact form and the generic form like this. σ_{mn} is equal to $l_{mi}, l_{ni}, \sigma_{ij}$, where in all these m, n, so also i and j take of values 1, 2, 3.

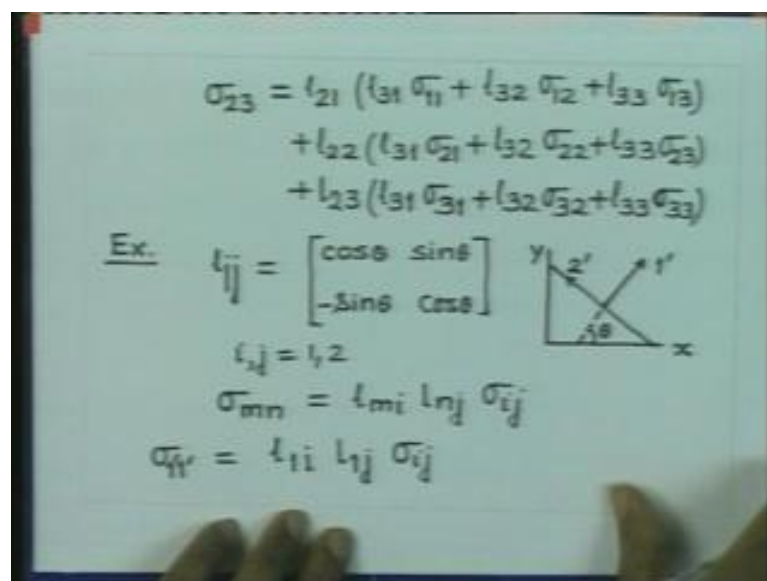
So, this is the transformation rule, for the second rank tensor. So, this rule for transforming the second rank tensor, σ_{mn} is equal l_{mi} into l_{ni} into σ_{ij} . Now, I would like to illustrate this by considering the examples later. But, if we want to

write any set, it will be like this. Suppose, if you are interested in writing. So, let us say will consider the example, we are interested in getting the new value of let us say, stress 23 and our original set is given by sigma 23. New is what, when these values are given.

So, again, we will make use of the fact that the new set will have the directed cosines given by l, m, n. Let us say, the direction cosines of the new set. So, I can now write the sigma 23 is equal to l_{2i} l_{3j} sigma_{ij}. Now, this can be further expanded like this. We will have values of i varying from 1, 2, 3, so also value of j varying from 1, 2, 3. So, this repetition of indices indicates that there is the summation implied.

So, we can now write, if i take the value of i to be first changing. So, therefore, I can write this thing as l_{2i}, I am sorry, l₂₁ into l_{3j} and this will become l₃₁ plus l₂₂, l_{3j}, sigma_{2j} plus l₂₃ into l_{3j} sigma_{3j}. So, you can link it up, that these are nothing but actually the tractions. These bracketed quantities are nothing but the tractions on the plane, which is perpendicular to the directions 2.

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$$\sigma_{23} = l_{21} (l_{31} \sigma_{11} + l_{32} \sigma_{12} + l_{33} \sigma_{13}) + l_{22} (l_{31} \sigma_{21} + l_{32} \sigma_{22} + l_{33} \sigma_{23}) + l_{23} (l_{31} \sigma_{31} + l_{32} \sigma_{32} + l_{33} \sigma_{33})$$

Ex. $l_{ij} = \begin{bmatrix} \cos \theta & \sin \theta \\ -\sin \theta & \cos \theta \end{bmatrix}$

$$l_{12} = l_{21}$$

$$\sigma_{mn} = l_{mi} l_{nj} \sigma_{ij}$$

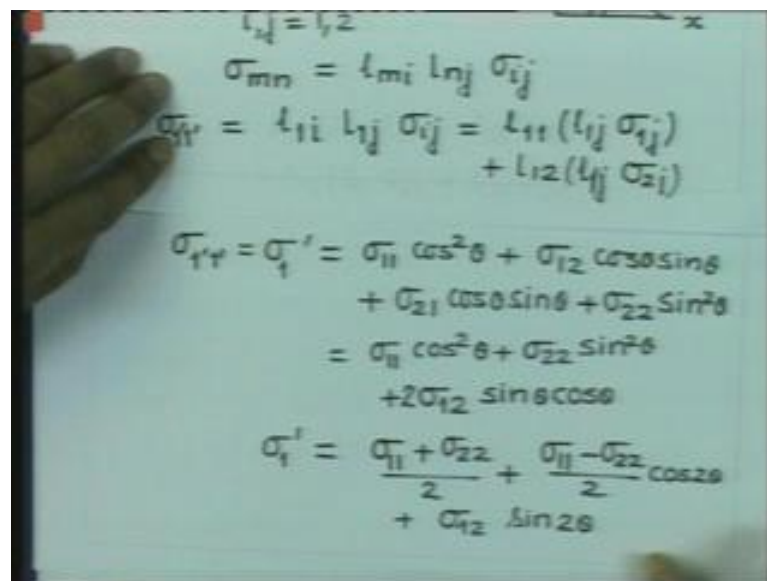
$$\sigma_{11'} = l_{1i} l_{1j} \sigma_{ij}$$

So, therefore, we can now expand it further. This is nothing but l₂₁ into now j can be taken as the given values 1, 2, 3. So, therefore, it will be l₃₁ sigma₁₁, l₃₂ sigma₁₂, plus l₃₃ sigma₁₃ plus l₂₂, l₃₁ sigma₂₁ plus l₃₂, sigma₂₂ plus l₃₃ sigma₂₃ plus l₂₃, l₃₁, sigma₃₁ plus l₃₂, sigma₃₂ plus l₃₃, sigma₃₃. So, you can go on writing the value for the stress components sigma 23. In the new system, from the old system, knowing the direction cosines of the new set of axis.

Now, we will like to consider another example to make you comfortable about the transformation. Let us consider the examples. Wherein, you are interested in considering the stresses in two dimensions. Let us say that, we have the stresses given in the system x, y and we will consider the new directions. Let us say the new directions are, this is the direction, new direction. And let us say, that this direction is 1 dash and this is the direction 2 dash. So, therefore, this direction 1 dash makes an angle θ with the x direction.

So, therefore, we have the direction cosines l_{ij} , l_{ij} is given by the direction cosines of this is nothing but $\cos \theta$ $\sin \theta$. And direction cosines of this direction 2 dash is $\sin \theta$, $\cos \theta$. So, therein, it is a two dimensional problem. So, therefore, i and j takes some value 1 to 2. Now, the transformation rule is going to be $\sigma_{mn} = l_{mi} l_{nj} \sigma_{ij}$. Let us find out the component σ_{11} . So, let us find out the component σ_{11} . So, σ_{11} dash, we indicate that with the σ_{11} dash. This is nothing but $l_{11} l_{11} \sigma_{11} + l_{12} l_{12} \sigma_{22} + 2 l_{11} l_{12} \sigma_{12}$.

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The image shows a handwritten derivation of the stress transformation equation for $\sigma_{1'}$. The equations are as follows:

$$l_{ij} = l_{ji}$$

$$\sigma_{mn} = l_{mi} l_{nj} \sigma_{ij}$$

$$\sigma_{1'1'} = l_{1'i} l_{1'j} \sigma_{ij} = l_{1'1} (l_{1'j} \sigma_{1j}) + l_{1'2} (l_{1'j} \sigma_{2j})$$

$$\sigma_{1'1'} = \sigma_{1'} = \sigma_{11} \cos^2 \theta + \sigma_{12} \cos \theta \sin \theta + \sigma_{21} \cos \theta \sin \theta + \sigma_{22} \sin^2 \theta$$

$$= \sigma_{11} \cos^2 \theta + \sigma_{22} \sin^2 \theta + 2 \sigma_{12} \sin \theta \cos \theta$$

$$\sigma_{1'} = \frac{\sigma_{11} + \sigma_{22}}{2} + \frac{\sigma_{11} - \sigma_{22}}{2} \cos 2\theta + \sigma_{12} \sin 2\theta$$

This is going to be given by $l_{11} l_{11} \sigma_{11} + l_{12} l_{12} \sigma_{22} + 2 l_{11} l_{12} \sigma_{12}$. So, we are taking the value of i to be 1. So, therefore, $l_{11} l_{11} \sigma_{11} + l_{12} l_{12} \sigma_{22} + 2 l_{11} l_{12} \sigma_{12}$. So, you can expand it further. So, this is equal to $l_{11} l_{11} \sigma_{11} + l_{12} l_{12} \sigma_{22} + 2 l_{11} l_{12} \sigma_{12}$. So, now we will try to substitute the values. So, we will get

now, we will simply represent the stress σ_1 as σ_1 . This is going to be now, we have σ_{11} , σ_{12} , σ_{21} and σ_{22} .

So, if we now substitute the values, then this is $\sigma_{11} \cos^2 \theta$ plus this is $\cos \theta$, this is $\sin \theta$, so $\sigma_{12} \cos \theta \sin \theta$. Similarly, we have this is $\sin \theta$, this is $\cos \theta$. So, therefore, it is $\sigma_{21} \cos \theta \sin \theta$. And this is nothing but $\sin^2 \theta$, so therefore, $\sigma_{22} \sin^2 \theta$. So, this finally is nothing but $\sigma_{11} \cos^2 \theta$ plus $\sigma_{22} \sin^2 \theta$.

Since, $\sigma_{12} = \sigma_{21}$, we can write this thing as $\sigma_{12} \sin 2\theta$. This can be written finally, that the σ_1 is nothing but if we write this $\cos^2 \theta$ is equal to $\frac{1 + \cos 2\theta}{2}$ and $\sin^2 \theta$ as $\frac{1 - \cos 2\theta}{2}$. Then, this will be given by σ_{11} , σ_{22} by 2 plus $\sigma_{11} - \sigma_{22}$ by 2 $\cos 2\theta$ plus $\sigma_{12} \sin 2\theta$. So, we find that σ_1 is given in this form.

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The image shows handwritten mathematical derivations for stress transformation. The top part shows the transformation for principal stress σ_1 in terms of σ_{11} , σ_{22} , and σ_{12} . The bottom part shows the transformation for normal stress σ_x in terms of σ_x , σ_y , and τ_{xy} .

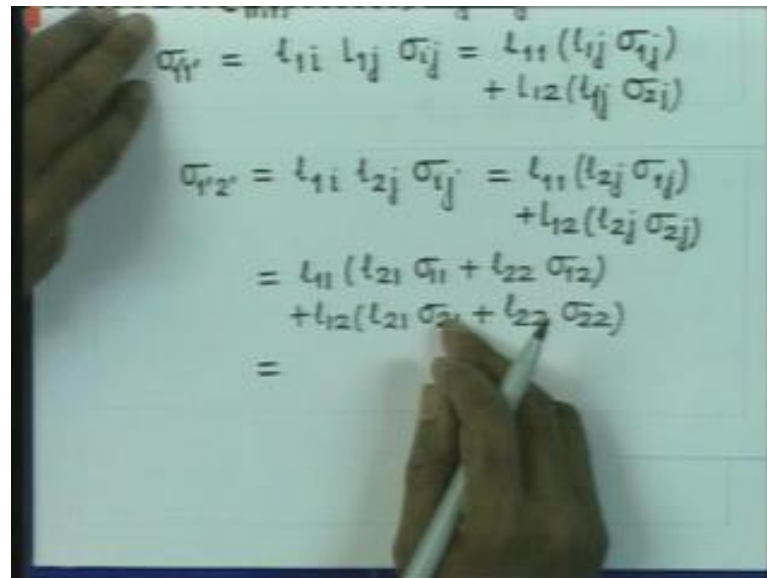
$$\begin{aligned} &= \sigma_{11} \cos^2 \theta + \sigma_{22} \sin^2 \theta + 2\sigma_{12} \sin \theta \cos \theta \\ \sigma_1 &= \frac{\sigma_{11} + \sigma_{22}}{2} + \frac{\sigma_{11} - \sigma_{22}}{2} \cos 2\theta + \sigma_{12} \sin 2\theta \end{aligned}$$

$$\begin{aligned} &\begin{matrix} 1 \rightarrow x' & 2 \rightarrow y' & 1 \rightarrow x & 2 \rightarrow y \end{matrix} \\ \sigma_x &= \frac{\sigma_x + \sigma_y}{2} + \frac{\sigma_x - \sigma_y}{2} \cos 2\theta + \tau_{xy} \sin 2\theta \end{aligned}$$

And if you want write in the x, y system. So, if represent 1 dash by x dash and 2 dash by y dash and similarly, 1 by x and 2 by y. Then, you can write the above expression in the form, that σ_x is equal to σ_x plus σ_y by 2 plus σ_x plus σ_y by 2 $\cos 2\theta$ plus $\tau_{xy} \sin 2\theta$. So, this is the transformation that you have considered in relation to Mohr's circles in two dimensions.

Similarly, if you now consider the other stress components. So, we would like to consider the component of the stress, which is going to be in the direction, 2 dash. So, it will be, this is the plane, we are considering and so this is the direction, which is outer normal. So, now we will get the component sigma 1 dash, 2 dash. This is now going to be equal to. So, it is again 1 i and n is 2 now, 2 j.

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The image shows a whiteboard with handwritten equations for stress transformation. The first equation is $\sigma_{1'1'} = l_{1i} l_{1j} \sigma_{ij} = l_{11}(l_{1j} \sigma_{1j}) + l_{12}(l_{1j} \sigma_{2j})$. The second equation is $\sigma_{1'2'} = l_{1i} l_{2j} \sigma_{ij} = l_{11}(l_{2j} \sigma_{1j}) + l_{12}(l_{2j} \sigma_{2j})$. The third equation is $= l_{11}(l_{21} \sigma_{11} + l_{22} \sigma_{12}) + l_{12}(l_{21} \sigma_{21} + l_{22} \sigma_{22})$. The fourth equation is $=$.

So, in fact, this dash has no meaning, we are just try to consider the new direction 1 and 2. So, this is how, we are going to get the link of the new system to the whole system. And again, we are going to have expansion of i first. So, it is l 11, l 2 j, sigma 1 j plus l 12, l 2 j, sigma 2 j. Now, we can write l 11, l 21 sigma 11, plus l 22 sigma 12, plus l 12, l 21 sigma 21, plus l 22 into sigma 22. So, all the four components of the stresses are now involved, sigma 11, sigma 12, sigma 21 and sigma 22.

So, if we now try to collate them, noting that the values of directed cosines. So, if we look into the values of the direction cosines, which we have given earlier. So, we will have now, this is cos theta, this is minus sine theta. So, therefore, it is the going to be cos theta, sine theta, sigma 11 plus this is cos theta, this is cos theta. So, therefore, cos square theta sigma 11 plus 2 cos theta sine theta sigma 12.

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$$\begin{aligned}
 &= l_{11}(l_{21}\sigma_{11} + l_{22}\sigma_{12}) \\
 &\quad + l_{12}(l_{21}\sigma_{21} + l_{22}\sigma_{22}) \\
 &= -\cos\theta\sin\theta\sigma_{11} + \sigma\cos^2\theta\sigma_{12} \\
 &\quad - \sin^2\theta\sigma_{21} + \cos\theta\sin\theta\sigma_{22} \\
 \\
 \sigma_{12}' = \tau_{x'y'} &= -\left(\frac{\sigma_{11} - \sigma_{22}}{2}\right)\sin 2\theta \\
 &\quad + \sigma_{12}\cos 2\theta \\
 \tau_{x'y'} &= -\frac{\sigma_x - \sigma_y}{2}\sin 2\theta \\
 &\quad + \tau_{xy}\cos 2\theta
 \end{aligned}$$

And this is again, $l_{12}\sin\theta$ and this is minus sine theta. So, therefore, it is minus sine square theta sigma 21 plus here this is $l_{12}\cos\theta$ and this is cosine theta. So, therefore, this is sine theta and this is cosine theta. So, therefore it is cosine theta, sine theta, sigma 22. So, we can now write it, sigma 1 dash, 2 dash, which is nothing but the stress.

And this is going to be minus sigma 11 minus sigma 22 by 2, sine 2 theta plus since sigma 12 is equal to sigma 21. So, you can write this thing as sigma 12 cosine 2 theta. So, this is the expression for the shear stress in the new coordinates. And if we want to write in terms of sigma x and sigma y, this is nothing but sigma x minus sigma y by 2 sine 2 theta plus tau x y cosine 2 theta.

So, the new shear stress is like this and this is what, again you have obtained in connection more circle diagram in the two dimensions. Just finally, to look at the expression for the transformed x dash and the shear stress is like this.

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$$\begin{aligned}\sigma_{x'} &= \frac{\sigma_x + \sigma_y}{2} + \frac{\sigma_x - \sigma_y}{2} \cos 2\theta \\ &\quad + \tau_{xy} \sin 2\theta \\ \sigma_{y'2'} = \tau_{x'y'} &= - \left(\frac{\sigma_{11} - \sigma_{22}}{2} \right) \sin 2\theta \\ &\quad + \tau_{12} \cos 2\theta \\ \tau_{x'y'} &= - \frac{\sigma_x - \sigma_y}{2} \sin 2\theta \\ &\quad + \tau_{xy} \cos 2\theta \\ \sigma_{x'2'} = \sigma_{y'} &= I_1 - \sigma_{x'} = \sigma_x + \sigma_y - \sigma_{x'} \\ &= \frac{\sigma_x + \sigma_y}{2} - \frac{\sigma_x - \sigma_y}{2} \cos 2\theta \\ &\quad - \tau_{xy} \sin 2\theta\end{aligned}$$

And if you are interested in finding out the stress $\sigma_{y'2'}$, which is nothing but $\sigma_{y'}$. You make use of the fact that, this is driven by the stress invariant. We can make use of I_1 minus $\sigma_{x'}$ and this is nothing but $\sigma_x + \sigma_y$, this is $\sigma_{x'}$ minus $\sigma_{x'}$. The stress invariant is same. So, therefore, one can use of the fact that $\sigma_{x'}$ equal to $\sigma_{y'}$, I_1 that gives these.

And finally, if you substitute the value of the $\sigma_{x'}$, it is finally, come to be like this, $\sigma_x + \sigma_y$ by 2 minus $\sigma_x - \sigma_y$ by 2 cosine 2 theta minus τ_{xy} sine 2 theta. So, therefore, that is the $\sigma_{y'}$. So, finally, your new stress are give in $\sigma_{x'}$ like this and $\sigma_{y'}$ is given here and shear stress. Now, we will like to see, how to draw more circle for stresses in three dimensions. You are familiar with the most circles in the two dimensions.

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MOHR CIRCLE IN 3-D

$$\sigma_{ij} \rightarrow \begin{matrix} \sigma_1, \sigma_2, \sigma_3 \\ l_{1j} \quad l_{2j} \quad l_{3j} \end{matrix}$$

$$\text{d.c.'s} \rightarrow l, m, n \quad (\text{Ref. } \sigma_1, \sigma_2, \sigma_3)$$

$$\sigma = l^2 \sigma_1 + m^2 \sigma_2 + n^2 \sigma_3$$

$$l^2 + m^2 + n^2 = 1$$

$$\sigma, \sigma = l^2 (\sigma_1 - \sigma_3) + m^2 (\sigma_2 - \sigma_3) + \sigma_3 \quad \dots (1)$$

So, we will like to consider the construction of Mohr circles in three dimensions. We have the stresses at the point i j starting from the stresses. You must calculate the principle stresses, sigma 1, sigma 2 and sigma 3; you can also calculate the directions. Now, we will like to consider an arbitrary plane with direction cosines l, m and n and the reference directions are sigma 1, sigma 2 and sigma 3. So, on this plane, we will have the normal stress given by $l^2 \sigma_1 + m^2 \sigma_2 + n^2 \sigma_3$.

We are also having the relation that $l^2 + m^2 + n^2 = 1$. So, if I substitute l^2 to $1 - m^2 - n^2$. That will give me sigma equal to $l^2 \sigma_1 - \sigma_3 + m^2 \sigma_2 - \sigma_3 + \sigma_3$. So, let us consider this relation to be 1.

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$\sigma_n = l(u_1 - u_3) + m(u_2 - u_3) + n \cdot 0$... (1)
 Total traction $S^2 = S_1^2 + S_2^2 + S_3^2$
 $S^2 = l^2 \sigma_1^2 + m^2 \sigma_2^2 + n^2 \sigma_3^2$
 $S^2 = \sigma^2 + \tau^2$
 ~~$\tau^2 = S^2$~~ $\tau^2 + \sigma^2 = l^2 \sigma_1^2 + m^2 \sigma_2^2 + n^2 \sigma_3^2$
 $\sigma, \tau^2 + \sigma^2 = l^2 (\sigma_1^2 - \sigma_3^2) + m^2 (\sigma_2^2 - \sigma_3^2) + \sigma_3^2$

Now, the tractions on the plane l, m and n are going to be given by sigma 1, m sigma 2 and n sigma 3. So, if we now consider the total traction, which is going to be related to S 1, S 2, S 3. The three tractions on the same plane and the relationship is nothing but S square is equal to S 1 square plus S 2 square plus S 3 square. Since, S 1 is l sigma 1 and S 2 is equal to m sigma 2 and S 3 is n sigma 3.

We can write the relation as S square is nothing but l square sigma 1 square plus m square sigma 2 square plus n square sigma 3 square. These types relationship, we are also derived earlier in connection with the introduction of the maximum shear stress and also the shear stress directions.

Now, since the normal stress on the plane is equal to sigma. And if you try to consider the shear stresses on the plane to be tau, then we write S square, which is also related to sigma square plus tau square. And therefore, we will have tau square is nothing but S square. We will continue writing tau square plus sigma square is equal to l square, sigma 1 square plus m square sigma 2 square plus n square sigma 3 square.

Again, we will like to substitute the n square in terms of 1 minus l square minus m square. So, we will get tau square plus sigma square is equal to, we have l 1 square into sigma 1 square minus sigma 3 square plus m square sigma 2 square minus sigma 3 square plus sigma 3 square.

So, this we can now rearranging, $\tau^2 + \sigma^2 - \sigma_3^2$ is equal to, this is not 1, this is 1 only. So, this is $1 - \sigma_1^2 - \sigma_3^2 + \sigma_2^2$ into $\sigma_2^2 - \sigma_3^2$. So, let us write this equation as equation number 2. I will show you, that these relations 1 and 2 are the basis to draw the Mohr circle in three dimensions in my next presentation.