

Advanced Strength of Materials
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Lecture – 4

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Handwritten equations on a whiteboard:

$$\left. \begin{aligned} l(\sigma_x - \sigma) + m\tau_{xy} + n\tau_{xz} &= 0 \\ l\tau_{xy} + m(\sigma_y - \sigma) + n\tau_{yz} &= 0 \\ l\tau_{xz} + m\tau_{yz} + n(\sigma_z - \sigma) &= 0 \end{aligned} \right\} \dots (1)$$

$$\det \begin{vmatrix} \sigma_x - \sigma & \tau_{xy} & \tau_{xz} \\ \tau_{xy} & \sigma_y - \sigma & \tau_{yz} \\ \tau_{xz} & \tau_{yz} & \sigma_z - \sigma \end{vmatrix} = 0 \dots (2)$$

Last time we have seen that if the principal plane has direction cosines l , m and n . Then, they will satisfy these set of equations and a magnitude of the principal stress is given by σ . To determine the principal stresses, we take the characteristic equation determinant of this expression equal to 0. You could get the principal stresses by solving this set of equations.

Now, the question is, how do you get the direction cosines corresponding to a particular value of σ ? You take the value of σ . Let us say σ_1 , substitute in these three equations, thereby, you will get again three equations involving l , m and n . Solving these equations; you are going to get the direction cosines corresponding to the stress σ . So, you will have three values of σ out of the characteristic equation. And therefore, you can solve these set 1, thrice and you can get all the direction cosines.

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EXAMPLE ON DETERMINATION OF σ_1, σ_2 & σ_3 AND THEIR DIRECTIONS

$$\sigma_{ij} = \begin{bmatrix} 12.31 & 4.20 & 0.84 \\ 4.20 & 8.96 & 5.27 \\ 0.84 & 5.27 & 4.34 \end{bmatrix} \text{ (MPa)}$$
$$I_1 = \sigma_x + \sigma_y + \sigma_z = 12.31 + 8.96 + 4.34 = 25.61 \text{ MPa}$$
$$I_2 = \sigma_y^2 + \sigma_z^2 + \sigma_x^2 - \sigma_x \sigma_y - \sigma_y \sigma_z - \sigma_z \sigma_x$$
$$= 4.20^2 + 5.27^2 + 0.84^2 - 12.31 \times 8.96 - 8.96 \times 4.34 - 4.34 \times 12.31 = -156.48 \text{ (MPa)}^2$$

We would like to illustrate these by taking some example today. It is given the stresses at a point are given by these tensor. These magnitudes are all in Mega Pascal. So, sigma x x is 12.31, tau x y is 4.20, tau x z is 0.84. Similarly, you have tau y x, tau sigma y y, tau y z and this is tau x z, tau y z and sigma z z. Now, we calculate first, the stress invariants, using the data I 1 is sigma x plus sigma y plus sigma z.

So, if you substitute the value, these are the three diagonal coefficients. It gives you 25.61 Mega Pascal. Similarly, if you consider the second stress invariant, I 2. It is given by tau x y square tau y z square plus tau z x square minus sigma x sigma y minus sigma y sigma z minus sigma z sigma x. So, this, when I substitute the values, this is 4.20, this stress. Then, this is 5.27. This is 0.84. Then, we have 12.31 into 8.96, 8.96 into 4.34 and then 4.34 multiplied by 12.31. So, on simplification, this will become 156.48 Mega Pascal square.

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The image shows a handwritten calculation on a piece of paper. The first line is the formula for the 3rd stress invariant I_3 :
$$I_3 = \sigma_x \sigma_y \sigma_z + 2 \tau_{xy} \tau_{yz} \tau_{zx} - \sigma_x \tau_{yz}^2 - \sigma_y \tau_{zx}^2 - \sigma_z \tau_{xy}^2$$
The next line shows the substitution of values:
$$= 12.31 \times 8.96 \times 4.34 + 2 \times 4.20 \times 5.27 \times 0.84 - (12.31 \times 5.27^2 - 8.96 \times 0.84^2 - 4.34 \times 4.20^2)$$
The result is:
$$= 91.12 \text{ (MPa)}^3$$
Then, the constant a is calculated:
$$a = - (I_2 + I_1^2/3) = - (-156.48 + \frac{25.61^2}{3})$$

$$= -62.14 \text{ (MPa)}^2$$
Finally, the constant b is calculated:
$$b = - (I_3 + 2 I_1^3/27 + I_1 I_2/3)$$

$$= - (91.12 + 2 \frac{25.61^3}{27} + \frac{25.61 \times (-156.48)}{3})$$

$$= +0.48 \text{ (MPa)}^3$$

Similarly, if you try to consider the 3rd stress invariant I_3 which is given by sigma x into sigma y into sigma z, plus 2 tau x y tau y z tau z x minus sigma x tau y z square minus sigma y tau z x square minus sigma z tau x y square. So, again, if you substitute the values of the respective stresses, you get these. And on simplification, this will become 91.12 Mega Pascal square.

Then, we try to calculate the constants, a is minus I_2 plus I_1 square by 3. And so if you substitute the value, we have got earlier. This is minus 156.48, that is I_2 , and then I_1 is 25.61. So, therefore, 25.61 square by 3 and on simplification, this will become 6 minus 62.14 MPa square. Similarly, if we calculate the constant b, which is related to minus I_3 plus 2 I_1 cube by 27 plus I_1 into I_2 by 3. And again, if we substitute the values, respective values, we get the simplification to be plus 0.48.

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$$\cos \phi = \frac{-b}{2\sqrt{\frac{a^2}{27}}} = \frac{-0.40}{2\sqrt{\frac{62.14}{27}}} \approx 0$$

$$\therefore \phi = 90^\circ$$

$$\sigma_1 = \frac{I_1}{3} + 2\sqrt{\frac{a}{3}} \cos \frac{\phi}{3} = \frac{25.61}{3} + 2\sqrt{\frac{62.14}{3}} \cos 30^\circ$$

$$= 16.41 \text{ MPa}$$

$$\sigma_2 = \frac{I_1}{3} + 2\sqrt{\frac{a}{3}} \cos(120^\circ + \frac{\phi}{3}) = 0.65 \text{ MPa}$$

Then, we calculate the angle phi, cos phi is equal to minus b by 2 into square root minus a cube by 27. And if we substitute the values, this comes out to be very small quantity, which is almost equal to 0. Therefore, it means that phi angle is equal to 90 degree. Now, the value of the principal stresses given by sigma 1 equal to I 1 by 3 plus 2 square root minus a by 3 cosine phi by 3.

So, substitute the value 25.61 by 3 plus 2 square root 62.14 by 3 into cosine phi is 90. So, therefore, this becomes 30 degree. And if you simplify this comes out to be 16.41 Mega Pascal. Similarly, if we consider the second principal stress, it is given by sigma 2 equal to I 1 by 3 plus 2 square root minus a by 3 cosine 120 degree plus phi by 3. Again, once you simplify, you get the sigma 2 stress as 0.65 MPa. If we now take the last one, it is going to be I 1 by 3 plus 2 square root minus a by 3 cosine 240 degree plus phi by 3. Again, on some simplification, this is equal to 8.53 MPa.

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$$\begin{aligned} \therefore \phi &= 30^\circ \\ \sigma_1 &= \frac{I_1}{3} + 2\sqrt{-\frac{J_2}{3}} \cos \frac{\phi}{3} = \frac{23.61}{3} + 2\sqrt{\frac{16.41}{3}} \cos 30^\circ \\ &= 16.41 \text{ MPa} \\ \sigma_2 &= \frac{I_1}{3} + 2\sqrt{-\frac{J_2}{3}} \cos(120^\circ + \frac{\phi}{3}) = 0.65 \text{ MPa} \\ \sigma_3 &= \frac{I_1}{3} + 2\sqrt{-\frac{J_2}{3}} \cos(240^\circ + \frac{\phi}{3}) = 8.53 \text{ MPa} \\ I_1 &= \sigma_x + \sigma_y + \sigma_z = \sigma_1 + \sigma_2 + \sigma_3 \\ \sigma_3 &= \sigma_x + \sigma_y + \sigma_z - (\sigma_1 + \sigma_2) = 8.53 \text{ MPa} \end{aligned}$$

Now, at this point itself, I would like to tell you, that sigma 3 can be calculated in a slightly different way. Since, we have the set stress invariant I 1. This I 1 is equal to, it is given by sigma x plus sigma y plus sigma z. And it can also be equal to the sum of the three principal stresses. If we consider our Cartesian system to coincide with the directions of the principal stresses. So, this must also be given by sigma 1 plus sigma 2 plus sigma 3.

Now, we have already got the values of the two stresses sigma 1 and sigma 2. And therefore, this sigma 3 is equal to sigma x plus sigma y plus sigma z minus sigma 1 plus sigma 2. So, this can be a simpler way of going about and if you calculate, this will also become 8.53 MPa. So, once, we have got the principal stresses, now we have to determine the direction cosines. That means, we are interested in finding out, what is the direction for the stress sigma 1. In other words, if you consider the direction cosines of the plane only sigma 1 is acting is given by l 1, m 1 and n 1. Then what are the magnitudes of these three direction cosines.

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$$\text{D.C.'s for } \sigma_1? \quad l_1, m_1, n_1 = ?$$

$$\left. \begin{aligned} l_1(\sigma_x - \sigma_1) + m_1 \tau_{yx} + n_1 \tau_{zx} &= 0 \\ l_1 \tau_{xy} + m_1(\sigma_y - \sigma_1) + n_1 \tau_{zy} &= 0 \\ l_1 \tau_{xz} + m_1 \tau_{yz} + n_1(\sigma_z - \sigma_1) &= 0 \end{aligned} \right\}$$

Now, these direction cosines are going to satisfy, equation number 1; that we wrote last time. And this will give us that l_1 into σ_x minus σ_1 . So, we have substituted σ_x by σ_1 plus m_1 into τ_{yx} plus n_1 into τ_{zx} equal to 0. Similarly, l_1 into τ_{xy} plus m_1 into σ_y minus σ_1 plus n_1 into τ_{zy} equal to 0. Similarly, l_1 into τ_{xz} plus m_1 into τ_{yz} plus n_1 into σ_z minus σ_1 equal to 0.

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$$\left. \begin{aligned} l_1(\sigma_x - \sigma_1) + m_1 \tau_{yx} + n_1 \tau_{zx} &= 0 \\ l_1 \tau_{xy} + m_1(\sigma_y - \sigma_1) + n_1 \tau_{zy} &= 0 \\ l_1 \tau_{xz} + m_1 \tau_{yz} + n_1(\sigma_z - \sigma_1) &= 0 \end{aligned} \right\}$$

$$\left. \begin{aligned} l_1(1235 - 1640) + m_1 420 + n_1 0.84 &= 0 \\ l_1 420 + m_1(836 - 1640) + n_1 527 &= 0 \\ l_1 0.84 + m_1 527 + n_1(434 - 1640) &= 0 \end{aligned} \right\}$$

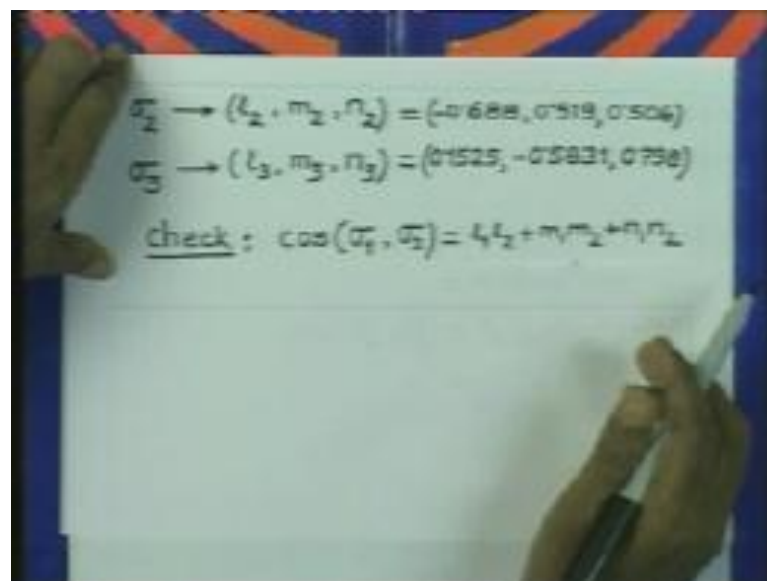
$$\rightarrow (l_1, m_1, n_1) = (0.706, 0.628, 0.324)$$

These three equations, once you substitute the value of the stresses Cartesian components of the stresses. Then, in that case, you get σ_x substituted by this, 4.20 is the value of

τ_{yx} and 0.84 is the value of τ_{zx} . Similarly, you have the second equation here, this is τ_{xy} plus m_1 into σ_{yy} minus 16.41 is the value of the first principal stress plus n_1 into 5.27. And last equation is l_1 into this is the value of τ_{xz} and m_1 into 5.27. That is the value of τ_{yz} plus n_1 into 4.34 will be the value of σ_{zz} minus 16.41. So, these three equations are the equations to solve for the constants l_1 , m_1 and n_1 .

Now, it is important to know, that the three direction cosines are going to satisfy the equation, which is nothing but l_1^2 plus m_1^2 plus n_1^2 equal to 1. So, you can make use of that. And after solving this set of equations, you get the direction cosines l_1 , m_1 , n_1 and 0.706, 0.628, 0.324. So, this is the direction a cosine of the plane on each σ_1 stress is acting.

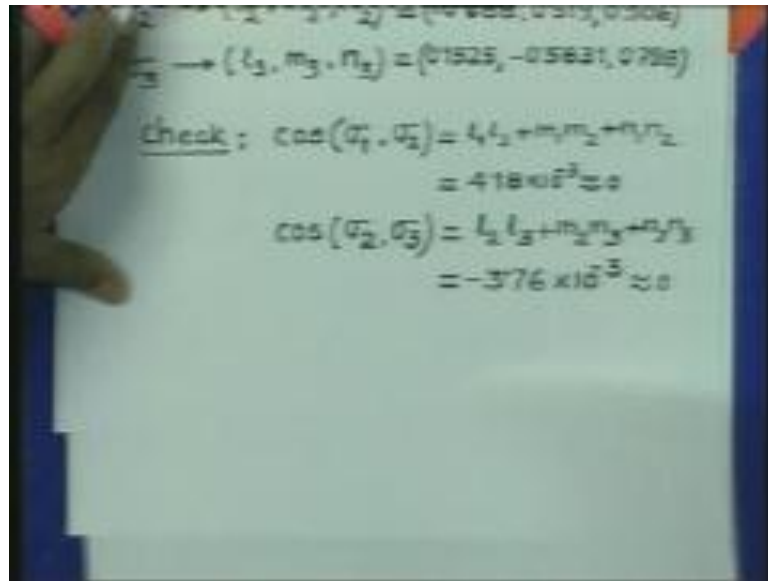
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$$\begin{aligned}\sigma_2 &\rightarrow (l_2, m_2, n_2) = (-0.688, 0.519, 0.506) \\ \sigma_3 &\rightarrow (l_3, m_3, n_3) = (0.1525, -0.5831, 0.798) \\ \text{check: } \cos(\sigma_1, \sigma_2) &= l_1 l_2 + m_1 m_2 + n_1 n_2\end{aligned}$$

Similarly, for the second stress, second principal stress σ_2 . If you go by the same procedure, we will find that the direction cosines l_2 , m_2 , n_2 will come out as minus 0.688, 0.519 and 0.506. In the similar way, σ_3 , the direction cosines of the plane on the σ_3 is acting are l_3 , m_3 , n_3 is going to be given by 0.1525 minus 0.5831, 0.798. Now, you have got the results, you can now check it. You can check whether the direction cosines are alright. We know that the directions σ_1 and σ_2 , they are orthogonal. And the angle between the two directions, cosine of the angle between the 2 direction σ_1 and σ_2 is going to be given by $l_1 l_2$ plus $m_1 m_2$ plus $n_1 n_2$.

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The image shows a whiteboard with handwritten mathematical work. At the top, a vector is defined as $\vec{e}_3 \rightarrow (l_3, m_3, n_3) = (0.1525, -0.5631, 0.798)$. Below this, a 'check' is performed for the cosine of the angle between $\vec{e}_1$ and $\vec{e}_2$:
$$\cos(\sigma_1, \sigma_2) = l_1 l_2 + m_1 m_2 + n_1 n_2$$
$$= 4.18 \times 10^{-3} \approx 0$$

Then, the cosine of the angle between $\vec{e}_2$ and $\vec{e}_3$ is calculated:
$$\cos(\sigma_2, \sigma_3) = l_2 l_3 + m_2 m_3 + n_2 n_3$$
$$= -3.76 \times 10^{-3} \approx 0$$

So, if we substitute the values that we have obtained earlier. We find that $l_1 l_2 + m_1 m_2 + n_1 n_2$ is equal to 4.18×10^{-3} , which is close to 0, corresponding to the values that we have got for the individual direction cosines. Similarly, you will find that for the directions σ_2 and σ_3 , cosine of the angle is given by $l_2 l_3 + m_2 m_3 + n_2 n_3$. This comes out to be -3.76×10^{-3} , which is again close to 0.

So, you can now go on considering the other two stresses σ_3 and σ_1 . And we will find that, the angles between the two directions are again going to be 90 degree. So, that gives you a final check on the calculations that you have done. Now, I would like to show you that, it is possible to consider. It is possible to prove that, the principal directions are mutually orthogonal. You have just now seen that, the principal directions direction cosines are such in the angle between the two directions are 90 degree.

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PRINCIPAL DIRECTIONS ARE MUTUALLY
ORTHOGONAL

σ_1	l_1	m_1	n_1
σ_2	l_2	m_2	n_2

$$\begin{aligned}
 l_2 \times l_1 \sigma_2 + m_1 \tau_{yz} + n_1 \tau_{zx} - l_1 \sigma_1 &= 0 \\
 m_2 \times l_1 \tau_{xy} + m_1 (\sigma_y) + n_1 \tau_{zy} - m_1 \sigma_1 &= 0 \\
 n_2 \times l_1 \tau_{xz} + m_1 \tau_{yz} + n_1 \sigma_z - n_1 \sigma_1 &= 0 \\
 l_1 \times l_2 \sigma_2 + m_2 \tau_{yz} + n_2 \tau_{zx} - l_2 \sigma_2 &= 0 \\
 m_2 \times l_2 \tau_{xy} + m_2 \sigma_y + n_2 \tau_{zy} - m_2 \sigma_2 &= 0 \\
 n_2 \times l_2 \tau_{xz} + m_2 \tau_{yz} + n_2 \sigma_z - n_2 \sigma_2 &= 0
 \end{aligned}$$

Now, we would like to show it as a general case. So, let us consider the proof for principal directions are mutually orthogonal. Let us consider that σ_1 is the principal stress 1 and the direction cosines are l_1, m_1, n_1 . Similarly, σ_2 is the second principal stress and the direction cosines are l_2, m_2, n_2 . Now, we have the equations, which we have written little while ago, that $l_1 \sigma_x + m_1 \tau_{yx} + n_1 \tau_{zx} - l_1 \sigma_1 = 0$. So, that will be first equation.

Similarly, will have $l_1 \tau_{xy} + m_1 \sigma_y + n_1 \tau_{zy} - m_1 \sigma_1 = 0$. Similarly, will have the third equation $l_1 \tau_{xz} + m_1 \tau_{yz} + n_1 \sigma_z - n_1 \sigma_1 = 0$. So, for the first principal stress will get these set of equations. Similarly, if we write the set for the second principal stress, it is going to be like this, $l_2 \sigma_x + m_2 \tau_{yx} + n_2 \tau_{zx} - l_2 \sigma_2 = 0$.

Again, $l_2 \tau_{xy} + m_2 \sigma_y + n_2 \tau_{zy} - m_2 \sigma_2 = 0$. And the last one, $l_2 \tau_{xz} + m_2 \tau_{yz} + n_2 \sigma_z - n_2 \sigma_2 = 0$. So, now we have two sets, we like to multiply this equation by l_2 , this equation by m_2 and this equation by n_2 . And we would like to add these three equations. Then, we multiply this equation by l_1 , this second equation by m_1 and this equation by n_1 .

Again, we add up these three relations, and then the resultant out of this sum is subtracted from the sum that we have got from this. So, you will find that, we are going

to have terms like $l_2, l_1 \sigma_x$, here also it will go to be $l_1, l_2 \sigma_x$. So, therefore, this term will cancel this 1. Similarly, $l_2 m_1 \tau_{yx}$ and you will find here, you are going to have l_2, m_1 . So, therefore, you will now find $l_2, m_1 \tau_{xy}$, because τ_{yx} is equal to τ_{xy} . So, all these terms are going to cancel.

However, you will find that, the term that you are going to get $l_2, l_1 \sigma_1$, it is not going to cancel. So, also m_2, m_1 into σ_1 , they are not going to cancel. So, the terms, which are going to come out of this and also of this, they are not going to cancel. And you will find that, the term here is $l_1, l_2 \sigma_1$. Here, also $l_1, l_2 \sigma_2$.

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$$(\sigma_2 - \sigma_1)(l_1 l_2 + m_1 m_2 + n_1 n_2) = 0$$

$$\sigma_1 \neq \sigma_2$$

$$l_1 l_2 + m_1 m_2 + n_1 n_2 = 0$$

$$\cos(\theta_{\sigma_1, \sigma_2}) = 0 \quad \sigma_1 \perp \sigma_2$$

$$\cos(\theta_{\sigma_2, \sigma_3}) = 0 \quad \sigma_2 \perp \sigma_3$$

$$\cos(\theta_{\sigma_3, \sigma_1}) = 0 \quad \sigma_3 \perp \sigma_1$$

So, you can immediately appreciate that. By this process, we are going to get relation like this. It is σ_2 minus σ_1 multiplied by l_1, l_2 plus $m_1 m_2$ plus n_1, n_2 . This is equal to 0. We know that, these two principal stresses are distinct and they are different. Therefore, σ_1 is not equal to σ_2 and this equation or rather this expression has got to be 0. So, this gives that l_1, l_2 plus $m_1 m_2$ plus n_1, n_2 is equal to 0.

This means that, cosine of the angle between the two directions, σ_1 and σ_2 is equal to 0. So, σ_1 and σ_2 directions are mutually orthogonal. By similar approach, you can again show that cosine of the angle between the direction σ_2 and σ_3 is equal to 0. Therefore, σ_2 is orthogonal to σ_3 and again cosine of the

angle between the directions sigma 3 and sigma 1 equal to 0. And therefore, sigma 3 is orthogonal to sigma 1. Thus the three directions are mutually orthogonal.

So, what it means is that, you are two directions, any two directions, if you consider that the direction cosines are l_{j1} and the other direction is l_{j2} and the third direction is l_{j3} . Then, these three directions are going to be already mutually orthogonal. Now, these expressions that we have written, it has taken more space to write. We can write in tensor notation and then it will be much more compact.

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The image shows a piece of paper with handwritten equations. At the top, there are labels for principal stresses σ_1 and σ_2 , and direction cosines l_{j1} and l_{j2} . The main equations are:

$$(l_{ji} \sigma_{ji} - l_{ji} \sigma \delta_{ji}) = 0$$

$$l_{j1}^2 - l_{j1}^i \sigma_{ji} - l_{j1}^i \delta_{ji} \sigma_1 = 0$$

$$l_{j2}^2 - l_{j2}^i \sigma_{ji} - l_{j2}^i \delta_{ji} \sigma_2 = 0$$

$$(l_{j1}^i l_{j1}^i + l_{j2}^i l_{j2}^i + l_{j3}^i l_{j3}^i) (\sigma_1 - \sigma_2) = 0$$

Now, let us consider, again that the sigma 1 is the principal stress and the corresponding direction cosines going to be given by l_{j1} . Again, we have sigma 2 and the corresponding direction cosines are l_{j2} . Now, the equations that we had obtained involving the principal stresses, they were of the form we are written earlier. That $l_{ji} \sigma_{ji} - l_{ji} \sigma \delta_{ji} = 0$.

So, now for the stress sigma 1, we can write that this is nothing but $l_{j1} \sigma_{ji} - l_{j1} \sigma_1 \delta_{ji} = 0$. So, this is the set, we get from the principal stress sigma 1. Similarly, for the principal stress sigma 2, without any difficulty we can write, that $l_{j2} \sigma_{ji} - l_{j2} \sigma_2 \delta_{ji} = 0$. Now, we will like to multiply this thing by l_{i2} , this equation. And this equation by l_{i1} , with the understanding that, i is also varying from 1 to 3.

Then, having multiplied, if you subtract this resultant equation from these, you will find that, you are going to get this terms canceling, only this terms are going to be not canceling. And at the same time, only when we have I equal to j , we will get some value. And therefore, by simplifying this, what you are going to find is this, σ_1^2 into super script 1, σ_1^2 into super script 2 plus σ_2^2 into super script 1, σ_2^2 into super script 2 plus σ_3^2 into super script 1, σ_3^2 into super script 2.

So, this is the term, which we are going to get out of this. This multiplied by σ_1 minus σ_2 equal to 0. Again, you find that, this is nothing but according to our earlier expression it is σ_1^2 plus σ_2^2 plus σ_3^2 . So, this is the advantage of using tensor notation, you get very compact relationship.

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$\sigma^3 - I_1 \sigma^2 - I_2 \sigma - I_3 = 0 \dots (3)$
ROOTS OF EQN. (3) ARE ALL REAL
 $\sigma_1, \sigma_2, \sigma_3$
 $\sigma_3 \rightarrow \text{Real} \quad \sigma_1 \& \sigma_2 \rightarrow \text{Complex}$
 $I_1 = \sigma_1 + \sigma_2 + \sigma_3 = \sigma_1 + \sigma_2 + \sigma_3$
 $\sigma_1 \& \sigma_2 \rightarrow \text{Complex conjugates}$
 $I_1 \sigma_j - I_2 \delta_{jk} \sigma = 0$

Now, we would like to show that, the principal stresses are real. They are not going to have complex quantity. Since, the equation, that we had gotten by writing the characteristic equation, it was nothing but σ^3 minus $I_1 \sigma^2$ minus $I_2 \sigma$ minus I_3 equal to 0. The roots of this equation are going to be all real. So, this equation, if we indicated as equation number, let us say 3. So, roots of equation 3 are all real.

Let us assume that, the three roots are again, σ_1 , σ_2 and σ_3 . Let us consider that, σ_3 is real, but σ_1 and σ_2 is complex. Since, I_1 is equal to $\sigma_1 + \sigma_2 + \sigma_3$. This is real quantity and at the same time, we have

seen that, this I_1 is also equal to $\sigma_1 + \sigma_2 + \sigma_3$. If sum of these three quantities, have got to be real. It demands that, σ_1 and σ_2 must be related.

In fact, one should be the complex conjugate of the other. So, it is σ_1 and σ_2 is complex to get the imaginary part of this sum of these two roots to cancel. We must have σ_1 and σ_2 must be conjugate complex conjugate. σ_1 and σ_2 are complex conjugates. Now, you would also recollect, that the equation, that we are going to be, that we satisfied by these are $\sum_j l_j \sigma_j - \sum_j l_j \delta_j = 0$. So, this is what, is the equation to be satisfied by these roots.

Now, we find that, this component yields complex one of the component is complex and therefore, this equation now going to involve complex coefficients. And hence it will indicate that the direction cosines have got to be complex. So, σ_1 and σ_2 are complex conjugates, and then this equation is to be satisfied, by all the roots.

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The image shows handwritten mathematical notes on a whiteboard. The notes are as follows:

$$\sigma_1 \rightarrow l_1^2 \quad \sigma_2 \rightarrow l_2^2$$

$$l_1^2 \sigma_{ji} - l_2^2 \delta_{ji} \sigma_1 = 0 \quad \dots (a)$$

$$l_2^2 \sigma_{ji} - l_1^2 \delta_{ji} \sigma_2 = 0 \quad \dots (b)$$

Below these equations, it is noted that l_1^2 and l_2^2 are complex.

$$l_1^2 \rightarrow \frac{l_1^2}{m_1^2} \quad l_2^2 \rightarrow \frac{l_2^2}{m_2^2}$$

Then, the direction cosines are defined as:

$$l_1 = a_1 + ib_1, \quad l_2 = c_1 + id_1, \quad l_3 = e_1 + if_1$$

Finally, the equations (a) and (b) are rewritten as:

$$(a) \rightarrow l_1^2 \sigma_{ji} - l_2^2 \delta_{ji} \sigma_1 = 0 \quad \dots (a')$$

Let us write that, the direction cosines corresponding to σ_1 is l_j super script 1. Similarly, for σ_2 , it is l_j super script 2. Now, this equation, you will have for each $j = 1, 2, 3$, $\sum_j l_j \sigma_j - \sum_j l_j \delta_j = 0$. Similarly, $\sum_j l_j \sigma_j - \sum_j l_j \delta_j = 0$. Now, these 2 equations a and b, these are all real quantities and here this is complex. So, also in this equation this is complex.

So, we are going to get a set of equations involving on complex coefficients. And it is natural to expect that, this l_{ij} and l_{ji} and l_{ij}^2 , roots of these two equations, must also be complex. So, l_{ij} and l_{ji} is complex; that means, if l_{ij} has components like $l_{ij} = l_{ij}^1 + i l_{ij}^2$. We can write this thing as l_{ij}^1 , l_{ij}^2 and n_{ij} . Then, each of these components must be complex. So, this could be that l_{ij}^1 is of the form $a_{ij} + i b_{ij}$, l_{ij}^2 is equal to $c_{ij} + i d_{ij}$ and this is let us say $e_{ij} + i f_{ij}$.

Now, if we try to take the equation 1 and take its complex conjugate. So, if you try to take complex conjugate of a, so it will look in this form $\overline{l_{ij}} \sigma_{ji} - \overline{l_{ji}} \delta_{ij} \sigma_2 = 0 \dots (a)$ $\delta_{ij} \sigma_1$ bar is equal to 0. So, that is the complex conjugate form of equation number 1.

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The image shows a whiteboard with handwritten mathematical derivations. The equations are as follows:

$$\overline{\sigma_1} = \sigma_2$$

$$\overline{l_{ij}} \sigma_{ji} - \overline{l_{ji}} \delta_{ij} \sigma_2 = 0 \dots (a)$$

$$\overline{l_{ij}} = l_{ij}^* \quad l_{ij}^* \text{ and } l_{ij}^2 \rightarrow \text{complex conjugate}$$

$$(a) \text{ \& } (b) \quad l_{ij}^* l_{ij}^2 (\sigma_1 - \sigma_2) = 0 \dots (c)$$

$$\sigma_1 = \alpha + i\beta \quad \sigma_2 = \alpha - i\beta$$

$$(c) \quad l_{ij}^* l_{ij}^2 2i\beta = 0$$

Now, this σ_1 bar is nothing but σ_2 . That is the stress σ_2 , because σ_2 is complex conjugate of σ_1 . So, equation a dash it now comes out to be of this form, $\overline{l_{ij}} \sigma_{ji} - \overline{l_{ji}} \delta_{ij} \sigma_2 = 0$. So, that is the form of a dash. Now, this is an equation, which is involving σ_2 here. So, it is natural to expect that, if you compare now these two equations a b and a dash. See, it is natural to expect that these two must be related.

So, this means that, $\overline{l_{ij}}$ must be equal to l_{ij}^2 . So, therefore, this comparison gives us that $\overline{l_{ij}}$ is equal to l_{ij}^2 . And hence, $\overline{l_{ij}}$ and l_{ij}^2 are again complex conjugates. Now, we go back to equation number a and b from a and b. If you

again multiply these equation by $l_j 2$ and these equation by $l_j 1$ and subtract. So, if you do the subtraction, so from a and b, we are going to get $l_j 1$, $l_j 2$ minus sigma 1 multiplied by sigma 1 minus sigma 2 equal to 0. So, this is equation.

So, therefore, we just consider these two, consider multiplying this thing by $l_j 2$ and this by $l_j 1$ and subtract the resultant equation out of here from the resultant equation from there, that will you give this. Now, if we assume that, sigma 1 is equal to alpha plus I beta. Since, it is complex and since, sigma 2 is complex conjugate of sigma 1. Therefore, sigma 2 can be written as alpha minus I beta. So, from c now, what we have is $l_j 1$, $l_j 2$ multiplied by 2 I beta equal to 0.

(Refer Slide Time: 43:44)

Handwritten derivation on a whiteboard:

$$\begin{aligned}
 (c) \quad l_j^1 l_j^2 &= 2i\beta = 0 \\
 l_j^1 &\Rightarrow (a_1 + ib_1) \quad (c_1 + id_1) \quad (e_1 + if_1) \\
 l_j^2 &\Rightarrow (a_1 - ib_1) \quad (c_1 - id_1) \quad (e_1 - if_1) \\
 a_1, c_1, e_1, b_1, d_1, f_1 &\rightarrow \text{real} \\
 l_j^1 l_j^2 &= l_1^1 l_1^2 + l_2^1 l_2^2 + l_3^1 l_3^2 \\
 &= (a_1^2 + b_1^2 + c_1^2 + d_1^2 + e_1^2 + f_1^2) \\
 &= \text{real} \\
 \beta &= 0 \quad \sigma_1 \text{ and } \sigma_2 \rightarrow \text{Cannot be complex} \\
 \text{Roots are real}
 \end{aligned}$$

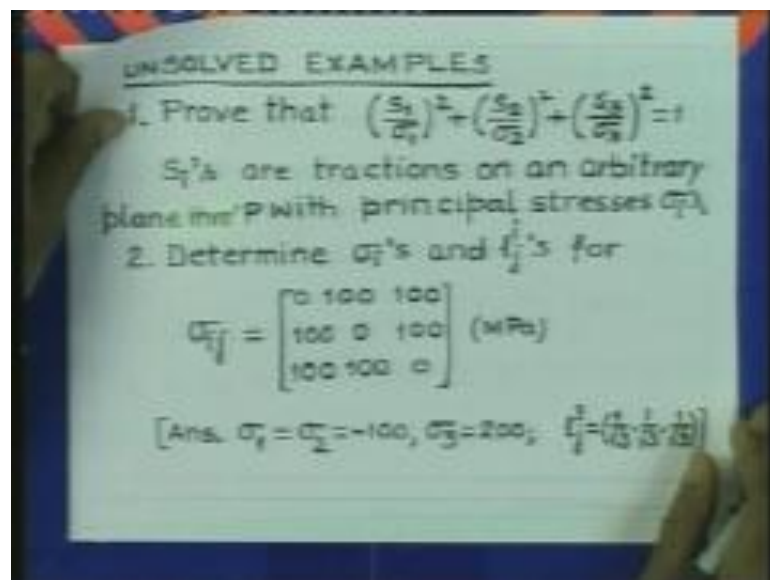
Now, since $l_j 1$ and $l_j 2$ are complex conjugates. Since, we can write now the coefficients of $l_j 1$ as a_1 plus I b_1 . And the components are these, this is c_1 plus I d_1 and this is e_1 plus I f_1 . Similarly, $l_j 2$ components are, since they are complex conjugates. So, therefore, it is going to be a_1 minus I b_1 , c_1 minus I d_1 and this is e_1 minus I f_1 . Now, these a_1 , b_1 , c_1 , d_1 , e_1 , f_1 they are all real. So, start a_1 , c_1 , e_1 . So, also b_1 , d_1 and f_1 are all real quantities.

So, now $l_j 1$ into $l_j 2$ should be equal to nothing but this multiplied by these plus this multiplied these plus this multiplied by these. So, which is nothing but $l_1 1$, $l_1 2$ plus, it is $l_2 1$, $l_2 2$, plus $l_3 1$, $l_3 2$. So, this will be, thus a_1^2 plus b_1^2 plus c_1^2

square plus d 1 square plus e 1 square plus f 1 square. Now, this expression is always some positive real value.

So, from equation c now, from equation c, that we have, it is quite clear. That beta, this product is real and this 2 I beta is product is 0. So, therefore, this is real and this simply means that, beta has got to be 0. So, this relation indicates that beta is equal to 0. And we assume that sigma 1 equal to alpha plus I beta sigma 2 equal to alpha minus I beta. It simply means that, sigma 1 and sigma 2 cannot be complex. So, the roots of the characteristic equation involving the principal stresses are real. So, roots are, thus we can show that, the other roots are also going to be real.

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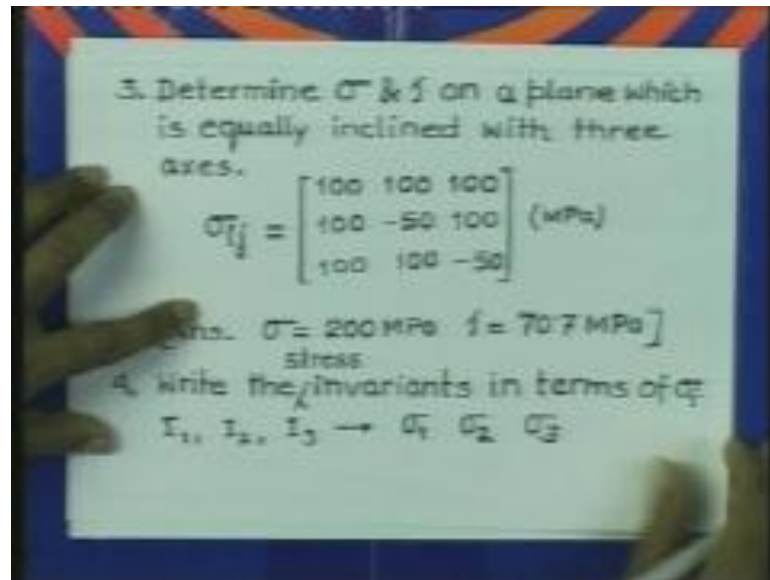


You can now consider, solving some problems yourself. You consider this example, prove that, $s_1^2/\sigma_1^2 + s_2^2/\sigma_2^2 + s_3^2/\sigma_3^2 = 1$. Where s_i 's are the traction on an arbitrary plane through the point p with principal stresses σ_i 's. So, $\sigma_1, \sigma_2, \sigma_3$ are the principal stresses at a point p. And s_1, s_2, s_3 is the tractions on an arbitrary plane inclined with the three principal directions; prove that, this relation is equal to unity.

Second example, you can consider determination of principal stresses and there direction cosines, where the stresses in Cartesian coordinates are given by this expression 0, 100, 100, 100, 0, 100, 100, 100, 0 MPa. And the answer for this case is going to be sigma 1 equal to sigma 2 equal to minus 100 MPa, sigma 3 equal to 200 MPa. And the direction

cosines connected with the sigma 3 stresses is 1 by root 3, 1 by root 3, 1 by root 3. Similarly, determine the stresses sigma and tau on a plane, which is equally inclined with the three Cartesian axes. And here the stress tensor is given by 100, 100, 100 MPa, 100 minus 50, 100 MPa 100 100 minus 50 MPa.

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And this answer for this case, the normal stress is given by 200 MPa and the shear stress is 70.7 MPa. You can consider the fourth example, you write the stress invariants, in terms of the principal stresses sigma I. That is you write the expressions for I 1, I 2, I 3; in terms of the principal stresses sigma 1 sigma 2 and sigma 3.

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stress
4. Write the invariants in terms of $\sigma_1, \sigma_2, \sigma_3$

$I_1, I_2, I_3 \rightarrow \sigma_1, \sigma_2, \sigma_3$

Ex 4

$$I_1 = \sigma_1 + \sigma_2 + \sigma_3$$

$$I_2 = \sigma_{xy}^2 + \sigma_{yz}^2 + \sigma_{zx}^2 - (\sigma_x \sigma_y + \sigma_y \sigma_z + \sigma_z \sigma_x)$$

$$= -(\sigma_1 \sigma_2 + \sigma_2 \sigma_3 + \sigma_3 \sigma_1)$$

$$I_3 = \sigma_x \sigma_y \sigma_z + 2\sigma_{xy} \sigma_{yz} \sigma_{zx} - \sigma_x^2 \sigma_{yz}^2 - \sigma_y^2 \sigma_{zx}^2 - \sigma_z^2 \sigma_{xy}^2$$

$$= \sigma_1 \sigma_2 \sigma_3$$

I would like to consider the solution for this example 4. So, example 4, we have I_1 , it is directly related to the three principal stresses. And their sum is going to be the first stress invariant. Now, second stress invariant, I would like you to consider the relationship. We had $\tau_{xy}^2 + \tau_{yz}^2 + \tau_{zx}^2$ and the third component minus $\sigma_x \sigma_y + \sigma_y \sigma_z + \sigma_z \sigma_x$.

Now, in this case, the shear stresses are all 0. Since, we are working in terms of the principal directions. Therefore, we are going to get directly σ_x to be σ_1 , σ_y to be σ_2 . And therefore, this is given by minus $\sigma_1 \sigma_2 + \sigma_2 \sigma_3 + \sigma_3 \sigma_1$. Similarly, the third stress invariant is given by $\sigma_x \sigma_y \sigma_z + 2\tau_{xy} \tau_{yz} \tau_{zx} - \sigma_x^2 \tau_{yz}^2 - \sigma_y^2 \tau_{zx}^2 - \sigma_z^2 \tau_{xy}^2$.

And since, in this case shear stresses are all absent. And the σ_x is nothing but σ_1 , this is σ_2 , this is σ_3 . Therefore, the third stress invariant is given by product of the three principal stresses, σ_1, σ_2 and σ_3 . So, the three principal, three stress invariants are going to be very simple expressions, I_1 is sum of the 3 principal stresses. I_2 is nothing but minus product of the two principal stresses, taken in succession and I_3 is nothing but product of the three principal stresses.