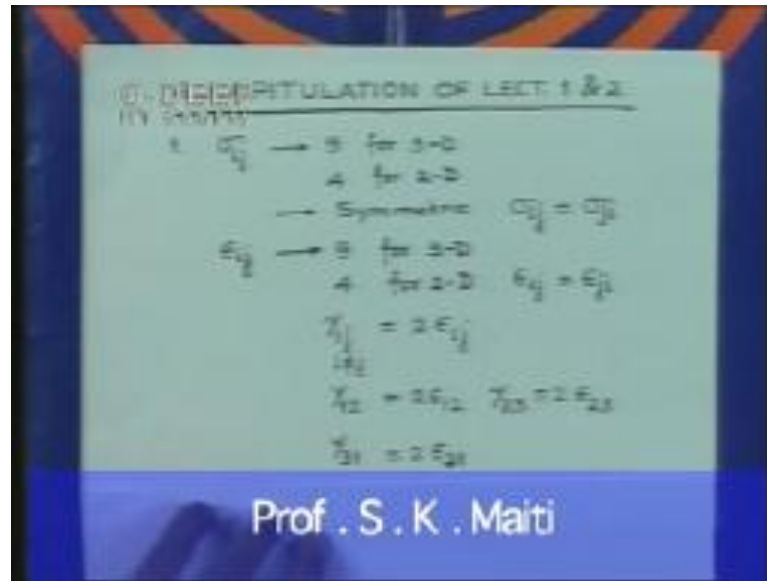


Advanced Strength of Materials
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Lecture – 3

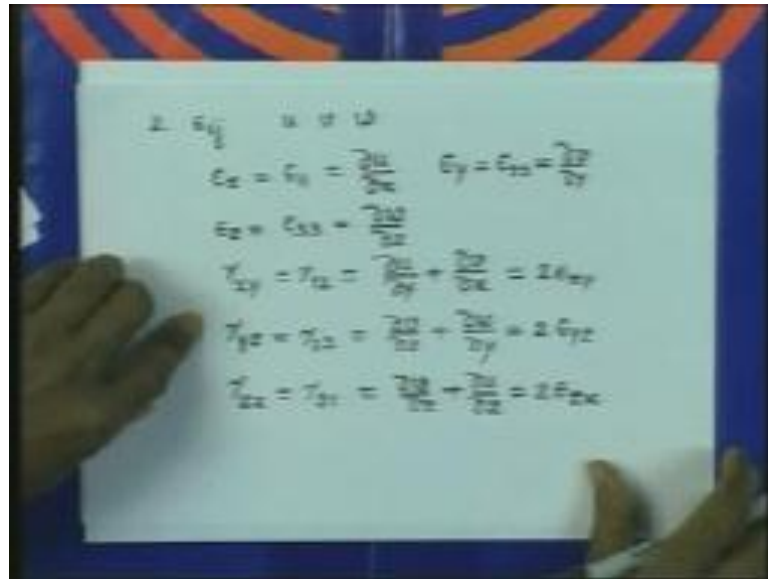
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We will recapitulate fast; what we have done in the earlier 2 lectures. We have first talked about the stress at a point. It is indicated by sigma ij, it is a 9 component tensor for 3 dimensions and it has 4 components for 2 dimensions. And we have also seen that this tensor is symmetric, thereby you can indicate it in the following form: sigma ij equal to sigma ji. Similarly, we have introduced the strain tensor which is also a tensor having 9 components for 3 dimensions and 4 components for 2 dimension. It is also symmetric tensor therefore, epsilon ij equal to epsilon ji.

We have also seen that, there is a difference between engineering shear strain and tensor shear strain. Engineering shear strain is twice the tensor shear strain. Specifically we have seen gamma 1 2 is equal to 2 times epsilon 1 2, similarly same relationship same type of relationship holds good for eps gamma 2 3 and gamma 3 1.

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Handwritten equations on a piece of paper:

$$\begin{aligned} \epsilon_x &= \epsilon_{11} = \frac{\partial u}{\partial x} & \epsilon_y &= \epsilon_{22} = \frac{\partial v}{\partial y} \\ \epsilon_z &= \epsilon_{33} = \frac{\partial w}{\partial z} \\ \gamma_{xy} &= \gamma_{12} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} = 2\epsilon_{xy} \\ \gamma_{yz} &= \gamma_{23} = \frac{\partial v}{\partial z} + \frac{\partial w}{\partial y} = 2\epsilon_{yz} \\ \gamma_{zx} &= \gamma_{31} = \frac{\partial w}{\partial x} + \frac{\partial u}{\partial z} = 2\epsilon_{zx} \end{aligned}$$

We have then define the strains at a point in terms of the 3 displacement components: u , v and w . These are the Cartesian components of displacements of a point on the body and we have seen that, ϵ_x is nothing but ϵ_{11} , it is given by the partial derivative $\frac{\partial u}{\partial x}$. Similarly, ϵ_y is nothing but ϵ_{22} , this is given by $\frac{\partial v}{\partial y}$ and ϵ_z which is nothing but ϵ_{33} , this is given by the partial derivative of w with respect to z . Then we have seen the shear strains; shear strain if it is γ_{xy} which is nothing but γ_{12} , this is given by $\frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}$.

Similarly, γ_{yz} is equal to γ_{23} which is nothing but $\frac{\partial v}{\partial z} + \frac{\partial w}{\partial y}$. Lastly we had γ_{zx} is equal to γ_{31} which is nothing but $\frac{\partial w}{\partial x} + \frac{\partial u}{\partial z}$. We have also noted that, this component is equal to 2 times ϵ_{xy} , this is 2 times ϵ_{yz} and this is also 2 times ϵ_{zx} .

(Refer Slide Time: 05:08)

3. $S_i = \sigma_{ji}$ Plane: l_j

$$S_i = l_j \sigma_{ji} \quad \begin{matrix} i=1,2,3 \\ j=1,2,3 \end{matrix}$$

Summation

$$S_i = l_1 \sigma_{i1} + l_2 \sigma_{i2} + l_3 \sigma_{i3}$$

Calculation of E_{ij} — Example

$P \rightarrow$ displacement

$$[i(x^2 - yz) + j(2y - xz^2) + k(2z - x^2y)] \times 10^{-3}$$

$$u = \frac{\partial}{\partial x}(x^2 - yz) = 2x \quad v = \frac{\partial}{\partial y}(2y - xz^2) = 2 - xz^2 \quad w = \frac{\partial}{\partial z}(2z - x^2y) = 2 - x^2y$$

Then we took up the calculations of tractions on an arbitrary plane, so traction component on an arbitrary plane, where the stresses are given by σ_{ji} . And if the arbitrary plane has direction cosines l_j then in that case, we have seen that S_i is nothing but $l_j \sigma_{ji}$ varying S_i of the components of traction and in this we have made use of the implied condition that i takes up value 1 to 3. So, also j takes up value 1 to 3. This indicates compact form of expanded expressions and herein we make use of the fact that, these 2 symbols are repeated therefore, a summation is indicated.

So, these 2 symbols are repeated and it indicates a summation. Just to the 1 case S_1 is nothing but $l_1 \sigma_{11}$ plus $l_2 \sigma_{21}$ plus $l_3 \sigma_{31}$, herein i is fixed here this is fixed and l is varying from 1 2 3. So, also the stress components are varying 1 2 3. Now, we would like to take 1 example of we would like to consider the calculation of strain at a point of ϵ_{ij} . It is given that in a body which is undergoing deformation, a point P has displacement given by the vector i into x^2 minus yz plus j 2 y minus xz^2 plus k into $3z$ minus x^2y square. This multiplied by 10^{-3} meter. So, the displacements are given in milli meter. Herein i, j and k are the unit vectors in the 3 Cartesian directions.

Therefore, you should have no difficulty in identifying that the x displacement of the typical point is given by x^2 minus yz and v displacement is equal to $2y$ minus xz^2 square and w is given by $3z$ minus x^2y square. And all this expressions are

multiplied by 10 to the power minus 3. So, these are the displacements. Now, using these displacements, we would like to now calculate the strains strain components particularly.

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The image shows a piece of paper with handwritten mathematical derivations for strain components. At the top, it states 'Displacement' and gives the displacement vector $\mathbf{u} = (x^2 - yz)\mathbf{i} + (xy - xz^2)\mathbf{j} + (yz - x^2y^2)\mathbf{k}$ in mm. Below this, the displacement components are listed: $u = (x^2 - yz) \times 10^{-3}$, $v = (xy - xz^2) \times 10^{-3}$, and $w = (yz - x^2y^2) \times 10^{-3}$. The normal strain components are then calculated as follows:

$$\epsilon_x = \frac{\partial u}{\partial x} = 2x \times 10^{-3} \text{ (mm/mm)}$$

$$\epsilon_y = \frac{\partial v}{\partial y} = 2 \times 10^{-3}$$

$$\epsilon_z = \frac{\partial w}{\partial z} = 2 \times 10^{-3}$$

The shear strain components are calculated using the formula $\gamma_{xy} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x}$:

$$\gamma_{xy} = \frac{\partial u}{\partial y} + \frac{\partial v}{\partial x} = (-z - xz)\mathbf{i} \times 10^{-3}$$

$$\gamma_{yz} = \frac{\partial v}{\partial z} + \frac{\partial w}{\partial y} = (-2xz - 2x^2y)\mathbf{j} \times 10^{-3}$$

$$\gamma_{zx} = \frac{\partial w}{\partial x} + \frac{\partial u}{\partial z} = (-2xy^2 - y)\mathbf{k} \times 10^{-3}$$

So, if I now consider epsilon x, this is given by delta u delta x. So, this will be given by twice x multiplied by 10 to the power minus 3. So, this is milli meter by milli meter. So, strains are dimensionless. So, hence forth I will not. So, the units epsilon y delta v delta y this is given by 2 10 to the power minus 3 and w is equal to w epsilon, the strain components epsilon z delta w delta z is equal to 3 multiplied by 10 to the power minus 3. Similarly, we can calculate the shear strains. Shear strains gamma xy delta u delta y plus delta v delta x, so this is given by minus z minus z square into 10 to the power minus 3 gamma yz delta v delta z plus delta w delta x.

So, this is going to be given by minus 2 xz minus 2 xy square, this should be delta w delta y, so it should be 2 x square y, so 2 x square y. And gamma zx is nothing but delta w delta x plus delta u delta z. This is given by minus 2 xy square and this is minus y into 10 to the power minus 3, here also you will have a multiplier 10 to the power minus 3. It is given specifically to calculate the strains at a point. The point P coordinates are given as 2 3 4 in milli meters unit.

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$$\begin{bmatrix} 4 & -20 & -39 \\ -20 & 2 & -40 \\ -39 & -40 & 3 \end{bmatrix} \times 10^{-3} \left(\frac{\text{mm}}{\text{mm}} \right)$$

$$\hookrightarrow \text{Engg Strains}$$

$$\begin{bmatrix} 4 & -10 & -19.5 \\ -10 & 2 & -20 \\ -19.5 & -20 & 3 \end{bmatrix} \times 10^{-3}$$

$$\hookrightarrow \epsilon_{ij}$$

So, if you now substitute the value of this coordinates, then we get the following strains; 4 minus 20 minus 39 minus 20 2 minus 40 minus 39 minus 40 3 and all these has a common multiplier 10 to the power minus 3, this is milli meter by milli meter. So, these are the engineering strains at the point P with coordinates 2 3 4. If you want the tensor components, it is going to be 4; this will be half of what we have here, so it is 10 minus 19.5 minus 10 2 minus 20 minus 19.5 minus 20 3 multiplied by 10 to the power minus 3, this is the strain tensor at point P.

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CALCULATION OF TRACTION - EXAMPLE

$$\vec{\sigma} = \begin{bmatrix} 100 & 0 & 0 \\ 0 & 200 & 0 \\ 0 & 0 & -100 \end{bmatrix} \text{ (MPa)}$$

$$S_i \text{ on plane with d.c.'s } (0.707, 0.707)$$

$$\text{Soln: } l_1 = 0.707 \quad l_2 = 0.707 \quad l_3 = 0$$

$$S_i = l_1 \sigma_{11} + l_2 \sigma_{22} + l_3 \sigma_{33}$$

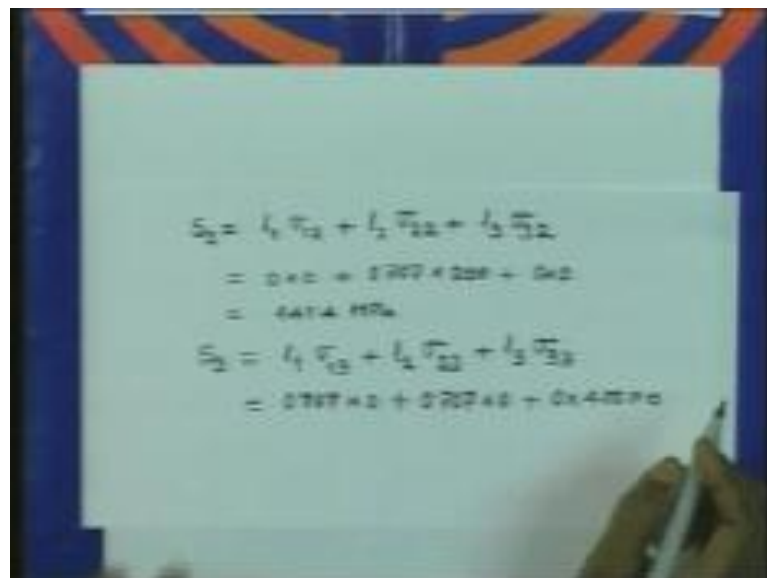
$$= 0.707 \times 100 + 0.707 \times 200 + 0 \times 0$$

$$= 70.7 \text{ MPa}$$

We would like to consider 1 example; calculation of tractions. Given the stress tensor at a point like this; 100 0 0 0 200 0 0 400, these are the stresses at a point. Now, it is necessary to calculate the tractions S_i on plane, with direction cosines 0.707 0.707 0. This means; so a solution can be done like this; we have l_1 is nothing but 0.707, l_2 0.707 and l_3 is equal to 0. So, these are the direction cosines.

Therefore, we can now calculate the traction component 1, which is given by $l_1 \sigma_{11} + l_2 \sigma_{21} + l_3 \sigma_{31}$, herein l_1 is 0.707 and the σ_{11} is 100, then l_2 is also 0.707 and σ_{21} is 0 and this l_3 0 and σ_{31} is 0. So, therefore, it gives us the value like this 70.7 MPa. So, that is the traction in the direction 1 and it has the same units as the stress.

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The image shows a person's hand holding a white marker, writing on a piece of paper. The paper has two equations written on it. The first equation is for S_1 and the second is for S_2 . The background is a blue and orange striped pattern.

$$S_1 = l_1 \sigma_{11} + l_2 \sigma_{21} + l_3 \sigma_{31}$$

$$= 0.707 \times 100 + 0.707 \times 200 + 0 \times 400$$

$$= 141.4 \text{ MPa}$$

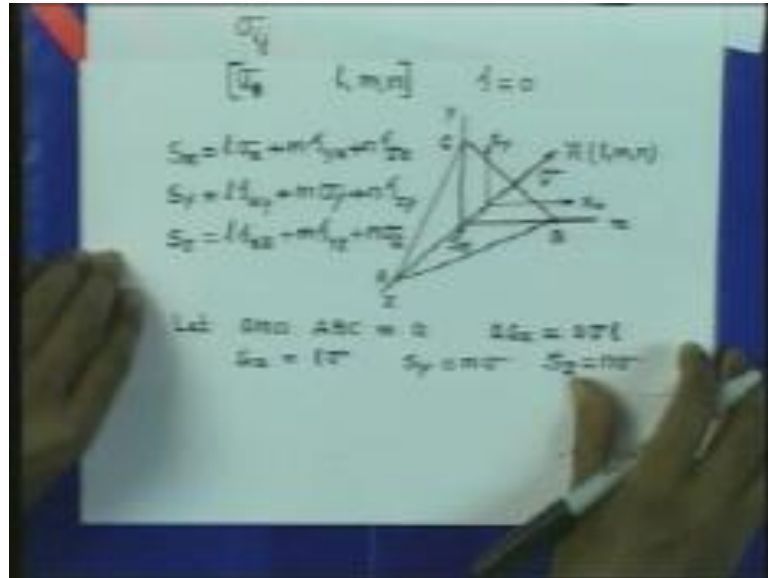
$$S_2 = l_1 \sigma_{12} + l_2 \sigma_{22} + l_3 \sigma_{32}$$

$$= 0.707 \times 0 + 0.707 \times 0 + 0 \times 400 = 0$$

You can calculate similarly the tractions for the other 2 directions. S_2 is given by $l_1 \sigma_{12} + l_2 \sigma_{22} + l_3 \sigma_{32}$ and if you substitute the value it becomes 0 0, this is 0.707 multiplied by 200 0 multiplied by 0. So, this gives us 141.4 MPa. Now, S_3 is given by $l_1 \sigma_{13} + l_2 \sigma_{23} + l_3 \sigma_{33}$, where l_1 0.707, σ_{13} 0, this is 0.707, this is 0 and this is 0 multiplied by 400. So, therefore, this is 0. Please note that, this value should not be 0; it should be 0.707. So, this amendment you can note, but the value of S_2 does not change.

So, the 3 traction components are; finally you have S_1 is 70.7 MPa, S_2 is 141.4 MPa, S_3 is 0. So, these are the 3 traction components on the plane, whose direction cosines are 0.707 0.707 and 0.

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Now, we would like to consider determination of principal planes and principal stresses. So, determination of principal stresses and principal planes; do you remember what is principal stress, these are the stresses which act on plane where there is no shear stress. So, we would like to now consider determination of some, determination of such stresses in 3 dimensions. So, we have the stresses at a point given by σ_{ij} and we would like to find out the stresses, which are acting on planes, where there is no shear stress.

So, we are interested in determine determining the stresses like; let us say σ_1 and its direction cosines l m and n . So, we will indicate the principal stresses by symbol; just a indicate by σ , so σ is the principal stress and it is acting on a plane with direction cosines l m and n . And this plane is such that, there is no shear stress, so shear stress on this plane equal to 0. So, we are really interested in determining this set, where there is a stress equal to 0.

Now, we try to consider 1 plane: A B C. Let us say this is the plane, whose outer normal is directed like this, this is the outer normal. Let us assume that, this is the plane which has direction cosines l m and n . And therefore, on this plane we have a stress σ acting and this plane A B C is divide of any shear stress.

Now, if I make use of the Cauchy formula then the traction components we can obtain, so we will like to guide the traction components here; these are the 3 traction components for this plane this is S_x , this is S_y and this component is S_z . So, we can write now S_x is equal to $l \sigma_x$ plus $m \tau_{yx}$ plus $n \tau_{zx}$ similarly, S_y is equal to $l \tau_{xy}$ plus $m \sigma_y$ plus $n \tau_{zy}$ S_z is equal to $l \tau_{xz}$ plus $m \tau_{yz}$ plus $n \sigma_z$ these are the tractions on these plane ABC.

Let, the area of this plane ABC be represented by the symbol a , now the force which is going to act in the x direction, because of the stress stresses on the plane is going to be nothing but a into S_x . And we can also, find out the component from the stress sigma the total force on this plane ABC is nothing but a into sigma and each component in the x direction is l sigma into a .

So therefore, we can write that the total force in the x direction is nothing but $a x$ and this should also be equal to a into sigma multiplied by a . So finally, we get S_x is equal to l into sigma similarly, you can have consideration of force in the y direction there by you will find that S_y is nothing but m into sigma and S_z is equal to n into sigma. Now, we will substitute the value of S_x S_y S_z in terms of sigma.

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$$\left. \begin{aligned} l(\sigma_x - \sigma) + m\tau_{yx} + n\tau_{zx} &= 0 \\ l\tau_{xy} + m(\sigma_y - \sigma) + n\tau_{zy} &= 0 \\ l\tau_{xz} + m\tau_{yz} + n(\sigma_z - \sigma) &= 0 \end{aligned} \right\} \text{--- (1)}$$

$$\det \begin{vmatrix} \sigma_x - \sigma & \tau_{yx} & \tau_{zx} \\ \tau_{xy} & \sigma_y - \sigma & \tau_{zy} \\ \tau_{xz} & \tau_{yz} & \sigma_z - \sigma \end{vmatrix} = 0 \quad \text{--- (2)}$$

$$\textcircled{3} \quad (I_j \sigma_j - I_i \delta_{ji} \sigma) = 0 \quad \text{--- (3)}$$

$\delta_{ji} = 0 \text{ if } i \neq j$
 $= 1 \text{ if } i = j$ $\swarrow$ Kronecker delta

So, if I do that we finally get the following that $l \sigma_x$ minus sigma plus $m \tau_{yx}$ plus $n \tau_{zx}$ equal to 0. So what I have done is that, I have coupled this relationships S_x and

these S_x and simplified similarly, we get from the second expression $m \sigma_y - n \tau_{zy} = 0$.

Similarly, $l \tau_{xz} + m \tau_{yz} + n \sigma_z - \sigma = 0$; this shape of equations let us represent it by set 1. Now, we say homogeneous set of equations with unknown l, m and n the direction cosines of the plane and if we have to have solutions for non 0 solutions for l, m and n it is necessary to have this condition satisfied determinant of the coefficient matrix which is nothing but $\sigma_x - \tau_{yx} \tau_{zx} \tau_{xy} \sigma_y - \tau_{zy} \tau_{xz} \tau_{yz} \sigma_z - \sigma$.

So this determinant must vanish, so it should be equal to 0 so this is the characteristic equation. This equation can be written in tensor notation particularly, we can use also the tensor notation to represent equation number 1 in the following form that $I_j \sigma_{ji} - I_j \delta_{ji} \sigma = 0$, where this δ_{ji} is nothing but kronecker delta. So, this kronecker delta has the characteristics, that δ_{ji} is equal to 0 when i is not equal to j and it is equal to unity when i equal to j . So, with those implications we can write the equation number 1 in this form, so this is really the form of equation number 1 so will indicate equation number 1 dash.

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$$(1) \rightarrow \det \left| \sigma_{ij} - \delta_{ij} \sigma \right| = 0 \dots (1)$$

$$(2) \rightarrow \sigma^3 - I_1 \sigma^2 - I_2 \sigma - I_3 = 0 \dots (2)$$

$$I_1 = \sigma_x + \sigma_y + \sigma_z$$

$$I_2 = \sigma_{xx}^2 + \sigma_{yy}^2 + \sigma_{zz}^2 - \sigma_x \sigma_y - \sigma_y \sigma_x - \sigma_x \sigma_z - \sigma_z \sigma_x - \sigma_y \sigma_z - \sigma_z \sigma_y$$

$$I_3 = \sigma_x \sigma_y \sigma_z + 2 \sigma_{xy} \sigma_{yz} \sigma_{zx} - \sigma_x \sigma_{yz}^2 - \sigma_y \sigma_{zx}^2 - \sigma_z \sigma_{xy}^2$$

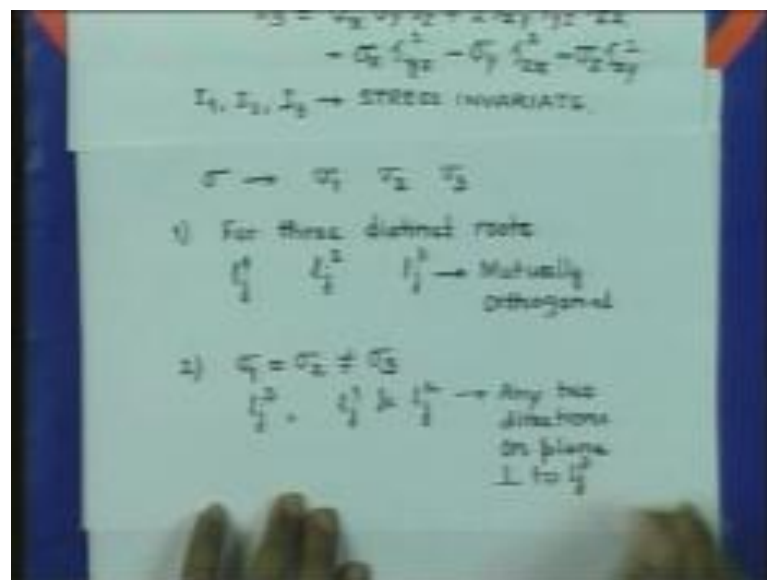
$I_1, I_2, I_3 \rightarrow$ STRESS INVARIANTS

Similarly the equation number 2 is nothing but determinant $\sigma_{ji} - \delta_{ji} \sigma = 0$, so this is in a compact form we can have. So we have seen that the equation number 2 can be written in tensor form like this. If we go back to the original

form of equation 2 this can be expanded and will get an equation in sigma, which is obvious form sigma cube minus I1 sigma square minus I2 sigma minus I3 equal to 0, where these coefficients I1 I2 and I3 are given in terms of the stresses I1 is equal to sigma x plus sigma y plus sigma z and I2 is tau xy square tau yz square plus tau zx square minus sigma x minus sigma y minus sigma y minus sigma z minus sigma z into sigma x.

Similarly, I3 is given by sigma x into sigma y into sigma z plus 2 into tau xy tau yz into tau zx minus sigma x into tau yz square minus sigma y tau zx square minus sigma z into tau xy square.

(Refer Slide Time: 37:18)



Let us, say the 3 roots are sigma 1 sigma 2 and sigma 3 if these 3 roots are distinct, then they are going to work on 3 mutually perpendicular planes. So far 3 distinct roots will have 3 directions which are going to be mutually perpendicular. So, if you consider the directions of the 3 stresses as l_1 for the sigma 1 stress l_2 for the sigma 2 stress and l_3 for the sigma 3 stress then these 3 directions are going to be mutually perpendicular.

So, these are 3 mutually orthogonal directions, if we have a situations like that sigma 1 there are 2 repeated roots sigma 1 equal to sigma 2, but then that is not equal to sigma 3. Then, you can find out this sigma 3 directions and then l_1 and 2 can be taken as any 2 directions on the plane perpendicular to l_3 . So, these are the coefficients these are called stress invariants I2 I3 are known as stress invariants.

So, this equation will have 3 roots for sigma and it can be shown that these roots are real. So I_1, I_2 which corresponds to sigma 1 and sigma 2 they will be any 2 directions any 2 directions on plane perpendicular to I_3 .

(Refer Slide Time: 40:20)

3) $\sigma_1 = \sigma_2 = \sigma_3$
 Any three orthogonal directions
 are principal directions.
Solution of Eqn. (2)
 $\sigma_1 = \frac{I_1}{3} + 2\sqrt{\frac{J_2}{3}} \cos(\phi/3)$
 $\sigma_2 = \frac{I_1}{3} + 2\sqrt{\frac{J_2}{3}} \cos(120^\circ + \frac{\phi}{3})$
 $\sigma_3 = \frac{I_1}{3} + 2\sqrt{\frac{J_2}{3}} \cos(240^\circ + \frac{\phi}{3})$
 $a = -\left(\frac{I_2}{3} + \frac{2I_3^2}{9}\right)$

If we have 3 roots repeated, if we have sigma 1 equal to sigma 2 equal to sigma 3 then any 3 orthogonal directions or principal directions. Now, we can find out the roots of this equation 3 easily so we can solve the equation. So, solution of equation 3 you can solve this cubic equation by the following steps in fact we can write sigma 1 to be given by $I_1/3 + 2\sqrt{J_2/3} \cos(\phi/3)$ sigma 2 is $I_1/3 + 2\sqrt{J_2/3} \cos(120^\circ + \phi/3)$ and sigma 3 is equal to $I_1/3 + 2\sqrt{J_2/3} \cos(240^\circ + \phi/3)$, where we have a is given by $-\left(\frac{I_2}{3} + \frac{2I_3^2}{9}\right)$.

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$$\cos \phi = -\frac{b}{2\sqrt{-a^3/27}}$$

$$b = -(I_3 + 2\frac{I_1^3}{27} - \frac{I_1 I_2}{3})$$

STRESS INVARIANTS (I_1, I_2, I_3)

$$I_1^3 - I_1^2 I_2 + I_1 I_2^2 - I_3 = 0$$

$$\sigma_{ij} \rightarrow x, y, z \quad \sigma_{ij}'$$

$$\sigma^3 - I_1 \sigma^2 + I_2 \sigma - I_3 = 0$$

$I_1, I_2, I_3 \rightarrow \text{Independent of coordinates}$

And we have cosine phi is equal to minus b by square root minus a cube by 27 and this b is given by minus I3 plus 2 I1 cube by 27 plus I1 I2 by 3. So, the solution for that equation cubic equation are given by this expressions wherein all the constants a phi and b are related to the stress invariants. Let us understand, why stress invariants? So, we have found that the quantities like I1, I2, and I3 are stress invariants.

So, you should have no difficulty in appreciating that when a body is loaded its deformation will take place and at a point there will be principal directions, there will at each point there will be 3 principal directions. And these principal directions are fixed so long as the loading on the body is fixed, the principal directions are fixed directions that means; the 3 directions l_1, l_2, l_3 which corresponds to stresses sigma 1, sigma 2 and sigma 3 they are fixed and they are dependent on the stresses stress tensor at the point.

Now, the direction of this principal stresses are fixed at a point. This will not depend on the selection of coordinates for defining the stress tensor. So, this stress tensor σ_{ji} is going to depend on selection of the coordinate ,x y, z, but the value of the principal stresses and there directions are fixed which implies that the cubic equation that we had retain little while ago that sigma cube I1 sigma square minus I2 sigma minus I3 equal to 0.

The roots of these equations must be unique, it will not depend on the selection of coordinate. If the 2 persons are selecting 2 different coordinates, it is going to probably

have values of σ_{ji} for 1 person and for another person it may have value σ_{ji} dash, but this 2 situations must lead to a equation in σ which should not depend on the coordinates. Therefore, this I_1 I_2 I_3 must be independent of the coordinate. So I_1 , I_2 , I_3 independent of coordinates; that is why they are called stress invariants.

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The image shows handwritten mathematical derivations for the first three stress invariants. At the top, it states: $I_1, I_2, I_3 \rightarrow$ Independent of coordinates. Below this, the derivations are as follows:

$$I_1 = \sigma_x + \sigma_y + \sigma_z$$

$$= \sigma_x' + \sigma_y' + \sigma_z'$$

$$I_2 = \sigma_{xy}^2 + \sigma_{yz}^2 + \sigma_{zx}^2 - (\sigma_x \sigma_y + \sigma_y \sigma_z + \sigma_z \sigma_x)$$

$$= \sigma_{xy}^2 + \sigma_{yz}^2 + \sigma_{zx}^2 - (\sigma_x' \sigma_y' + \sigma_y' \sigma_z' + \sigma_z' \sigma_x')$$

$$I_3 = \sigma_x \sigma_y \sigma_z + 3\sigma_{xy} \sigma_{yz} \sigma_{zx} - (\sigma_x^3 + \sigma_y^3 + \sigma_z^3)$$

$$= \sigma_x' \sigma_y' \sigma_z' + 3\sigma_{xy}' \sigma_{yz}' \sigma_{zx}' - (\sigma_x'^3 + \sigma_y'^3 + \sigma_z'^3)$$

And you can therefore, finally write that the value of I_1 which is given by σ_x plus σ_y plus σ_z it must be also equal to σ_x dash σ_y dash plus σ_z dash. So also, we will have I_2 given by τ_{xy}^2 plus τ_{yz}^2 plus τ_{zx}^2 minus $\sigma_x \sigma_y$ minus $\sigma_y \sigma_z$ plus $\sigma_z \sigma_x$ this should also be equal to τ_{xy} dash square τ_{yz} dash square plus τ_{zx} dash square plus σ_x dash σ_y dash plus σ_y dash σ_z dash plus σ_z dash σ_x dash.

Similarly, we can write the third stress invariant also along similar line $\sigma_x \sigma_y \sigma_z$ plus $\tau_{xy} \tau_{yz} \tau_{zx}$ minus $\sigma_x \tau_{yz}^2$ plus $\sigma_y \tau_{zx}^2$ plus $\sigma_z \tau_{xy}^2$ it is also, equal to σ_x dash σ_y dash σ_z dash plus τ_{xy} dash τ_{yz} dash τ_{zx} dash minus σ_x dash square τ_{yz} dash square plus σ_y dash into τ_{yz} dash square plus σ_z dash into τ_{xy} dash square. So, the expression for the quantities I_1 , I_2 , I_3 are going to come to the same value; that is why they are not dependent on the coordinates and it is because of this reason they are called stress invariants.