

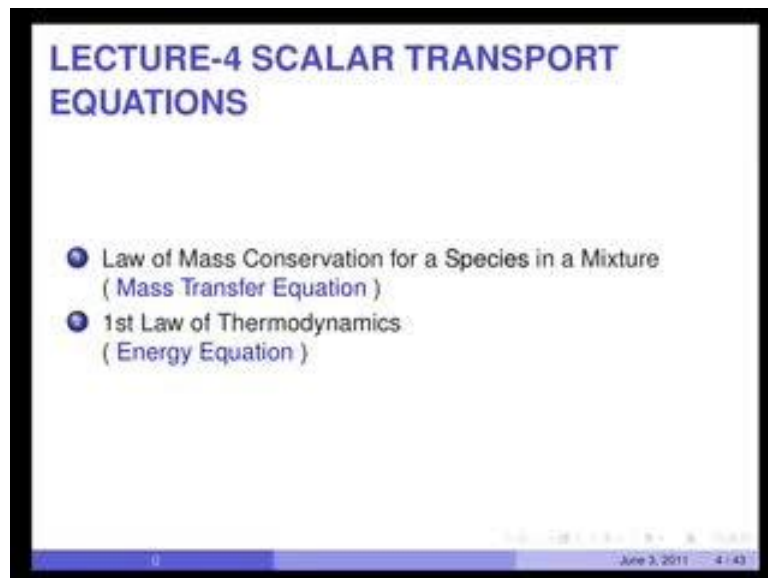
**Convective Heat and Mass Transfer**  
**Prof. A.W. Date**  
**Department of Mechanical Engineering**  
**Indian Institute of Technology, Bombay**

**Module No. # 01**  
**Lecture No. # 04**  
**Scalar Transport Equations**

In the last lecture, we saw equations of fluid motion namely the law of mass conservation applied to the bulk fluid as well as the second law of motion also called the Navier-Stokes equations.

Moving fluids also carry with them scalar quantities. Today, we are going to look at equations that govern transport of scalar quantities.

(Refer Slide Time: 00:47)



The first scalar quantity is the specie in a mixture. For example, if you are dealing with a combustion problem then the species carried are oxygen, carbon dioxide, fuel, carbon monoxide and so on and so forth.

The law of mass conservation for a specie in a mixture is called the mass transfer equation. We will also invoke the first law of thermodynamics which transports energy in a moving fluid. Both energy and the concentration of the specie are scalar quantities.

(Refer Slide Time: 01:31)

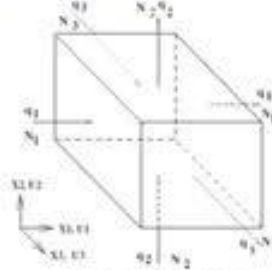
### Mass Transfer Equation - L4( $\frac{1}{17}$ )

For Specie k in a Mixture  
 Rate of accumulation of mass ( $M_{k,acc}$ ) =  
 Rate of mass in ( $M_{k,in}$ )  
 - Rate of mass out ( $M_{k,out}$ )  
 + Rate of generation within CV ( $R_{k,cv}$ )

$$M_{k,acc} = \frac{d(\rho_k \Delta V)}{dt}$$

$$M_{k,in} = N_{1,k} \Delta A_1 |_{x_1} + N_{2,k} \Delta A_2 |_{x_2} + N_{3,k} \Delta A_3 |_{x_3}$$

$$M_{k,out} = N_{1,k} \Delta A_1 |_{x_1+\Delta x_1} + N_{2,k} \Delta A_2 |_{x_2+\Delta x_2} + N_{3,k} \Delta A_3 |_{x_3+\Delta x_3}$$



$N$  = TOTAL MASS FLUXES  
 $q$  = TOTAL HEAT FLUXES

$\rho_k$  = Specie Density  
 Substitute, Divide each term by  $\Delta V$  and Let  $\Delta x_1, \Delta x_2, \Delta x_3 \rightarrow 0$

March 16, 2010 21/33

What does the law say for a specie? For a specie k in a mixture, the rate of accumulation of mass of the specie k that is given by  $M \cdot k$  accumulation equals rate of mass in minus rate of mass out plus rate of generation of specie within the control volume.

As a result of chemical reaction, as you know some species are generated and while others are destroyed. As such each specie will have either a generation or a destruction rate associated with it. What is  $M \cdot k$  accumulation? That is simply the mass of the specie k within the control volume  $\Delta v - d$  by  $dt$  of  $\rho_k \Delta v$ .

The mass of specie in will be the flux of specie k in direction 1 multiplied by the area  $\Delta A_1$  which is this area plus  $N_{2,k}$  which is coming from the bottom multiplied by the area  $\Delta A_2$  plus  $N_{3,k}$  multiplied by  $\Delta A_3$  at  $x_3$  and likewise the same quantities at the outgoing phases at  $x_1 + \Delta x_1$ ,  $x_2 + \Delta x_2$  and  $x_3 + \Delta x_3$ ;  $\rho_k$  is the specie density.

If we now substitute these expressions in this verbal statement and divide each term by  $\Delta v$  which is the product of  $\Delta x_1$ ,  $\Delta x_2$  and  $\Delta x_3$  and let each of these increments tend to 0; this is a procedure we have gone through before.

(Refer Slide Time: 03:23)

**Mass Transfer Equation - I L4(2/17)**

$$\frac{\partial(\rho_k)}{\partial t} + \frac{\partial(N_{1,k})}{\partial x_1} + \frac{\partial(N_{2,k})}{\partial x_2} + \frac{\partial(N_{3,k})}{\partial x_3} = R_k \quad (1)$$

Now, the total mass transfer flux  $N_{i,k}$  in direction  $i$  is the sum of *Convective Flux* ( $\rho_k u_i$ ) due to bulk fluid motion and *Diffusion Flux* ( $m''_{i,k}$ ) due to density difference. Thus,

$$N_{i,k} = \rho_k u_i + m''_{i,k} \quad (2)$$

where

$$m''_{i,k} = -D \frac{\partial \rho_k}{\partial x_i} \quad (3)$$

where  $D$  ( $m^2/s$ ) is the mass-diffusivity

Then you will see you will get an equation of this type  $d\rho_k$  by  $dt$  plus net transport of specie  $k$  in direction 1, plus net transport in direction 2, plus net transport in direction 3 equal to rate of generation of specie  $k$ .

Now, the total mass flux  $N_{i,k}$  in direction  $i$  is the sum of the convective flux due to  $\rho_k u_i$  due to bulk fluid motion and diffusion flux  $m''_{i,k}$  due to density difference. Thus  $N_{i,k}$  is represented as  $\rho_k u_i$  plus  $m''_{i,k}$ .

Writing convective flux in this manner is indicative of the fact that we are assuming that each specie is traveling or is being carried at the same velocity as the bulk fluid. The diffusion flux on the other hand  $m''_{i,k}$  arises simply due to the differences in density at neighbouring locations.

The expression for mass flux by diffusion is given as  $m''_{i,k}$  equal to minus  $D \rho_k$  by  $dx_i$ . This equation is analogous to the conduction heat transfer and is like a Fourier's law of heat conduction.

In mass transfer, it is called the Fick's law of mass diffusion. So, just as in convective heat transfer you know that the total flux of energy is given by the convecting flux plus conduction flux. In mass transfer, we say it is the convective flux plus diffusion flux.

D is called the mass-diffusivity; it has units of meter square per second. So, if I have to substitute for  $N_{i,k}$  for each of these terms then I can rewrite this equation in the following manner.

(Refer Slide Time: 05:56)

**Mass Transfer Equation - II L4(3/17)**  
 Substituting for  $N_{i,k}$  gives

$$\frac{\partial(\rho_k)}{\partial t} + \frac{\partial(\rho_k u_1)}{\partial x_1} + \frac{\partial(\rho_k u_2)}{\partial x_2} + \frac{\partial(\rho_k u_3)}{\partial x_3} = \frac{\partial}{\partial x_1} \left( D \frac{\partial \rho_k}{\partial x_1} \right) + \frac{\partial}{\partial x_2} \left( D \frac{\partial \rho_k}{\partial x_2} \right) + \frac{\partial}{\partial x_3} \left( D \frac{\partial \rho_k}{\partial x_3} \right) + R_k \quad (4)$$

Define Mass Fraction  $\omega_k$

$$\omega_k = \frac{\rho_k}{\rho_m} \quad \sum_{\text{all species}} \omega_k = 1 \quad \text{and} \quad \sum \rho_k = \rho_m \quad (5)$$

$$\frac{\partial(\rho_m \omega_k)}{\partial t} + \frac{\partial(\rho_m u_j \omega_k)}{\partial x_j} = \frac{\partial}{\partial x_j} \left( \rho_m D \frac{\partial \omega_k}{\partial x_j} \right) + R_k \quad (6)$$

Summation gives Bulk Mass Conservation because  $\sum \omega_k = 1$  and hence,  $\sum R_k = 0$ . Also,  $\sum N_{i,k} = \rho_m u_i + \sum \dot{m}_{i,k} = \rho_m u_i$

It would read as  $\frac{d \rho_k}{dt} + \frac{d \rho_k u_1}{dx_1}$  and so on and so forth is equal to  $\frac{d}{dx_1} \left( D \frac{d \rho_k}{dx_1} \right)$  and likewise in direction 2 and 3 plus  $R_k$ .

$\rho_k$  has units of density. It is customary to define mass fraction  $\omega_k$  as the specie density divided by the mixture density and therefore,  $\sum \omega_k = 1$  by Dalton's law. Another way of saying is sum of densities is equal to the mixture density.

The notion behind Dalton's law is that each specie behaves as though it occupies the volume of the total mixture.

If I were to substitute now,  $\rho_k$  as  $\rho_m$  multiplied by  $\omega_k$ , this equation would be written in tensor notation in this fashion this is called the mass transfer equation. It has a transient term, a convection term, a diffusion term and a source or a generation term.

Now, let us sum each of these terms over all species - that is sum over all  $k$  as done here. (Refer Slide 07:38) Then clearly here this term would reduce to  $\frac{d \rho_m}{dt}$ ; this term would reduce to  $\frac{d \rho_m u_j}{dx_j}$  because  $\omega_k$  is equal to 1.

$\sigma_{\omega k}$  will simply become unity or constant. Therefore, this term would simply vanish. Just see now that this equation would be  $\frac{d\rho}{dt} + \frac{d\rho u_j}{dx_j} = \sigma R_k$  and that term is here. You will recognize readily then in the absence of  $\omega k$ , the left hand side is simply the bulk mass conservation and it equals 0; it follows therefore the  $\sigma R_k$  must be 0.

These two last deductions are very important.  $\sigma R_k = 0$  says that whenever there is a chemical reaction it is true that some species will be generated, but there would be others that would be destroyed and the total mass cannot be generated or destroyed; that idea is expressed in  $\sigma R_k = 0$ .

What does summation of diffusion equal to 0 imply? It implies that when some species are being diffused in a certain direction, other species are being diffused in the opposite direction and that stands to reason.

For example, if fuel was decreasing then product would increase. The net result is that summation of all diffusive quantities would be 0 and therefore,  $\sum N_i k$  would be  $\sum \rho u_i + \sum m''_{i,k}$ , but that quantity is 0 and therefore, sum of the mass fluxes in direction  $i$  are simply the bulk mass flux  $\rho u_i$ . We shall refer to this equation much later in the course.

(Refer Slide Time: 09:55)

**Mass Diffusivity L4(4/17)**

- 1 The mass diffusivity is *defined* only for a Binary Mixture of two fluids 1 and 2 as  $D_{12}$ .
- 2 In **Multicomponent Gaseous Mixtures**, however, diffusivities for pairs of species are nearly equal and a single symbol  $D$  suffices for all species.
- 3 Incidentally, in **turbulent flows**, this assumption of equal ( effective ) diffusivities holds even greater validity as will be shown later

6 March 15, 2010 8:39

A word about Mass Diffusivity: strictly speaking mass diffusivity is defined only for a Binary Mixture of two fluids 1 and 2 as  $D_{12}$ .

But in a combusting product mixture for example, there are several species present and diffusivities of pairs of species are truly different. So, **diffusion of carbon dioxide in nitrogen** or the diffusivity of carbon dioxide in nitrogen is different from diffusivity of oxygen in nitrogen and vice versa, but in gaseous mixtures these diffusivity pairs between species tend to be very nearly equal and therefore, we discard  $D_{12}$  or  $d_{ij}$  as the symbol for diffusivity and replace it by single symbol  $D$  and that suffices for combustion calculations.

Incidentally in turbulent flow, this assumption of equal effective diffusivity holds even greater validity; this we shall appreciate a little later.

(Refer Slide Time: 11:23)

**1st Law of Thermodynamics L4( $\frac{5}{17}$ )**  
In Rate Form ( $W / m^3$ ), the 1st Law of Thermodynamics reads as

$$\dot{E} = \dot{Q}_{conv} + \dot{Q}_{cond} + \dot{Q}_{gen} - \dot{W}_s - \dot{W}_b \quad (7)$$

where

- $\dot{E}$  = Rate of Change of Energy of the CV
- $\dot{Q}_{conv}$  = Net Rate of Energy transferred by Convection
- $\dot{Q}_{cond}$  = Net Rate of Energy transferred by Conduction
- $\dot{Q}_{gen}$  = Net Rate of Volumetric Heat Generation within CV
- $\dot{W}_s$  = Net Rate of Work Done by Surface Forces
- $\dot{W}_b$  = Net Rate of Work Done by Body Forces

Each term will now be represented by a mathematical expression.

March 15, 2012 7:13

We now turn to the first law of thermodynamics in rate form that is Watts per cubic meter. The first law of thermodynamics reads as  $dE$  by  $dt$  equal to  $dQ$  convection by  $dt$  plus  $dQ$  conduction by  $dt$  plus rate of generation of energy minus work done by shear forces minus work done by body forces.

Each of these terms is defined here and we shall seek mathematical representations for each one of them.

(Refer Slide Time: 12:09)

**Rate of Change L4(6/17)**

$$\dot{E}_D = \frac{d(\rho_m e^o)}{dt} \quad (8)$$
$$e^o = e_m + \frac{V^2}{2} = h_m - \frac{p}{\rho_m} + \frac{V^2}{2} \quad (9)$$

where

- $e_m$  = Mixture Specific Energy ( J / kg )
- $h_m$  = Mixture Specific Enthalpy ( J / kg )
- $V^2 = u_1^2 + u_2^2 + u_3^2$  ( J / kg )
- $\rho_m = \sum \rho_k$  = Mixture Density ( kg / m<sup>3</sup> )

Contributions to  $e^o$  from other forms ( Potential, Electro-Magnetic etc ) of energy are neglected.

March 16, 2012 8/19

So, let us consider the first term  $\dot{E}_D$ ; that would simply be  $d$  by  $dt$   $\rho_m e^o$  or the total energy. The total energy is the sum of static energy plus kinetic energy and the static energy by a thermodynamic relation is nothing, but enthalpy minus  $p$  into specific volume or  $p$  divided by  $\rho_m$  plus  $V$  square by 2.

So,  $e_m$  is mixture specific energy,  $h_m$  is mixture specific enthalpy,  $V$  square is the kinetic energy - sum of  $u_1$  square,  $u_2$  square and  $u_3$  square and  $\rho_m$  as we saw before is  $\sum \rho_k$ , the Mixture Density.

In general,  $e^o$  will have contributions from many other **force** effects like potential energy associated with rise or fall in elevation, electro-magnetic energies and so on and so forth, but these we will neglect because practical equipment's are fairly small.

**But in which** Yes, kinetic energy would be of some interest, but not these energies that I mentioned at the moment.

(Refer Slide Time: 13:35)

### Thermodynamic Convention L4( $\frac{7}{17}$ )

Convention:

- Heat Flow **into CV** is Positive whereas Heat Flow **out of CV** is Negative.
- Also, all species are transported at **Mixture Velocity**

N = TOTAL MASS FLUXES  
q = TOTAL HEAT FLUXES

March 16, 2010 9/19

Now, in order to represent the heat transfer and the work transfer across a control volume we shall follow the thermodynamic conventions. We will say that heat flow into the control volume is positive whereas, the heat flow out of the control volume is negative.

As we said before, all species are transported at mixture velocity. Here, I have shown all the net mass transfers and heat transfers that could take place across a control volume phase. N is total mass flux and q is the total heat flux across the control volume phase.

(Refer Slide Time: 14:25)

### Net Convection L4( $\frac{8}{17}$ )

Since  $\sum N_{j,k} = \rho_m u_j$

$$\dot{Q}_{conv} = -\frac{\partial \sum (N_{j,k} \theta_k^0)}{\partial x_j} = -\frac{\partial}{\partial x_j} \left[ \left( \sum N_{j,k} h_k \right) + \rho_m u_j \left( -\frac{p}{\rho_m} + \frac{V^2}{2} \right) \right] \quad (10)$$

Now, following definition of  $N_{j,k}$  and noting  $\sum_{-k} h_k = h_m$ ,

$$\begin{aligned} \sum N_{j,k} h_k &= \sum (\rho_m u_j \omega_k + m_{j,k}^* ) h_k \\ &= \rho_m u_j h_m + \sum m_{j,k}^* h_k \end{aligned} \quad (11)$$

Hence, after some algebra

$$\dot{Q}_{conv} = -\frac{\partial (\rho_m u_j \theta^0)}{\partial x_j} - \frac{\partial (\sum m_{j,k}^* h_k)}{\partial x_j} \quad (12)$$

March 16, 2010 10/19



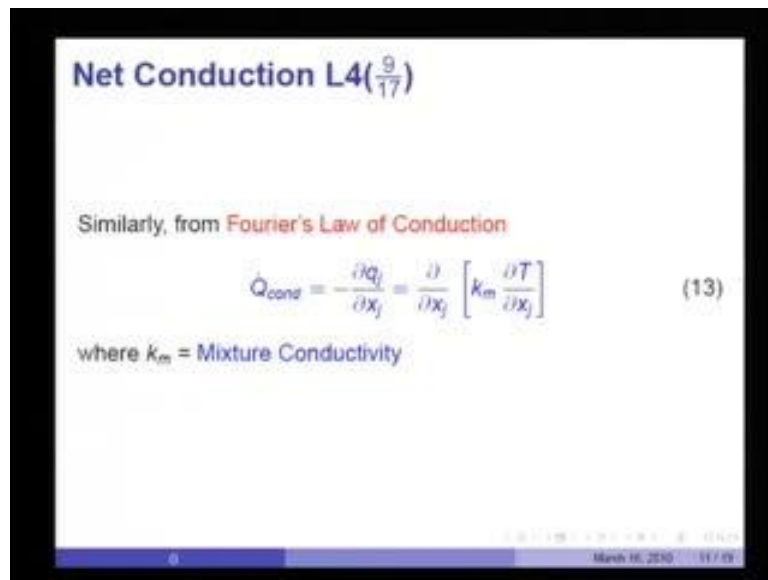
Net convection then by following the convention that the heat in is positive whereas, heat out is negative, we would have convection as simply  $d$  by  $d x_j$  of  $\sum N_j k e_{naught k}$  and if I replace  $e_{naught k}$  by  $h_k$  minus  $p$  by  $\rho_m$  plus  $V$  square by  $2$ .

Then you will see  $N_j k$  would multiply with  $h_k$  into the bracket, but the sum of  $N_j k$  would simply be  $\rho_m u_j$  as we saw before minus  $p$  by  $\rho_m$  plus  $V$  square by  $2$ . If we now note that  $\omega_k$  into  $h_k$ , that is the mass fraction of specie  $k$  multiplied by its specific enthalpy must add up to the mixture enthalpy; then  $\sum N_j k h_k$  would be  $\sum \rho_m u_j \omega_k$  plus  $m$  double  $\dot{j} k$ , the diffusion flux into  $h_k$ .

Since  $\sum \omega_k h_k$  is equal to  $h_m$ , it will add up to  $\rho_m u_j h_m$  plus  $\sum m$  double prime  $j k h_k$ .

(Refer Slide Time: 15:37) Therefore, if I replace this term here then you will see you have  $\rho_m u_j h_m$  minus  $p$  by  $\rho_m$  plus  $V$  square by  $2$  which would be written in this fashion and then the left over term, the diffusion flux term that could be given by that term. So, that is the expression for the convective flux.

(Refer Slide Time: 16:06)



Net conduction again by Fourier's law of conduction:  $Q$  conduction will be simply minus  $d q_j$  by  $d x_j$  where  $q_j$  is  $k_m$  multiplied by  $d T$  by  $d x_j$  and  $k_m$  is the mixture conductivity.

(Refer Slide Time: 16:22)

**Volumetric Generation L4(10/17)**

$$\dot{Q}_{gen} = \dot{Q}_{chem} + \dot{Q}_{rad} \quad (14)$$

- 1 The Chemical Energy is positive for exothermic reactions and negative for endothermic reactions.
- 2 Evaluation of  $\dot{Q}_{chem}$  depends on the chemical reaction model employed in a particular situation.
- 3 The  $\dot{Q}_{rad}$  term represents the Net Radiation Exchange between the CV and its surroundings. Evaluation of this term, in general, requires solution of integro-differential equations
- 4 When absorptivity ( $a$ ) and scattering coefficient ( $s$ ) are large,

$$\dot{Q}_{rad} = \frac{d}{dx_j} \left[ k_{rad} \frac{dT}{dx_j} \right] \quad k_{rad} = \frac{16 \sigma T^3}{a + s} \quad (15)$$

where  $\sigma$  is the Stefan-Boltzmann constant.

Net volumetric generation: Now, volumetric generation in a moving fluid typically would comprise of the generation due to chemical energy because some reactions are exothermic whereas, some other chemical reactions are endothermic. So,  $\dot{Q}_{chem}$  would be positive in case of exothermic reactions and it would be negative in terms of endothermic reaction.

If you want to evaluate  $\dot{Q}_{chem}$  in a moving fluid then one needs to postulate first of all the chemical reaction model that is being employed and there are variety of chemical reaction mechanisms of different levels of complexity. We shall see all these when we come to study of mass transfer.

$\dot{Q}_{rad}$  represents the net radiation exchange between the control volume and its surroundings. Usually evaluation of this term would require what is called the radiation transfer equation and it happens to be an integro-differential equation.

Treatment of that equation is really beyond the scope of the present lectures and it would require usually numerical calculations in a real practical equipment, but suppose the mixture that we are carrying has very high absorptivity may be because of soot, may be because of particulates that are present in the gaseous mixture.

Then the absorptivity and scattering of the radiation would be very high. When absorptivity and scattering coefficients are large,  $\dot{Q}_{rad}$  can actually be represented in

a manner similar to the conduction equation. The radiation conductivity can be defined as  $\frac{16}{3} \frac{\sigma T^3}{\kappa}$  where  $\sigma$  is the Stefan-Boltzmann constant that you are all familiar with.

But remember this kind of representation is justified only when the gaseous mixture has very high absorption and scattering coefficients.

(Refer Slide Time: 18:48)

The slide contains the following text and equations:

**Work-Done by Forces-I L4(11/17)**

Convention: The work done **on the CV** is negative,

$$\begin{aligned}
 -W_s &= \frac{\partial}{\partial x_1} [\sigma_1 u_1 + \tau_{12} u_2 + \tau_{13} u_3] \\
 &+ \frac{\partial}{\partial x_2} [\tau_{21} u_1 + \sigma_2 u_2 + \tau_{23} u_3] \\
 &+ \frac{\partial}{\partial x_3} [\tau_{31} u_1 + \tau_{32} u_2 + \sigma_3 u_3] \quad (16)
 \end{aligned}$$

$$-W_b = \rho_m (B_1 u_1 + B_2 u_2 + B_3 u_3) \quad (17)$$

where  $W_s =$  **Stress Work** and  $W_b =$  **Body-Force Work**  
 Further use of differentiation of product gives ( see next slide )

Now, we come to the work done terms and the derivation to follow is somewhat longish and I would request you to play good attention to how the derivation progresses.

Firstly there are shear forces and normal forces. Here is the stress; stress is force per unit area multiplied by velocity gives you the work done and d by d x 1 of that gives you the net work done.

So,  $\sigma_1$  which acts in the direction 1 multiplied by  $u_1$  plus  $\tau_{12}$  into  $u_2$  is the **work is the shear work in** net shear work in x 1 direction and so on and so forth.

You will get shear stress multiplied by the associated velocity in the two directions and the net work done by stresses  $W \cdot s$  would be given as that.

The body forces work is  $\rho_m B_1$ , the force multiplied by the velocity  $u_1$  is the work done in direction 1, 2 and 3 likewise. So,  $\dot{W}_s$  is the stress work;  $\dot{W}_B$  is the Body-Force Work.

Now we shall write these things as differentiation of a product. So, this  $\frac{d}{dx_1} \sigma_{11} u_1$  will be written as  $\sigma_{11} \frac{du_1}{dx_1} + u_1 \frac{d\sigma_{11}}{dx_1}$  and so on and so forth.

(Refer Slide Time: 20:22)

**Work-Done by Forces-II L4( $\frac{12}{17}$ )**

$$-(\dot{W}_s + \dot{W}_B) = u_1 \left[ \frac{\partial \sigma_{11}}{\partial x_1} + \frac{\partial \tau_{21}}{\partial x_2} + \frac{\partial \tau_{31}}{\partial x_3} + \rho_m B_1 \right] \quad (1)$$

$$+ u_2 \left[ \frac{\partial \tau_{12}}{\partial x_1} + \frac{\partial \sigma_{22}}{\partial x_2} + \frac{\partial \tau_{32}}{\partial x_3} + \rho_m B_2 \right] \quad (2)$$

$$+ u_3 \left[ \frac{\partial \tau_{13}}{\partial x_1} + \frac{\partial \tau_{23}}{\partial x_2} + \frac{\partial \sigma_{33}}{\partial x_3} + \rho_m B_3 \right] \quad (3)$$

$$+ \sigma_{11} \frac{\partial u_1}{\partial x_1} + \sigma_{22} \frac{\partial u_2}{\partial x_2} + \sigma_{33} \frac{\partial u_3}{\partial x_3} \quad (4)$$

$$+ \tau_{12} \left( \frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1} \right) + \tau_{13} \left( \frac{\partial u_1}{\partial x_3} + \frac{\partial u_3}{\partial x_1} \right) \quad (5)$$

$$+ \tau_{23} \left( \frac{\partial u_2}{\partial x_3} + \frac{\partial u_3}{\partial x_2} \right)$$

Complementary stress  $\tau_{12} = \tau_{21}$  etc are recognised ( see next slide )

On the next line you will see the result. So, the total work-done will be  $u_1$  multiplied by all these terms,  $u_2$  multiplied by all these terms,  $u_3$  multiplied by all these terms, plus  $\sigma_{11} \frac{du_1}{dx_1}$ ,  $\sigma_{22} \frac{du_2}{dx_2}$ ,  $\sigma_{33} \frac{du_3}{dx_3}$ ,  $\tau_{12}$  into the strain rate associated with  $\tau_{12}$ ,  $\tau_{13}$  into the strain rate associated with  $\tau_{13}$ , and  $\tau_{23}$  into the strain rate associated with  $\tau_{23}$  and here I have already used the idea of complementarities of stresses that is  $\tau_{12}$  is equal to  $\tau_{21}$ . Now I draw your attention to this term.

If you recall, when we wrote the Navier-Stoke's equations before or the Newton's second law of motion, these are simply the net forces acting in direction 1 and therefore, these are nothing, but the right hand sides of the Newton's second law of motion.

They can be replaced by the left hand side of the Newton's second law of motion and that is what I will do on the next slide plus all the stresses can be replaced in terms of viscosity multiplied by strain rates.

(Refer Slide Time: 21:48)

**Work-Done by Forces-III L4(13/17)**

Multipliers of  $u_1, u_2, u_3$  in equations 18, 19 and 20 are simply RHS of Momentum equations ( See Lecture 3, slides 13-14-15 ). They are replaced by LHS of Momentum equations. Hence,

$$\begin{aligned} \text{Equations 18, 19, 20} &= \rho m \left[ u_1 \frac{Du_1}{Dt} + u_2 \frac{Du_2}{Dt} + u_3 \frac{Du_3}{Dt} \right] \\ &= \rho m \frac{D}{Dt} \left( \frac{V^2}{2} \right) \end{aligned} \quad (23)$$

Let us do this first. So, multipliers of  $u_1, u_2,$  and  $u_3$  in equations 18, 19 are simply right hand sides of momentum equations, if you recall that on lectures 3 slides 13-14-15.

We have replaced them by left hand side of momentum equations. Left hand side was  $u_1 \frac{Du_1}{Dt}$  into  $\rho m u_2 \frac{Du_2}{Dt}$  and  $u_3 \frac{Du_3}{Dt}$ .

This would simply be  $\frac{Du_1^2}{2Dt}, \frac{Du_2^2}{Dt}$  and  $\frac{Du_3^2}{Dt}$  which we write as  $\frac{DV^2}{2Dt}$  into  $\rho m$ . (Refer Slide Time: 22:30) These are the first 3 terms that we have expressed; now we look at the remaining terms.

If I replace  $\sigma_1$  by  $-\rho p + \tau_{11}$ ,  $\sigma_2$  as  $-\rho p + \sigma_2 + \tau_{22}$  and  $\sigma_3$  as  $-\rho p + \tau_{33}$  and  $\tau_{12}$  as  $\mu \frac{Du_1}{Dx_2}$  into  $\frac{Du_2}{Dx_2}$  and so on and so forth.

(Refer Slide Time: 23:06)

**Work-Done by Forces-IV L4( $\frac{14}{17}$ )**

Similarly, using Stokes's Laws, equations 21, and 22 can be written as

$$\text{Equations 21, 22} = \mu \Phi_v - p \nabla \cdot V \quad (24)$$

where Viscous Dissipation Function is

$$\Phi_v = 2 \left[ \left( \frac{\partial u_1}{\partial x_1} \right)^2 + \left( \frac{\partial u_2}{\partial x_2} \right)^2 + \left( \frac{\partial u_3}{\partial x_3} \right)^2 \right] + \left( \frac{\partial u_1}{\partial x_2} + \frac{\partial u_2}{\partial x_1} \right)^2 + \left( \frac{\partial u_1}{\partial x_3} + \frac{\partial u_3}{\partial x_1} \right)^2 + \left( \frac{\partial u_2}{\partial x_3} + \frac{\partial u_3}{\partial x_2} \right)^2 \quad (25)$$

March 16, 2010 16/19

You will see that **the term would be** equations 21 and 22 would be simply  $\mu \phi_v$  minus  $p \nabla \cdot V$  where  $\phi_v$  is called the Viscous Dissipation Function.

It would take the form 2 times  $\frac{\partial u_1}{\partial x_1}$  squared which arises from  $\tau_{11}$ ,  $\frac{\partial u_2}{\partial x_2}$  squared,  $\frac{\partial u_3}{\partial x_3}$  squared and so on and so forth and  $\nabla \cdot V$  as we recall is simply,  $\frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2} + \frac{\partial u_3}{\partial x_3}$ .

(Refer Slide Time: 23:43)

**Work-Done by Forces-V L4( $\frac{15}{17}$ )**

Hence, from equations 23 and 24, we have

$$-(W_s + W_b) = \rho \frac{D}{Dt} \left( \frac{V^2}{2} \right) + \mu \Phi_v - p \nabla \cdot V \quad (26)$$

where

$$\nabla \cdot V = \frac{\partial u_1}{\partial x_1} + \frac{\partial u_2}{\partial x_2} + \frac{\partial u_3}{\partial x_3} \quad (27)$$

March 16, 2010 17/19

Therefore, the total work done terms are simply  $\rho_m \int D \cdot D \, t \, V \, \text{square by 2 plus } \mu \nabla \cdot \nabla \phi \, v \, \text{minus } p \, \text{del dot } V$ . You will notice that  $\nabla \cdot \nabla \phi \, v$  will always be positive because it is a sum of all the gradient square.

So, that sum would always be positive and in general and its purpose is to increase the energy level of the mixture. (Refer Slide Time: 24:08) This is the kinetic energy term and this is the minus  $p \, d \, V$  term or sometimes also called the pressure work term.  $\text{Del dot } v$  as you know is this.

(Refer Slide Time: 24:20)

**Summary L4(16/17)**

$$\dot{E} = \dot{Q}_{\text{conv}} + \dot{Q}_{\text{cond}} + \dot{Q}_{\text{gen}} - \dot{W}_s - \dot{W}_b \quad (28)$$

where

$$\dot{E} = \frac{d(\rho_m e^o)}{dt} \quad (29)$$

$$\dot{Q}_{\text{conv}} = -\frac{d(\rho_m u_j e^o)}{dx_j} - \frac{d(\sum_k m_k^o h_k)}{dx_j} \quad (30)$$

$$\dot{Q}_{\text{cond}} = \frac{d}{dx_j} \left[ k_m \frac{dT}{dx_j} \right] \quad (31)$$

$$\dot{Q}_{\text{gen}} = \dot{Q}_{\text{chem}} + \dot{Q}_{\text{rad}} \quad (32)$$

$$-(\dot{W}_s + \dot{W}_b) = \rho_m \frac{D}{Dt} \left( \frac{V^2}{2} \right) + \mu \phi_v - p \nabla \cdot V \quad (33)$$

By way of summary: we wrote expression for  $\dot{E}$  as  $d \rho_m e \, \text{naught by } d \, t$ ; expression for convection was written as  $\text{minus } d \rho_m u_j e \, \text{naught by } d x_j \text{ minus diffusion heat transfer } d \, \text{by } d x_j \text{ of sum of } m \, \text{dot } j \, k \, h \, k$ .

$\dot{Q}$  dot conduction was written in this fashion;  $\dot{Q}$  gen was  $\dot{Q}$  dot chem plus  $\dot{Q}$  dot rad and lastly we saw this term is this kinetic energy term, viscous dissipation term and so on and so forth.

You will see that  $\mu \nabla \cdot \nabla \phi \, v$  is positive and therefore, it tends to increase the rate of energy with time, conduction and if I were to replace  $\dot{Q}$  dot rad as  $k \, \text{rad } d \, T \, d x_j$  and then it would simply get added to that term. (Refer Slide Time: 25:17) This is the  $p \, \text{del dot } V$ ; this can be both positive and negative and so can this be positive and negative.

(Refer Slide Time: 25:22)

**Final Energy Equation L4(17/17)**

$$\frac{\partial(\rho m e^o)}{\partial t} + \frac{\partial(\rho m u_j e^o)}{\partial x_j} = \frac{\partial}{\partial x_j} \left[ k_m \frac{\partial T}{\partial x_j} \right] - \frac{\partial(\sum m_{j,k} h_k)}{\partial x_j} + \frac{D}{Dt} \left[ \frac{V^2}{2} \right] - \rho \nabla \cdot V + \mu \Phi_v + \dot{Q}_{chem} + \dot{Q}_{rad} \quad (34)$$

or, using definition of  $e^o$ ,

$$\rho m \frac{Dh}{Dt} = \frac{\partial}{\partial x_j} \left[ k_m \frac{\partial T}{\partial x_j} \right] - \frac{\partial(\sum m_{j,k} h_k)}{\partial x_j} + \mu \Phi_v + \frac{Dp}{Dt} + \dot{Q}_{chem} + \dot{Q}_{rad} \quad (35)$$

The final energy equation then reads as follows. What I have done is simply transported this negative term to the left hand side so that it becomes positive and we have  $\frac{d}{dt}(\rho m e) + \frac{d}{dx_j}(\rho m u_j e) = \text{net conduction heat transfer}$ .

Net conduction heat transfer due to mass diffusion, net kinetic energy generation minus  $\rho \nabla \cdot V$  plus Viscous Dissipation, chemical energy generation and radiation generation.

Now, if I replace  $e$  as  $h$  minus  $\frac{V^2}{2}$  then this term would become simply  $\rho m \frac{Dh}{Dt}$ .

The kinetic energy term on the left hand side would cancel with this kinetic energy term and the differential of  $\rho m u_j p$  by  $\rho m$  which is essentially  $u_j p$  would simply give me  $\frac{Dp}{Dt}$  would cancel partly with minus  $\rho \nabla \cdot V$  and I will have a term called  $\frac{Dp}{Dt} + \dot{Q}_{chem} + \dot{Q}_{rad}$  and so, this is called the Energy Equation written in enthalpy form.

(Refer Slide Time: 26:56) This is the energy equation written in internal energy form. It is this equation which is largely used in flows with or without chemical reaction.

Notice that all it says is that rate of change of enthalpy is equal to the net conduction heat transfer, net mass diffusion heat transfer and total viscous dissipation.



$\frac{Dp}{Dt}$  - this term is first of all the  $\frac{Dp}{Dt}$  will be  $\frac{\partial p}{\partial t}$  which would be of importance in unsteady flows such as one in which explosions are studied plus  $u \frac{Dp}{dx}$   $u_1 \frac{Dp}{dx_1}$ ,  $u_2 \frac{Dp}{dx_2}$  plus  $u_3 \frac{Dp}{dx_3}$  - these terms are important particularly when shocks are present where very large pressure gradients occur in a given direction, but in low speed flow really this term is relatively unimportant and we usually tend to ignore it.

$\dot{Q}_{chem}$  would arise out of the combustion model that one I specified and radiation is to be included in very high temperature reacting flows.

I will conclude here our equations of energy. (Refer Slide Time: 28:35) I would leave you with the **simple** recall of the mass transfer equation that we derived that was  $\rho \omega \frac{dk}{dt}$  and so on and so forth. This is the mass transport equation and then the last one is the energy transport equation. It is these equations that we shall develop further.