

Partial Differential Equations - 1
Professor A. K. Nandakumaran
Department of Mathematics
Indian Institute of Science Bengaluru
Professor P. S. Datt
Former Faculty, Tata Institute of Fundamental Research - Centre for Applicable
Mathematics, Bengaluru
Lecture 38
One Dimensional Wave Equation

(Refer Slide Time: 0:47)

D'Alembert's formula generalises to:

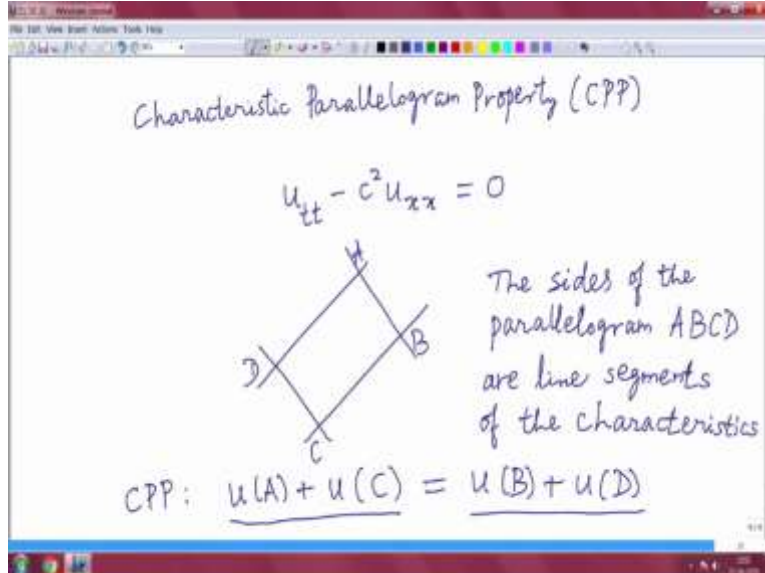
$$u(x,t) = \left(\frac{c_1}{c_1 - c_2} u_0(x + c_2 t) - \frac{c_2}{c_1 - c_2} u_0(x + c_1 t) \right) + \frac{1}{c_1 - c_2} \int_{x + c_2 t}^{x + c_1 t} u_1(\eta) d\eta$$

If $c_2 = -c_1 = -c$, this reduces to D'Alembert's formula

Welcome back. In today's lecture again we begin with general second-order equation with two different speeds. Last time we started that and we derived this formula for the solution of the second-order equation with two different speeds, c_1 and c_2 . So it generalizes the D'Alembert's formula, known D'Alembert's formula.

Of course, when we take the c_2 equal to minus c_1 equal to minus c , this reduces to D'Alembert's formula. And now you asked the question, what happens if c_1 equal to c_2 ?

(Refer Slide Time: 1:16)



So before coming to that, so let me just mention one more important property of the wave equation, so this is called Characteristic Parallelogram Property, which we are going to use later. In fact, this property characterizes the wave equation as we shall see. Characteristic Parallelogram Property, for short I will use the CPP to denote this, so CPP.

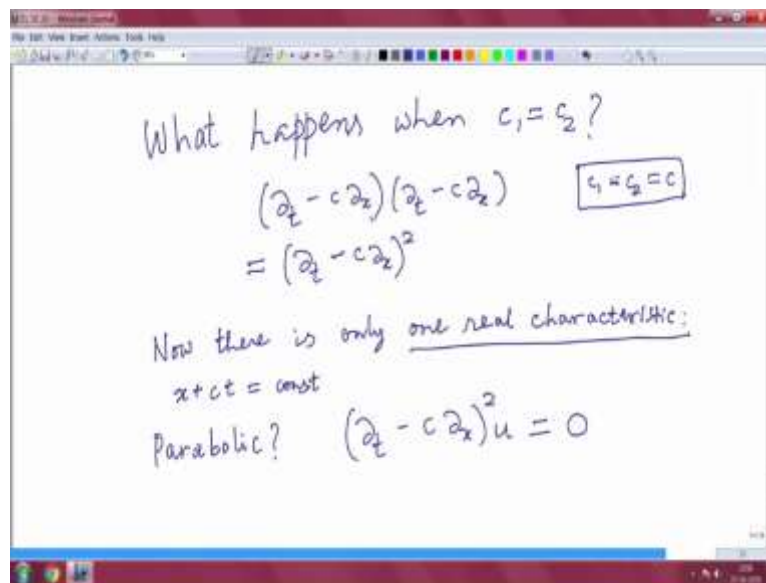
So again, we consider the wave equation, so here the initial conditions are not explicitly used, but they are in the background, so this is the wave equation. Now you consider a parallelogram, not any parallelogram, but a specific one parallelogram, so here I denote this A, B, C, D. So the sides of the parallelogram, A, B, C, D are line segments of the characteristics. So that is why it is called characteristic parallelogram.

Then the important property, CPP states that U of A plus U of C equal to U of B plus U of D. So you consider any solution of the wave equation and then you evaluate the solution at the opposite vertices, A and C, and B and D, then you add them, so this is sum of the solutions at A and C and this is sum of the solutions at B and D.

And the characteristic parallelogram property states that U_A plus U_C is equal to U_B plus U_D . So if you know the values of the solution at three points, any three points of the characteristic parallelogram, then the value at the fourth point is determined by this property. So this is an important property which we are going to use later.

So conversely, if that is the characterization, conversely if U is any C^2 , so I am not necessarily I am not referring to solution of the wave equation, if conversely, if U is any C^2 function satisfying CPP for all characteristic parallelograms then U satisfy the wave equation. So that is the characterization. So we are going to use it later.

(Refer Slide Time: 6:26)



So again back to this second-order equation with only one real characteristic. So this equation, I will write, $\partial_t^2 u - 2c\partial_t\partial_x u + c^2\partial_x^2 u = 0$, so homogeneous equation I am considering. So there is only one real characteristic, but this should not be called as parabolic.

(Refer Slide Time: 7:10)

The eqn $(\partial_t - c\partial_x)^2 u = 0$ double characteristic
is classified as weakly hyperbolic.
Characteristic variable: $\xi = x + ct$
 $\tau = x - ct$ ($c \neq 0$)
 $(x, t) \mapsto (\xi, \tau)$ is non-singular
The eqn $\Rightarrow \partial_\tau^2 u = 0$ ← This does not look like heat eqn
To solve the eqn, we need 2 initial conditions

So the equation, $\partial_t - c\partial_x$ squared u equal to 0 is classified as weakly hyperbolic. So perhaps you are first time hearing this class of equations. I will give you two reasons why this equation should not be classified as parabolic. So the first reason is, so consider the characteristic variable, so there is only one characteristic, so ξ is equal to x plus ct .

And now we choose another variable so that the change of variables is again non-singular, so Jacobian should not be a 0, so the easiest choice is x minus ct . But remember this time this is not a characteristic variable, so we are choosing a different variable, so that this xt change of variable to $\xi \tau$ is non-singular.

So we are assuming of course, $c \neq 0$, if c is 0 then the equation reduces to OD, so it is not even a PD. So again, you do the computation, so you see that the equation implies, so some computation, you show that $\partial_\tau^2 u = 0$. So in these new variables, you show that $\partial_\tau^2 u$ is equal to 0, and this does not resemble the heat equation.

So this is more like a OD, because it is only ∂_τ^2 , so there is no ∂_ξ there, so it is more like an OD, so far from the heat equation, which is a typical example of parabolic equation. So this is one reason why this equation should not be called as parabolic. The another reason to solve this equation we need two initial conditions, whereas heat equation requires only one initial condition.

So for these two reasons, again I stress that this equation should not be classified as parabolic, so it is classified as weakly hyperbolic. Now let us quickly solve the initial value problem or Cauchy problem associated with this equation. So since there the c , c repeats here, so such we can call this double characteristic equation.

So if the multiplicity of this characteristic is more than 1, so generally it is called multiple characteristic, and in this case since there it only repeats twice, so it is a double characteristic equation. These are more difficult than, then in contrast, the equation so let me show you that.

(Refer Slide Time: 12:50)

Then,
 $(\partial_t^2 - c_1^2 \partial_x^2)(\partial_t^2 - c_2^2 \partial_x^2)u = 0 \Rightarrow \partial_{\xi\tau}^2 u = 0$
 General solution: $u = F(\xi) + G(\tau)$
 $= F(x+c_1 t) + G(x+c_2 t)$
 Prescribe initial condns
 $u(x,0) = u_0(x)$
 $u_t(x,0) = u_1(x)$
 simple characteristics

This equation is called equation with simple characteristics. So we are assuming c_1 is different from c_2 , so this, there are two real and distinct characteristics here, so it is called simple characteristics.

(Refer Slide Time: 13:22)

IVP or Cauchy problem

$$\boxed{(\partial_t - c\partial_x)^2 u = 0}$$

$$u(x,0) = u_0(x), \quad u_t(x,0) = u_1(x)$$

$$\Rightarrow \partial_t^2 u = 0 \Rightarrow u(\xi, \tau) = f_1(\xi) + f_2(\xi)\tau$$

Gen. soln $u(x,t) = f_1(x+ct) + (x-ct)f_2(x+ct)$

Write $x-ct = (x+ct) - 2ct$

$$\text{Gen. soln: } \boxed{u(x,t) = F_1(x+ct) + t F_2(x+ct)}$$

The eqn $(\partial_t - c\partial_x)^2 u = 0$ double characteristic
 is classified as weakly hyperbolic.

Characteristic variable: $\xi = x+ct$
 $\tau = x-ct$ ($c \neq 0$)

$(x,t) \mapsto (\xi, \tau)$ is non-singular

The eqn $\Rightarrow \boxed{\partial_\tau^2 u = 0}$ ← This does not look like heat eqn

To solve the eqn, we need 2 initial conditions

Now let us quickly consider the initial value problem or Cauchy problem of this double characteristic equation. So $\partial_t - c\partial_x$ square equal to 0, as I said, it requires two initial conditions. In case of wave equation, we require u_0 to be a C^2 function and u_1 to be a C^1 function and let us see what difference we will observe here.

So for the time being assume u_0 and u_1 are such that we can obtain a C^2 solution of the equation. So now since there is only one characteristic, we already done it here, let us go back, so we get

this. With this change of variable in the new variable ξ and τ , this equation reduces to $\frac{\partial^2 u}{\partial \tau^2} = 0$. So that we can immediately integrate and obtain the solution.

So this implies $\frac{\partial^2 u}{\partial \tau^2}$, so let me not repeat again. So which again in terms implies u of, remember, u is a function of ξ and τ . So this is a linear function in τ , so we can write this as, so the constants can be functions of ξ . So this is the general solution of this equation. So let us go back to the original variables, so general solution, $u(x, t)$, so now again I write in terms of the original variables x and t , so this $F_1(x + ct)$, so τ is $x - ct$, $F_2(x - ct)$.

So we can arrange the term little bit and write it in a neat form, so we write $x - ct$ as $x + ct - 2ct$, and you absorb this $x + ct$ term in this function, and so we have only a multiplication by t . So general solution we can write it as, one more time I will write it, so $u(x, t)$, so let me use different notation, $F_1(x + ct) + t F_2(x - ct)$.

So we can easily verify that the expression given by this for u satisfy the given equation, this is the equation. And now in order to determine F_1 , F_2 , again just recall what we did in the case of wave equation to arrive at D'Alembert's formula. So I leave it as an easy exercise, so you plug in the initial conditions in this general form and finally you obtain a solution.

(Refer Slide Time: 18:28)

Solution

$$u(x, t) = u_0(x + ct) - ct \underline{u_0'(x + ct)} + t \underline{u_1(x + ct)}$$

For u to be a C^2 function, we require

$u_0 \in C^3, u_1 \in C^2$

Loss of regularity: More smoothness of the initial values is required.

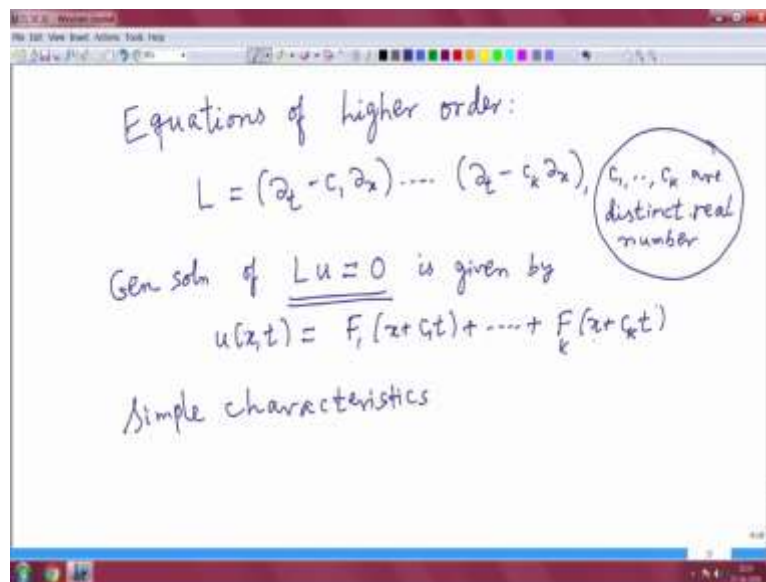
So I am skipping few steps, can easily point them and work out $x + ct - 2ct$, u_0' , so this derivative of u_0 , plus $t u_1(x + ct)$. So if you compare this form of the solution with the

D'Alembert's formula, you see a big change here. First of all, in the D'Alembert's formula no derivative of the initial conditions appeared, but here, the derivative of the initial condition appears, there is no integration of u_1 , but u_1 appears as it is.

So now we require so far u to be a C^2 function and that is what we need, at least two derivatives, u to be a C^2 function of x and t . We require u_0 to be a C^3 function, because there is one derivative in the formula of the solution, so if you want u to be C^2 , so this has to be C^3 and similarly u_1 has to be C^2 . So we require more smoothness in the initial values to obtain a less smooth u , u_0 is C^3 , but that gives u only a C^2 function.

So this is referred to as loss of regularity and this is typical of weakly hyperbolic equations and systems, loss of regularity. So more smoothness of the initial values required.

(Refer Slide Time: 22:27)



But now we can simply combine the two cases, so let me just mention as a general remark, so equations of higher order. So this is just follows from an induction argument. So first consider equation with simple characteristics, so the operator L , I write it, so this is C^1 del x , so a k -th order differential operator, which can be factored into linear factors, so that is the form of the higher order equation, where c_1, c_2 all distinct real numbers.

So we did for it in the case of two values, but the same procedures are real. So what we can show is this general solution now, so this is just followed by an induction argument, so I will not do

the details, you can do it yourself, general solution of $Lu = 0$ is given by u of x, t is equal to $F_1 x + c_1 t + F$ of $x + c_1 t$ plus where F_1, F_2, F_k are k times differential.

So here it is important there are distinct, otherwise you will not get this one. Of course, in order to determine this F_1, F_2, F_k , we have to prescribe k initial conditions on the solution, but it is more algebra, you have to solve k by k linear system. But in principle, it can be done. So that is with simple characteristics, so this is simple characteristics.

(Refer Slide Time: 26:00)

Multiple characteristics

$$(\partial_t - c \partial_x)^k u = 0, \quad k > 1$$

The general solution is given by

$$u(x, t) = F_1(x+ct) + t F_2(x+ct) + \dots + \frac{t^{k-1}}{(k-1)!} F_k(x+ct)$$

Equations of higher order:

$$L = (\partial_t - c_1 \partial_x) \dots (\partial_t - c_k \partial_x), \quad c_1, \dots, c_k \text{ are distinct real number}$$

Gen soln of $Lu = 0$ is given by

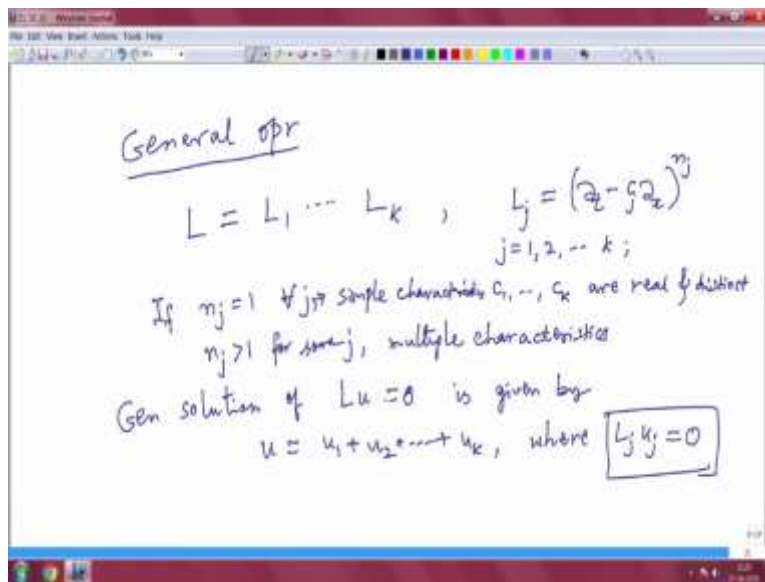
$$u(x, t) = F_1(x+c_1 t) + \dots + F_k(x+c_k t)$$

Simple characteristics

What about multiple characteristics? So it is similar thing. So this is the equation, $\text{del } x^k, u$ equal to 0. So k is bigger than 1, earlier we did for k equal to 2. So again, by an induction argument, this general solution is given by u of x is equal to F_1 of x plus $c_1 t$ plus t^2 , this we saw for k equal to 2 and now you do an induction argument and you get this.

And you see this factor t , which was not there in D'Alembert's formula and even in this simple characteristic problem, c_1 in simple characteristic problem, there is no t factor at all but when you go to multiple characteristics so you see this factor of t . So now we can combine these two simple characteristic equation and multiple characteristic equation into a single one.

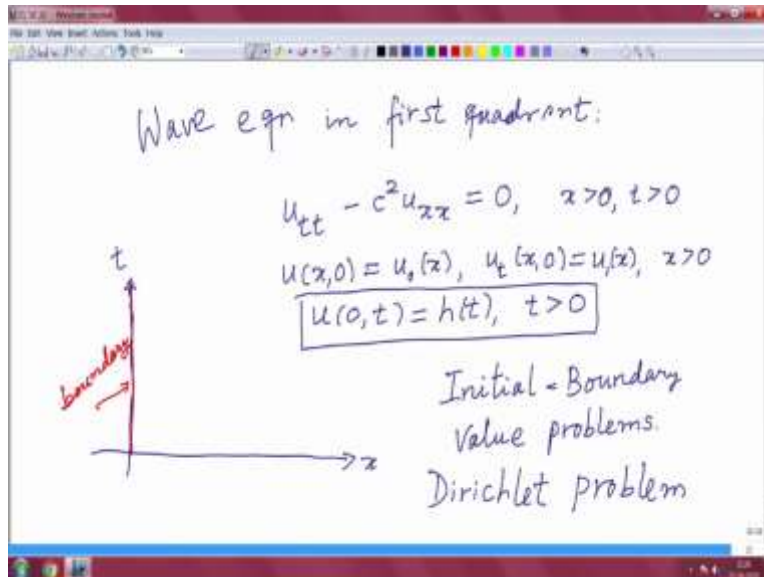
(Refer Slide Time: 28:26)



So, a more general equation, general operator. So L is equal to L_1, L_k , so L_k is this operator L_j , let me with $L_j, c_j \text{ del } x, I \text{ put } n_j$. So j equal to 1, 2, k and c_1, c_2, c_k are real and distinct. So if n_j equal to 1 for all j , then we have simple characteristics, so that leads to simple characteristics. And if n_j is bigger than 1 for some j , even for $1j$, it is multiple characteristic problem.

So the general solution of Lu equal to 0 is given by u is equal to u_1 plus u_2 plus u_k , where L_j, u_j is 0. And this we have now learned how to solve that equation. So later on we will provide with some example just to list this general equations of higher order. There is more algebra in it, so that is I just want to make remark of that.

(Refer Slide Time: 31:52)



So with this thing now we will move on to again wave equation, but now wave equation in first quadrant. So, so far consider on the real line, but now we are considering only on the positive axis. So here is the problem, u_{tt} , again wave equation, let me again just consider the homogeneous one, so we can always use Duhamel's principle in order to get the formula for the solution of the inhomogeneous equation.

So now we consider this equation only in x positive, and as usual, t is always positive. So like the previous case, first we provide the initial conditions at time t equal to 0, as I remarked in the case of wave equation, you can provide on any t equal to t_0 line, so there is absolutely no problem, by just translation the solution formula can be derived. But now only x is positive.

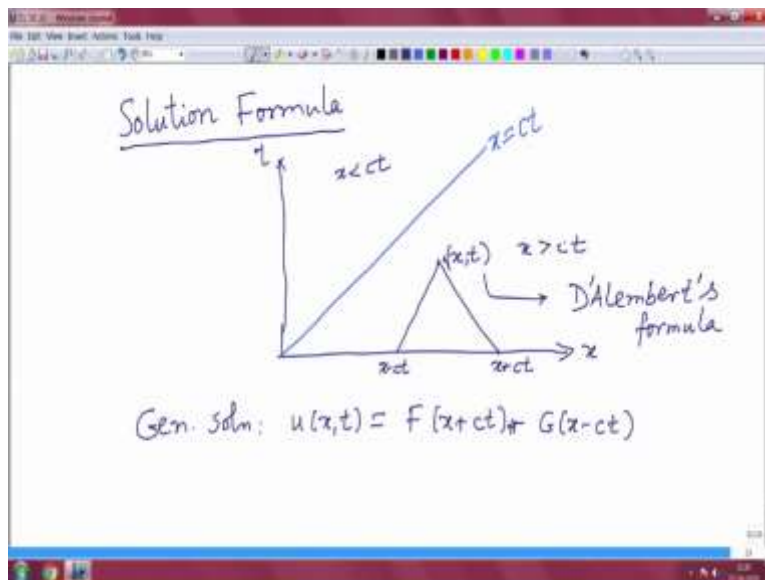
So let me draw this domain where we are interested, so this is the x axis and this is positive t axis. So we are interested in this solution, finding the solution in this first quadrant, x positive, t positive. And now, see earlier this was not there, but now there is a boundary here, so we distinguish the initial line and the boundary line, so let me use a different color here, a boundary. So what is the boundary, boundary is the t axis, so that is x equal to 0.

So we have to provide some conditions on the boundary, so for simplicity again let me just begin with, so x equal to 0 now, t , I put the h of t , and t positive. So such problems are referred to as

initial boundary value problems. And in this case, since we are prescribing the value of the solution on the boundary, it is referred to as Dirichlet problem.

There are other kinds of boundary conditions one can prescribe, so I will mention as you go along, but let us now try to find the formula for this wave equation with this initial and boundary conditions. Let us try to do that. So since most of the work is already done in the case of wave equation on the entire real line, namely D'Alembert's formula and other things, so here it is much easier, as we shall see.

(Refer Slide Time: 37:25)



So solution formula. And let us see how the boundary condition influences the formula. So again, it is very easy to explain using diagrams. So this is the interface, there is only point here, the origin, and the origin which is the intersection of the initial line and the boundary line. Let us draw the characteristic here, so they play a role there, let me use different color here for the time being, let me use blue.

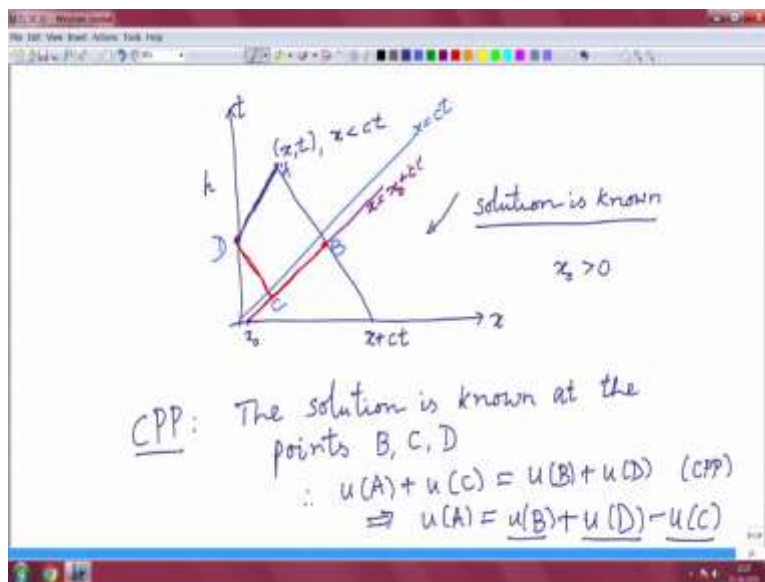
So that is the characteristic x equal to ct passing through the origin. While we are doing that one, let me just see that. So if I take any point x, t here, so below this characteristic, that means x is bigger than ct here and x is less than ct here. So the geometry changes, that is what is important, and if you look at the, so again, general solution of the wave equation, we already know that.

So, general solution in whatever domain is of this form, F of x plus ct , there is no change here, we can use just characteristic variable and obtain this. So F and G are two C^2 functions. So again, if you look at the domain of dependence property, so this value of the solution at this point is determined by the initial values in this interval, that is what is the domain of dependence.

And since x is bigger than ct , this x minus ct lies on the positive axis, thus the value of the solution at such a point is determined only by the initial values, because this will not see the boundary, so the solution simply obtained by D'Alembert's formula.

So geometrically also we can see that, so the value of the solution at this point below the characteristic is not influenced by the boundary value. But situation is quite different, so let draw again one more picture, so a different picture.

(Refer Slide Time: 41:40)



So if I take xt here, so now remember, x is less than ct . So again if I draw the characteristics through this xt , so one characteristic intersects, see the earlier case it was not doing that and one characteristic there is no problem, so this is x plus ct is there. But this characteristic does not meet the initial axis in the positive x axis and here we do not have any data, but we have data here, wherever it hits, we pick up the value h there.

So this is okay, but how to incorporate this value to find the value of the solution at this xt . So here we use, this is one way, so you can again use the general solution and try to determine G ,

but let me use the characteristic parallelogram property to find the solution of u at this point. So we have construct a characteristic parallelogram, so there are already two sides, we see that.

So this is also a characteristic and that is also a characteristic, and two points, so one to find the solution there and the data is given here, so we also know the solution at this point. So we need another two points where solution is known. But in this region we already know the solution, the solution is known here and that is given by the D'Alembert's formula.

So if we can choose two appropriate points in this domain, namely x we get then ct , and which form a characteristic parallelogram, then we are able to find the solution u at the point xt . And for that, so let me just use different color again, so you pick as point x_0 here, x_0 small, so we will see how small it is, x_0 is positive, small.

And now you draw a characteristic, so this is the characteristic, so this, you can write this x equal to x_0 plus ct . So this blue one is x equal to ct , that is passing through the origin and now this is passing through the point x_0 . So that is a characteristic, so there is now third side and so we know the value here because that is in the region where solution is known.

So we need one more point, so just you draw a line, a characteristic parallel to this. So this is one characteristic and this, so it does not look like a parallelogram, but we can just, so this is A, this is B, this is C, this is D. The solution is known at the points B, C, D, so therefore u_A plus u_C equal to u_B plus u_D . This is CPP that implies u of A is equal to u of B, u of D minus u of C.

So we have to find out what is u of B and what is u of D, u of D is just very easy, so that is the boundary data and u of C we have to find. So this we will continue in the next class.