

First Course on Partial Defferential Equations – 1

Professor A. K. Nandakumaran

Department of Mathematics,

Indian Institute of Science, Bengaluru

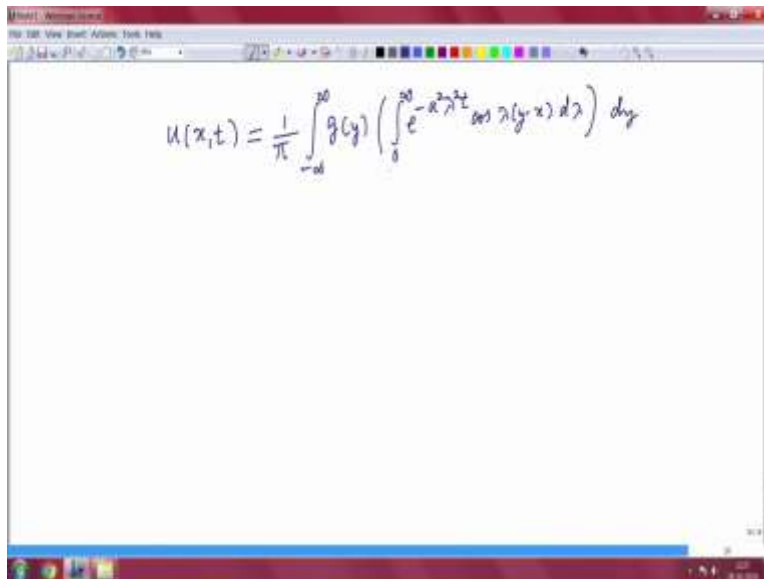
Professor P. S. Datti

Former Faculty, Tata Institute of Fundamental Research - Centre for Applicable
Mathematics

Lecture 30

One dimensional heat equation-3

(Refer Slide Time: 1:03)


$$u(x,t) = \frac{1}{\pi} \int_{-\infty}^{\infty} g(y) \left(\int_0^{\infty} e^{-\lambda^2 t} \cos(\lambda y - \lambda x) d\lambda \right) dy$$

So welcome back, we will continue discussion with the solution of the heat equation we started in the previous lecture. Now we complete the formula, so we are in the process of analyzing this integral. So again let me write it, so we had obtained the solution in this form $g(y) e^{-\lambda^2 t} \cos(\lambda y - \lambda x) d\lambda$ then dy . So we analyze this thing separately.

(Refer Slide Time: 1:49)

Concentrate on $\int_0^{\infty} e^{-\lambda^2 x t} \cos \lambda(y-x) d\lambda = \frac{1}{\sqrt{\pi t}} \int_0^{\infty} e^{-z^2} \cos(\eta z) dz$

Put $\eta = \frac{y-x}{\sqrt{\pi t}}$; $\sqrt{\pi t} \lambda = z \Rightarrow d\lambda = \frac{dz}{\sqrt{\pi t}}$

Define $H(\eta) = \int_0^{\infty} e^{-z^2} \cos(\eta z) dz$

$\frac{dH}{d\eta} = \int_0^{\infty} e^{-z^2} (-z \sin(\eta z)) dz$; integrate by parts

$= -\frac{\eta}{2} \int_0^{\infty} e^{-z^2} \cos(\eta z) dz = -\frac{\eta}{2} H(\eta)$

$\Rightarrow H(\eta) = \text{const.} \cdot e^{-\eta^2/4}$ Since $H(0) = \frac{\sqrt{\pi}}{2}$,
 $H(\eta) = \frac{\sqrt{\pi}}{2} e^{-\eta^2/4}$

And we obtained this function H eta and there is 1 by a root t and again we substitute back this variable eta. In the original variables

(Refer Slide Time: 2:05)

$u(x,t) = \frac{1}{\pi} \int_{-\infty}^{\infty} g(y) \left(\int_0^{\infty} e^{-\lambda^2 x t} \cos \lambda(y-x) d\lambda \right) dy$

$= (4\pi a^2 t)^{-1/2} \int_{-\infty}^{\infty} \exp\left(-\frac{(y-x)^2}{4a^2 t}\right) g(y) dy$

$u(x,t) = \int_{-\infty}^{\infty} K(x-y, t) g(y) dy$, $K = \text{Fundamental Solution of the heat eqn}$

$\cdot K$ is symmetric: $K(x,t) = K(-x,t), t > 0$

So that is simplifies into so let me write 1 by 4 pi a square t to the minus half minus infinity to infinity exponential, so since expression is little big. So, let me write exponential so y minus x square divided by 4 a square t into g y dy.

And you recognize this multiplied by this is precisely the fundamental solution of the heat equation. So this is a, write it here K of x minus y t g y dy. So K is the fundamental solution of

the heat operator which we first obtain as a special solution to the heat equation. And now it also makes appearance in the formula we are trying to derive for the heat equation. So just let me write $u(x, t)$.

So we also see that K is apart from the properties I listed earlier, K is also symmetry meaning K of x, t is equal to K of $-x, t$ positive. So if t is positive so that you can, because in the exponential term this sits as an x square so x and x square equal to minus x square so there is no change there. And now you look at for just a few comments on this formula. So let me just introduce some notations.

(Refer Slide Time: 5:24)

If, $f, g: \mathbb{R} \rightarrow \mathbb{R}$, their convolution is defined by

$$(f * g)(x) = \int_{-\infty}^{\infty} f(x-y)g(y) dy = \int_{-\infty}^{\infty} f(y)g(x-y) dy = g * f(x)$$

$$u(x, t) = (K(\cdot, t) * g)(x), \quad t > 0 \quad (2)$$

If g is bounded, u is well-defined.

Theorem If g is cont. & bdd, then u given by (2) indeed satisfies the heat eqn for $t > 0$.
 u also satisfies the initial condn in the sense that $\lim_{\substack{x \rightarrow \eta \\ t \rightarrow 0}} u(x, t) = g(\eta), \quad \eta \in \mathbb{R}$

$$u(x, t) = \frac{1}{\pi} \int_{-\infty}^{\infty} g(y) \left(\int_0^{\infty} e^{-\lambda^2 t} \cos \lambda(y-x) d\lambda \right) dy$$

$$= (4\pi a^2 t)^{-1/2} \int_{-\infty}^{\infty} \exp\left(-\frac{(y-x)^2}{4a^2 t}\right) g(y) dy$$

$$u(x, t) = \int_{-\infty}^{\infty} K(x-y, t) g(y) dy$$

$K =$ Fundamental Solution of the heat opr

$\cdot K$ is symmetric: $K(x, t) = K(-x, t), \quad t > 0$

So if f and g are functions from \mathbb{R} to \mathbb{R} , \mathbb{R} to \mathbb{C} their convolution is defined by so convolution of two functions is again a function so defined by $f \star g$ of x minus infinity to infinity f of x minus y g of y dy which we also can written as by changing the variable. So this is f of y g of x minus y dy , so this $g \star f$. So provided the integral is finite and other so this f the star operation the convolution operator is a commutative operator.

So in this sense we can write with this notation $u(x, t)$ as $K \star g$ so I will put the function here $t \star g$ of x . So the u we obtain by heuristic arguments can be written as convolution of the fundamental solution and g . That is what I want to comment. I want to stress that. So now so we obtain this formula by assuming several steps in between but the final formula is very neat, very neat.

And it make sense for example if g is bounded because of the exponential factor in K u is well defined so there is no problem at all. So the integral adjust for all t positive. And now try to show that u is indeed the solution of the heat equation. We cannot claim uniqueness so you just say so let me state it this as a theorem so if g is continuous and bounded we have to assume continuity does not imply boundedness.

Because we are on the real line that is not a compact check and bounded continuity is also not require if you are using Lebesgue integral let me just state that g is continuous and bounded then so let me put this as 2 then u given by 2 indeed satisfies the heat equation for t positive u also satisfies the initial condition in the sense that so limit u of x, t as x tends to η and t tends to 0 is equal to g of η so η is real variable.

So in general this limit is satisfied at all the points of continuity of g . So if you are not assuming g is continuous everywhere. So this limit exists wherever g is continuous so the prove of this theorem depends on one more important property of the fundamental solution.

(Refer Slide Time: 12:23)

Proof: Recall

- 1) $K \in C^\infty$ for $x \in \mathbb{R}, t > 0$; $K(x,t) = K(-x,t)$
- 2) $K_t = \alpha^2 K_{xx}$, $x \in \mathbb{R}, t > 0$
- 3) $\lim_{t \rightarrow 0^+} \int_{|x-y| \geq \delta} K(x-y,t) dy = 0$ unif in x
($\delta > 0$ fixed)

Proof of (3) $\int_{|x-y| \geq \delta} K(x-y,t) dy$; $\frac{x-y}{2\alpha\sqrt{t}} = z$

$$= \frac{2}{\sqrt{\pi}} \int_{\delta/2\alpha\sqrt{t}}^{\infty} e^{-z^2} dz = \frac{2}{\sqrt{\pi}} \int_{\frac{\delta}{2\alpha\sqrt{t}}}^{\infty} \frac{z}{z} e^{-z^2} dz < \frac{1}{\sqrt{\pi}} \frac{2\alpha\sqrt{t}}{\delta} \int_0^{\infty} z e^{-z^2} dz$$

$\rightarrow 0$ as $t \rightarrow 0^+$

If, $f, g: \mathbb{R} \rightarrow \mathbb{R}$, their convolution is defined by

$$(f * g)(x) = \int_{-\infty}^{\infty} f(x-y)g(y) dy = \int_{-\infty}^{\infty} f(y)g(x-y) dy = g * f(x)$$

$$u(x,t) = (K(\cdot,t) * g)(x), \quad t > 0 \quad (2)$$

If g is bounded, u is well-defined.

Theorem If g is cont & bdd, then u given by (2) indeed satisfies the heat eqn for $t > 0$.
 u also satisfies the initial condn in the sense that

$$\lim_{\substack{x \rightarrow \eta \\ t \rightarrow 0}} u(x,t) = g(\eta), \quad \eta \in \mathbb{R}$$

So let me just so proof, so recall the again properties of the fundamental solution. K is a C^∞ function for x in \mathbb{R} and t positive K is symmetric we saw that and K satisfy the heat equation so this is we also saw this K_{xx} for x in \mathbb{R} and t positive. So one more important property so I write it here so if you integrate $K(x-y,t)$ with respect to y not on the whole real line but you just leave stay away from this x at a positive distance.

And then you take this limit as t tends to 0 this is 0 uniformly in x that is important uniformly in x . So usual epsilon delta if you take that definition of the uniformity means that delta does not depend on x . So this is crucial in verifying this part of the theorem. This part so we will come to

that just so proof of three that is not very difficult. So again let us write that what is in terms of K . So you consider this x minus y bigger than or equal to δ , δ is positive.

So for any δ positive so if you stay away from B x and then this limit is 0. So make the substitution x minus y by 2 a root t is equal to z . And then this integral is changed to 2 by root π δ to infinity, so there is a portion coming from minus infinity to minus δ and δ to infinity but because of symmetry we can convert that minus infinity to minus δ integral to δ to infinity integral and that is why that 2 comes.

So it is very simple one, there is no complication at all. And sorry the δ so this become δ by 2 a root t because we are making the substitution so if x minus y is in absolute value bigger than δ then $\text{mod } z$ is bigger than or equal to δ by 2 a t^2 a square root t . So what you do here is now its simple thing 2 by π , now let me write that you add a z and divide by z e to minus z square.

And this we do can do because z is away from the origin and for 1 by z you replace it by z is bigger than or equal δ I but δ divided by 2 a root t . So 1 by z is less than or equal to 1 by root π 2 a root t by δ and the remaining integral you can just write it as 0 to infinity now there is no problem 2 z e to minus z square dz .

And this integral is just 1 and since root t sits in the numerator so that goes to 0 as t goes to 0 . Of course we are always considering t positive so let me stress that also there is a square root there. So there is this prove property 3, so now it is very easy to prove the stated theorem.

(Refer Slide Time: 19:33)

Differentiating under the integral, we obtain

$$u_t = \int_{-\infty}^{\infty} K_t(x-y, t) g(y) dy$$

$$u_{xx} = \int_{-\infty}^{\infty} K_{xx}(x-y, t) g(y) dy$$

$$\Rightarrow \boxed{u_t = a^2 u_{xx}}$$

If $b > 0$, then for every $n > 0$, we have $x^n e^{-bx^2} \leq C_n e^{-\frac{1}{2}x^2} \forall x > 0$

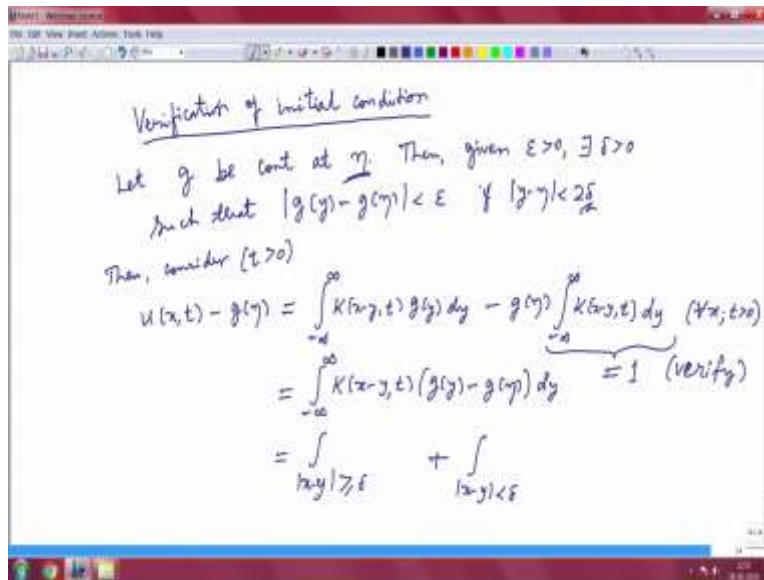
Since $K \in C^\infty$ for $t > 0$, it follows that $u(\cdot, t) \in C^\infty$ for $t > 0$
Something proper

So again differentiating under the integral sign so the this is let me make a comment here differentiating under the integral sign. So we obtain u of t same as minus infinity to infinity K of t x minus y t g y dy and u_{xx} is K_{xx} x minus y t g y dy . So the justification of taking the differentiation under the integral sign the important thing we use here for the exponential functions.

So if b is positive so this is a small result in analysis then for every n positive can be integer it can be anything we have x to the n e to the minus b x square is less than or equal to some constant which depends on n e to minus b by 2 x square for all x positive. So this crucial property helps us in taking the differentiation under the integral sign and that in the immediate implies because the fundamental solution satisfies the heat equation u_{xx} .

So the third part is proving the initial condition is also satisfied. So, before doing that thing since K is a C infinity function for t positive and just like we did it here taking the differentiation under the integral sign it follow that u dot t is also a function, as a function of x for t positive. So though initial at t equal to 0 g may be just a bounded function but this heat Kernel makes the solution as C infinity function for any t positive. And this is refer to as smoothing effect smoothing property.

(Refer Slide Time: 24:03)



Now we will come to the verification of initial condition. So let g be continuous at η then given $\epsilon > 0$ there exist $\delta > 0$ such that $|g(y) - g(\eta)| < \epsilon$ if $|y - \eta| < 2\delta$ for technical reasons otherwise you simply write δ there then consider $u(x, t) - g(\eta)$.

So let us simplify this so I forgot to mention one more property which I am writing now $\int_{-\infty}^{\infty} K(x-y, t) dy = 1$. This is an easy verification this is simply 1. And you just go back and see the expression for K you this is sorry making mistakes. So this is for all x and t positive.

So there is x here so this is just constant one. So verify this, so this is one more property of one more important property. So this we combining the two integrals and write it as $\int_{-\infty}^{\infty} K(x-y, t) (g(y) - g(\eta)) dy$. And now you break the integral into two parts. So, this mod $x - y$ bigger than or equal to δ and mod $x - y$ less than δ . So δ is coming from here, so the same δ we take we break this integral into two parts. So now we rewrite so let me just again write it.

(Refer Slide Time: 28:52)

$$|u(x,t) - g(\eta)| \leq \int_{|x-y| < \delta} K(x,y,t) |g(y) - g(\eta)| dy + \int_{|x-y| \geq \delta} K(x,y,t) |g(y) - g(\eta)| dy$$

$\leq 2M$ (':g' is odd)

$$|y-\eta| \leq |y-x| + |x-\eta| < \delta + \delta = 2\delta \text{ if } |x-\eta| < \delta$$

$$< \epsilon + 2M \int_{|x-y| \geq \delta} K(x,y,t) dy \rightarrow 0 \text{ as } t \rightarrow 0^+$$

$$< 2\epsilon \text{ if } |x-\eta| < \delta \text{ and } t \text{ suff small}$$

This completes the proof.

Verification of initial condition

Let g be cont at η . Then, given $\epsilon > 0$, $\exists \delta > 0$ such that $|g(y) - g(\eta)| < \epsilon$ if $|y-\eta| < \frac{2\delta}{\alpha}$

Then, consider ($t > 0$)

$$u(x,t) - g(\eta) = \int_{-\infty}^{\infty} K(x,y,t) g(y) dy - g(\eta) \int_{-\infty}^{\infty} K(x,y,t) dy \quad (\forall x, t > 0)$$

$$= \int_{-\infty}^{\infty} K(x,y,t) (g(y) - g(\eta)) dy \quad \int_{-\infty}^{\infty} K(x,y,t) dy = 1 \text{ (verify)}$$

$$= \int_{|x-y| \geq \delta} + \int_{|x-y| < \delta}$$

$$u(x,t) = \frac{1}{\pi} \int_{-\infty}^{\infty} g(y) \left(\int_0^{\infty} e^{-\lambda^2 t} \cos \lambda(y-x) d\lambda \right) dy$$

$$= (4\pi t)^{-1/2} \int_{-\infty}^{\infty} \exp\left(-\frac{(y-x)^2}{4t}\right) g(y) dy$$

Fourier-Poisson formula: $u(x,t) = \int_{-\infty}^{\infty} K(x-y, t) g(y) dy$, $K =$ Fundamental Solution of the heat eqn
 • K is symmetric: $K(x,t) = K(-x,t), t > 0$

So $|x - y| < \delta$ now we take absolute value here less than or equal to $|x - y| < \delta$. K is a positive function so that is also, K is a positive function because it is exponential, so $\int_{-\infty}^{\infty} g(y) dy$ plus $|x - y| > \delta$, K of $|x - y| > \delta$ of y minus $\int_{-\infty}^{\infty} g(y) dy$. Now concentrate on the first integral, so here $|x - y| < \delta$.

So, what about $|y - x|$, so $|y - x|$ by triangle inequality is $|y - x| \leq |y - x| + |x - x|$. So if I also take x within δ neighborhood x . So this is less than δ another δ so that will be 2δ . If $|x - x| < \delta$. And that is anyhow we are going to take the limit as $\delta \rightarrow 0$. So this assumption is fine. So that makes $|y - x| < 2\delta$.

And by the assumed continuity property $\int_{-\infty}^{\infty} g(y) dy$ is less than ϵ so the first integral. So it just this is separate thing so this is less than ϵ . Because the integral of K is 1 so the first integral the second one so we are assuming g is bounded so this is less than or equal to say $2M$ because g is bounded. So M bound is the bound on g . So this is $2M$ and $|x - y| > \delta$ $\int_{-\infty}^{\infty} K(x-y, t) dy$.

And the third property of the fundamental solution since δ is positive say that this goes to 0 as $t \rightarrow 0$. So we make it less than 2ϵ if $|x - x| < \delta$ and t sufficiently small. And that prove the required so finally what we have assume that absolute value of $|x - y| < \delta$ is less than 2ϵ if $|x - x| < \delta$ and t is sufficiently small. And this completes the prove.

So though is started with some heuristic arguments we have arrived at a formula for a function u and which we are able to show that it is a solution of the heat equation and also satisfying the initial condition as described in the theorem. So with we will stop here and in the next lecture we continue from here and discuss uniqueness and some other results. So just remember this and this is refer to as Fourier Poisson formula.

So I conclude here this talk with so though we started with some heuristic arguments but to arrive at this Fourier Poisson formula and under boundedness and continuity assumption on g we did verify that it satisfies the heat equation and the initial condition. And in the next lecture we will discuss uniqueness and non-uniqueness results. And also heat equation in a finite rod. Thank you.