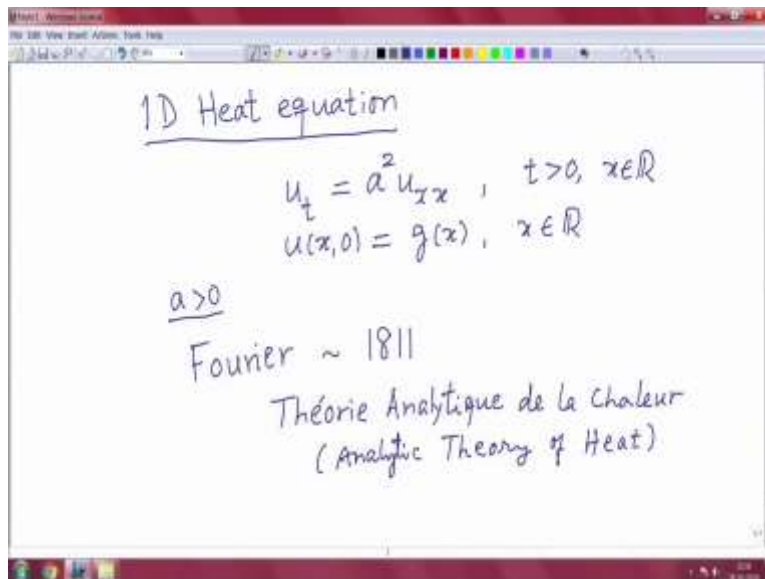


First Course on Partial Defferential Equations – 1
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Lecture 28
One dimensional heat equation-1

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1D Heat equation

$$u_t = a^2 u_{xx}, \quad t > 0, x \in \mathbb{R}$$
$$u(x, 0) = g(x), \quad x \in \mathbb{R}$$

$a > 0$

Fourier ~ 1811

Théorie Analytique de la Chaleur
(Analytic Theory of Heat)

Hello everyone, today we will study the one D heat equation. So this is a equation given by u_t equal to a square u_{xx} again t is positive x is in real line and initial value g of x . So here is a is positive a constant. So physically this represents the evolution of temperature in an in a finite rod. So though not physically meaningful but we will see that just like wave equation. This also helps us in analyzing later the temperature evolution in a finite rod.

So, before we go further few historical remarks. So this equation first appears in the work of Josef Fourier a French scientist. So, somewhere in the beginning of 19 century. There he derives this equation using Newton's law of cooling. So most of the things we study now like fourier series fourier transform they all appear in this works and this is called is in French, so let me just write it *Théorie Analytique*, this is in French, *de la chaleur*.

So simply translated this is analytic theory of heat. So whatever he Fourier in this work it influence a lot especially in the 19 century mathematical physicist, mathematicians so and in fact, now there is separate branch called harmonic analysis which originated in this work Fourier. So he was also successful in deciding to construct chateur to escape extreme weathers, so extreme cold condition would like to remain warm and in extreme heat conditions would like to remain somewhat cool.

So he was successful in determining so remember this is in 19 century so of course now we have air conditioners heaters and what not but in those, so this is more than 200 years. So he was successful in determining the depth at which chateur must be constructed in order to escape extreme weather. So and there he exploited the linearity of this heat equation and this linear super imposition that we all aware of now.

But it in this work he so there will be temperature variation daily temperature variations and yearly temperature variations. So he took all that into account and successful arrived at the depth at which a chateur must be constructed. So a brief account of this I will give you reference.

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T.W. Körner, Fourier Analysis, Cambridge Univ. Press (1988)

$$u_t = a^2 u_{xx}$$

Prototype of a parabolic eqns: has only one real characteristic family: $t = \text{constant}$

Characteristic family: $t = \text{constant}$

Characteristic IVP

- Existence
- Uniqueness
- Cont. dependence on the initial data

Some additional hypothesis may be required

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Théorie Analytique de la Chaleur
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So you can just there is a book by T.W.K Korner title is Fourier Analysis so you will also find a good account of fourier series fourier transform which we are not covering in this course. So this is published by Cambridge University Press now there are many books. So I am just mentioning one. So today what we will do is this try to find a formula for the solution again in the fourier spirit.

So before going further let me mention one thing about this heat equation. So it is a prototype of parabolic equation. So it has only one real characteristic, has only one family when I say one and you can easily see that these characteristics, the characteristic family are given by, characteristic family t equal to constant. And remember we are giving prescribing initial condition at t equal to 0 though physically meaning full.

So at some t equal to 0 initial time u prescribe the temperature in a rod and then you see how it evolves though physically meaningful but as far as mathematician is concerned. So what we are dealing here is so characteristic initial value problem. So I just press this thing characteristic initial value. Because we are prescribing the initial values on a characteristic t equal to 0 is also characteristic.

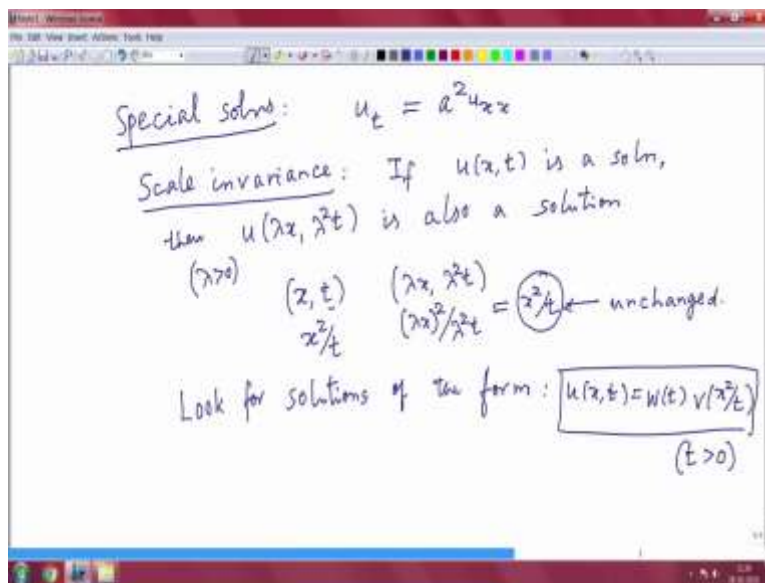
So in mathematically such characteristic initial value problem they may fail. So there are three things we should be interested existence of a solution and then uniqueness this is again important as far as the physics is concerned, uniqueness and then the continuous dependence on the initial

value. So together they are known as Hard mass well pose conditions initially. So when we are dealing with characteristic initial value problem there could be problem with any of these three things one or two.

So that means so even for existence or uniqueness some additional conditions may be prescribe. So some additional hypothesis may be required. So in fact there is no uniqueness in general but if you restrict the initial value to certain class then there is uniqueness. As we see existence will be there again this continuous depended on the initial data is there only if you restrict the initial data to a certain class.

So it is a very difficult to do all this in this set of lectures. So we will just concentrate on certain things but this is an important equation which makes it appearance in several branches of mathematics. So and it is also a typical parabolic equation and it has many properties interesting properties some of them will be discussion.

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So first we will obtain some special solutions and you just let me write it. So we see first you observe that there is a scale in variance. What does this mean? Let me explain that so if $u(x,t)$ is a solution then $u(\lambda x, \lambda^2 t)$ is also a solution. This easily follows by direct differentiation. So if we scale the variable x to λx so λ is a positive number any positive number. And t by $\lambda^2 t$ you see that. So what does this do?

So remember w and v are at our choice so we make first choice, so choose v such that this term in the bracket is 0. And that immediately gives, so one solution, so that is a though it is a second order equation for v . It is a first order equation for v prime, so you do two integrations and you get one particular solution, which is so $v z$ is equal to e to minus z by $4 a$ square that is one solution. And now you plug in this solution v in this thing and you obtain a equation for w prime.

So this is gone and now you differentiate this one and plug in here. So what you get is, so let me just then w satisfies, so you have to work it out. I am just writing w prime plus 1 by $2t$ w equal to 0 and again a particular solution of this first order ODE. Remember this is ODE, so this is $w t$ is t to the minus half. So thus we have obtain a particular solution of the heat equation.

So thus a particular solution of the heat equation is given by u of $x t$ is equal to $w t v$ of x square by t and this one. Now you replace z by x square by t , so let me write this is expression in the exponential is big so let me write it as exponential. So this is x square by $4 a$ square t . A constant multiple of this solution is going to play an important role. So that is called so let us put some name.

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The fundamental ^{solution} of the heat eqn (or opr) is defined by

$$K(x,t) = \begin{cases} (4\pi a^2 t)^{-1/2} \exp\left(-\frac{x^2}{4a^2 t}\right), & t > 0, x \in \mathbb{R} \\ 0, & t < 0, x \in \mathbb{R} \end{cases}$$

- K has a singularity at $t=0$
- For $t > 0$, $K \in C^\infty$ for $x \in \mathbb{R}$
- For $t > 0$, $K_t = a^2 K_{xx}$

The diagram shows a graph of K versus x for three different times t_1, t_2, t_3 where $t_1 < t_2 < t_3$. The curves are bell-shaped and centered at $x=0$. As time increases, the peak of the curve decreases in height and the curve spreads out horizontally, illustrating the diffusion process.

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So, the fundamental solution of the heat equation or heat operator fundamental solution is defined by or operator is defined by denoted by $K(x,t)$. So there is a constant multiple here. So $4\pi a^2 t$ to the minus half exponential so this is t positive x in \mathbb{R} and 0 where t less than 0. So what we obtain here, so I am just adding putting a constant term there. So this $4\pi a^2 t$ we will see the reason for putting that particular constant how that plays an important role.

So immediate observations about this K . So for t positive exponential is infinitely differentiable also this t to the minus half. So immediately just by looking at the expression K has a singularity at t equal to 0 for t positive K is an infinitely differentiable function for x in \mathbb{R} . And just now we have seen that, so again for t positive K satisfies K_t equal to a square K_{xx} .

So K satisfies the heat equation in the region t positive and it is a very smooth function it has only singularity at t equal to 0. Some of them, some profiles of K I can show you, so this is also a since this a Gaussian curve so it is also refer to as Gaussian. So it is, so yes t becomes smaller and smaller you will notice that the exponential term becomes closer to 1. Because t is infinity so this will become 0.

So exponential become 1 but there is a t to the minus half factor so as t becomes larger and larger and we will see that space is less here, so this safe, so t equal to t^4 equal to t^3 sometimes. So this is t equal to t^2 and t equal to t^1 and you drawing x and K . So here t^4 is less than t^3 less than t^2 . So important thing we will notice that that we are going to use that the fundamental solutions

satisfy the heat equation in the region $t > 0$ that is important and we will also see that this K is a very smooth function. So it is infinitely differentiable.

So these two things will help us in finding a formula for the initial value problem. So this occurs in the formula for the solution of the initial value problem. So anything else yeah with this we close this lecture and in the next lecture we will obtain the solution of the initial value problem for the heat equation. And there again we see the role played this fundamental solution. Thank you.