

Advanced Engineering Mathematics
Prof. P. N. Agrawal
Department of Mathematics
Indian Institute of Technology - Roorkee

Lecture – 04
Applications to the Problems of Potential Flow - I

Hello friends, welcome to my lecture on applications to the problems of potential flow, the Laplace equation $\nabla^2 u = 0$ is one of the most important partial differential equations and engineering mathematics because it occurs in connection with gravitational fields, electrostatic fields, a steady state heat conduction in incompressible fluid flow etc.

(Refer Slide Time: 00:50)

Laplace equation $\nabla^2 u = 0$ is one of the most important partial differential equations in engineering mathematics because it occurs in connection with gravitational fields, electrostatic fields, steady state heat conduction, incompressible fluid flow etc. The theory of the solutions of this equation is called Potential theory.

In the two dimensional case, Laplace equation becomes

$$u_{xx} + u_{yy} = 0.$$

We shall see that its solutions are closely related to complex analytic functions.

The theory of solutions of the Laplace equation is called the potential theory. Now, in the 2 dimensional case, the Laplace equation $\nabla^2 u = 0$ becomes $u_{xx} + u_{yy} = 0$, we shall see that its solutions are closely related to complex analytic functions, when we will look at the solution of 2 dimensional Laplace equation, we are going to see that they are closely related to complex analytic functions.

(Refer Slide Time: 01:24)

Coulomb's law

Let 'r' be the distance between two point electric charges q_1 and q_2 . Then the force of attraction and repulsion between them is given in magnitude by

$$F = \frac{q_1 q_2}{kr^2}$$

where the constant k is called the dielectric constant and depends on the medium; in a vacuum $k = 1$, in other cases $k > 1$.

q_1 q_2

Let us first discuss Coulomb's law; let r be the distance between 2 point electric charges q_1 and q_2 , let us say 2 point electric charges q_1 and q_2 , okay then the force of attraction and repulsion between them is given in magnitude by $F = q_1 q_2$ over kr square. Here, k is the dielectric constant and it depends on the medium in a vacuum, the value of k is 1 and in other cases, the value of k is > 1 .

(Refer Slide Time: 01:56)

Electrostatic Potential

Consider a discrete or continuous or a combination of both, charge distribution. This charge distribution sets up an electric field. If a unit positive charge (small enough so as to not affect the field appreciably) is placed at any point 'A' not already occupied by charge then the force acting on this charge is called the electric field intensity at 'A' and is denoted by ϵ . This force is derivable from a potential ϕ which is called the electrostatic potential. We have

$$\epsilon = -\nabla\phi.$$

If the charge distribution is two dimensional then

$$\epsilon = E_x + iE_y = -\phi_x - i\phi_y$$

$$\Rightarrow E_x = -\phi_x, E_y = -\phi_y$$

Now, let us consider a discrete or continuous or a combination of both charge distribution, this charge distribution sets up an electric field, if a unit positive charge is small enough so as to not affect the field appreciably is placed at any point A not already occupied by charge then the force acting on this charge is called the electric field intensity at A and is denoted by this symbol, okay.

This force is derivable from a potential ϕ which is called the electrostatic potential, okay and we have this; we can call it as big epsilon = - del phi, okay. If the charge distribution is 2 dimensional than this big epsilon is $E_x + i E_y$, okay and this big epsilon is - del phi, so this is = - phi x - i phi y, when we write the; this is the notation in the complex plane, okay as such del phi is i phi x + j phi y in 2 dimensions.

But when we write it corresponding representation in the complex, okay it will be written as $-5x - i \phi y$, so we have - phi x - i phi y and here also you can see this force, epsilon is written as $E_x + i E_y$, this is the complex representation of the field intensity which is $E_x i + E_y j$, okay, this iota is the complex number, this iota is not the unit vector, okay, this iota is; here this iota is square root -1, so $E_x =$ then - phi x and $E_y = - \phi y$.

(Refer Slide Time: 04:12)

Electrostatic Potential cont...

It follows that in any region, free of charge,

$$\frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} = 0$$

$\Rightarrow \phi_{xx} + \phi_{yy} = 0$

$\Rightarrow \phi$ is harmonic at all points not occupied by charge.

$E_x = -\phi_x$
 $E_y = -\phi_y$
 $E_{xx} = -\phi_{xx}$
 $E_{yy} = -\phi_{yy}$
 Since $\frac{\partial E_x}{\partial x} + \frac{\partial E_y}{\partial y} = 0$
 we have $\phi_{xx} + \phi_{yy} = 0$
 $\Rightarrow \phi$ is a harmonic function

Now, it follow that in any region free of charge, the partial derivatives of E_x + partial derivative of E_y that is = 0. Now, $E_x =$; we have seen $E_x = - \phi x$ and $E_y = - \phi y$, so $E_{xx} = - \phi_{xx}$ and $E_{yy} = - \phi_{yy}$, which implies that we have the partial derivative of E_x with respect to x + partial derivative of E_y with respect to y , okay, this is what we have, okay. So, since we have; since partial derivative of E_x with respect to x + partial derivative of E_y with respect to y is 0.

We have $\phi_{xx} + \phi_{yy} = 0$, okay which means that ϕ is a harmonic function, at all points which are not occupied by a charge.

(Refer Slide Time: 05:39)

The Complex Electrostatic Potential

Since ϕ is harmonic there exists a harmonic function ψ conjugate to ϕ such that $\Omega(z) = \phi(x, y) + i\psi(x, y)$ is analytic in any region not occupied by charge. The function $\Omega(z)$ is called the complex electrostatic potential. The curve $\phi(x, y) = \alpha$, $\psi(x, y) = \beta$ are called the equipotential lines and the lines of force respectively. The lines of force cut the equipotential lines orthogonally. Note that $\psi(x, y) = \beta$ are called the lines of force because the electrical force being $-\nabla\phi$ acts in a direction perpendicular to the equipotential lines i.e. along the lines $\psi(x, y) = \beta$.

$\underline{\underline{\epsilon}} = -\nabla\phi$ $\nabla\phi$ is a vector perpendicular to $\phi(x, y) = \alpha$

Now, we have harmonic function ϕ with us, okay, so we can find its corresponding conjugate harmonic functions ψ , okay such that the complex function, $\omega(z) = \phi(x, y) + i\psi(x, y)$, $\phi(x, y)$ is the real part of $\omega(z)$ and $\psi(x, y)$ is the imaginary part of $\omega(z)$. Now, this function is analytic in any region not occupied by the charge, the function $\omega(z)$ is called the complex electrostatic potential, okay, this function is called the complex electrostatic potential.

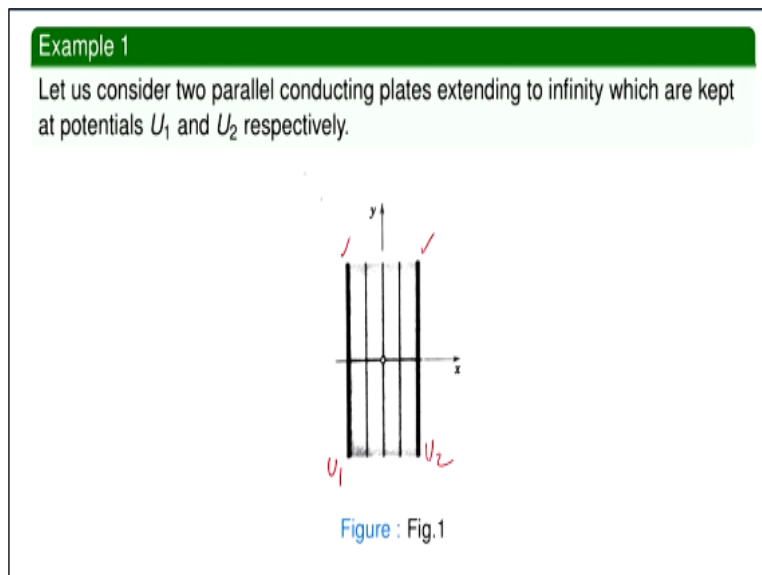
The curve $\phi(x, y) = \alpha$, okay and $\psi(x, y) = \beta$, find 2 level curves, family of level curves, okay since ϕ is the potential, okay, we are writing this big epsilon as $-\nabla\phi$, so that is the ϕ is the potential there, so $\phi(x, y) = \alpha$ means, the curves $\phi(x, y) = \alpha$, the level curve $\phi(x, y) = \alpha$ are called as equipotential lines or you can call them as equipotential curves because on every curve, the potential is the same.

And then $\psi(x, y) = \beta$ are known as the lines of force, okay respectively, lines; we write lines but they can be called as curves, so the lines of force cut the; now we have seen that the family of level curves $\phi(x, y) = \alpha$ and $\psi(x, y) = \beta$, they cut each other orthogonally, okay so, the lines of force cut, the equipotential line orthogonally. Note that $\psi(x, y)$; why we call $\psi(x, y) = \beta$ as lines of force.

Because $\psi = \beta$, the direction of the electrical force, okay we had the electrical force $-\nabla\phi$, okay and this $\nabla\phi$ is a vector which is perpendicular to $\phi = \alpha$, we know this that $\nabla\phi$ is a vector, a perpendicular to $\phi = \text{constant}$, okay. So, $\nabla\phi$ is a vector perpendicular to $\phi = \alpha$ and a $\psi = \beta$ is also perpendicular to this $\phi = \alpha$, therefore the direction of $\psi = \beta$ and the vector; electrical force, this electrical force are same.

And therefore, we call $\psi = \beta$ s, lines of force, so let us note that $\psi = \beta$ are called the lines of force because the electrical force which is $-\nabla\phi$, it reacts in a direction perpendicular to the equipotential lines $\phi = \alpha$, okay, that is along the lines, $\psi = \beta$.

(Refer Slide Time: 08:33)



Now, let us consider 2 parallel conducting plates; one conducting plate is this one, the other conducting plate is this one, they are extending to infinity and they are capped at potentials; u_1 and u_2 , so this is having the potential u_1 throughout and these having potential u_2 throughout and they are infinitely long.

(Refer Slide Time: 08:54)

Example cont...

Then from the shape of the plates it follows that u depends only on x . Hence from $\nabla^2 u = 0$, we have

$$u_{xx} = 0 \text{ or } u(x) = ax + b.$$

Then ② - ①
 $\Rightarrow U_2 - U_1 = 2a$
 $\Rightarrow a = \frac{U_2 - U_1}{2}$

$$b = \frac{U_1 + U_2}{2}, a = \frac{-U_1 + U_2}{2}.$$

Hence, the potential of the field between the given plates is

$$u(x) = \frac{1}{2}(-U_1 + U_2)x + \frac{1}{2}(U_1 + U_2)$$

\Rightarrow the equipotential surfaces are parallel lines.

$$\begin{aligned} u_{xx} + u_{yy} &= 0 \\ \Downarrow \\ u_{xx} &= 0 \\ \Downarrow \\ u_x &= a \\ \Downarrow \\ u(x) &= ax + b \\ U_1 = -a + b &\text{--- (1)} \\ U_2 = a + b &\text{--- (2)} \\ \text{①} + \text{②} &\Rightarrow \frac{U_1 + U_2}{2} = b \end{aligned}$$

Then from the shape of the throughout; shape of the plates, it follow that u depends only on x , you can see that U is constant on this plate, okay, it does not depend on y , it is $E = U_1$, okay throughout this plate, and this is let us say this plate is $x =$; this plate is parallel to y axis, so the equation of this plate is $x =$ some constant, okay and so for all values of y , you can see $U =$; if you denote the potential then $U = U_1$.

And similarly, here this is; this plate is given by another equation $h =$ some constant, okay, so along this plate also for all values of y , okay we have $U = U_2$, so u depends only on x , it does not depend on y , okay. So, since U depends only on x , the Laplace equation del square $U = 0$ gives us $U_{xx} = 0$, del square u is $U_{xx} + U_{yy} = 0$, okay. This implies that $U_{xx} = 0$ because U is independent of y , okay.

Now, $U_{xx} = 0$ is; when we integrate it twice, if you integrate it once, what you get; $U_x =$ some constants say, a okay and when you integrate it again with respect to x , you get $U_x = ax + b$, so U is the function of x , it is linear function, $ax + b$. Now, let us use the condition that when $x =$; we have the condition that and this plate $U = U_1$, on this plate $U = U_2$, so when we use that we get the values of a and b , okay.

So, let us put them as $x = 1$ and -1 , so this is $x = -1$, okay and this is $x = 1$, okay, so when you take $x = -1$, $x = 1$, what you get; $U_x = U_1$, so $U_1 = -a + b$, okay, U_1 is $-a + b$ and $u_2 = a + b$,

okay, so when you add the 2 equations, $U_1 + U_2 = b$, okay we get the value of U_1 ; value of b as $U_1 + U_2/2$ and when we subtract the coefficient 1 from 2, so $2 - 1$ gives you what?

$U_2 - U_1 = 2a$ so that gives $a = (U_2 - U_1)/2$ okay, so we get the values of a and b by taking the 2 plates at left plate at $x = -1$ and the other plate at $x = 1$, okay we get the values of a and b and then putting the values of a and b then $U_x = 1/2 - U_1 + U_2 \cdot x$, $1/2 U_1 + U_2 =$, okay. Now, U_x is potential, okay U_x is potential, so if you take u_x potential to be constant then you see that $U_x =$ constant gives $-1/2 U_2 - U_1 + 1/2 U_2 - U_1 \cdot x + 1/2 U_1 + U_2 =$ constant, okay.

Since $U_1 + U_2/2$ is the constant, $U_2 - U_1/2$ is the constant, this implies that $x =$ some constant, okay, some constant c , okay. So, when we take U_{xy} ; $U_x =$ some constant we get $x =$ some constant, so means for different values of the constant will get different lines which are parallel to y axis, so equipotential lines are surface lines; equipotential surfaces are parallel lines, they are parallel to y axis.

(Refer Slide Time: 13:37)

Example cont...

If $v(x, y)$ is the conjugate harmonic function of $u(x, y)$, then from (1), we have

$$v(x, y) = \frac{1}{2}(U_2 - U_1)y + c$$

Thus, the complex potential is

$$\Omega(z) = u(x, y) + iv(x, y) = az + (b + ic)$$

The lines of force are given by

$$v(x, y) = \text{constant} \Rightarrow ay + c = \text{constant} \Rightarrow y = \text{constant}$$

\Rightarrow the lines of force are parallel to x -axis.

Handwritten notes:

we have $u(x) = ax + b$
 where $a = \frac{U_2 - U_1}{2}$
 $b = \frac{U_1 + U_2}{2}$

$\frac{\partial u}{\partial x} = \frac{\partial v}{\partial y}$ and $\frac{\partial u}{\partial y} = -\frac{\partial v}{\partial x}$
 $a = \frac{\partial v}{\partial y}$
 $v = ay + c(x)$
 $\frac{\partial v}{\partial x} = c'(x) = 0 \Rightarrow c(x) = \text{constant}$
 i.e. $c(x) = c$

If v_{xy} is the conjugate harmonic function of u_{xy} then we can find the value of v_{xy} , why using Cauchy- Riemann equations, we have $u_{xy} = v_{yx}$; we have u_x ; u is not depending on y , so $u_x = ax + b$ we have, we have $a = (U_2 - U_1)/2$, $b = (U_1 + U_2)/2$, okay. Now, let us use the Cauchy- Riemann equations, so $u_x = v_y$ and $u_y = -v_x$, so let us first use this equation, so $u_x = v_y = a$, okay, so $a = v_y$,

so this implies $v = ay + \text{some constant}$ let us say, some constant of integration which will be a function of x , so let us write $c(x)$, okay.

Now, we will make use of this to determine this unknown function of x , okay, so this gives you; let us differentiate this with respect to x , so derivative of b with respect to x gives you y is independent effect, so this will give you $c'(x)$, okay but from here this $v_x = -u_y$, so this $= -u_y$ is how much; $u_y = 0$, okay, u_y is not depending on y , so this implies $c'(x) = 0$ and which implies $c(x)$ is a constant, okay, it is a constant.

I have denoted this $c(x)/c$, okay, so that is $c(x)$ is some constant c , okay, so what do we get; by $v = ay + c$, where a is this, $\frac{1}{2}(u_2 - u_1)$, okay, so $v_{xy} = \text{this one}$, $\frac{1}{2}(u_2 - u_1) \cdot y + c$ and now let us find the complex potential because we have u and v with us, so $\omega(z) = u_{xy} + i v_{xy}$, $u_{xy} = ax + b$, so $\omega(z) = u_{xy}$ is $ax + b + i$ times v_{xy} is; this is a , okay, so $ay + c$, so what do we get; $a(x + iy) + b + i \cdot c$.

Now, this is $az + b + ic$, okay, so this $az + b + ic$; a is $\frac{1}{2}(u_2 - u_1)$ and b is $\frac{1}{2}(u_1 + u_2)$, c is a real constant, the lines of force are now, what are the lines of force? Lines of force are given by $u_{xy} = \text{constant}$ gave us the equipotential lines, the $v_{xy} = \text{constant}$ gives us the lines of force, so $v_{xy} = \text{constant}$ means $v_{xy} = \text{constant}$, implies $ay + c = \text{constant}$, okay, now c is the constant, a is the constant, so this implies $y = \text{constant}$ okay.

So, the lines of force are parallel to x axis, okay, so lines of force are parallel to x axis the equipotential lines are parallel to y axis, so they cut each other at right angles, they define 2 orthogonal families.

(Refer Slide Time: 17:40)

Example 2

Let us consider two conducting coaxial cylinders which extend to infinity on both sides and are kept at potentials U_1 and U_2 respectively.

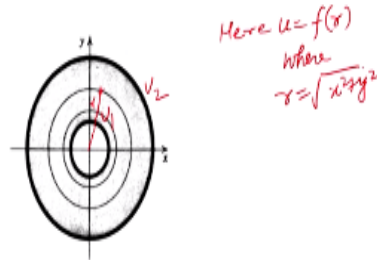


Figure : Fig.2

Now, let us consider two conducting coaxial cylinders, so we have 2 coaxial cylinder, they have the same axis, okay they are coaxial which extend to infinity on both sides okay, so on both sides of the axis they are extending to infinity and are kept at potentials U_1 and U_2 respectively, so let us say this is U_1 okay, this is U_1 , this is U_2 , okay, they are kept at potentials U_1 and U_2 respectively.

So, here what is happening is that the potential U okay at any point is given by at any let us say at this point, okay, it is given by the distance of this point from the origin, let us say this distance of this point from origin is r , then here u is the function of r , okay, we have the 2 conducting coaxial cylinders which are kept at potentials U_1 at; so that mean that on the coaxial cylinders, the constant; the potential is remains constant.

So, potential varies when r varies, okay, u is the function only of r , so u is the function of r , and r where $r = \text{under root } x^2 + y^2$.

(Refer Slide Time: 19:06)

A and b are constants, okay, $a \ln r + b = \text{constant}$ means $\ln r$ is constant and which implies that r is constant, okay, r is constant means what; are the equipotential curves, u is the potential, okay, so equipotential curves are concentric circle with centre at 00, $r = \text{constant}$ means it defines a circle with centre at 00 and $v_{xy} = \text{constant}$; v_{xy} is what; v_{xy} , let us find v_{xy} , we can find by using the Cauchy- Riemann equations in polar form are $u_r = 1/r$ v_{θ} .

So, using Cauchy-Riemann equations, $u_r = 1/r$ v_{θ} , what do we get; $u_r =$; we have found f_r , $f_r = a \ln r + b$, so $u = a \ln r + b$ gives partial derivative of u with respect to r ; $u_r = a/r$, okay, C-R equations in polar form are $u_r = 1/r$ b_{θ} $u_{\theta} = -r v_r$ okay. Now, $u = a \ln r + b$, so $u_r = a/r$, so $a/r = 1/r$ v_{θ} , first equation gives us this which implies that $v_{\theta} = a$ and when we integrate with respect to θ what we get?

$V = a \theta + \text{some function of } r$, I can write that as $\phi(r)$, now here $u_{\theta} = 0$ because u does not depend on θ , so this second equation gives us $0 = -r v_r$ which implies that $v_r = 0$, okay and this mean that when you differentiate this v with respect to r , what do we get; $v = a \theta + \phi(r)$, when we differentiate with respect to r , so we have this, θ is independent of r , so we get $\phi'(r)$.

But this is $= 0$, let us use this, so this is $= 0$ implies $\phi'(r) = \text{arbitrary constant say, } c$ okay, so what do we get then, okay hence, $v =$; let us we call this; $v = a \theta + \text{some arbitrary constant } c$, so if $v_{xy} = \text{constant}$, okay, $v_{xy} = \text{constant}$ means $a \theta + c = \text{constant}$ and this will implies that $\theta = \text{constant}$, okay, so the lines of force are a straight lines which pass through 00, okay because θ is constant.

So, you can see the lines of force are passing through origin, okay and equipotential curves are concentric circles with centre at origin, so at the point of intersection they are; they cut each other at right angles, with this I would like to end my lecture, thank you very much for your attention.