

Advanced Engineering Mathematics
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Lecture – 38
Fourier Integral and Fourier Transforms

Hello friends let's welcome to my lecture on Fourier integral and Fourier transforms. Let us first discuss Fourier integral. We have seen Fourier series are powerful tools in treating various problems involving periodic functions but in many practical problems we do not have periodic functions. Hence it is desirable to generalize the method of Fourier series to include non-periodic functions. Let us consider 2 simple examples of periodic functions of period T .

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The Fourier integral

Fourier series are powerful tools in treating various problems involving periodic functions. But many practical problems do not involve periodic functions hence it is desirable to generalize the method of Fourier series to include non-periodic functions. Let us consider two simple examples of periodic functions of period T and see what happens if we let $T \rightarrow \infty$.

And see what happens if we let T go to infinity.

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Example 1

Consider the function

$$f_T(x) = \begin{cases} 0, & \text{when } -\frac{T}{2} < x < -1 \\ 1, & \text{when } -1 < x < 1 \\ 0, & \text{when } 1 < x < \frac{T}{2}. \end{cases}$$

having period $T > 2$.

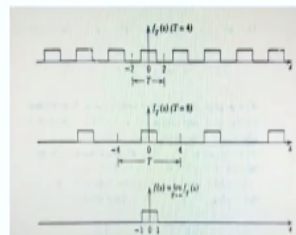
So, consider this function $f_T(x) = 0$ / interval- $T/2 - 1$ 1 when $-1 < x < 1$ 0 when $1 < x < T/2$ and let us assume that $T > 2$.

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Example 1 cont...

When $T \rightarrow \infty$, we obtain a function $f(x)$ which is not periodic

$$f(x) = \lim_{T \rightarrow \infty} f_T(x) = \begin{cases} 1 & \text{when } -1 < x < 1 \\ 0, & \text{otherwise} \end{cases}$$



So, we have this graph okay this is the graph of $f_T(x)$ when we take $T=4$ and we are taking this example for values of T more than 2 let us 1 st consider this special case when $T=4$ then the graph of the function $f_T(x)$ which is 0 when $-x$ belongs to the open interval $-T/2 - 1$ when $1 < x$ belongs to the open interval -1 1 and 0 when x belongs to open interval 1 $T/2$ is this 1. This is the graph of $f_T(x)$ when $T=4$ and this graph changes to this graph.

When we take $T = 8$ so this is the graph for $T=8$ now when T goes to infinity here when in this example when T goes to infinity what happens we see that $f_T x$ goes to let us say f_x which is 0/ interval – infinity to -1/ the interval -1 to 1 and 0 when $1 < x < \infty$ we get the definition of f_T f_x will be = limit of $f_T x$ goes to infinity it is 1 when $-1 < x < 1$ and 0 other wise so when we consider the limit of $f_T x$ as T goes to infinity.

The limit in function f_x assumes value 1/ the open interval -1 1 else where it is 0 so this is the graph of the limiting function f_x you can see. And you can also see that this function is not a periodic function so we get a non-periodic function when we let T go to infinity now let us consider another example.

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Example 2

Let $f_T(x) = e^{-|x|}$, for $-\frac{T}{2} < x < \frac{T}{2}$ and $f_T(x+T) = f_T(x)$.

When $T \rightarrow \infty$ we get the function $f(x) = e^{-|x|}$ which is no longer periodic

$\lim_{T \rightarrow \infty} f_T(x) = f(x)$ $f(x) = \lim_{T \rightarrow \infty} f_T(x) = e^{-|x|}$, $-\infty < x < \infty$.

$f(x) = e^{-|x|}$, $-\infty < x < \infty$

Suppose $f_T x = e$ to the power – mod of x for the interval $-T/2 < x < T/2$ so this is the interval – $T/2$ and this is the graph of $f_T x$ this graph of $f_T x$ this one / the interval $-T/2$ $T/2$ and we are assuming it to be T periodic so by using T periodicity or with period T we can extend this graph / the whole real line and this is how the graph of $f_T x$ looks like for all real values of x now let us take T goes to infinity $f_T x$ let us say goes to f_x limit of $f_T x$ as T goes to infinity.

We are taking $s f_x$ so then what will happen f_x will be = e to the power – mod of x / the interval – infinity to infinity and this is how we get the graph of f_x and you can see that this function f_x is not a periodic function so we get a non-periodic function so these are 2 examples where we have

tried to show that a periodic function $f_T(x)$ when T goes to infinity does not lead us to a periodic function. So, how we will treat non-periodic functions that we will see through Fourier integral

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Fourier integral cont...

Let us consider any periodic function $f_T(x)$ of period T which can be represented by the Fourier series

$$f(x) = \frac{a_0}{2} + \sum_{n=1}^{\infty} (a_n \cos(\omega_n x) + b_n \sin(\omega_n x)) \quad \text{where } \omega_n = \frac{2\pi n}{T}.$$

If we let $T \rightarrow \infty$ and assume that the resulting non-periodic function $f(x) = \lim_{T \rightarrow \infty} f_T(x)$ is absolutely integrable i.e. $\int_{-\infty}^{\infty} |f(x)| dx$ exists then we obtain

$$f(x) = \frac{1}{\pi} \int_0^{\infty} \{A(\omega) \cos(\omega x) + B(\omega) \sin(\omega x)\} d\omega \quad (1)$$

where

$$A(\omega) = \int_{-\infty}^{\infty} f(v) \cos(\omega v) dv$$

So, let us consider any periodic function $f_T(x)$ of period T which can be represented by this Fourier series $f(x) = a_0/2 + \sum_{n=1}^{\infty} (a_n \cos(\omega_n x) + b_n \sin(\omega_n x))$ where ω_n we are writing as $2\pi n/T$ because $f_T(x)$ has period T and in the case of period T we know that we get the terms in \sin and $\sin x$ $\sin 2\pi n x/T$ $\sin 2\pi n x/T$ so for convenience we are writing $2\pi n/T$ as ω_n so we get this Fourier series of the function $f_T(x)$.

Now let us take T goes to infinity so if we take T goes to infinity and assume that the resulting non-periodic function $f(x)$ is absolutely integrable so this condition we put on $f(x)$ that $f(x)$ is absolutely integrable meaning that $\int_{-\infty}^{\infty} |f(x)| dx$ exists then we obtain $f(x) = 1/\pi \int_0^{\infty} (A(\omega) \cos(\omega x) + B(\omega) \sin(\omega x)) d\omega$ where $A(\omega)$ is given by this integral $\int_{-\infty}^{\infty} f(v) \cos(\omega v) dv$.

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and

$$B(\omega) = \int_{-\infty}^{\infty} f(v) \sin \omega v \, dv$$

The integral on right hand side of (1) is called as the Fourier integral. The sufficient condition for the validity of (1) are given in the following result:

Theorem

If $f(x)$ is piecewise continuous in every finite interval and has a left and right hand derivatives at every point and f is absolutely integrable, then $f(x)$ can be represented by a Fourier integral. At a point of discontinuity, the value of the Fourier integral is equal to the average of the left and right hand limits at that point.

And $B(\omega)$ given by $\int_{-\infty}^{\infty} f(v) \sin \omega v \, dv$ the integral on the right hand side of (1) this integral is known as the Fourier integral. The sufficient condition for the validity of Fourier integral is given in the following theorem. If $f(x)$ is piecewise continuous in every finite interval and has a left and right hand derivative at every point and f is absolutely integrable then $f(x)$ can be represented by a Fourier integral.

At a point of discontinuity the value of the Fourier integral is equal to the average of the left hand and right hand limits at that point. So here when we write this Fourier integral to be equal to $f(x)$ we are assuming that x is a continuity point of $f(x)$. If x is the discontinuity point of f then this $f(x)$ will be replaced by the average of left and right hand limits of f .

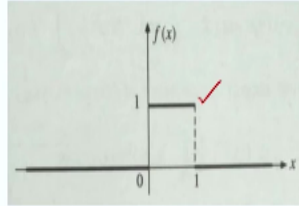
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Example 3

Find the Fourier integral representation of the function

$$f(x) = \begin{cases} 0, & \text{when } x < 0 \\ 1, & \text{when } 0 \leq x \leq 1 \\ 0, & \text{when } x > 1. \end{cases}$$

$$\begin{aligned} B(w) &= \int_{-\infty}^{\infty} f(v) \sin wv \, dv \\ &= \int_0^1 1 \sin wv \, dv \\ &= \left(-\frac{\cos wv}{w} \right)_0^1 \\ &= \frac{1 - \cos w}{w} \end{aligned}$$



$$\begin{aligned} \int_{-\infty}^{\infty} |f(x)| \, dx &= \int_0^1 1 \, dx = 1 < \infty \\ A(w) &= \int_{-\infty}^{\infty} f(v) \cos wv \, dv \\ &= \int_0^1 \cos wv \, dv \\ &= \left(\frac{\sin wv}{w} \right)_0^1 \\ &= \frac{\sin w}{w} \end{aligned}$$

Hence show that $\int_0^{\infty} \frac{\sin(x/2)}{x} \, dx = \frac{\pi}{2}$.

Now let us look at this example suppose if we have the function $f(x) = 0$ when x is $>$ than 0 1 0 is $0 \leq x \leq 1$ and 0 when $x > 1$ so this is the graph of this function we can let us try to find the Fourier integral representation of this we can see that it satisfies all the conditions which are stated here in the theorem it is piecewise condition in a every finite interval it has left hand and right hand derivative at every point and it is absolutely integrable.

We can see absolutely integrable integral $-\infty$ to ∞ mod of $f(x) \, dx$ will reduce to integral 0 to 1 because otherwise it is 0 and it is everywhere else other than the interval is 0 so 0 to 1 mod of $f(x)$ is $f(x) = 0$ to 1 mod of $1 \, dx$ which is $= 1$ so the integral converges integral $-\infty$ to ∞ mod of $f(x) \, dx$ is $< \infty$ so it is absolutely integrable so therefore we can find its Fourier integral representation.

In order to find the Fourier integral representation we need to calculate a ω and b ω here so let us find the value of $A(\omega)$ will be $=$ let us go to the definition $-\infty$ to ∞ $f(v) \cos \omega v \, dv$ so using the definition of $f(x)$ it is 0 to 1 $f(v) = 1$ we get $\cos \omega v \, dv$ so this is $\sin \omega v / \omega$ let us put the limits and we get $\sin \omega / \omega$ let us find the value of $B(\omega)$ ω is integral $-\infty$ to ∞ $f(v) \sin \omega v \, dv$.

So, this will reduce to integral 0 to 1 $f(v) \sin \omega v \, dv$ but $f(v) = 1$ / the interval 0 1 so 0 to 1 $\sin \omega v \, dv$ with integrate with respect to v we get $-\cos \omega v / \omega$ 0 1 this is $1 - \cos$

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So, we get the following we get the value of $A \omega \sin \omega x / \omega$ we get the value of $B \omega (1 - \cos \omega x) / \omega$ then the Fourier integral representation of $f(x)$ given by $\frac{1}{\pi} \int_0^\infty \int_{-\infty}^\infty f(x) \cos \omega(x-t) d\omega$ we have so this is $\frac{1}{\pi} \int_0^\infty \int_{-\infty}^\infty f(x) \cos \omega(x-t) d\omega$ we get $A \omega = \sin \omega x / \omega \cos \omega x + B \omega (1 - \cos \omega x) / \omega$.

That is $\cos \omega x - \frac{1}{2}d\omega$ we get the Fourier integral representation of f and this will be $=f(x)$ at each point of continuity of f this will be $=f(x)$ when x is a continuity point of f and when you have taken any discontinuity point then it will be average of left hand and right hand limit now we have to find the value of this integral 0 to $\infty \sin x/2/x$ so let us see we can consider $x = 1/2$ so let us take $x = 1/2$.

Then what we would notice $x=1/2$ is here and that $1/2 f(x)$ is a continuous function so what we have $x = 1/2$ is the continuity point and $f(1/2) = 1$ so we get $1=2/\pi \int_0^\infty \sin \omega x/2 dx = \pi/2$ and we get this result so $\int_0^\infty \sin x/2 /x dx = \pi/2$.

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Note

If f is an even function in $(-\infty, \infty)$ then $B(\omega) = 0$, hence

$$f(x) = \frac{1}{\pi} \int_0^\infty A(\omega) \cos(\omega x) d\omega, \quad (2)$$

where

$$A(\omega) = 2 \int_0^\infty f(v) \cos(\omega v) dv$$

and if f is an odd function, $A(\omega) = 0$, hence

$$f(x) = \frac{1}{\pi} \int_0^\infty B(\omega) \sin(\omega x) d\omega, \quad (3)$$

where

$$B(\omega) = 2 \int_0^\infty f(v) \sin(\omega v) dv$$

Now if f is an even function in the interval $-\infty$ to ∞ then you can see $B(\omega)$ will be $= 0$ because f is an even function and $\sin \omega v$ is an odd function of v so their product is an odd function and therefore $\int_{-\infty}^\infty f(v) \sin \omega v dv$ will be 0 so $B(\omega)$ is 0 and therefore $f(x) = \frac{1}{\pi} \int_0^\infty A(\omega) \cos \omega x d\omega$.

Where $A(\omega)$ now $A(\omega)$ will become twice $\int_0^\infty f(v) \cos \omega v dv$ this is $A(\omega)$ so if f is even $\cos \omega v$ is also even then this $A(\omega)$ will become 2 times $\int_0^\infty f(v) \cos \omega v dv$ so $f(x)$ will be given by this integral where $A(\omega)$ will be given by twice $\int_0^\infty f(v) \cos \omega v dv$ if f is an odd function then $B(\omega)$ will become double $B(\omega)$ will become twice $\int_0^\infty f(v) \sin \omega v dv$.

While $A(\omega)$ will become 0 because $A(\omega)$ is integral $-\infty$ to ∞ $f(v) \cos \omega v dv$ so when f is odd $\cos \omega v$ is even their product will be odd and therefore their integral $-\infty$ to ∞ will be 0 so $f(x) = \frac{1}{\pi} \int_0^{\infty} B(\omega) \sin \omega x d\omega$ where $B(\omega)$ is $\frac{1}{\pi}$ this twice $\int_0^{\infty} f(v) \sin \omega v dv$.

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The integrals in (2) and (3) are called Fourier cosine and sine integrals respectively.

Remark

Just as in the case of half range Fourier series, a function $f(x)$ defined over the interval $(0, \infty)$ may be expressed as a Fourier sine or cosine integral. Further, we observe that the representation of a non-periodic function $f(x)$ given by (1) is similar to the representation of a function by a Fourier series except that the range is now $(-\infty, \infty)$ and the summation has been replaced by integration.

The integrals in 2 and 3 are called Fourier cosine and sine integral so this is called as Fourier cosine integral this 1 is called Fourier cosine integral this 1 is called Fourier sin integral now just as in the case of 1/2 range Fourier series a function defined over the interval suppose the function there in the 1/2 range Fourier series what we had we had function defined over the interval 0 to 1 so over the 1/2 interval $-\infty$ to ∞ 0.

We were free to define in any way so we define it by odd and even periodic extensions when we used odd periodic extension we found 1/2 range Fourier sine series when we define the function f as an even extension / the interval -1 to 1 then we found Fourier sine series so just like that so just like that in the case of 1/2 Fourier series a function defined / the interval 0 to ∞ may be expressed as a Fourier sine or cosine integral/ interval $-\infty$ to 0.

We can define by taking even extension or by odd extension then the function will be even or odd and so we will get the corresponding Fourier cosine or Fourier sine integral so we observe that the presentation of a non-periodic function $f(x)$ given by 1 is similar and now we can see here

this integral you can see Fourier cosine integral this Fourier cosine integral is similar to the Fourier sine Fourier series so given by 1 is similar to the representation of a function.

By a Fourier series except that the range is now – infinity to infinity and the summation has been replaced by integration so there also we had similar situation we had an $\cos nx + B_n \sin nx$ and so it is similar expression here there we had summation here we have integration and there we had the interval – π to π here we have the interval – infinity to infinity in the Fourier integral so it is similar to the Fourier series so this is what we have observed here

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Inversion formulae for Fourier transforms

We may write (1) as

$$f(x) = \frac{1}{\pi} \int_0^{\infty} \left[\int_{-\infty}^{\infty} f(v) \cos \omega(x-v) dv \right] d\omega, \quad (4)$$

by substituting the values of $A(\omega)$ and $B(\omega)$.
 Since $\cos \omega(x-v)$ is an even function of ω , (4) may be written as

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \left[\int_{-\infty}^{\infty} f(v) \cos \omega(x-v) dv \right] d\omega, \quad (5)$$

Since $\sin \omega(x-v)$ is an odd function of ω , we have

$$0 = \frac{i}{2\pi} \int_{-\infty}^{\infty} \left[\int_{-\infty}^{\infty} f(v) \sin \omega(x-v) dv \right] d\omega. \quad (6)$$

Let us now move to inversion formulae for Fourier transforms now this is equation 1 can be written in the following manner this equation $f(x) = \frac{1}{\pi} \int_0^{\infty} A(\omega) \cos \omega x + B(\omega) \sin \omega x$ let us put the value of $A(\omega)$ and $B(\omega)$ here what we get $\frac{1}{\omega} \int_{-\infty}^{\infty} f(v) \cos \omega(x-v) dv$ so $f(x) = \frac{1}{\pi} \int_0^{\infty} \left[\int_{-\infty}^{\infty} f(v) \cos \omega(x-v) dv \right] d\omega$ so we can write like this now $\cos \omega(x-v)$ is an even function now ω we can make it $\frac{1}{2\pi} \int_{-\infty}^{\infty} \left[\int_{-\infty}^{\infty} f(v) \cos \omega(x-v) dv \right] d\omega$ this is because $\cos \omega(x-v)$ is an even function of ω so let us now go there so this is how.

So, I can write it as $\frac{1}{\pi} \int_0^{\infty} \left[\int_{-\infty}^{\infty} f(v) \cos \omega(x-v) dv \right] d\omega$ so I can write it as $\frac{1}{\pi} \int_0^{\infty} \left[\int_{-\infty}^{\infty} f(v) \cos \omega(x-v) dv \right] d\omega$ so we can write like this now $\cos \omega(x-v)$ is an even function now ω we can make it $\frac{1}{2\pi} \int_{-\infty}^{\infty} \left[\int_{-\infty}^{\infty} f(v) \cos \omega(x-v) dv \right] d\omega$ this is because $\cos \omega(x-v)$ is an even function of ω so let us now go there so this is how.

We can use the fact that $\cos \omega x - v$ is an even function of ω so we can write $f(x)$ like this
 now $\sin \omega x - v$ is an odd function of ω therefore $\int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(v) \sin \omega(x-v) dv d\omega$ will be $= 0$. I can multiply by $i/2\pi$ it will be also $0 = i/2\pi \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(v) \sin \omega(x-v) dv d\omega$ now let us add this quantity to this equation. Let us say 5 and 6 what we will get.

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Inversion formulae for Fourier transforms cont...

Adding (5) and (6), we have

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \left[\int_{-\infty}^{\infty} f(v) e^{i\omega(x-v)} dv \right] d\omega. \quad (7)$$

This is known as complex form of the Fourier integral. We may write (7) as

$$f(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} \underbrace{e^{i\omega x}}_{\text{}} \left[\int_{-\infty}^{\infty} \underbrace{f(v) e^{-i\omega v}}_{\text{}} dv \right] d\omega = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \underbrace{F(\omega) e^{i\omega x}}_{\text{}} d\omega, \quad (8)$$

where

$$F(\omega) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(v) e^{-i\omega v} dv. \quad (9)$$

The function $F(\omega)$ given by (9) is called the Fourier transform of the function $f(x)$.
 Also, the function $f(x)$, given by (8) is known as inverse Fourier transform of $F(\omega)$.

We will get $f(x) = 1/2\pi \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(v) \cos \omega(x-v) dv d\omega$ and here $\cos \theta + i \sin \theta = e^{i\theta}$ so you will write $e^{i\omega(x-v)}$ to the power $i\omega x - i\omega v$ $d\omega$ this is known as complex form of the Fourier integral. We may write this as $1/2\pi \int_{-\infty}^{\infty} e^{i\omega x} \left[\int_{-\infty}^{\infty} f(v) e^{-i\omega v} dv \right] d\omega$. I can write $i\omega x$ here and the integral $\int_{-\infty}^{\infty} f(v) e^{-i\omega v} dv d\omega$.

This I can write as $1/\sqrt{2\pi} \int_{-\infty}^{\infty} F(\omega) e^{i\omega x} d\omega$ this I write as $1/\sqrt{2\pi} \int_{-\infty}^{\infty} F(\omega) e^{i\omega x} d\omega$ we take from this coefficient and define the integral $\int_{-\infty}^{\infty} f(v) e^{-i\omega v} dv = \sqrt{2\pi} F(\omega)$ so this is $F(\omega)$ actually so this $F(\omega)$ and then we get $1/\sqrt{2\pi} \int_{-\infty}^{\infty} F(\omega) e^{i\omega x} d\omega$ now the function $F(\omega)$ given by this equation 9 is called the Fourier transform of the function $f(x)$.

The function F given by this equation $f(x) = \dots$ this is known as inverse Fourier transform of $f(\omega)$. Some people take 1 here in $f(\omega)$ and while taking the inversion formulae they take $1/2\pi$ here. We have taken $1/\sqrt{2\pi}$ here so there is no problem because when we use this Fourier transform in the solution of partial differential equations first we apply the transform and then we take the inverse Fourier transform so this coefficient is taken into account there.

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Inversion formulae for Fourier transforms cont...

In a similar fashion, we obtain from equations (2) and (3)

$$f(x) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} F_c(\omega) \cos(\omega x) d\omega, \quad (10)$$

where

$$F_c(\omega) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} f(v) \cos(\omega v) dv, \quad (11)$$

and

$$f(x) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} F_s(\omega) \sin(\omega x) d\omega, \quad (12)$$

where

$$F_s(\omega) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} f(v) \sin(\omega v) dv, \quad (13)$$

So, we now in a similar manner let us see how we define the Fourier sine and cosine transforms in their inversion formulae let us see how we find so from equation 2 and 3 let us look at question 2 and 3 this is equation 2 and this is equation 3 you put the value of $A(\omega)$ here but you get $2/\pi \int_0^{\infty} A(\omega) \cos(\omega x) d\omega$ and let us write it as $\int_0^{\infty} f(v) \cos(\omega v) dv \cos(\omega x) d\omega$.

So, this can be written as $2/\pi$ we have 2 integrals $\int_0^{\infty} f(v) \cos(\omega v) \cos(\omega x) dv d\omega$ now let us see how we define Fourier sine and cosine transform so we can write it as $\int_0^{\infty} \sqrt{2/\pi} F_c(\omega) \cos(\omega x) d\omega$ so we can write it like this this = $\sqrt{2/\pi} \int_0^{\infty} f(v) \cos(v-x) dv$ this is because $\cos(\omega v-x)$ is an even function of ω okay.

So, let us now go there okay this is how we can use the fact that $\cos(\omega f-b)$ is an even function. So, we can write $f(x)$ like this. Now $\sin(\omega x-b)$ is an odd function of ω okay

therefore $\int_{-\infty}^{\infty} f(v) \sin \omega x - v \, dv \, d\omega$ will be $=0$. So I can multiply $i/2\pi$ it will be also 0. So, $0 = i/2\pi \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(v) \sin \omega x - v \, dv \, d\omega$.

Now. Let us add this quantity to this equation to this equation. Let us take 5 and 6 what we will get we will get $f(x) = 1/2\pi \int_{-\infty}^{\infty} \int_{-\infty}^{\infty} f(v) \cos \theta + i \sin \theta$ so e to the power $i \theta$ we will get you write e to the power $i \omega x - v \, dv \, d\omega$ this is known as complex form of the Fourier integral. Okay we may write this as $1/2\pi \int_{-\infty}^{\infty}$.

e to the power $i \omega x$ e to the power $i \omega v$ we have here. So, I can write it for $i \omega x$ here $\int_{-\infty}^{\infty} f(v) e$ to the power $-i \omega v \, dv \, d\omega$ okay. This I can write as $1/\sqrt{2\pi} \int_{-\infty}^{\infty} f(\omega) \, d\omega$. This I write as $1/\sqrt{2\pi}$ we take this from this coefficient and define the integral $-\infty$ to ∞ $f(v) e$ to the power $-i \omega v \, dv$ as $\sqrt{2\pi} * f(\omega)$.

So, this is $f(\omega)$ actually. Okay so this is $f(\omega)$ and then we get $1/\sqrt{2\pi} \int_{-\infty}^{\infty} f(\omega) e$ to the power $i \omega x \, d\omega$. Now the function $f(\omega)$ given by this equation 9 is called the Fourier transform of the function $f(x)$. The function f given by this equation. $F(x) =$ this okay this is known as inverse Fourier transforms of $F(\omega)$ okay. Some people take one here okay $f(\omega)$.

And while taking the inversion formula they take $1/2\pi$ here we have taken $1/\sqrt{2\pi}$ here $1/\sqrt{2\pi}$ here. So, there is no problem because when we use Fourier transform in the solution of partial differential equations there first, we apply the Fourier transform then we take the inverse Fourier transform. So, this coefficient is taken into account there. So, we now in a similar manner let us see how we define the Fourier sin transforms.

And their inverse formula. Let us see how we find okay. So, from equation 2 and 3 let us look at equation 2 and 3 okay this is equation 2 and this is equation 3. So, put the value of ω here what value you get $2/\pi \int_0^{\infty} f(\omega) \cos \omega x \, d\omega$ okay $2/\pi \int_0^{\infty}$ to

infinity. And let us write it as $\int_0^\infty \int_0^\infty f(b) \cos(\omega v) dv \cos(\omega x) d\omega$ okay so this can be written as $\frac{2}{\pi}$.

We have 2 integrals $\int_0^\infty \int_0^\infty f(b) \cos(\omega v) dv d\omega$. Now let us see how we define Fourier sin and cosine transforms. So, we can write it as $\frac{2}{\pi} \int_0^\infty f(\omega) \cos(\omega x) d\omega$ okay. So, we can write it like this this is $= \frac{2}{\pi} \int_0^\infty f(v) \cos(\omega v) dv$ we interchange the order of integration. We have Instead of $dv d\omega$ we have $d\omega dv$.

Okay so $\cos(\omega v) d\omega$ we have and outside we have $\cos(\omega x) d\omega$. Okay then we define it write like this. $\frac{2}{\pi} \int_0^\infty f(\omega) \cos(\omega x) d\omega$. Where $f(\omega) = \frac{2}{\pi} \int_0^\infty f(v) \cos(\omega v) dv$. Okay so this is defined as Fourier sin transform of f and this formula which gives the $f(x)$ okay this formula $f(x) = \frac{2}{\pi} \int_0^\infty f(\omega) \cos(\omega x) d\omega$ is called the inverse formula for the Fourier cosine transform. Similarly, we have the formula for Fourier sin transform.

You put the value for v ω here $2 \times \int_0^\infty f(v) \sin(\omega v) dv$ interchange the order of integration. So, after inter changing order of integration $\frac{2}{\pi} \int_0^\infty f(v) \sin(\omega v) dv \sin(\omega x) d\omega$. Okay $\sin(\omega v) d\omega$ okay $f(\omega)$ will be $= \frac{2}{\pi} \int_0^\infty f(v) \sin(\omega v) dv$ this will be Fourier sin transforms and inverse formula will be $f(x) = \frac{2}{\pi} \int_0^\infty f(\omega) \sin(\omega x) d\omega$ in a similar manner we do it okay just like in the case of Fourier cosine transform.

We put the value of v ω from here in this integral and we get the Fourier sin transform and inverse formula for Fourier sin transform okay. So, this is formula for inversion formula for Fourier cosine transform this is inversion formula for Fourier sin transform this is Fourier formula for Fourier sin transform.

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Inversion formulae for Fourier transforms cont...

The equations (10) and (12) give the inversion formulae for the Fourier cosine and sine transforms of f given by (11) and (13) respectively.

Note that some authors take the coefficients in the Fourier sine and cosine transforms as 1 and in their inversion formulae $\frac{2}{\pi}$ instead of $\sqrt{\frac{2}{\pi}}$ each taken by us.

This does not make any difference in applications to the solutions of physical problems.

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Now let us look at this the equations 10 and 12 give the inversion formulae for the Fourier cosine and sine transforms this one and this one they give the formulae for the inversion formula for the Fourier cosine and sine transforms of f which are given 11 and 13. Note that some authors take the coefficients in the Fourier Sine and cosine transform as 1 and in their inversion formulae $2/\pi$ instead of root $2/\pi$ each taken by us. This does not make any difference in applications to the solutions of physical problems.

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Example 4

If

$$\int_0^{\infty} f(x) \sin ax \, dx = \begin{cases} 1, & 0 < a < 1, \\ 0, & a > 1, \end{cases}$$

then find $f(x)$.

ANS:

$$f(x) = \frac{2}{\pi} \left(\frac{1 - \cos x}{x} \right).$$

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Now let us take this example suppose we have an integral equation because we have here an unknown function $f(x)$ it occurs inside the integral sine. So, integral equations can also be solved using the Fourier cosine sine transform so 0 to infinity $f(x) \sin ax \, dx = 1$ when $0 < a < 1$; 0 when

$a > 1$, we have to this unknown function $f(x)$ okay. So, let us compare this with the Fourier sine transform okay.

So, this is the Fourier sine transform okay 0 to infinity $f(\omega) = \frac{1}{\sqrt{2\pi}} \int_0^\infty f(x) \sin \omega x \, dx$ okay. So, we have given $f(\omega)$ okay. So, here we are given Fourier sine instead of ω we have here a okay we do not have $\sin \omega x$ we have $\sin x$ so we have $f(a)$ and $\sqrt{2\pi}$ I am taking as 1 instead of $\sqrt{2\pi}$ here for convenience I am taking it as this coefficient as 1. So, while taking this inverse Fourier sine transform, I will take here $2/\pi$.

So, $f(a)$ given to be 0 to infinity $f(x) \sin ax \, dx = 1$ $0 < a < 1$ and 0 when $a > 1$. Okay now by then inversion formula $f(x)$ will be $2/\pi \int_0^\infty f(\omega) \sin \omega x \, d\omega$ okay. Why I have to use sine inverse Fourier sine transform because we have sine function here. If here we have here cos as function ax then we will use Fourier cosine transform formula. So, this is $= 2/\pi$ and over the interval 0 it is 1 otherwise 0.

So, 0 to 1 $f(\omega) = 1 \sin \omega x \, d\omega$, so this is $2/\pi \int_0^1 \sin \omega x \, d\omega$ because we are integrating the two ω 0 to 1 okay. So, this is $2/\pi [1 - \cos \omega x]_0^1$ so this is 1. We have 0 1 interval okay $1 - \cos x/dx$ because these are the limits for ω . So, $1 - \cos x/x$ so we have $2/\pi \int_0^1 (1 - \cos \omega x) \, d\omega$.

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Application to partial differential equations

Integral transforms are used to find the solution of a partial differential equation. The choice of particular transform to be used for the solution of the differential equation depends upon the nature of the boundary conditions of the equation and the facility with which the transform $F(\omega)$ can be inverted to give $f(x)$.

Let $u(x, t)$ and $\frac{\partial u}{\partial x}$ be functions which tend to zero as $x \rightarrow \pm\infty$. Then Fourier transform of $\frac{\partial^2 u}{\partial x^2}$ is given by

$$F\left(\frac{\partial^2 u}{\partial x^2}\right) = -\omega^2 U(\omega, t), \text{ where } U(\omega, t) = F(u(x, t)).$$

$$F(u_{xx}) = \int_{-\infty}^{\infty} u_{xx} e^{-i\omega x} dx = \left[u_x e^{-i\omega x} \right]_{-\infty}^{\infty} - \int_{-\infty}^{\infty} u_x (-i\omega) e^{-i\omega x} dx$$

$$= 0 - (-i\omega) \int_{-\infty}^{\infty} u_x e^{-i\omega x} dx = i\omega \int_{-\infty}^{\infty} u_x e^{-i\omega x} dx$$

Now application to partial differential equations integral transforms can be used to find the solution of a partial differential equation. The Choice of a particular transform to be used for the solution of the differential equation depends upon the nature of the boundary conditions of the equation and the facility with which the transform $F(\omega)$ can be inverted to give $f(x)$ let $u(x)$ and u_x be functions which tend to 0 as x goes to $\pm\infty$.

Then the Fourier transform of u_{xx} is given by $F(u_{xx}) = -\omega^2 F(u)$ we can see this. Fourier transform of u_{xx} okay let us go to the formula for Fourier transform so Fourier transform is $1/\sqrt{2\pi} \int_{-\infty}^{\infty} u(x) e^{-i\omega x} dx$ I will take here 1 and then while we take the inversion formula, we will take $1/2\pi \int_{-\infty}^{\infty} f(\omega) e^{i\omega x} d\omega$ here, so $f(\omega)$ is $-\infty$ to ∞ $\int_{-\infty}^{\infty} f(\omega) e^{i\omega x} d\omega$ okay let us use this formula.

So, this will be integral $-\infty$ to ∞ $u_{xx} e^{-i\omega x} dx$ okay. This will be now let us integrate it by parts so we will have u_x when we integrate with respect to x $u_x e^{-i\omega x}$ to the power $-i\omega x$ $-\infty$ to ∞ $- \int_{-\infty}^{\infty} u (-i\omega) e^{-i\omega x} dx$ okay now u_x goes to 0 we have assumed u_x goes to 0 as x goes to $\pm\infty$. So, this will go to 0 when x goes to $\pm\infty$.

And this will also go to 0 when x goes to $-\infty$. $e^{-i\omega x}$ is a bounded quantity we get mod of $e^{-i\omega x} = 1$. Mod of $e^{-i\omega x}$ is under $\cos^2 \omega x + \sin^2 \omega x$. So, that is 1 so bounded quantity $\times 0$ so we get 0 here and this becomes $i\omega \int_{-\infty}^{\infty} u e^{-i\omega x} dx$. Okay and we integrate it again so when you integrate it again.

What do you get $\int_{-\infty}^{\infty} u_x e^{-i\omega x} dx$ is now x so $u_x \times e^{-i\omega x}$ $-\int_{-\infty}^{\infty} u (-i\omega) e^{-i\omega x} dx$ okay so this will again go to 0 because we have assumed that u_x goes to 0 as x goes to $+$ or $-\infty$ and $e^{-i\omega x}$ is a bounded quantity it is bounded by 1. So, this will go to 0 and we will have then $i\omega \times i\omega \int_{-\infty}^{\infty} u e^{-i\omega x} dx = -\omega^2 \int_{-\infty}^{\infty} u e^{-i\omega x} dx$.

So, we will get $-\omega^2$ and we will get integral from $-\infty$ to ∞ of $u(x,t)$ to the power $-i$ ωx dx which is Fourier transform of $u(x,t)$ and we are denoting it by $U(\omega, t)$. Fourier transform of $u(x,t)$ we are denoting by $U(\omega, t)$. So, Fourier transform of u_{xx} is $-\omega^2 U(\omega, t)$.

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Application to partial differential equations cont...

If $U_s(\omega, t)$ and $U_c(\omega, t)$ are Fourier sine and cosine transforms of $u(x, t)$ then

$$U_s(\omega, t) = \int_0^\infty u(x, t) \sin(\omega x) dx$$

hence

$$F_s\left(\frac{\partial^2 u}{\partial x^2}\right) = -\omega u(x, t)|_{x=0} - \omega^2 U_s(\omega, t)$$

Similarly,

$$U_c(\omega, t) = \int_0^\infty u(x, t) \cos(\omega x) dx$$

Hence

$$F_c\left(\frac{\partial^2 u}{\partial x^2}\right) = -\left(\frac{\partial u}{\partial x}\right)\bigg|_{x=0} - \omega^2 U_c(\omega, t)$$

If $U_s(\omega, t)$ and $U_c(\omega, t)$ are Fourier Sine and Cosine transform of $u(x, t)$ then $U_s(\omega, t) = \int_0^\infty u(x, t) \sin \omega x$. This is how we denote the Fourier by this notation we are denoting the Fourier sine transform of $u(x, t)$ okay Fourier sine transform of u_{xx} then will be given by this formula we can easily see that. So, Fourier sine transform of u_{xx} will be then u_{xx} will be integral from 0 to infinity $u_{xx} \sin \omega x$ dx .

So, integrate by parts so $u \sin \omega x$ from 0 to infinity $-u$ we are integrating with respect to x . So, $\omega \cos \omega x$ dx from 0 to infinity. Now u goes to 0 when x goes to infinity and $\sin \omega x$ is a bounded quantity. It is bounded by 1 so it is 0. So, when x is 0 $\sin \omega x$ is 0. So, this quantity becomes 0 and what we get $-\omega \int_0^\infty u \cos \omega x$ dx . Let us integrate by parts again so we get $u \sin \omega x$ function.

Now integral of u is $u \sin \omega x$ from 0 to infinity -0 to infinity $u \sin \omega x$ dx . Okay when x goes to infinity, we have assumed u goes to 0 so this goes to 0 when x goes to 0 this becomes $u - u_0 \cos 0$ becomes 1 and here we get $+\omega$ Fourier sine transform of $u(x, t)$. Okay

so because we are getting this ω times 0 to infinity $u(x,t) \sin \omega x \, dx$ okay so what we get.

Let us see we get $-\omega$ we got outside and then we get here $-\omega^2$ okay $-\omega^2$ square $f(u(x,t))$ $f(u(x,t))$. We denote by $u(x,t)$ okay $u(x,t)$ and this will be $\omega u(x,t)$ at $x=0$. Okay so $\omega u(x,t)$ at $x=0$ $-\omega^2 f(u(x,t))$ which is $u(x,t)$. In a similar manner we can derive for the Fourier Cosine transform the formula $f(u_{xx}) = -u(x,t)$ at $x=0$ $-\omega^2 u(x,t)$.

Now when we apply these sin transform, or Cosine transform to the partial differential equation we have to see the boundary condition. If the boundary condition is given at $x=0$ we are given $u(0,t)$ then we apply Fourier Sine transform. If we are given the gradient $u_x(0,t)$ that is at $u_{xx}(0,t)$ we apply Fourier Cosine transform. So, we can decide from the boundary condition which transform we have to apply in case of a semi-infinite domain.

Okay here we are using $\int_{-\infty}^{\infty}$ $-\int_0^{\infty}$ to infinity if we use $\int_{-\infty}^{\infty}$. If we are given the problem over $-\infty < x < \infty$, then we will need to use Fourier transform. So, let us apply the Fourier transforms to solve partial differential equation in one dimensional problem the pd is transformed in to an of by applying a suitable transform if In a problem $u(x,t)$ at $x=0$ is given then we use Fourier sine transform to remove u_{xx} .

In case, u_x at $x=0$ is given then the Fourier Cosine transform is applied to remove u_{xx} .

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Example 5

Determine the distribution of temperature in the semi-infinite medium $x \geq 0$ when the end $x = 0$ is maintained at zero temperature and the initial distribution of temperature is $f(x) = e^{-x}$.

Then

$$u(x, t) = \frac{2}{\pi} \int_0^{\infty} \frac{\omega}{1 + \omega^2} e^{-c^2 \omega^2 t} \sin \omega x d\omega.$$

One dimensional heat equation is given by $\frac{\partial u}{\partial t} = c^2 \frac{\partial^2 u}{\partial x^2}$

$\frac{d}{dt} U_{\beta}(w, t) = -c^2 w^2 U_{\beta}(w, t)$
 $U_{\beta}(w, t) = A(w) e^{-c^2 w^2 t}$
 $\frac{w}{1+w^2} = A(w)$

$\int_0^{\infty} \sin \omega x \frac{\partial u}{\partial t} dx = c^2 \int_0^{\infty} \frac{\partial^2 u}{\partial x^2} \sin \omega x dx$
 $\frac{d}{dt} \left(\int_0^{\infty} u(x, t) \sin \omega x dx \right) = c^2 \left(-w^2 U_{\beta}(w, t) \right)$
 $U_{\beta}(w, t) = \int_0^{\infty} u(x, t) \sin \omega x dx$

$u(0, t) = 0$
 $u(x, 0) = f(x) = e^{-x}$
 $U_{\beta}(w, 0) = \int_0^{\infty} e^{-x} \sin \omega x dx = \frac{w}{1+w^2}$

Let us cosine determine the distribution of temperature in the semi-infinite medium $x \geq 0$ when the end $x=0$ is maintained at 0 temperature so we are given $u(0, t)$ is given as 0 and initial distribution of temperature is e^{-x} . So, we have $u(x, 0) = f(x) = e^{-x}$ okay now we know that the heat equation in 1 dimensional heat equation is given by 1 dimensional heat equation is given by $u_t = c^2 u_{xx}$.

Now we are given $u(x, t)$ at $x=0$ we will use Fourier sine transform. So, we will apply Fourier Sine transform here okay we multiply this equation by Fourier Sine transform we have to apply. Okay so we multiply this equation by $\sin \omega x$ and integrate with respect to x okay. So, let us multiply by $\sin \omega x$ and integrate with respect to x over the interval 0 to infinity okay. so, then x and t are independent of each other.

I can write it as $\frac{d}{dt}$ of okay derivative of with respect to t of integral to infinity $u(x, t) \sin \omega x dx$ and this is c^2 times $4a$ sine transform of u_{xx} . So we have seen Fourier sine transform of u_{xx} is given by ω^2 times $U_{\beta}(w, t)$ at $x=0$ $-\omega^2 U_{\beta}(w, t)$ when $x=0$ so it reduces to $-\omega^2 U_{\beta}(w, t)$. So, we get $-\omega^2 U_{\beta}(w, t)$ okay. So, left hand side is this is $U_{\beta}(w, t)$ okay.

So, we get $\frac{d}{dt} U_{\beta}(w, t) = -c^2 \omega^2 U_{\beta}(w, t)$. Okay so this is $\frac{dy}{dt} = -c^2 \omega^2 y$ so we can integrate it easily now okay so we get $U_{\beta}(w, t) = e^{-c^2 \omega^2 t}$ some constant A times $e^{-c^2 \omega^2 t}$.

to the power $-c^2 \omega^2 t$ we can integrate with respect to t keeping ω as constant. So, this will depend on ω and will be a function of ω . Now let us take this boundary condition.

Okay so we have given that $\int_0^\infty u(x,t) \sin \omega x \, dx$ this is $= u(x,0)$ this is what we have assumed. So, put $t=0$ in this so $u(x,0)$ will be $= \int_0^\infty e^{-c^2 \omega^2 t} \sin \omega x \, dx$ and its value is $\omega/(1+\omega^2)$ okay. So, put $t=0$ here so $u(x,0)$ becomes $\omega/(1+\omega^2)$ and this is $A(\omega)$, and this is 1. Okay we got the value of $A(\omega)$ okay thus we get.

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$$U_p(\omega, t) = \frac{\omega}{1+\omega^2} e^{-c^2 \omega^2 t}$$

Taking inverse Fourier sine transform

$$u(x,t) = \frac{2}{\pi} \int_0^\infty \frac{\omega}{1+\omega^2} e^{-c^2 \omega^2 t} \sin \omega x \, d\omega$$

$u(x,t) = \omega/(1+\omega^2) e^{-c^2 \omega^2 t}$. Now we take the inversion formula use the inversion formula, so we take the inverse Fourier sine transform of this. Taking inverse Fourier sine transform we get $u(x,t)$ we multiply by $\sin \omega x$ and integrate with respect to ω okay. Now we have taken while taking Fourier Sine transform, we have taken the Fourier coefficient as 1.

So, here we will take the coefficient $2/\pi$. So, $\int_0^\infty \omega/(1+\omega^2) e^{-c^2 \omega^2 t} \sin \omega x \, d\omega$. Okay so this is the temperature distribution in the semi-infinite medium for $x > 0$ $t > 0$. Okay so this is what we have $2/\pi \int_0^\infty$

$\frac{\omega}{1 + \omega^2 c^2 t^2} e^{-\frac{\omega^2 c^2 t^2}{2}} \sin \omega x$. With this
I have come to the end of this lecture thank you very much for your attention.