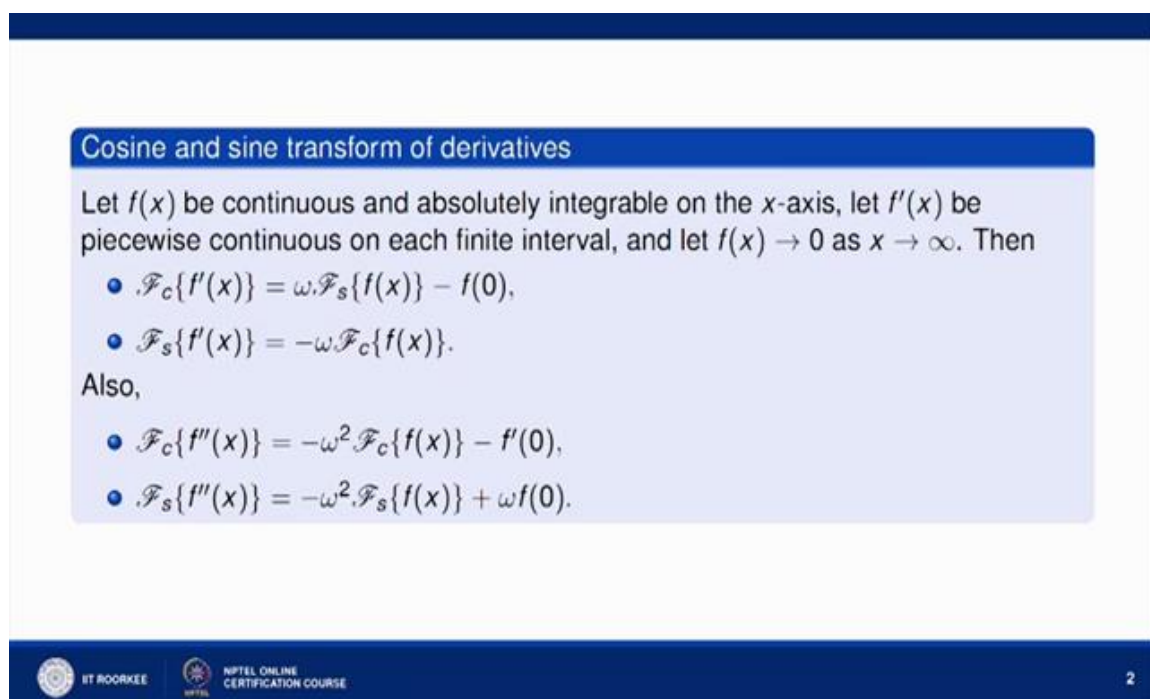


Mathematical methods and its applications
Dr. S. K. Gupta
Department of Mathematics
Indian Institute of Technology, Roorkee

Lecture - 60
Applications of Fourier transforms to BVP – III

Welcome to the lecture series on Mathematical Methods and its Applications. Let us discuss some more application of Fourier transforms. So, we have already discussed some application in last lectures; Fourier series, Fourier integral, Fourier transforms. Now this we have already discussed that how we can find out Fourier derivative of sin or cosine transform of derivatives.

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Cosine and sine transform of derivatives

Let $f(x)$ be continuous and absolutely integrable on the x -axis, let $f'(x)$ be piecewise continuous on each finite interval, and let $f(x) \rightarrow 0$ as $x \rightarrow \infty$. Then

- $\mathcal{F}_c\{f'(x)\} = \omega \mathcal{F}_s\{f(x)\} - f(0)$,
- $\mathcal{F}_s\{f'(x)\} = -\omega \mathcal{F}_c\{f(x)\}$.

Also,

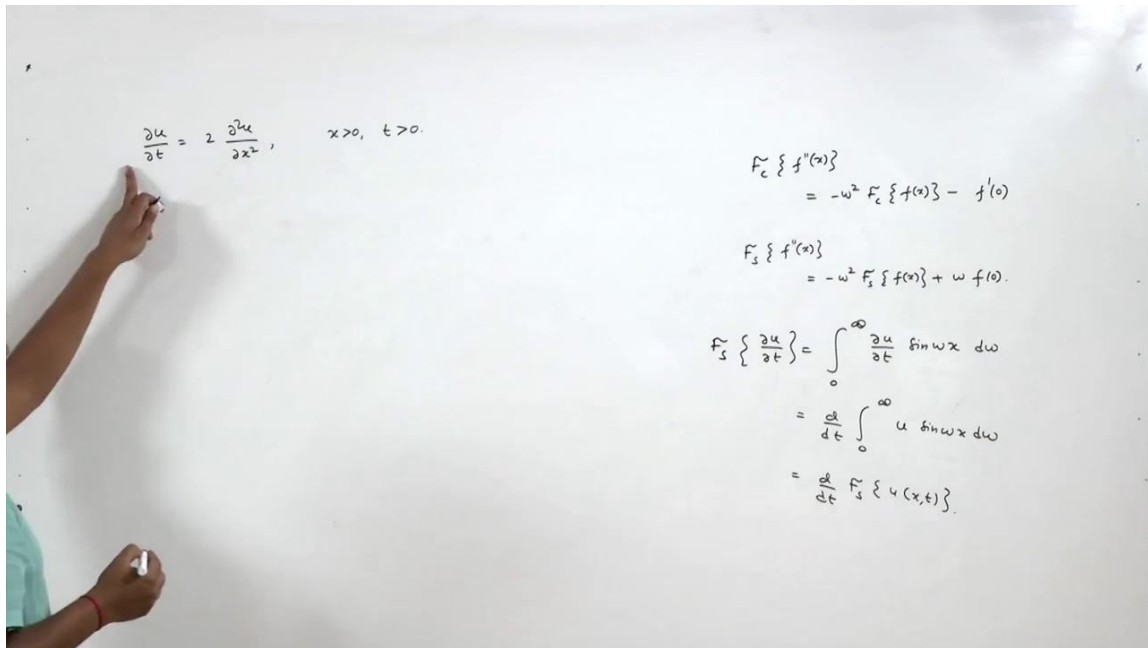
- $\mathcal{F}_c\{f''(x)\} = -\omega^2 \mathcal{F}_c\{f(x)\} - f'(0)$,
- $\mathcal{F}_s\{f''(x)\} = -\omega^2 \mathcal{F}_s\{f(x)\} + \omega f(0)$.

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This we have already discussed, that if we want to find out Fourier cosine of f dash x . So, that will be nothing but ω Fourier sin of f x minus $f(0)$. And similarly if you want to find out Fourier sin of f dash x it will nothing but minus ω Fourier cosine of f x .

So, when you replace f by f dash in the first expression and f by f dash in the second expression you will get back to these two expressions. So, what are these two expressions? So, let us this.

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This we have already derived in the last lectures. So, what is Fourier cosine of $f''(x)$? Fourier cosine of $f''(x)$ will be nothing but x dash we have already derived it is minus omega square Fourier cosine of $f(x)$ minus $f'(0)$. And Fourier sine of $f''(x)$ is nothing but omega square Fourier sine of $f(x)$; it is minus omega square Fourier sine of this plus omega $f(0)$.

So, these we have already discussed in the last lectures, now what how they are important, and why I am discussing this these things over here let us see. Now in this problem: for the equation u_t equal to $\frac{\partial u}{\partial t}$ is equal to $2 \frac{\partial^2 u}{\partial x^2}$, this equation we have to solve, now x is greater than 0 and t is greater than 0. Now x is varying from 0 to infinity; that means, we have either to apply a Fourier cosine or Fourier sine transform, because if this x is not varying from minus infinity to plus infinity; if x is varying if any one of variable is varying from minus infinity to plus infinity then we have to apply Fourier transform.

But this x is varying only in half range from 0 to infinity; that means we have to apply either Fourier cosine or Fourier sine.

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Problem

Solve the equation $\frac{\partial u}{\partial t} = 2 \frac{\partial^2 u}{\partial x^2}$, $(x > 0, t > 0)$, subject to the conditions

- (i) $u = 0$, when $x = 0, t > 0$
- (ii) $u = e^{-x}$, when $t = 0$
- (iii) $u(x, t)$ is bounded.

Solution. Given

$$\frac{\partial u}{\partial t} = 2 \frac{\partial^2 u}{\partial x^2}, \quad x > 0, t > 0. \quad (1)$$

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Now, which we have to apply Fourier sin or cosine how we will decide? We will decide by seeing the conditions. If we apply a Fourier cosine in f double dash; here f double dash if this term we are applying Fourier sin or cosine transform respect to x. So, we have a double derivative. When we have a double derivative we apply suppose Fourier cosine. So, in this we need f dash 0. So, here we need del u by del x at t equal to 0; at x equal to 0 sorry. We need the derivative condition which is not here, it is only u equal to 0 when x equal to 0 the first condition. So, if we now observe Fourier sin transform they are we need only f 0, when x is 0. So, in this way we conclude that we have to apply Fourier sin transform in this expression.

Basically, whenever we have a half range; we have to apply either Fourier sin or Fourier cosine. When in the initial condition we have the partial derivative at x equal to 0, so then we have to apply Fourier cosine otherwise we have to apply Fourier sin transform. So, in this particular problem we have to apply Fourier sin transform. So, take Fourier sin transform both the sides, because we have the condition when x equal to 0 u equal to 0. If we have derivative condition at x equal to 0 then we apply Fourier cosine transform.

Now take Fourier sin transform both the side, what will be the Fourier sin transform of del u by del t? It will be nothing but 0 to infinity del u by del t of sin omega x into d

omega. So, that will be nothing but d by d t, because this term is free from t so 0 to infinity u into sin omega x into d omega which is nothing but d by d t of Fourier sin transform of u x t.

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$$\frac{\partial u}{\partial t} = 2 \frac{\partial^2 u}{\partial x^2}, \quad x > 0, \quad t > 0.$$

$$\frac{d}{dt} f_s \{ u(x, t) \} = 2 \left[-\omega^2 f_s \{ u(x, t) \} + \omega x^0 \right].$$

$$\text{Let } f_s \{ u(x, t) \} = f_s(\omega, t)$$

$$\frac{d f_s}{dt} + 2\omega^2 f_s = 0.$$

$$\frac{d f_s}{f_s} = -2\omega^2 dt \Rightarrow f_s = k e^{-2\omega^2 t}, \quad k = k(\omega).$$

$$f_s(\omega, 0) = k(\omega) = \frac{\omega}{1 + \omega^2}.$$

$$f_s(\omega, t) = \frac{\omega}{1 + \omega^2} e^{-2\omega^2 t}$$

$$u(x, 0) = e^{-x}$$

$$f_s \{ u(x, 0) \} = f_s \{ e^{-x} \}$$

$$= \int_0^{\infty} e^{-x} \sin \omega x \, dx$$

$$= \frac{1}{1 + \omega^2} \left[-e^{-x} \sin \omega x - e^{-x} \omega \cos \omega x \right]_0^{\infty}$$

$$= \frac{1}{1 + \omega^2} \left[\omega \right] = \frac{\omega}{1 + \omega^2}$$

$$= f_s(\omega, 0)$$

So, now when you take Fourier sin transform both the side it will be nothing but d by d t of Fourier sin transform of u x t which is equal to 2 times minus omega square; apply this result, so Fourier transform of sin transform of u x t plus omega into 0 because u is 0 when x is 0.

Now, let Fourier sin transform of u x t as suppose you take it 4 f of s it is omega t, because we are applying Fourier transform with respect to x. So, that is it is d upon f x upon d t which is minus minus plus 2 omega square f s which is equals to 0. So that implies; d f s upon f s is equals to minus 2 omega square d t, when w e integrate both sides f x is nothing but you obtain as k times e k power minus 2 omega square d.

Again, this k will be nothing but function of omega. So, now we have and the next condition where when t equal to 0 u equal to e k power minus x where this condition. Now apply this condition over here. What we obtain? So, u x 0 is e k power minus x. So, apply Fourier sin transform both the sides. So, Fourier sin transform of this expression will be nothing but Fourier sin transform of e k power minus x, which is nothing but 0 to infinity e k power minus x sin omega x into d omega into d x sorry it will be d x. So, it

will be equal to $\frac{1}{1 + \omega^2}$ and it is e^{-kx} to integrate by parts and $e^{-kx} \omega \cos \omega x$ from 0 to infinity, so this is nothing but $\frac{1}{1 + \omega^2}$.

So, when x tends to infinity both will tend to 0 when x tends to 0 it will be nothing but $\frac{1}{1 + \omega^2}$. So, that will be nothing but $\frac{1}{1 + \omega^2}$. Now when you take t equal to 0 here, so Fourier sin transform of $u(x, 0)$ will be nothing but $f(\omega)$; this is nothing but basically $f(x)$. When you take $f(x)$ here; $f(x)$, when you substitute t equal to 0 here basically it will be $k\omega$. And $k\omega$ is nothing but from here will be nothing but $\frac{1}{1 + \omega^2}$.

Therefore, what will be $f(\omega, t)$? It will be $\frac{1}{1 + \omega^2} e^{-k\omega t}$. Now take inverse Fourier sin transform both the side to find out the solution. So, this implies $u(x, t)$ will be nothing but $\frac{2}{\pi} \int_0^\infty \frac{1}{1 + \omega^2} e^{-k\omega t} \sin \omega x \, d\omega$. So, this will be the final solution of this problem.

So, here is a here is a solution, the same solution I have discussed here. So, this is a solution $\frac{2}{\pi} \int_0^\infty \frac{1}{1 + \omega^2} e^{-k\omega t} \sin \omega x \, d\omega$.

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Problem

The steady state temperature distribution $u(x, y)$ in a thin, homogeneous semi-infinite plate is governed by the boundary value problem

$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0, \quad 0 < x < l, \quad 0 < y < \infty$$
$$u(0, y) = e^{-2y}, \quad u(l, y) = 0, \quad y > 0; \quad \left(\frac{\partial u}{\partial y}\right)(x, 0) = 0, \quad 0 < x < l.$$

Find the temperature distribution $u(x, y)$, $0 < x < l$, $y > 0$.

Solution. Since the domain of x is finite and the domain of y is $0 < y < \infty$, we can use Fourier cosine or sine transform (w.r.t the variable y).

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Now, the next problem: it is again in half range. Now here the half range is in y , y is varying from 0 to infinity. That means, we have to apply Fourier sin or Fourier cosine transform with respect to y . First we have to decide that for which variable we have to apply Fourier sin or cosine transform and that depends on the range. Here x is in 0 to 1 and y is varying from 0 to infinity. So, we have to apply Fourier sin or cosine transform with respect to y , number 1. Number 2: we have to apply Fourier sin or Fourier cosine that will decide by the conditions.

Now, here we will apply Fourier sin or cosine transform with respect to y and the partial derivative with respect to y when y is 0 is 0, it is given to us. That means we will apply a Fourier cosine transform here. Because, when you see this expression in Fourier cosine transform the second derivative here we have the derivative at t equal to at x equal to 0 that means we have to apply Fourier cosine transform when we have the derivative term in the initial conditions; in the conditions.

So, here $\frac{\partial u}{\partial y}$ at y equal to 0 is given; that means, we have to apply Fourier cosine transform. Now, in this problem apply Fourier cosine transform both the sides. So, what is the problem basically?

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$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0,$$
 Apply Fourier cosine transform w.r.t 'y':

$$\frac{d^2}{dx^2} F_c(x, \omega) + (-\omega^2 F_c(x, \omega) - \left(\frac{\partial u}{\partial y}\right)_{(x,0)}) = 0$$

$$\frac{d^2 F_c}{dx^2} - \omega^2 F_c = 0$$

$$(D^2 - \omega^2) F_c = 0, \quad D \equiv \frac{d}{dx}$$

$$F_c = c_1 \sinh \omega x + c_2 \cosh \omega x, \quad c_1 = c_1(\omega)$$

$$F_c(0, \omega) = c_2 = \frac{2}{1 + \omega^2}, \quad c_2 = c_2(\omega)$$

$$F_c \{ u(x, y) \} = F_c(x, \omega)$$

$$u(0, y) = e^{-2y}$$

$$F_c \{ u(0, y) \} = F_c \{ e^{-2y} \}$$

$$= \int_0^{\infty} e^{-2y} \cos \omega y \, dy$$

$$= \frac{1}{1 + \omega^2} \left[-2 e^{-2y} \cos \omega y + e^{-2y} \omega \sin \omega y \right]_0^{\infty}$$

$$= \frac{2}{1 + \omega^2} = F_c(0, \omega)$$

Now, the problems here is del square u upon del x square plus del square u upon del y square equal to 0. So, you will apply Fourier cosine transform that is clear and with respect to variable y that you have to apply. Apply Fourier cosine transform with respect to y.

So, what we will obtain? So, x will be 3, so it is d square f upon d x square. Now let us suppose Fourier cosine transform of u x y will be nothing but suppose it is f c x omega, because we are apply with respect to y so x we will treat as it is. So, it will be f c x omega plus. Now when we apply Fourier cosine transform here we will apply that result that is nothing but minus omega square Fourier sin cosine transform of x omega plus and minus it is del u by del y when y equal to at x comma 0 equal to 0. This condition we have applied for Fourier cosine transform.

So, what we obtain finally? Sorry, so basically, now this term is 0. So, what we obtained? D square f c upon d x square minus omega square f c will be 0 because this is 0, so this gives d square minus omega square f c is equal to 0 where d is nothing but d by d x. So, f c will be nothing but c 1 e k power or you can break in terms of hyperbolic functions. So, it is sin hyperbolic omega x plus c 2 sin hyperbolic omega y. And again c 1 and 2 are the arbitrary functions and f c is a function of x and omega, so c 1 and c 2 are the nothing but functions of omega.

So, it is $c_1 \sin$ hyperbolic ωy , sorry; it is \cos hyperbolic ωy . Now use the conditions; conditions are $u(0, y) = e^{-k y}$ is given to you, you apply a Fourier cosine transform both the sides. So, it is Fourier cosine transform of $e^{-k y}$ which is given as 0 to infinity $e^{-k y}$ and it is $\cos \omega y$ into $d y$. So, when you integrate it what we will obtain? It is $\frac{1}{4 + \omega^2}$ and again you integrate. So, it is $\frac{1}{4 + \omega^2} \cos \omega y$ from 0 to infinity.

At infinity both will tend to 0 and at 0 this is 0 and this will be 1 and minus minus plus, so it is $\frac{2}{4 + \omega^2}$. So, this we will obtain here as a Fourier cosine transform of this term. And again when you substitute x equal to 0 both the sides this will be nothing but is equal to $f_c(0, \omega)$. So, when you take $f_c(0, \omega)$ here, so \sin hyperbolic 0 is 0 , and \cos hyperbolic 0 is 1 that is nothing but c_2 ; it is it is also actually x , we are putting x equal to 0 . When you put x equal to 0 it is 0 and it is 1 , so it is c_2 . And c_2 will be nothing but $\frac{2}{4 + \omega^2}$.

So, in this way c_2 can be obtained. Now for c_1 we apply another condition.

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$$\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} = 0,$$
 Apply Fourier cosine transform w.r.t 'y'.

$$\frac{d^2}{dx^2} f_c(x, \omega) + (-\omega^2 f_c(x, \omega) - \left(\frac{\partial u}{\partial y}\right)_{(x, 0)}) = 0$$

$$\frac{d^2 f_c}{dx^2} - \omega^2 f_c = 0$$

$$(D^2 - \omega^2) f_c = 0, \quad D \equiv \frac{d}{dx}$$

$$f_c = c_1 \sinh \omega x + c_2 \cosh \omega x, \quad c_1 = c_1(\omega), \quad c_2 = c_2(\omega)$$

$$f_c(0, \omega) = c_2 = \frac{2}{4 + \omega^2}$$

$$f_c(l, \omega) = c_1 \sinh \omega l + c_2 \cosh \omega l = 0$$

$$c_1 = -c_2 \frac{\cosh \omega l}{\sinh \omega l}$$

$$= -\frac{2}{4 + \omega^2} \frac{\cosh \omega l}{\sinh \omega l}$$

$$f_c(x, \omega) = c_1 \sinh \omega x + c_2 \cosh \omega x$$

$$u(x, y) = \frac{2}{\pi} \int_0^{\infty} f_c(x, \omega) \cos \omega y d\omega$$

$$= \frac{2}{\pi} \int_0^{\infty} \left[\left(-\frac{2}{4 + \omega^2} \frac{\cosh \omega l}{\sinh \omega l} \right) \sinh \omega x + \left(\frac{2}{4 + \omega^2} \right) \cosh \omega x \right] \cos \omega y d\omega$$

Another condition is u at l is 0 . Now again take Fourier cosine transform both the sides. So, Fourier cosine transform of u at l will be nothing but Fourier cosine transform of 0

which is nothing but 0. And this is also equal to; where you substitute x equal to 1 both the sides this is nothing but $f c_1 \omega$. So, when you find $f c_1 \omega$ from here, $f c_1 \omega$ so that will be nothing but $c_1 \sin \text{hyperbolic } \omega l$ plus $c_2 \sin c_2 \cos \text{hyperbolic } \omega l$ which is equal to 0 from here.

Now, we know $c_2; c_1$ can be find out. So, what will be c_1 ? C_1 will be nothing but minus $c_2 \cos \text{hyperbolic } \omega l$ upon $\sin \text{hyperbolic } \omega l$. And c_2 is nothing but this term, so it is minus 2 upon 4 plus ω square into $\cos \text{hyperbolic } \omega l$ upon $\sin \text{hyperbolic } \omega l$. So, in this way once we obtain c_1 and c_2 which is here, then the only thing is you take inverse Fourier cosine transform both the sides to find out the final solution; that is the only thing remain now.

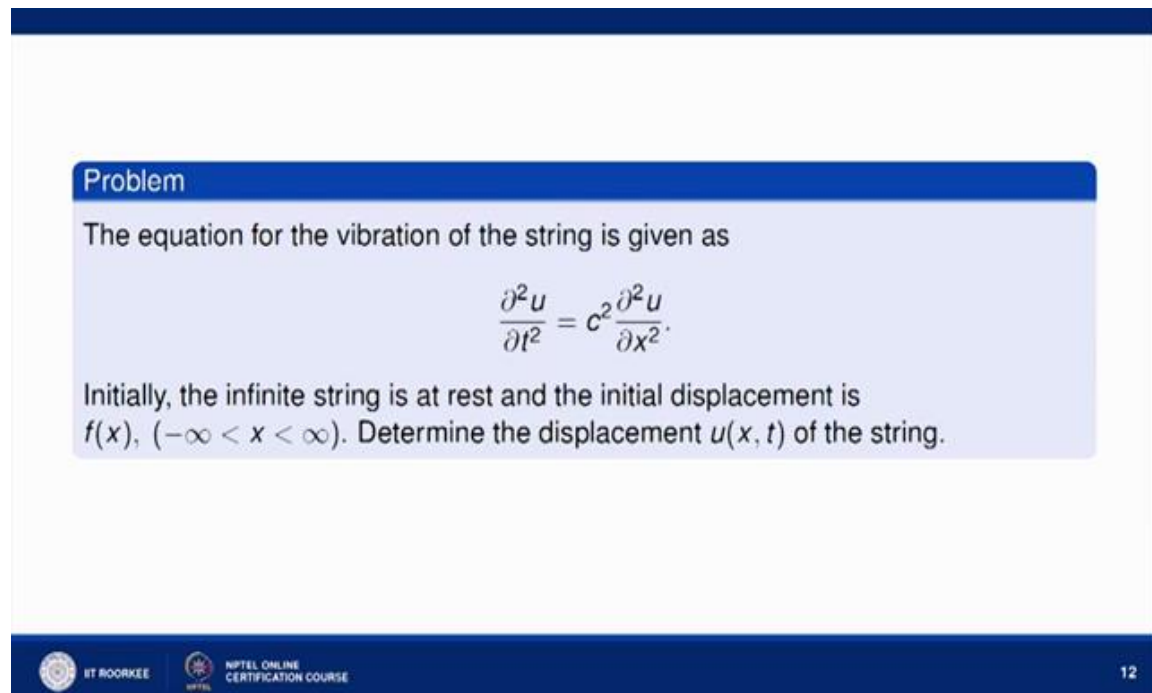
So, what is $f c \omega t$? $F c \omega t$ here is nothing but $f c x \omega$ sorry, because $f c$ is a function of x in ω . So, x and ω is nothing but $c_1 \sin \text{hyperbolic } \omega x$ plus $c_2 \cos \text{hyperbolic } \omega x$, here c_1 and c_2 are given by these two expressions so the case are inverse cosine transform both the sides. So, $u x y$ will be nothing but 2 upon π integral 0 to infinity and that will be $f c x \omega$ into $e^{-k} \text{power } \eta \omega y$ into $d \omega$. So, that will be 2 upon π integral 0 to infinity and $f c$ is nothing but c_1 times, c_1 is this term. So, this is minus 2 upon 4 plus ω square $\cos \text{hyperbolic } \omega l$ upon $\sin \text{hyperbolic } \omega l$ c_1 at this term into $\sin \text{hyperbolic } \omega x$ and plus $c_2; c_2$ is 2 upon 4 plus ω square into $\cos \text{hyperbolic } \omega x$ and whole multiplied by $e^{-k} \text{power } \eta \omega y$ in to $d y d \omega$.

So, we can simplify this expression and that will give the final solution $u x y$ which is a solution of this partial differential equation. So, hence we can solve such type of problems. So, whenever the half range is given to us you first make sure that respect to which variable you will have to apply Fourier transform x or y and that will be decided by the conditions given to us. Now you apply either Fourier sin or Fourier cosine transforms and for finding the arbitrary constants c_1 and c_2 you have the conditions given the problem. And final take initial sin or cosine transforms to find out the final solution of the given partial differential equation.

So, this is what I have done here also. This is when you simplify this; when you simplify this expression so we will get this form of answer, this one can easily solve. So, this is the final answer which we can obtain after solving this. Ok sorry, one thing I have

missed here, it is not e k power eta y it is basically, because we are taking inverse cosine transform it will be nothing but cos omega y and it is d y, d omega. Similarly it will be cos omega y into d y into d omega, because we are taking inverse Fourier cosine transform; it will be cosine term will be here. So, that we have missed, this way we will get the final answer.

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Problem

The equation for the vibration of the string is given as

$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}.$$

Initially, the infinite string is at rest and the initial displacement is $f(x)$, $(-\infty < x < \infty)$. Determine the displacement $u(x, t)$ of the string.

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Now, the last problem the equation for the vibration of the string is given as this and initially the string is at rest and the initial displacement is $f(x)$. Now here x is ranging from minus infinity to plus infinity; that means we have to apply Fourier transform. So, initially string is at rest; these are the conditions given to us so simply apply Fourier transform both the sides to solve this problem.

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$$\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}, \quad \left(\frac{\partial u}{\partial t} \right)_{(x,0)} = 0$$

$$u(x,0) = f(x)$$

$$\frac{d^2 F}{dt^2} = c^2 (-i\omega)^2 F$$

$$\frac{d^2 F}{dt^2} + c^2 \omega^2 F = 0$$

$$(D^2 + c^2 \omega^2) F = 0, \quad D \equiv \frac{d}{dt}$$

$$f = c_1 \cos(c\omega t) + c_2 \sin(c\omega t), \quad c_1 = c_1(\omega)$$

$$\frac{df}{dt} = 0 \text{ at } t=0 \Rightarrow c_2 c\omega = 0, \quad c_2 = c_2(\omega)$$

$$\Rightarrow c_2 = 0$$

$$f = c_1 \cos(c\omega t)$$

$$F = \mathcal{F}\{u(x,t)\} = F(\omega, t)$$

$$\left(\frac{\partial u}{\partial t} \right)_{(x,0)} = 0$$

$$\left(F \left\{ \frac{\partial u}{\partial t} \right\} \right)_{(x,0)} = 0$$

$$\left(\frac{d}{dt} F \right)_{(x,0)} = 0$$

So, what is the problem del square u upon del t square is equals to c square del square u upon del x square. And initially the infinity string is at rest; that means, del by del t of u at x comma 0 is 0. And initial displacement is f x; that means x comma 0 is f x. So, these are the conditions given to you. And x is varying from minus infinity to plus infinity; you have to solve this equation.

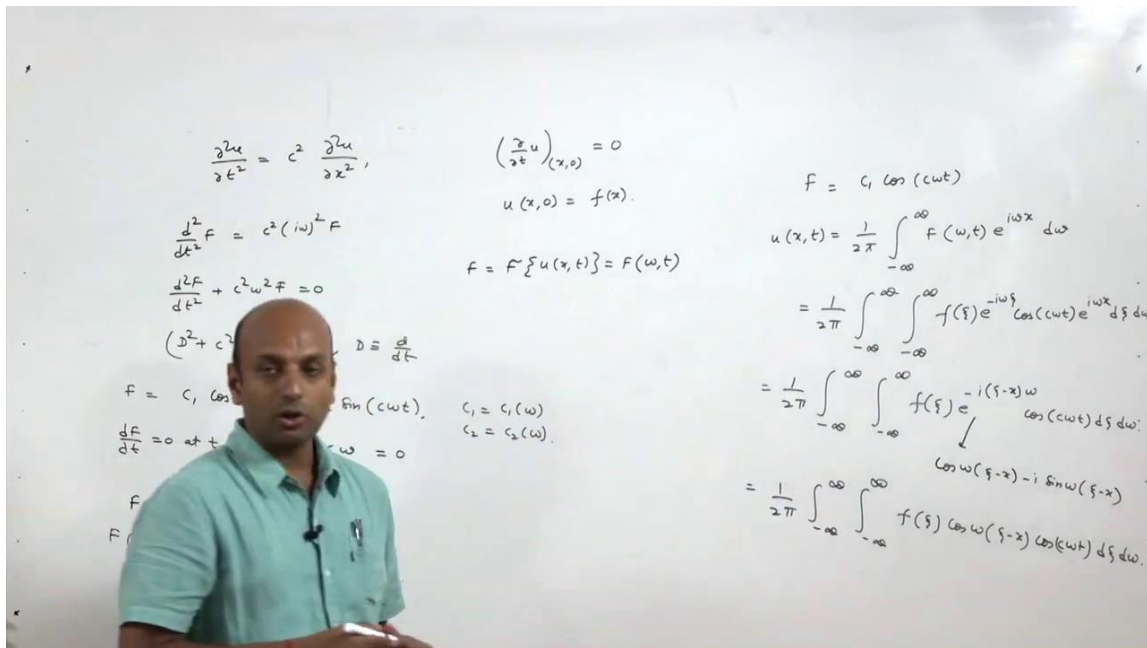
So, again we will apply Fourier transform technique. You will take Fourier transform both the sides respect to x, because x is varying from minus infinity to plus infinity. That means, it is d square upon d t square of f which is equals to c square eta omega square f; here f is nothing but Fourier transform of u x t which is a function of basically omega and t. When you simply d square f upon d t square and it is minus 1, so it is plus c square omega square f equal to 0, so it is d square plus c square omega square into f equal to 0. So, where d is nothing but d by d t. When you simply this; so f we can obtain as c 1 e k power, it is cos; so it is cos c omega t plus c 2 sin c omega t. Again c 1 and c 2 are the functions of omega.

Now, we will apply these conditions. So, we will take the first condition is del u upon d t at x comma 0 is 0. When you take Fourier transform both the sides of this expression you will obtain Fourier transform of del u by del t at x comma 0 is 0. That means, d by d t of capital F at x comma 0 is 0. A capital F is nothing but Fourier transform of u x t; it is

Fourier transform of $u(x, t)$. That means, when you take $\frac{d}{dx}$ by $\frac{d}{dt}$ that is 0 at t equal to 0.

So, when you apply this condition over here. So the derivative of this will be \sin , minus \sin ; when t equal to 0 it will be 0 and it will be \cos when t equal to 0 that will be 1. So, we will left with $c_2 \cos c\omega t$ which is equal to 0; $c_2 \cos c\omega t$ which is equal to 0. Of course, this implies c_2 equal to 0, because c_1 c and ω are constant; so c_2 cannot be 0 and ω cannot be 0, c_2 is the only thing which can be 0. So, f is nothing but $c_1 \cos c\omega t$.

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Now we have one more condition to find out the value of c_1 ; that is $u(x, 0) = f(x)$. Now $u(x, 0) = f(x)$, so take Fourier transform both the sides. So, Fourier transform of $u(x, 0)$ will be nothing but Fourier transform of $f(x)$ and that will be nothing but minus infinity to plus infinity $f(x)$ into e^{ikx} power minus η ω x into $d\omega$ into dx sorry; so that will be nothing but $f(\omega)$, because you are putting t equal to 0 so that will be $f(\omega)$. So, when you take $f(\omega)$ that will be c_1 and that c_1 will be nothing but this expression. So, that will be nothing but minus infinity to plus infinity $f(x)$ e^{ikx} power minus η ω x into dx so.

Now we have the value of c_1 and c_2 in our hand, so we will simply take inverse Fourier transform to find out the final solution. So, what is f ? F is nothing but $c_1 \cos c \omega t$ where c is given by this expression take inverse Fourier transforms both the sides, so it is $u(x, t)$ will be nothing but $\frac{1}{2\pi} \int_{-\infty}^{+\infty} f(\omega) t e^{k \text{ power } \eta \omega x} dx$; that will give the final answer.

It is nothing but $\frac{1}{2\pi} \int_{-\infty}^{+\infty} f(\omega) t e^{k \text{ power } \eta \omega x} dx$. And $f(\omega) t$ is $c_1 \cos c \omega t$ and c_1 is given by this expression. So, we can substitute c_1 also; it is $f(\psi) e^{k \text{ power } \eta \omega \psi} \cos c \omega t e^{k \text{ power } \eta \omega x} d\psi d\omega$. So, that will be here. Now when you further simplify it, so it is $\frac{1}{2\pi} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} f(\psi) e^{k \text{ power } \eta \omega \psi} \cos c \omega t e^{k \text{ power } \eta \omega x} d\psi d\omega$.

Again when you write $e^{k \text{ power } \eta \omega x}$ as $\cos \theta + i \sin \theta$ then \cos ; $\theta \cos \theta$ will contain ω and it is also even in ω even into even remain even. However, when you have a \sin term now, this can be written as $\cos \omega \psi - i \sin \omega \psi$. And this into this is an even function in ω and this into this is an odd function in ω . So, this will be 0 in terms of ω , so we will left with only one term that is $\frac{1}{\pi} \int_{-\infty}^{+\infty} \int_{-\infty}^{+\infty} \cos \omega \psi \cos c \omega t e^{k \text{ power } \eta \omega x} d\psi d\omega$. So, that will be the final solution of the given differential equation.

In this way we can solve several problems, several partial differential equation using Fourier series, Fourier integrals, or Fourier transforms. I hope in this course you have learnt lot of things about mathematical matters, lot of techniques to solve ordinary differential equation, partial differential equation, Laplace transforms its applications, Fourier series, Fourier transform, and Fourier integrals their applications that all we have seen. So, I hope you might have enjoyed the course. I wish you all the best for your bright career ahead.

Thank you very much.