

**Sobolev Spaces and Partial Differential Equations**  
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**Trace theory - Part 2**

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$\gamma_0$  is called the trace map of order 0.

Theorem: If  $v \in H^1(\mathbb{R}_+^N) \Rightarrow \gamma_0(v) \in H^{\frac{1}{2}}(\mathbb{R}^{N-1})$ .

and  $\gamma_0: H^1(\mathbb{R}_+^N) \rightarrow H^{\frac{1}{2}}(\mathbb{R}^{N-1})$  is cont.

So, now the  $\gamma_0$  is called the trace map of order 0. So, our next theorem states that the range is not all of  $L^2(\mathbb{R}^{N-1})$ . So, the




**Theorem:** If  $v \in H^1(\mathbb{R}_+^N)$  this implies that  $\gamma_0(v) \in H^{\frac{1}{2}}(\mathbb{R}^{N-1})$  and

$$\gamma_0: H^1(\mathbb{R}_+^N) \rightarrow H^{\frac{1}{2}}(\mathbb{R}^{N-1})$$




is continuous. So, this is what we would like to show. And in fact, later we will show, I will not prove it fully but we it is known that in fact this is a subjective map. That will come to it at the end.

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Pp: Step 1. Let  $v \in \mathcal{D}(\mathbb{R}^N)$   $w(x') = v(x', 0) \in L^2(\mathbb{R}^{N-1})$ .  
 Let  $\tilde{w}(\xi')$  be the F.T. (in  $\mathbb{R}^{N-1}$ ) of  $w$ .  $\xi = (\xi', \xi_N)$   
 $\xi' \in \mathbb{R}^{N-1}$ .  
 Let  $\hat{v}(\xi)$  be the F.T. (in  $\mathbb{R}^N$ ) of  $v$ .  
 Fourier inversion  $\Rightarrow$ .  
 $v(x', 0) = \int_{\mathbb{R}^N} e^{2\pi i x' \cdot \xi'} \hat{v}(\xi) d\xi$ . ( $x_N = 0$ )  
 $= \int_{\mathbb{R}^{N-1}} e^{2\pi i x' \cdot \xi'} \left( \int_{-\infty}^{\infty} \hat{v}(\xi) d\xi_N \right) d\xi'$   
 $v(x', 0) = w(x') = \int_{\mathbb{R}^{N-1}} e^{2\pi i x' \cdot \xi'} \tilde{w}(\xi') d\xi'$ .

Let  $\hat{v}(\xi)$  be the F.T. (in  $\mathbb{R}^N$ ) of  $v$ .  
 Fourier inversion  $\Rightarrow$ .  
 $v(x', 0) = \int_{\mathbb{R}^N} e^{2\pi i x' \cdot \xi'} \hat{v}(\xi) d\xi$ . ( $x_N = 0$ )  
 $= \int_{\mathbb{R}^{N-1}} e^{2\pi i x' \cdot \xi'} \left( \int_{-\infty}^{\infty} \hat{v}(\xi) d\xi_N \right) d\xi'$   
 $v(x', 0) = w(x') = \int_{\mathbb{R}^{N-1}} e^{2\pi i x' \cdot \xi'} \tilde{w}(\xi') d\xi'$ .  
 $\tilde{w}, \int_{\mathbb{R}^{N-1}} d\xi' \in \mathcal{D}(\mathbb{R}^{N-1})$  By uniqueness of F.T. inversion  
 $\tilde{w}(\xi') = \int_{-\infty}^{\infty} \hat{v}(\xi) d\xi_N$ .

So proof: Step one, so let  $v$  belong to  $\mathcal{D}(\mathbb{R}^N)$  that is always our starting point and you say  $w(x') = v(x', 0)$ . So, now we are going to we can it is in  $L^2(\mathbb{R}^{N-1})$ , therefore we can apply the Fourier transform to this so let  $\tilde{w}$  of  $\psi$  dash be the Fourier transform in  $\mathbb{R}^{N-1}$  of  $w$ . So, again  $x_i$ , I am going to write as  $x_i', x_i'$  in  $\mathbb{R}^N_-$ .

Let  $\hat{v}(x_i)$  be the Fourier transform in  $\mathbb{R}^N$  of  $v$ . So, tilde usually I have been reserving for an extension by 0, but now there is no confusion because we are working in all of  $\mathbb{R}^N$  and therefore extension does not arise, so I am using the symbol tilde for the Fourier transform in

$\mathbb{R}^n$  minus 1. So, we want to connect these two what is the connection between these two things?

By the Fourier inversion formula

$$v(x', 0) = \int_{\mathbb{R}^N} e^{2\pi i x' \cdot \xi'} \hat{v}(\xi') d\xi' = \int_{\mathbb{R}^{N-1}} e^{2\pi i x' \cdot \xi'} \left( \int_{-\infty}^{+\infty} \hat{v}(\xi') d\xi_N \right) d\xi'$$

$$v(x', 0) = w(x') = \int_{\mathbb{R}^{N-1}} e^{2\pi i x' \cdot \xi'} \tilde{w}(\xi') d\xi' \quad ; \quad \tilde{w}(\xi') = \int_{-\infty}^{+\infty} \hat{v}(\xi') d\xi_N$$

Now, apply the Fourier inversion formula in  $\mathbb{R}^{N-1}$ . So, this is nothing but  $w$  of  $x$  dash and that is equal to integral over  $\mathbb{R}^{N-1}$   $e^{2\pi i x' \cdot \xi'} \tilde{w}(\xi') d\xi'$ . So, now  $\tilde{w}$  and integral  $\hat{v} d\xi_N$  all belong to  $\mathcal{S}'(\mathbb{R}^{N-1})$ , this is a very easy thing to check and therefore by uniqueness of Fourier inversion we have the  $\tilde{w}(\xi')$  is nothing but integral, minus infinity to infinity we had  $\hat{v} d\xi_N$ . So, that is the first step where we have related the two.

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$\mathbb{R}^n$

Step 2. To show  $w \in H^{1/2}(\mathbb{R}^{N-1})$ .

To show  $\xi' \mapsto (1+|\xi'|^2)^{1/2} |\tilde{w}(\xi')|^2$  integrable.

$$\int_{\mathbb{R}^{N-1}} (1+|\xi'|^2)^{1/2} |\tilde{w}(\xi')|^2 d\xi' = \int_{\mathbb{R}^{N-1}} (1+|\xi'|^2)^{1/2} \left| \int_{-\infty}^{\infty} \hat{v}(\xi') d\xi_N \right|^2 d\xi'$$

$$= \int_{\mathbb{R}^{N-1}} (1+|\xi'|^2)^{1/2} \left| \int_{-\infty}^{\infty} \hat{v}(\xi') (1+|\xi'|^2)^{1/2} d\xi_N \right|^2 d\xi'$$

Cauchy-Schwarz  $\leq \int_{\mathbb{R}^{N-1}} (1+|\xi'|^2)^{1/2} \left[ \int_{-\infty}^{\infty} \hat{v}(\xi')^2 (1+|\xi'|^2) d\xi_N \right] \left[ \int_{-\infty}^{\infty} (1+|\xi'|^2)^{-1} d\xi_N \right] d\xi'$

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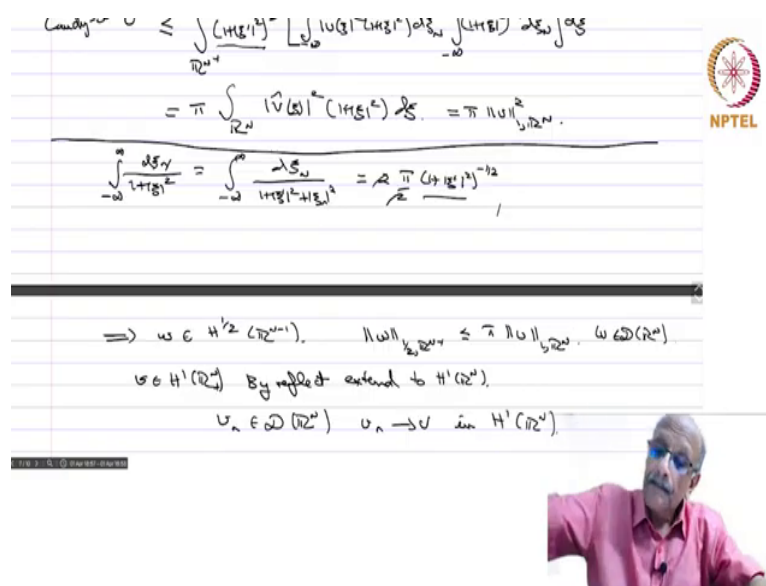
So, now we go to step two. So, what is the claim? We claim that it is in  $H^{\frac{N-1}{2}}(\mathbb{R}^{N-1})$ .

So, we want to show that to... so to show  $W$  belongs to  $H^{\frac{N-1}{2}}(\mathbb{R}^{N-1})$ , so in other words we need to show that  $\int_{\mathbb{R}^{N-1}} (1 + |\xi|^2)^{\frac{N-1}{4}} |W(\xi)|^2 d\xi < \infty$ . So, we just go ahead and try to compute that integral. So,  $\int_{\mathbb{R}^{N-1}} (1 + |\xi|^2)^{\frac{N-1}{4}} |W(\xi)|^2 d\xi = \int_{\mathbb{R}^{N-1}} (1 + |\xi|^2)^{\frac{N-1}{4}} |\int_{\mathbb{R}} \psi(x) e^{ix \cdot \xi} dx|^2 d\xi$  that is equal to  $\int_{\mathbb{R}^{N-1}} \int_{\mathbb{R}} \int_{\mathbb{R}} \psi(x) \psi(y) e^{ix \cdot \xi} e^{-iy \cdot \xi} dx dy d\xi$ .

$\int_{\mathbb{R}^{N-1}} (1 + |\xi|^2)^{\frac{N-1}{4}} |\int_{\mathbb{R}} \psi(x) e^{ix \cdot \xi} dx|^2 d\xi$  into modulus integral, from minus infinity to infinity we had  $\int_{\mathbb{R}^{N-1}} (1 + |\xi|^2)^{\frac{N-1}{4}} |\int_{\mathbb{R}} \psi(x) e^{ix \cdot \xi} dx|^2 d\xi$ ; that is equal to  $\int_{\mathbb{R}^{N-1}} \int_{\mathbb{R}} \int_{\mathbb{R}} \psi(x) \psi(y) e^{ix \cdot \xi} e^{-iy \cdot \xi} dx dy d\xi$  half into modulus integral minus infinity to infinity, we had  $\int_{\mathbb{R}^{N-1}} \int_{\mathbb{R}} \int_{\mathbb{R}} \psi(x) \psi(y) e^{ix \cdot \xi} e^{-iy \cdot \xi} dx dy d\xi$  half into  $\int_{\mathbb{R}^{N-1}} (1 + |\xi|^2)^{\frac{N-1}{4}} d\xi$ .

Now, for the term inside the modulus I will apply Cauchy-Schwarz inequality, so Cauchy-Schwarz, this is less than equal  $\int_{\mathbb{R}^{N-1}} (1 + |\xi|^2)^{\frac{N-1}{4}} |\int_{\mathbb{R}} \psi(x) e^{ix \cdot \xi} dx|^2 d\xi$  half into  $\int_{\mathbb{R}^{N-1}} (1 + |\xi|^2)^{\frac{N-1}{2}} d\xi$  times  $\int_{\mathbb{R}} \int_{\mathbb{R}} \psi(x) \psi(y) e^{ix \cdot \xi} e^{-iy \cdot \xi} dx dy d\xi$  half, sorry, into  $\int_{\mathbb{R}^{N-1}} (1 + |\xi|^2)^{\frac{N-1}{2}} d\xi$  times  $\int_{\mathbb{R}} \int_{\mathbb{R}} \psi(x) \psi(y) e^{ix \cdot \xi} e^{-iy \cdot \xi} dx dy d\xi$  half. So, the power half which should have come is taken up by the square here and therefore we have this.

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Cauchy-Schwarz 
$$\left| \int_{\mathbb{R}^{N-1}} \frac{(1+|\xi|^2)^{-\frac{N-1}{2}}}{2} \int_{\mathbb{R}} \psi(x) e^{ix \cdot \xi} dx d\xi \right| \leq \left( \int_{\mathbb{R}^{N-1}} \frac{(1+|\xi|^2)^{-\frac{N-1}{2}}}{2} d\xi \right)^{\frac{1}{2}} \left( \int_{\mathbb{R}^{N-1}} \left| \int_{\mathbb{R}} \psi(x) e^{ix \cdot \xi} dx \right|^2 d\xi \right)^{\frac{1}{2}}$$

$$= \pi \int_{\mathbb{R}^{N-1}} \frac{(1+|\xi|^2)^{-\frac{N-1}{2}}}{2} d\xi = \pi \int_{\mathbb{R}^{N-1}} \frac{1}{(1+|\xi|^2)^{\frac{N-1}{2}}} d\xi$$

$$\int_{\mathbb{R}^{N-1}} \frac{1}{(1+|\xi|^2)^{\frac{N-1}{2}}} d\xi = \int_{\mathbb{R}^{N-1}} \frac{1}{(1+|\xi|^2)^{\frac{N-1}{2}}} d\xi = \frac{2\pi}{\sqrt{2}} \int_{\mathbb{R}^{N-1}} \frac{1}{(1+|\xi|^2)^{\frac{N-1}{2}}} d\xi$$

$$\Rightarrow W \in H^{\frac{N-1}{2}}(\mathbb{R}^{N-1}), \quad \|W\|_{H^{\frac{N-1}{2}}(\mathbb{R}^{N-1})} \leq \pi \|\psi\|_{L^2(\mathbb{R})}, \quad W \in \mathcal{D}(\mathbb{R}^{N-1})$$

$$W \in H^1(\mathbb{R}^N) \text{ By reflect extend to } H^1(\mathbb{R}^N).$$

$$W_n \in \mathcal{D}(\mathbb{R}^N) \quad W_n \rightarrow W \text{ in } H^1(\mathbb{R}^N).$$



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integral here  $d\xi_N$  and the  $d\xi'$  and therefore you get integral of  $\mathbb{R}_+^N$ , sorry, mod  $v$  hat  $z$  square into  $1 + \text{mod } \xi \text{ square } dx, d\xi$

That is not here,  $\pi$  times integral over  $\mathbb{R}^N$  of  $\text{mod } v \text{ hat } \xi \text{ square into } 1 + \text{mod } \xi \text{ square } d\xi$  and what is that that is nothing but  $\pi$  into  $\text{norm } v \text{ square } 1 \mathbb{R}^N$  so this implies that  $w$  belongs to  $H^{1/2}(\mathbb{R}^{N-1})$ , and you have  $\text{norm } W \text{ in half } \mathbb{R}^{N-1}$  is less than equal to  $\pi$  times  $\text{norm } v \text{ in } 1 \mathbb{R}^N$  and this is for  $v$  in  $D$  of  $\mathbb{R}^N$ . So now if you take  $v$  in  $H^1(\mathbb{R}_+^N)$  then by reflection extend to  $H^1(\mathbb{R}^N)$ .

Now, if you extend it to  $H^1(\mathbb{R}^N)$ , then you know  $D$  of  $\mathbb{R}^N$  is dense there, so you take  $W_n$ , sorry,  $v_n$  in  $D$  of  $\mathbb{R}^N$ ,  $v_n$  converging to  $v$  in  $H^1$  of  $\mathbb{R}^N$ . Then by what we have just see, we saw that if you take  $W_n - W_m$  in half  $\mathbb{R}^{N-1}$ , this is less than equal to  $\pi$  times  $\text{norm } v_n - v_m \text{ in } 1 \mathbb{R}^N$  and therefore this implies that  $W_n$  is Cauchy in  $H^{1/2}(\mathbb{R}^{N-1})$ , so  $W_n$  converges to some  $W$  in  $H^{1/2}(\mathbb{R}^{N-1})$ .

But  $v_n$  also converges to  $v$  in  $H^1$  of  $\mathbb{R}^N$  plus and that implies that  $\gamma v_n$  converges to  $\gamma v$  in  $L^2$  of  $\mathbb{R}^{N-1}$ , but  $\gamma v_n$  is nothing but  $W_n$  and that goes to some  $W$  and therefore this implies that  $W = \gamma v$  and this implies that  $\gamma v$  belongs to  $H^{1/2}(\mathbb{R}^{N-1})$ ; also you have that  $\text{norm } \gamma v \text{ in } H^{1/2}(\mathbb{R}^{N-1})$  is less than or equal to  $\pi$  times  $\text{norm } v$  in  $H^1$  of  $\mathbb{R}^N$ ...

We have this 1,  $\mathbb{R}^N$  and therefore  $\text{norm } \gamma v \text{ in half } \mathbb{R}^{N-1}$  is less than equal to  $\pi$  times  $\text{norm } v \text{ in } 1 \mathbb{R}^N$ , but  $\gamma$  is nothing but the reflection thing so that is a prolongation operator, so that is  $C$  times  $\text{norm } v \text{ in } \mathbb{R}^N$  and therefore you have that the map  $\gamma$  from  $H^1$  of  $\mathbb{R}_+^N$  to  $H^{1/2}$  of  $\mathbb{R}^{N-1}$  is continuous.

So, in fact, we can show that  $\gamma$  is surjective; that is, the range of  $\gamma$  equals  $H^{1/2}$  of  $\mathbb{R}^{N-1}$ . I am not going to prove this, you can check the proof in the

book Topics in Functional Analysis and Applications which I mentioned to you and you have that in fact you have, you can actually construct a pre image for this one.

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Handwritten notes on a slide with an NPTEL logo. The notes are as follows:

$\mathcal{R}(\gamma_0) = H^{1/2}(\mathbb{R}^{N-1})$

Remark With very minor modifications in the computations above, we can show that  $\gamma_0: H^m(\mathbb{R}_+^N) \rightarrow H^{m-1/2}(\mathbb{R}^{N-1})$

Also if  $u \in D(\mathbb{R}^N)$  we can show  $-\Delta u \in L^2(\mathbb{R}^{N-1})$  and that, in fact it lies in  $H^{1/2}(\mathbb{R}^{N-1})$  and that this extends as a cont. lin. map  $\gamma_1: H^2(\mathbb{R}_+^N) \rightarrow H^{1/2}(\mathbb{R}^{N-1})$ .

More generally, we have a sequence of maps  $\{\gamma_j\}$  s.t.  $\gamma = (\gamma_0, \dots, \gamma_{m-1}): H^m(\mathbb{R}_+^N) \rightarrow \prod_{j=0}^{m-1} H^{m-j-1/2}(\mathbb{R}^{N-1})$  for any  $m$ .

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So, now, so remarks, so with very minor modifications in the computations above, we can show that

$$\gamma_0: H^m(\mathbb{R}_+^N) \rightarrow H^{m-1/2}(\mathbb{R}^{N-1}).$$

So, it is very easy to show this. Also if  $u$  belongs to  $D(\mathbb{R}^N)$ , we can show minus  $-\frac{du}{dx_n}$

belongs to  $L^2(\mathbb{R}^{N-1})$  and in that in fact it lies in  $H^{1/2}(\mathbb{R}^{N-1})$  and that this extends as a continuous linear map  $\gamma_1$  from  $H^2(\mathbb{R}_+^N)$  to  $H^{1/2}(\mathbb{R}^{N-1})$

More generally, we have a sequence of maps  $\gamma_j$  such that  $\gamma = (\gamma_0, \dots, \gamma_{m-1})$  for any  $m$  positive integer maps  $H^m(\mathbb{R}_+^N)$  to  $\prod_{j=0}^{m-1} H^{m-j-1/2}(\mathbb{R}^{N-1})$ . So, all this can be done. This just modifications of the calculations which we have been doing. Our next aim is to look at the kernel of this trace maps  $\gamma$  and we will see that subsequently.