

Transform Calculus and its Applications in Differential Equations
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Lecture – 48
Solution of Partial Differential Equations using Fourier Transform – 1

Now let us see how to find out the solution of a PDE using Fourier Transform.

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Example

Solve the following PDE using F.C.T.

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} \quad x > 0, t > 0$$

with $u_x(0, t) = 0$ when $t > 0$,

$$u(x, 0) = \begin{cases} x & , \text{ for } 0 \leq x \leq 1 \\ 0 & , \text{ for } x > 1 \end{cases}$$

$u, \frac{\partial u}{\partial x} \rightarrow 0$ as $x \rightarrow \infty$.

The problem is

$$\frac{\partial u}{\partial t} = \frac{\partial^2 u}{\partial x^2} \quad x > 0, t > 0$$

with, $u_x(0, t) = 0$ when $t > 0$ and $u(x, 0) = \begin{cases} x, & 0 \leq x \leq 1 \\ 0, & x > 1 \end{cases}$

$$u, \frac{\partial u}{\partial x} \rightarrow 0 \text{ as } x \rightarrow \infty$$

We apply Fourier cosine transform with respect to x on the given equation. Therefore, we obtain,

$$\begin{aligned} \mathcal{F}_c \left[\frac{\partial u}{\partial t} \right] &= \mathcal{F}_c \left[\frac{\partial^2 u}{\partial x^2} \right] \\ \Rightarrow \sqrt{\frac{2}{\pi}} \int_0^\infty \frac{\partial u}{\partial t} \cos ax dx &= \sqrt{\frac{2}{\pi}} \int_0^\infty \frac{\partial^2 u}{\partial x^2} \cos ax dx \end{aligned}$$

$$\begin{aligned}
& \Rightarrow \frac{d}{dt} \left[\sqrt{\frac{2}{\pi}} \int_0^\infty u \cos \alpha x \, dx \right] = \sqrt{\frac{2}{\pi}} \left[\left[\frac{\partial u}{\partial x} \cos \alpha x \right]_0^\infty + \alpha \int_0^\infty \frac{\partial u}{\partial x} \sin \alpha x \, dx \right] \\
& \Rightarrow \frac{d\bar{u}_c}{dt} = \alpha \sqrt{\frac{2}{\pi}} \left[[u \sin \alpha x]_0^\infty - \alpha \int_0^\infty u \cos \alpha x \, dx \right] \\
& \quad \left[\because \frac{\partial u}{\partial x} \rightarrow 0 \text{ as } x \rightarrow \infty, \quad u_x(0, t) = 0 \right] \\
& \Rightarrow \frac{d\bar{u}_c}{dt} = -\alpha^2 \int_0^\infty u \cos \alpha x \, dx \quad [\because u \rightarrow 0 \text{ as } x \rightarrow \infty] \\
& \Rightarrow \frac{d\bar{u}_c}{dt} = -\alpha^2 \bar{u}_c \quad \text{where } \bar{u}_c(\alpha, t) = \mathcal{F}_c[u(x, t)]
\end{aligned}$$

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Apply F.C.T. w.r.t. x

$$\begin{aligned}
& \int_{-\infty}^{\infty} \int_0^\infty \frac{\partial u}{\partial x} \cos \alpha x \, dx \, dz = \int_{-\infty}^{\infty} \int_0^\infty \frac{\partial u}{\partial z} \cos \alpha z \, dz \, dx \\
& \frac{d}{dt} (\bar{u}_c) = \int_{-\infty}^{\infty} \left[\left[\frac{\partial u}{\partial z} \cos \alpha z \right]_0^\infty + \alpha \int_0^\infty \frac{\partial u}{\partial z} \sin \alpha z \, dz \right] \\
& = \int_{-\infty}^{\infty} \left[\alpha \int_0^\infty \frac{\partial u}{\partial z} \sin \alpha z \, dz \right] \quad u_z(0, t) = 0 \\
& \frac{d}{dt} (\bar{u}_c) = \int_{-\infty}^{\infty} \left[\alpha \underbrace{[u \sin \alpha z]}_0^\infty - \alpha^2 \int_0^\infty u \cos \alpha z \, dz \right] \\
& = -\alpha^2 \bar{u}_c \quad u \rightarrow 0 \text{ as } z \rightarrow \infty
\end{aligned}$$

Thus, the given PDE is reduced to a first order ODE. The obtained ODE can be easily solved to get the solution as

$$\bar{u}_c(\alpha, t) = A e^{-\alpha^2 t} \quad (1)$$

where, A is the constant of integration.

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$$\begin{aligned}
 \frac{d\bar{u}_c}{dt} + \alpha^2 \bar{u}_c &= 0 \Rightarrow \bar{u}_c = A e^{-\alpha^2 t} \\
 \text{At } t=0, \bar{u}_c(\alpha, 0) &= \sqrt{\frac{2}{\pi}} \int_0^\infty u(x, 0) \cos \alpha x dx \\
 &= \sqrt{\frac{2}{\pi}} \int_0^1 x \cos \alpha x dx \\
 &= \sqrt{\frac{2}{\pi}} \left[\frac{x \sin \alpha x}{\alpha} - \frac{1}{\alpha^2} (-\cos \alpha x) \right]_0^1 \\
 &= \sqrt{\frac{2}{\pi}} \left[\frac{\sin \alpha}{\alpha} + \frac{\cos \alpha - 1}{\alpha^2} \right]
 \end{aligned}$$

We are given the following initial condition:

$$u(x, 0) = \begin{cases} x, & 0 \leq x \leq 1 \\ 0, & x > 1 \end{cases}$$

Now, at $t = 0$, we have,

$$\begin{aligned}
 \bar{u}_c(\alpha, 0) &= \sqrt{\frac{2}{\pi}} \int_0^\infty u(x, 0) \cos \alpha x dx \\
 &= \sqrt{\frac{2}{\pi}} \int_0^1 x \cos \alpha x dx \\
 &= \sqrt{\frac{2}{\pi}} \left[\left[\frac{x}{\alpha} \sin \alpha x \right]_0^1 - \frac{1}{\alpha} \int_0^1 \sin \alpha x dx \right] \\
 &= \sqrt{\frac{2}{\pi}} \left[\frac{\sin \alpha}{\alpha} + \frac{1}{\alpha^2} [\cos \alpha x]_0^1 \right] \\
 &= \sqrt{\frac{2}{\pi}} \left[\frac{\sin \alpha}{\alpha} + \frac{1}{\alpha^2} (\cos \alpha - 1) \right]
 \end{aligned}$$

Therefore, (1) implies

$$A = \bar{u}_c(\alpha, 0) = \sqrt{\frac{2}{\pi}} \left[\frac{\sin \alpha}{\alpha} + \frac{1}{\alpha^2} (\cos \alpha - 1) \right]$$

Thus, from (1), we have,

$$\bar{u}_c(\alpha, t) = \sqrt{\frac{2}{\pi}} \left[\frac{\sin \alpha}{\alpha} + \frac{\cos \alpha - 1}{\alpha^2} \right] e^{-\alpha^2 t}$$

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The image shows a handwritten derivation of the inverse Fourier cosine transform formula. It starts with the expression for A and then derives the formula for $\bar{u}_c(\alpha, t)$. The final result is:

$$u(x, t) = \int_0^\infty \bar{u}_c(\alpha) e^{-\alpha^2 t} \cos \alpha x d\alpha$$

$$= \sqrt{\frac{2}{\pi}} \left[\sqrt{\frac{2}{\pi}} \int_0^\infty \left[\frac{\sin \alpha}{\alpha} + \frac{\cos \alpha - 1}{\alpha^2} \right] e^{-\alpha^2 t} \cos \alpha x d\alpha \right]$$

Now, taking the inverse Fourier cosine transform, we have,

$$\begin{aligned} u(x, t) &= \mathcal{F}_c^{-1}[\bar{u}_c(\alpha, t)] \\ &= \sqrt{\frac{2}{\pi}} \mathcal{F}_c^{-1} \left[\left[\frac{\sin \alpha}{\alpha} + \frac{\cos \alpha - 1}{\alpha^2} \right] e^{-\alpha^2 t} \right] \\ &= \sqrt{\frac{2}{\pi}} \sqrt{\frac{2}{\pi}} \int_0^\infty \left[\frac{\sin \alpha}{\alpha} + \frac{\cos \alpha - 1}{\alpha^2} \right] e^{-\alpha^2 t} \cos \alpha x d\alpha \\ &= \frac{2}{\pi} \int_0^\infty \left[\frac{\sin \alpha}{\alpha} + \frac{\cos \alpha - 1}{\alpha^2} \right] e^{-\alpha^2 t} \cos \alpha x d\alpha \end{aligned}$$

Evaluation of the above integral will give the required solution for $u(x, t)$.

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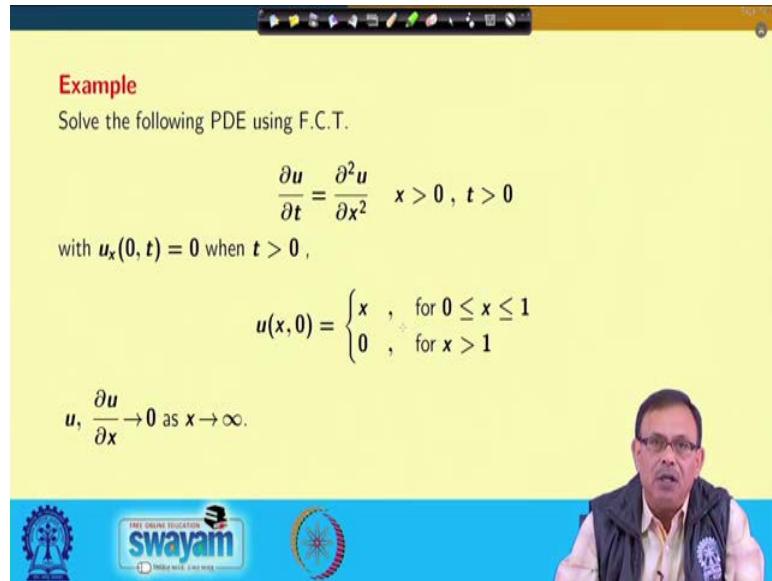
Example
Solve the following PDE using F.C.T.

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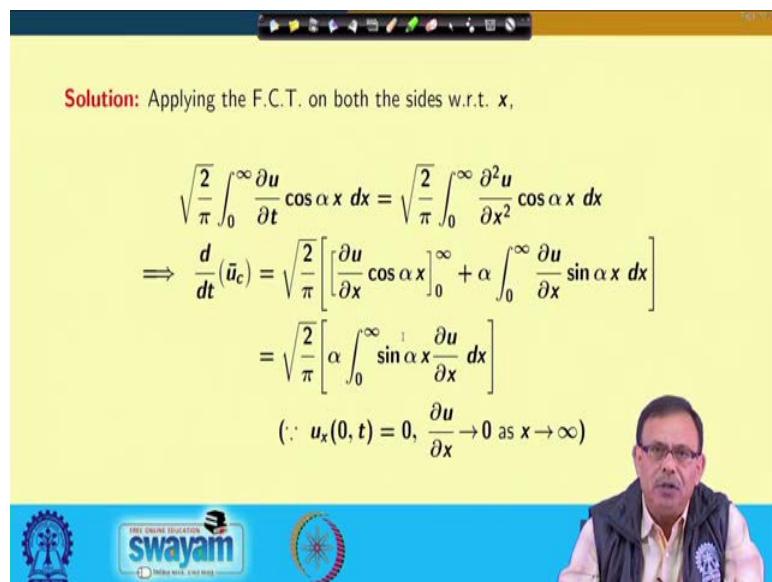
$$u(x, 0) = \begin{cases} x & , \text{ for } 0 \leq x \leq 1 \\ 0 & , \text{ for } x > 1 \end{cases}$$

$u, \frac{\partial u}{\partial x} \rightarrow 0$ as $x \rightarrow \infty$.



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Solution: Applying the F.C.T. on both the sides w.r.t. x ,

$$\sqrt{\frac{2}{\pi}} \int_0^\infty \frac{\partial u}{\partial t} \cos \alpha x \, dx = \sqrt{\frac{2}{\pi}} \int_0^\infty \frac{\partial^2 u}{\partial x^2} \cos \alpha x \, dx$$
$$\Rightarrow \frac{d}{dt}(\tilde{u}_c) = \sqrt{\frac{2}{\pi}} \left[\left[\frac{\partial u}{\partial x} \cos \alpha x \right]_0^\infty + \alpha \int_0^\infty \frac{\partial u}{\partial x} \sin \alpha x \, dx \right]$$
$$= \sqrt{\frac{2}{\pi}} \left[\alpha \int_0^\infty \sin \alpha x \frac{\partial u}{\partial x} \, dx \right]$$
$$(\because u_x(0, t) = 0, \frac{\partial u}{\partial x} \rightarrow 0 \text{ as } x \rightarrow \infty)$$


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$$\begin{aligned}\implies \frac{d}{dt}(\bar{u}_c) &= \sqrt{\frac{2}{\pi}} \left[\alpha [u \sin \alpha x]_0^\infty - \alpha^2 \int_0^\infty u \cos \alpha x \, dx \right] \\ &= -\alpha^2 \bar{u}_c \quad (\because u \rightarrow 0 \text{ as } x \rightarrow \infty) \\ \therefore \frac{d\bar{u}_c}{dt} + \alpha^2 \bar{u}_c &= 0 \\ \therefore \bar{u}_c &= Ae^{-\alpha^2 t}\end{aligned}$$

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$$\begin{aligned}\text{At } t = 0, \bar{u}_c(\alpha, 0) &= \sqrt{\frac{2}{\pi}} \int_0^\infty u(x, 0) \cos \alpha x \, dx \\ &= \sqrt{\frac{2}{\pi}} \int_0^1 x \cos \alpha x \, dx \\ &= \sqrt{\frac{2}{\pi}} \left[\frac{x \sin \alpha x}{\alpha} - \frac{1}{\alpha^2} (-\cos \alpha x) \right]_0^1 \\ &= \sqrt{\frac{2}{\pi}} \left[\frac{\sin \alpha}{\alpha} + \frac{\cos \alpha - 1}{\alpha^2} \right]\end{aligned}$$

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The slide contains the following mathematical steps:

$$\therefore A = \sqrt{\frac{2}{\pi}} \left[\frac{\sin \alpha}{\alpha} + \frac{\cos \alpha - 1}{\alpha^2} \right]$$
$$\implies \bar{u}_c = \sqrt{\frac{2}{\pi}} \left[\frac{\sin \alpha}{\alpha} + \frac{\cos \alpha - 1}{\alpha^2} \right] e^{-\alpha^2 t}$$
$$\implies u(x, t) = \sqrt{\frac{2}{\pi}} \int_0^\infty \bar{u}_c \cos \alpha x \, d\alpha$$
$$= \frac{2}{\pi} \int_0^\infty \left(\frac{\sin \alpha}{\alpha} + \frac{\cos \alpha - 1}{\alpha^2} \right) e^{-\alpha^2 t} \cos \alpha x \, d\alpha$$

At the bottom of the slide, there is a blue footer bar with the Indian Space Research Organisation (ISRO) logo, the text "FREE ONLINE EDUCATION", the "swayam" logo, and the text "India Me, Jai Me".

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The slide features a red "Example" header. Below it, the text reads "Solve the following PDE using F.T." followed by the PDE equation $\frac{\partial \theta}{\partial t} = k \frac{\partial^2 \theta}{\partial x^2}$ with the condition $-\infty < x < \infty, t > 0$. It also includes the boundary condition $\theta = f(x)$ at $t = 0$ and $\theta, \frac{\partial \theta}{\partial x} \rightarrow 0$ as $x \rightarrow \infty$.

On the right side of the slide, there is a video feed of a man wearing glasses and a dark vest, speaking to the camera.

At the bottom of the slide, there is a blue footer bar with the Indian Space Research Organisation (ISRO) logo, the text "FREE ONLINE EDUCATION", the "swayam" logo, and the text "India Me, Jai Me".

Now, let us solve another problem using Fourier transform. The problem is given as:

$$\frac{\partial \theta}{\partial t} = k \frac{\partial^2 \theta}{\partial x^2} \quad -\infty < x < \infty, t > 0$$

with $\theta = f(x)$ at $t = 0$ and

$$\theta, \frac{\partial \theta}{\partial x} \rightarrow 0 \text{ as } x \rightarrow \pm\infty$$

We have two variables here, x and t ; t is given as greater than 0 and x lies between $-\infty$ to ∞ .

And we know that the range for Fourier transform is from $-\infty$ to ∞ . Also it has been given that, $\theta, \frac{\partial \theta}{\partial x} \rightarrow 0$ as $x \rightarrow \pm\infty$. Therefore, from the given criteria, we can tell that, we have to use Fourier transform on this given equation with respect to the variable x , not with respect to the variable t because range of t is from 0 to ∞ , not from $-\infty$ to ∞ .

Here, since x varies from $-\infty$ to ∞ , we can take Fourier transform.

We apply Fourier transform with respect to x on the given equation. Therefore, we obtain,

$$\begin{aligned}\mathcal{F}\left[\frac{\partial \theta}{\partial t}\right] &= k \mathcal{F}\left[\frac{\partial^2 \theta}{\partial x^2}\right] \\ \Rightarrow \frac{d}{dt} \mathcal{F}[\theta] &= k(-i\alpha)^2 \mathcal{F}[\theta] \\ \Rightarrow \frac{d\bar{\theta}}{dt} &= -k\alpha^2 \bar{\theta} \quad \text{where, } \bar{\theta}(\alpha, t) = \mathcal{F}[\theta(x, t)]\end{aligned}$$

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The handwritten derivation shows the application of Fourier transform to the partial differential equation. It starts with the PDE $\frac{\partial \theta}{\partial t} = k \frac{\partial^2 \theta}{\partial x^2}$. The Fourier transform of the second derivative is $\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{\partial \theta}{\partial x} e^{i\alpha x} dx = k \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{\partial^2 \theta}{\partial x^2} e^{i\alpha x} dx$. This leads to the transformed equation $\frac{d\bar{\theta}}{dt} = k \frac{1}{\sqrt{2\pi}} \left[\left[\frac{\partial \theta}{\partial x} e^{i\alpha x} \right]_{-\infty}^{\infty} - i\alpha \int_{-\infty}^{\infty} \frac{\partial \theta}{\partial x} e^{i\alpha x} dx \right]$. Simplifying, we get $= -\frac{i\alpha k}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{i\alpha x} \frac{\partial \theta}{\partial x} dx$. As $\theta \rightarrow 0$ at $x \rightarrow \infty$, the boundary term vanishes. This results in $= -\frac{i\alpha k}{\sqrt{2\pi}} \left[\left[\theta e^{i\alpha x} \right]_{-\infty}^{\infty} - i\alpha \int_{-\infty}^{\infty} \theta e^{i\alpha x} dx \right]$. Finally, as $\theta \rightarrow 0$ at $x \rightarrow \infty$, the boundary term vanishes, leaving us with $= -k\alpha^2 \bar{\theta}$.

Therefore, the given PDE is reduced to a first order ODE whose solution is given as:

$$\bar{\theta}(\alpha, t) = Ae^{-k\alpha^2 t} \tag{2}$$

where, A is the constant of integration.

We are given the initial condition as

$$\theta = f(x) \text{ at } t = 0$$

Now,

$$\begin{aligned}\bar{\theta}(\alpha, 0) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \theta(x, 0) e^{i\alpha x} dx \\ &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{i\alpha x} dx \\ &= \bar{f}(\alpha)\end{aligned}$$

where, $\bar{f}(\alpha)$ denotes the Fourier transform of $f(x)$.

Therefore, (2) implies

$$A = \bar{\theta}(\alpha, 0) = \bar{f}(\alpha)$$

Thus, from (2), we have,

$$\bar{\theta}(\alpha, t) = \bar{f}(\alpha) e^{-k\alpha^2 t}$$

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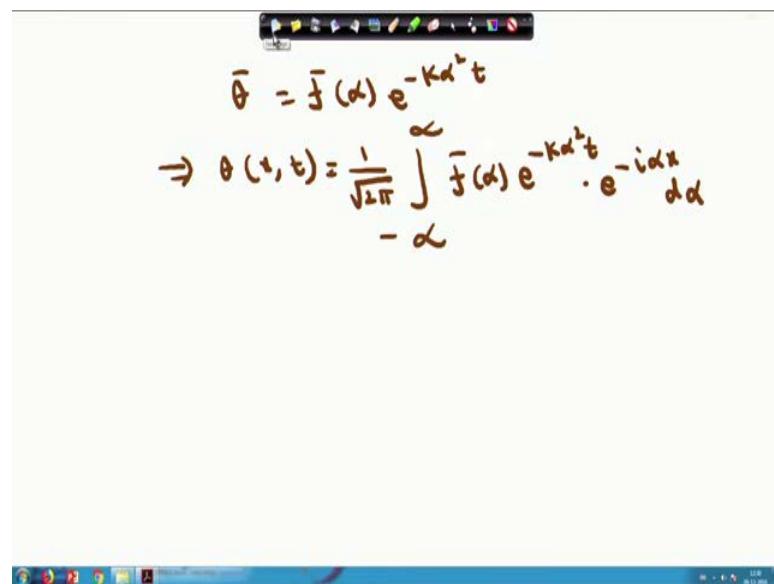
The image shows a handwritten derivation of the Fourier transform of the heat kernel. It starts with the heat equation $\frac{d\bar{\theta}}{dt} + k\alpha^2 \bar{\theta} = 0$, leading to the solution $\bar{\theta} = A e^{-k\alpha^2 t}$. At $t = 0$, $\bar{\theta}(\alpha, 0) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \theta(x, 0) e^{i\alpha x} dx$. This integral is shown as $= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{i\alpha x} dx$, where $f(x)$ is indicated by a bracket under the integral. The result is $= \bar{f}(\alpha)$. Below this, it is noted that $A = \bar{f}(\alpha)$.

Now, taking the inverse Fourier transform, we have,

$$\begin{aligned}\theta(x, t) &= \mathcal{F}^{-1}[\bar{\theta}(\alpha, t)] \\ &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \bar{f}(\alpha) e^{-k\alpha^2 t} e^{-i\alpha x} d\alpha\end{aligned}$$

If $f(x)$ is known to us, then we can calculate $\bar{f}(\alpha)$ also. And hence, we can evaluate the integral to obtain the required solution for $\theta(x, t)$.

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Example

Solve the following PDE using F.T.

$$\frac{\partial \theta}{\partial t} = k \frac{\partial^2 \theta}{\partial x^2} \quad -\infty < x < \infty, t > 0$$

with $\theta = f(x)$, at $t = 0$ and $\theta, \frac{\partial \theta}{\partial x} \rightarrow 0$ as $x \rightarrow \infty$.

A video player interface is visible at the top of the slide, showing a progress bar and control buttons. The bottom of the slide features a blue footer with the Indian Space Research Organisation (ISRO) logo, the Swayam logo, and the text "FREE ONLINE EDUCATION" and "India's first, only and largest". A man in a grey vest and glasses is visible on the right side of the slide.

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Solution: Applying the F.T. on both the sides w.r.t. x ,

$$\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{\partial \theta}{\partial t} e^{i\alpha x} dx = k \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \frac{\partial^2 \theta}{\partial x^2} e^{i\alpha x} dx$$
$$\Rightarrow \frac{d}{dt}(\bar{\theta}) = k \frac{1}{\sqrt{2\pi}} \left[\left[\frac{\partial \theta}{\partial x} e^{i\alpha x} \right]_{-\infty}^{\infty} - i\alpha \int_{-\infty}^{\infty} \frac{\partial \theta}{\partial x} e^{i\alpha x} dx \right]$$
$$\Rightarrow \frac{d}{dt}(\bar{\theta}) = -\frac{i\alpha k}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{i\alpha x} \frac{\partial \theta}{\partial x} dx$$
$$\left(\because \frac{\partial \theta}{\partial x} \rightarrow 0 \text{ as } x \rightarrow \infty \right)$$

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$$\Rightarrow \frac{d}{dt}(\bar{\theta}) = -\frac{i\alpha k}{\sqrt{2\pi}} \left[\left[\theta e^{i\alpha x} \right]_{-\infty}^{\infty} - i\alpha \int_{-\infty}^{\infty} \theta e^{i\alpha x} dx \right]$$
$$= -k\alpha^2 \bar{\theta} \quad (\because \theta \rightarrow 0 \text{ as } x \rightarrow \infty)$$
$$\therefore \frac{d\bar{\theta}}{dt} + k\alpha^2 \bar{\theta} = 0$$
$$\therefore \bar{\theta} = Ae^{-k\alpha^2 t}$$

At $t = 0$, $\bar{\theta}(\alpha, 0) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \theta(x, 0) e^{i\alpha x} dx$

$$= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(x) e^{i\alpha x} dx$$
$$= \bar{f}(\alpha)$$

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$$\begin{aligned}\therefore A &= \tilde{f}(\alpha) \\ \Rightarrow \bar{\theta} &= \tilde{f}(\alpha)e^{-k\alpha^2 t} \\ \Rightarrow \theta(x, t) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \tilde{f}(\alpha)e^{-k\alpha^2 t} e^{-i\alpha x} d\alpha \\ &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \tilde{f}(\alpha)e^{-k\alpha^2 t - i\alpha x} d\alpha\end{aligned}$$

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Example
Solve the following PDE using F.S.T.

$$\frac{\partial v}{\partial t} = k \frac{\partial^2 v}{\partial x^2} \quad x > 0, t > 0$$

with $v = v_0$, when $x = 0, t > 0$, $v = 0$, when $t = 0, x > 0$, and
 $v, \frac{\partial v}{\partial x} \rightarrow 0$ as $x \rightarrow \infty$.

Now, let us take another problem. We have solved this particular problem earlier using Fourier cosine transform. Now we want to solve the same problem using Fourier sine transform. The problem is given as:

$$\frac{\partial v}{\partial t} = k \frac{\partial^2 v}{\partial x^2} \quad x > 0, t > 0$$

with, $v(0, t) = v_0$ when $t > 0$ and $v(x, 0) = 0$ when $x > 0$

$$v, \frac{\partial v}{\partial x} \rightarrow 0 \text{ as } x \rightarrow \infty$$

We apply Fourier sine transform with respect to x on the given equation. Therefore, we obtain,

$$\begin{aligned}
 \mathcal{F}_s \left[\frac{\partial v}{\partial t} \right] &= k \mathcal{F}_s \left[\frac{\partial^2 v}{\partial x^2} \right] \\
 \Rightarrow \sqrt{\frac{2}{\pi}} \int_0^\infty \frac{\partial v}{\partial t} \sin \alpha x \, dx &= k \sqrt{\frac{2}{\pi}} \int_0^\infty \frac{\partial^2 v}{\partial x^2} \sin \alpha x \, dx \\
 \Rightarrow \frac{d}{dt} \left[\sqrt{\frac{2}{\pi}} \int_0^\infty v \sin \alpha x \, dx \right] &= k \sqrt{\frac{2}{\pi}} \left[\left[\frac{\partial v}{\partial x} \sin \alpha x \right]_0^\infty - \alpha \int_0^\infty \frac{\partial v}{\partial x} \cos \alpha x \, dx \right] \\
 \Rightarrow \frac{d\bar{v}_s}{dt} &= -k\alpha \sqrt{\frac{2}{\pi}} \left[[v \cos \alpha x]_0^\infty + \alpha \int_0^\infty v \sin \alpha x \, dx \right] \quad [\because \frac{\partial v}{\partial x} \rightarrow 0 \text{ as } x \rightarrow \infty] \\
 \Rightarrow \frac{d\bar{v}_s}{dt} &= -k\alpha \sqrt{\frac{2}{\pi}} \left[-v_0 + \alpha \int_0^\infty v \sin \alpha x \, dx \right] \quad [\because v \rightarrow 0 \text{ as } x \rightarrow \infty, v(0, t) = v_0] \\
 \Rightarrow \frac{d\bar{v}_s}{dt} &= \sqrt{\frac{2}{\pi}} k\alpha v_0 - k\alpha^2 \bar{v}_s \quad \text{where } \bar{v}_s(\alpha, t) = \mathcal{F}_s[v(x, t)]
 \end{aligned}$$

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The image shows a handwritten derivation of the Fourier sine transform application. It starts with the equation:

$$\sqrt{\frac{2}{\pi}} \int_0^\infty \frac{\partial v}{\partial t} \sin \alpha x \, dx = k \sqrt{\frac{2}{\pi}} \int_0^\infty \frac{\partial^2 v}{\partial x^2} \sin \alpha x \, dx$$

Then it differentiates with respect to t :

$$\frac{d}{dt} \left(\sqrt{\frac{2}{\pi}} \int_0^\infty v \sin \alpha x \, dx \right) = k \sqrt{\frac{2}{\pi}} \left[\left[\frac{\partial v}{\partial x} \sin \alpha x \right]_0^\infty - \alpha \int_0^\infty \frac{\partial v}{\partial x} \cos \alpha x \, dx \right]$$

This results in:

$$= -k\alpha \sqrt{\frac{2}{\pi}} \left[\int_0^\infty \frac{\partial v}{\partial x} \cos \alpha x \, dx \right]$$

$$= -k\alpha \sqrt{\frac{2}{\pi}} \left[[v \cos \alpha x]_0^\infty + \alpha \int_0^\infty v \sin \alpha x \, dx \right]$$

Finally, it is simplified to:

$$= k\alpha v_0 \sqrt{\frac{2}{\pi}} - k\alpha^2 \bar{v}_s$$

Notes at the bottom right indicate $v \rightarrow 0$ and $v = v_0$ when $x \rightarrow \infty$.

This reduces the given PDE to a first order ODE which can be easily solved and the solution is given as:

$$\bar{v}_s e^{k\alpha^2 t} = C + \sqrt{\frac{2}{\pi}} \frac{v_0}{\alpha} e^{k\alpha^2 t} \quad (3)$$

where, C is the constant of integration.

We have the initial condition as $v(x, 0) = 0$. Therefore,

$$\bar{v}_s(\alpha, 0) = 0$$

Thus, (3) implies

$$C = -\sqrt{\frac{2}{\pi}} \frac{v_0}{\alpha}$$

Therefore, from (3), we have,

$$\begin{aligned} \bar{v}_s e^{k\alpha^2 t} &= -\sqrt{\frac{2}{\pi}} \frac{v_0}{\alpha} + \sqrt{\frac{2}{\pi}} \frac{v_0}{\alpha} e^{k\alpha^2 t} \\ \Rightarrow \bar{v}_s(\alpha, t) &= \sqrt{\frac{2}{\pi}} \frac{v_0}{\alpha} (1 - e^{-k\alpha^2 t}) \end{aligned}$$

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The image shows a handwritten derivation of the solution for \bar{v}_s over time t . It starts with the differential equation:

$$\frac{d\bar{v}_s}{dt} + k\alpha^2 \bar{v}_s = kdv_0 \sqrt{\frac{2}{\pi}}$$

Then, it separates the variables and integrates both sides:

$$\int \bar{v}_s e^{k\alpha^2 t} dt = C + \sqrt{\frac{2}{\pi}} \frac{v_0}{\alpha} e^{k\alpha^2 t} \quad (1)$$

At $t = 0$, $\bar{v}_s = 0$ (underlined), which implies $v = 0$ at $t = 0$.

From (1),

$$0 = C + \sqrt{\frac{2}{\pi}} \frac{v_0}{\alpha}$$

$$C = -\sqrt{\frac{2}{\pi}} \frac{v_0}{\alpha}$$

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The derivation shows the following steps:

$$\bar{v}_s = \sqrt{\frac{2}{\pi}} \frac{v_0}{a} (1 - e^{-ka^2 t})$$

$$v(x, t) = \frac{2}{\pi} \int_0^\infty \frac{1 - e^{-ka^2 t}}{a} \sin ax da$$

$$= \frac{2v_0}{\pi} \left[\left(\int_0^\infty \frac{\sin ax}{a} da \right) - \int_0^\infty \frac{e^{-ka^2 t}}{a} \sin ax da \right]$$

$$= \cancel{2v_0} \left[1 - \frac{2}{\pi} \int_0^\infty \frac{e^{-ka^2 t}}{a} \sin ax da \right]$$

Taking the inverse Fourier sine transform, we have,

$$v(x, t) = \mathcal{F}_s^{-1}[\bar{v}_s(\alpha, t)]$$

$$\Rightarrow v(x, t) = \sqrt{\frac{2}{\pi}} v_0 \mathcal{F}_s^{-1} \left[\frac{1}{\alpha} (1 - e^{-ka^2 t}) \right]$$

$$\Rightarrow v(x, t) = \frac{2v_0}{\pi} \int_0^\infty \frac{1}{\alpha} (1 - e^{-ka^2 t}) \sin ax d\alpha$$

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The derivation shows the following steps:

$$f(t) = \int_0^\infty \frac{e^{-ka^2 t}}{a} \sin ax da$$

$$f(t) = \int_0^\infty e^{-ka^2 t} \cos ax da$$

$$= \frac{\sqrt{\pi}}{2\sqrt{t}} e^{-\frac{x^2}{4kt}}$$

$$f(t) = \frac{1}{2} \sqrt{\frac{\pi}{t}} \int_0^x e^{-\frac{x^2}{4kt}} dx$$

$$= \frac{1}{2} \sqrt{\frac{\pi}{t}} 2\sqrt{kt} \int_0^0 e^{-\frac{x^2}{4kt}} d\left(\frac{x}{2\sqrt{kt}}\right)$$

Therefore, we have,

$$\begin{aligned}
 v(x, t) &= \frac{2v_0}{\pi} \int_0^\infty \frac{1}{\alpha} (1 - e^{-k\alpha^2 t}) \sin \alpha x d\alpha \\
 \Rightarrow v(x, t) &= \frac{2v_0}{\pi} \int_0^\infty \frac{\sin \alpha x}{\alpha} d\alpha - \frac{2v_0}{\pi} \int_0^\infty \frac{\sin \alpha x}{\alpha} e^{-k\alpha^2 t} d\alpha \\
 \Rightarrow v(x, t) &= \frac{2v_0 \pi}{\pi} \frac{1}{2} - \frac{2v_0}{\pi} \int_0^\infty \frac{\sin \alpha x}{\alpha} e^{-k\alpha^2 t} d\alpha \\
 \Rightarrow v(x, t) &= v_0 - \frac{2v_0}{\pi} f(x)
 \end{aligned}$$

where,

$$\begin{aligned}
 f(x) &= \int_0^\infty \frac{\sin \alpha x}{\alpha} e^{-k\alpha^2 t} d\alpha \\
 \Rightarrow f'(x) &= \int_0^\infty e^{-k\alpha^2 t} \cos \alpha x d\alpha \\
 \Rightarrow f'(x) &= \frac{\sqrt{\pi}}{2\sqrt{kt}} e^{-\frac{x^2}{4kt}} \\
 \Rightarrow f(x) &= \frac{\sqrt{\pi}}{2\sqrt{kt}} \int_0^x e^{-\frac{u^2}{4kt}} du \\
 \Rightarrow f(x) &= \sqrt{\pi} \int_0^x e^{-\frac{u^2}{4kt}} d\left(\frac{u}{2\sqrt{kt}}\right)
 \end{aligned}$$

(Refer Slide Time: 31:37)

$$\begin{aligned}
 \int v(x, t) = & \frac{1}{2} \sqrt{\frac{\pi}{t}} \cdot 2 \sqrt{kt} \int e^{-\frac{u^2}{4kt}} du \quad u = \frac{x}{2\sqrt{kt}} \\
 & = \frac{\pi}{2} e^{-\frac{x^2}{4kt}}
 \end{aligned}$$

$$\begin{aligned}
 v(x, t) &= v_0 \left[1 - e^{-\frac{x^2}{4kt}} \right] \\
 &= v_0 e^{-\frac{x^2}{4kt}}
 \end{aligned}$$

Put

$$u = \frac{x}{2\sqrt{kt}} \quad \text{so that} \quad du = \frac{dx}{2\sqrt{kt}}$$

Therefore,

$$\begin{aligned} f(x) &= \sqrt{\pi} \int_0^{\frac{x}{2\sqrt{kt}}} e^{-u^2} du \\ \Rightarrow f(x) &= \sqrt{\pi} \frac{\sqrt{\pi}}{2} \operatorname{erf}\left(\frac{x}{2\sqrt{kt}}\right) \\ \Rightarrow f(x) &= \frac{\pi}{2} \operatorname{erf}\left(\frac{x}{2\sqrt{kt}}\right) \end{aligned}$$

Therefore,

$$\begin{aligned} v(x, t) &= v_0 - \frac{2v_0}{\pi} \frac{\pi}{2} \operatorname{erf}\left(\frac{x}{2\sqrt{kt}}\right) \\ \Rightarrow v(x, t) &= v_0 - v_0 \operatorname{erf}\left(\frac{x}{2\sqrt{kt}}\right) \\ \Rightarrow v(x, t) &= v_0 \left[1 - \operatorname{erf}\left(\frac{x}{2\sqrt{kt}}\right) \right] = v_0 \operatorname{erfc}\left(\frac{x}{2\sqrt{kt}}\right) \end{aligned}$$

which is the required solution.

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Solution: Applying the F.S.T. on both the sides w.r.t. x ,

$$\begin{aligned} \sqrt{\frac{2}{\pi}} \int_0^\infty \frac{\partial v}{\partial t} \sin \alpha x \, dx &= k \sqrt{\frac{2}{\pi}} \int_0^\infty \frac{\partial^2 v}{\partial x^2} \sin \alpha x \, dx \\ \Rightarrow \frac{d}{dt} (\tilde{v}_s) &= k \sqrt{\frac{2}{\pi}} \left[\left[\frac{\partial v}{\partial x} \sin \alpha x \right]_0^\infty - \alpha \int_0^\infty \frac{\partial v}{\partial x} \cos \alpha x \, dx \right] \\ \Rightarrow \frac{d}{dt} (\tilde{v}_s) &= -\alpha k \sqrt{\frac{2}{\pi}} \left[\int_0^\infty \frac{\partial v}{\partial x} \cos \alpha x \, dx \right] \\ &\quad \left(\because \frac{\partial v}{\partial x} \rightarrow 0 \text{ as } x \rightarrow \infty \right) \end{aligned}$$

(Refer Slide Time: 33:11)

The slide shows the derivation of the equation for velocity \bar{v}_s over time t . It starts with the differential equation $\frac{d}{dt}(\bar{v}_s) = -\alpha k \sqrt{\frac{2}{\pi}} \left[[v \cos \alpha x]_0^\infty + \alpha \int_0^\infty v \sin \alpha x \, dx \right]$, which simplifies to $= k\alpha v_0 \sqrt{\frac{2}{\pi}} - k\alpha^2 \bar{v}_s$ ($\because v \rightarrow 0$ as $x \rightarrow \infty$, $v = v_0$ at $x = 0$). This leads to the homogeneous differential equation $\frac{d\bar{v}_s}{dt} + k\alpha^2 \bar{v}_s = k\alpha v_0 \sqrt{\frac{2}{\pi}}$. The solution is given as $\bar{v}_s e^{k\alpha^2 t} = c + \sqrt{\frac{2}{\pi}} \frac{v_0}{\alpha} e^{k\alpha^2 t}$, where c is a constant. The slide is titled "FREE ONLINE EDUCATION" and features the "swayam" logo.

$$\begin{aligned}\Rightarrow \frac{d}{dt}(\bar{v}_s) &= -\alpha k \sqrt{\frac{2}{\pi}} \left[[v \cos \alpha x]_0^\infty + \alpha \int_0^\infty v \sin \alpha x \, dx \right] \\ &= k\alpha v_0 \sqrt{\frac{2}{\pi}} - k\alpha^2 \bar{v}_s \quad (\because v \rightarrow 0 \text{ as } x \rightarrow \infty, v = v_0 \text{ at } x = 0) \\ \therefore \frac{d\bar{v}_s}{dt} + k\alpha^2 \bar{v}_s &= k\alpha v_0 \sqrt{\frac{2}{\pi}} \\ \therefore \bar{v}_s e^{k\alpha^2 t} &= c + \sqrt{\frac{2}{\pi}} \frac{v_0}{\alpha} e^{k\alpha^2 t} \quad (1)\end{aligned}$$

(Refer Slide Time: 33:21)

The slide shows the solution for velocity \bar{v}_s at $t = 0$. It states that at $t = 0$, $\bar{v}_s = 0$ ($\because v = 0$ at $t = 0$). From equation (1), $0 = c + \sqrt{\frac{2}{\pi}} \frac{v_0}{\alpha}$, it follows that $c = -\sqrt{\frac{2}{\pi}} \frac{v_0}{\alpha}$. Substituting c back into the general solution, we get $\bar{v}_s e^{k\alpha^2 t} = \sqrt{\frac{2}{\pi}} \frac{v_0}{\alpha} (e^{k\alpha^2 t} - 1)$, and finally $\bar{v}_s = \sqrt{\frac{2}{\pi}} \frac{v_0}{\alpha} (1 - e^{-k\alpha^2 t})$. The slide features the "FREE ONLINE EDUCATION" and "swayam" logos.

$$\begin{aligned}\text{At } t = 0, \bar{v}_s &= 0 \quad (\because v = 0 \text{ at } t = 0) \\ \therefore \text{From (1), } 0 &= c + \sqrt{\frac{2}{\pi}} \frac{v_0}{\alpha} \\ \Rightarrow c &= -\sqrt{\frac{2}{\pi}} \frac{v_0}{\alpha} \\ \therefore \bar{v}_s e^{k\alpha^2 t} &= \sqrt{\frac{2}{\pi}} \frac{v_0}{\alpha} (e^{k\alpha^2 t} - 1) \\ \Rightarrow \bar{v}_s &= \sqrt{\frac{2}{\pi}} \frac{v_0}{\alpha} (1 - e^{-k\alpha^2 t})\end{aligned}$$

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$$\begin{aligned}\Rightarrow v(x, t) &= \frac{2}{\pi} v_0 \int_0^{\infty} \frac{1 - e^{-k\alpha^2 t}}{\alpha} \sin \alpha x \, d\alpha \\ &= \frac{2v_0}{\pi} \left[\int_0^{\infty} \frac{\sin \alpha x}{\alpha} \, d\alpha - \int_0^{\infty} \frac{e^{-k\alpha^2 t}}{\alpha} \sin \alpha x \, d\alpha \right] \\ &= \frac{2v_0}{\pi} \left[\frac{\pi}{2} - \int_0^{\infty} \frac{e^{-k\alpha^2 t}}{\alpha} \sin \alpha x \, d\alpha \right] \\ &= v_0 \left[1 - \frac{2}{\pi} \int_0^{\infty} \frac{e^{-k\alpha^2 t}}{\alpha} \sin \alpha x \, d\alpha \right]\end{aligned}$$


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$$\begin{aligned}\text{Let } f(x) &= \int_0^{\infty} \frac{e^{-k\alpha^2 t}}{\alpha} \sin \alpha x \, d\alpha \\ \therefore f'(x) &= \int_0^{\infty} e^{-k\alpha^2 t} \cos \alpha x \, d\alpha \\ &= \frac{\sqrt{\pi}}{2\sqrt{t}} e^{-\frac{x^2}{4kt}} \\ \Rightarrow f(x) &= \frac{1}{2} \sqrt{\frac{\pi}{t}} \int_0^x e^{-\frac{s^2}{4kt}} \, ds \\ &= \frac{1}{2} \sqrt{\frac{\pi}{t}} 2\sqrt{kt} \int_0^{\frac{x}{2\sqrt{kt}}} e^{-\frac{s^2}{4kt}} \, d\left(\frac{x}{2\sqrt{kt}}\right)\end{aligned}$$


(Refer Slide Time: 34:35)

$$\begin{aligned}\implies f(x) &= \frac{1}{2} \sqrt{\frac{\pi}{t}} 2\sqrt{kt} \int_0^{\frac{x}{2\sqrt{kt}}} e^{-u^2} du \quad [\text{put } u = \frac{x}{2\sqrt{kt}}] \\ &= \frac{\pi}{2} \operatorname{erf} \frac{x}{2\sqrt{kt}} \\ \therefore v(x, t) &= v_0 \left[1 - \operatorname{erf} \frac{x}{2\sqrt{kt}} \right] \\ &= v_0 \operatorname{erfc} \frac{x}{2\sqrt{kt}}\end{aligned}$$

Thank you.