

Advanced Engineering Mathematics
Prof. Jitendra Kumar
Department of Mathematics
Indian Institute of Technology, Kharagpur

Lecture No # 31
Applications of Fourier Transform to PDEs

Welcome back to the lectures on transform calculus. And in the last lecture, we have studied Fourier transform, and its properties. So, today we will continue this lecture for the application part, and we will discuss its application to partial differential equations. So, the procedure is very similar, what we have done in the case of Laplace transform. So, for a given partial differential equation, we will apply the Fourier transform to both the side of the equation. And then, this partial differential equation will be transformed to a simpler ordinary differential equation that we will solve, and at the end by taking the inverse Fourier transform of the solution of this ordinary differential equation. We will arrive for the solution of the original partial differential equation.

(Refer Slide Time: 01:22)

© CEE
IIT KGP

FOURIER COSINE AND SINE TRANSFORM

$F_c(f) = \hat{f}_c(\alpha) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} f(u) \cos \alpha u \, du$	$F_c^{-1}(\hat{f}_c) = f(x) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} \hat{f}_c(\alpha) \cos \alpha x \, d\alpha$
$F_s(f) = \hat{f}_s(\alpha) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} f(u) \sin \alpha u \, du$	$F_s^{-1}(\hat{f}_s) = f(x) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} \hat{f}_s(\alpha) \sin \alpha x \, d\alpha$
$F(f) = \hat{f}(\alpha) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} f(u) e^{i\alpha u} \, du$	$F^{-1}(\hat{f}(\alpha)) = f(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(\alpha) e^{-i\alpha x} \, d\alpha$
PROPERTIES:	
$F_c\{f'(x)\} = \alpha F_s\{f(x)\} - \sqrt{\frac{2}{\pi}} f(0)$ <small>(f(0) = f(0+), f(0-) = 0)</small>	$F_c\{f''(x)\} = -\sqrt{\frac{2}{\pi}} f'(0) - \alpha^2 F_c\{f(x)\}$ <small>(f'(0) = f'(0+) - f'(0-), f'(0-) = 0)</small>
$F_s\{f'(x)\} = -\alpha F_c\{f(x)\}$ <small>(f(0) = 0, f'(0) = f'(0+))</small>	$F_s\{f''(x)\} = \sqrt{\frac{2}{\pi}} \alpha f(0) - \alpha^2 F_s\{f(x)\}$ <small>(f(0) = 0, f'(0) = f'(0+))</small>
$F\{f'(x)\} = -i\alpha F\{f(x)\}$ <small>(f(-\infty) = f(\infty) = 0)</small>	$F\{f''(x)\} = -\alpha^2 F\{f(x)\}$ <small>(f(-\infty) = f(\infty) = 0)</small>

So, before I go for the example let me just summarize some formulas, and the properties of the Fourier transform. So, we have here the Fourier Cosine, and sine transform. So, Fourier Cosine transform was just denoted by F_c hat alpha, and the definition was

square root 2 over pi and integral 0 to infinity $f(u) \cos \alpha u \, du$, and its inverse of Fourier Cosine **Cosine** inverse of this \hat{f}_c will be just $f(x)$, and this is given by square root 2 over pi 0 to infinity $\hat{f}_c \cos \alpha x \, d\alpha$.

Similarly, we have for the sine transform instead of this $\cos \alpha u$, we have $\sin \alpha u$; and similarly here this in place of $\cos \alpha x$, we have $\sin \alpha x \, d\alpha$, and then we studied this Fourier transform. So, in this case, the Fourier transform of f , we will also denote by this \hat{f}_s , because is a function of α now; $1/\sqrt{2\pi}$ and the integral over the whole axis, and we have $f(u) e^{-i\alpha u} \, du$. And for the inverse, we have again the same constant there, $1/\sqrt{2\pi}$ and minus infinity to plus infinity, and then we have \hat{f}_s and this will be $e^{i\alpha x}$ on the integral over this $d\alpha$.

So, these are the main properties are mainly the derivative theorem for the Fourier transform, and Fourier Cosine and sine transform. So, first for the Fourier Cosine transform, we will use this for the first derivative it will be α , and Fourier sine transform of f and minus square root 2 over pi $f(0)$. And this is the **the** condition under which we got this result that was that $f(x)$ approaches to 0 as x approaches to infinity; for the second derivatives the Fourier Cosine transform of the second derivative of f , we have minus square root 2 over pi $f'(0) - \alpha^2$ Fourier Cosine transform of $f(x)$. And in this case, we have this result under the conditions that $f(x)$, and f' as x approaches to 0, the first derivative, and the function itself approaches to 0 as x approaches to 0.

So, in **in** these cases... So, basically we will have solve today the partial differential equations mainly, and in that case the function depends on 2 variables, that we will get this Fourier Cosine or sine or Fourier transform with respect to one variable. So, the same formulas will whole the other variable will be treated as constant. One more point we should mention here, before we go for the **the** examples that we have to see for a particular problem that which one is applicable, whether we should apply the Fourier, Cosine transform or Fourier sine transform or the Fourier transform.

So, one is clear that when the limit of the variable, where we will be taking the Fourier transform is from minus infinity to plus infinity, then we will of course, apply this Fourier transform, but if our range for the variable is given from 0 to infinity. Then we have these two **(())** either Fourier Cosine transform or Fourier sine transform. So, in

these cases, if we just look at the properties - the derivative properties here, so for the Fourier Cosine transform; for example, here in the double derivative will be using for the second order partial differential equations.

So, here this $f'(0)$ appears whereas, in the sine transform this $f(0)$ appears with function value at x is equal to 0, and this is the first derivative of that function at x is equal to 0. So, if this condition is given, and the range of x is 0 to infinity we will apply the sine transform, and if the first derivative - this condition is given a first derivative of 0 is whatever this is given, then we will apply the Fourier Cosine transform. And Fourier transform, when the limit when the range of that variable x from minus infinity to plus infinity.

So, with this information, we continue now for the different partial differential equations, short introduction to partial differential equation, I have already given in Laplace transform case.

(Refer Slide Time: 07:32)

Ex: Heat equation

$$k \frac{\partial^2 u}{\partial x^2} = \frac{\partial u}{\partial t} \quad -\infty < x < \infty, t > 0$$

Bcs: $u(x,t)$ and $u_x(x,t)$ both $\rightarrow 0$ as $|x| \rightarrow \infty$

IC: $u(x,0) = f(x) \quad -\infty < x < \infty$

Sol: Taking Fourier transform w.r.t x :

$$-k\alpha^2 \hat{u}(\alpha, t) = \frac{d}{dt} \hat{u}(\alpha, t)$$

$$\Rightarrow \frac{d\hat{u}}{dt} + k\alpha^2 \hat{u}(\alpha, t) = 0$$

The solution of the ODE:

$$\hat{u}(\alpha, t) = C e^{-k\alpha^2 t} \quad \text{--- (1)}$$

Fourier transform w I.C.

$$\hat{u}(\alpha, 0) = \hat{f}(\alpha) \Rightarrow C = \hat{f}(\alpha)$$

So, we will directly go to the application to be these. So, we will solve first the heat equation **heat equation**, and that is $k \frac{\partial^2 u}{\partial x^2} = \frac{\partial u}{\partial t}$, and our limits are minus infinity to plus infinity t as given positive. The boundary conditions are given $u(x,t)$ and $u_x(x,t)$; both goes to 0 as absolute value of x goes to infinity, and the initial conditions are $f(x,0)$ as $f(x)$. So, we have only the first derivative here is only one initial condition, and for x minus infinity to plus infinity.

So, here the choice of the Fourier transform is clear. So, we will apply here the Fourier transform, because our variable is from minus infinity to plus infinity; just remember for the Laplace transform, **we** we applied Laplace transform with respect to t , because t is always from 0 to infinity, t can vary from 0 to infinity, but now we will apply for the **for** **the** x variable from minus infinity to plus infinity.

So, now taking Fourier transform with respect to x . So, what we will get here, with respect to x is the double derivative, and if we look at the table. So, the double derivatives minus π^2 , and the Fourier transform of $f(x)$. So, we have minus k^2 , and Fourier transform I will denote by this α , and t will remain as it is, because we have taken with respect to this x , we have $u(x,t)$. So, with respect to x we have taken, so this x is replaced by α . And the right hand side this d over dt will remain as it is, and the Fourier transform of this u will be \hat{u} and αt .

So, what we get $d \hat{u} / dt + k^2 \hat{u} = 0$; note that this boundary conditions for this problem, we have already used here, because the Fourier transform of this double derivative minus $\alpha^2 \hat{u}$. We have use these 2 boundary conditions. So, the solution of this ODE will be just the characteristic **(())**, here directly we can have this. So, $d \hat{u} / dt$ we can separate the variable, and we will get this L and \hat{u} . So, clear the **the** solution is $e^{-k^2 L}$, and the integral this side will be t , and we have some constant of integration. So, we now know the Fourier transform of the initial condition given $u(x,0)$ is $f(x)$.

So, Fourier **Fourier** transform **transform** of the initial condition, what we will get. So, \hat{u} at $x=0$, and the Fourier transform of $f(x)$. So, \hat{f} at α . Now, we use this condition to get this constant. So, t is equal to 0, we have \hat{f} at α . So, c is simply. So, this implies that our c is \hat{f} at α , and then we apply this.

(Refer Slide Time: 11:45)

$$\Rightarrow \hat{u}(\alpha, t) = \hat{f}(\alpha) e^{-k\alpha^2 t}.$$

Taking Inverse Fourier transform:

$$u(x, t) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(\alpha) e^{-k\alpha^2 t} e^{-i\alpha x} d\alpha$$

Recall: $F\{f * g\} = \sqrt{2\pi} \hat{f}(\alpha) \hat{g}(\alpha)$

Let $e^{-k\alpha^2 t}$ be the F.T. of $g(x)$:

$$g(x) = \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-k\alpha^2 t} e^{-i\alpha x} d\alpha \quad \text{--- (2)}$$

Consider the integral:

$$I = \int_{-\infty}^{\infty} e^{-a x^2 - 2bx} dx$$

$$= \int_{-\infty}^{\infty} e^{-\left(\sqrt{a}x + \frac{b}{\sqrt{a}}\right)^2 + \frac{b^2}{a}} dx$$

So, we get $\hat{u}(\alpha, t)$ is $\hat{f}(\alpha)$, and $e^{-k\alpha^2 t}$. Now, we take the inverse Laplace transform taking the inverse Fourier **sorry** Fourier transform not the Laplace Fourier transform, what we get $u(x, t)$; if we directly apply the definition of the inverse, what we will get $\hat{f}(\alpha) e^{-k\alpha^2 t}$, and $e^{-i\alpha x}$. So, now we note that in this solution, if we leave this **(())**, it is not a close form at all, because in this integral we are using this Fourier transform of f of the initial condition. So, it is better to have a solution which does not have this Fourier transform, we may have the initial, condition because that is given all already, but we should not expect to have this Fourier transform.

So, in order to avoid this, what we do just recall the convolution theorem. So, we had the Fourier transform of $f * g$ was square root 2π , and $\hat{f}(\alpha)$ and $\hat{g}(\alpha)$. So, if we take the Fourier inverse transform here. So, f inverse of **of** this the multiplication of 2 Fourier transform like we have here. So, if we know the inverse Fourier transform of this, then we can get $u(x, t)$ just by the convolution of f and g . So, we let now that $e^{-k\alpha^2 t}$ be the Fourier transform of $g(x)$. So then, by the definition what we have $1/\sqrt{2\pi}$, and minus infinity to plus infinity $e^{-k\alpha^2 t}$, and $e^{-i\alpha x}$.

So, now we need to evaluate this integral. So, for this we consider a simplified form. So, consider the integral I , and then we will come back to this integral again minus infinity to

plus infinity e. So, here we have with respect to alpha. So, this is alpha square, so we take here some constant times alpha square, this constant will be k t in our case; and minus just for the simplicity we take 2 b x and dx. So, here we have x here also we have this alpha. So, what we do now, we try to put this in the square form of a so, minus infinity to plus infinity e minus square root a x, this is the whole square, and $\frac{b}{\sqrt{a}}$ b over square root a well, k was the multiplication of these 2 times; 2 x b with minus so, we have this extra. What extra we have here b square over a with minus, so we have 2 add here b square over a and then d x. So, this e power b square by a, we can take out of this integral.

(Refer Slide Time: 15:37)

The image shows a handwritten derivation on a blue background. It starts with the integral $I = e^{\frac{b^2}{a}} \int_{-\infty}^{\infty} e^{-\left(\sqrt{a}x + \frac{b}{\sqrt{a}}\right)^2} dx$. A substitution is made: $\sqrt{a}x + \frac{b}{\sqrt{a}} = t \Rightarrow dx = \frac{dt}{\sqrt{a}}$. This leads to $I = e^{\frac{b^2}{a}} \int_{-\infty}^{\infty} e^{-t^2} \cdot \frac{dt}{\sqrt{a}} = e^{\frac{b^2}{a}} \cdot \frac{1}{\sqrt{a}} \cdot \sqrt{\pi}$. Then, the integral $\int_{-\infty}^{\infty} e^{-ax^2 - 2bx} dx = \sqrt{\frac{\pi}{a}} e^{\frac{b^2}{a}}$ is shown. Next, it identifies $a = kt$ and $b = ix/2$. This gives $\int_{-\infty}^{\infty} e^{-ktx^2 - ix x} dx = \sqrt{\frac{\pi}{kt}} e^{-\frac{x^2}{4kt}}$. Finally, the function $g(x) = \frac{1}{\sqrt{it}} \cdot \frac{\sqrt{\pi}}{\sqrt{kt}} e^{-\frac{x^2}{4kt}} = \frac{1}{\sqrt{2kt}} e^{-\frac{x^2}{4kt}}$ is derived. Logos for NPTEL and CCEET are visible in the bottom left and top right corners respectively.

So, we have this integral e b square over a minus infinity to plus infinity minus square root a x plus b over square root a square and dx. Now, we substitute this new variable a x plus b over square root a to t. So, that we have dx is equal to d t over square root a, and this implies that I is e b square over a, and we have minus infinity to plus infinity e minus this is t. So, we have t square, and dx is dt over square root a. So, we have then this square root a, we can take out of this integral. So, e b square over a, and 1 over square root a, and this integral minus infinity to plus infinity minus t square dt that is the Gaussian standard integral, we have a square root pi the value. So, we got this integral we had minus infinity to plus infinity e minus a x square minus 2 b x dx is equal to pi over a with square root, and e b square over a. So, if we let now, because you want to go back to this integral.

So, we will choose our a is $k t$, and this $2 b$ is $i x$. So, a as $k t$, and this $2 b$ are $b i x$ by 2 . So, we will get the our required integral, and we change this integral variable to α . So, e minus $k t$ α square for this x square, and minus $2 b$; so, b is $2 b$ is $i x$. So, minus $i x$, and this x we replace to α . So, $d \alpha$ and we have π over a is $k t$, and $e b$ square i square x square by four. So, here we have then b square. So, that is minus x square over π by a is $k t$. So, $a b$ square over a is $k t$. So, we have $k t$, and this 4 comes from here. So, we have minus x square over $4 k t$ is square root π over $k t$. So, now we go back to this integral $g(x)$ to equation 2. So, this we have evaluated, and 1 over square root 2π will come.

So, our $g(x)$ is 1 over square root 2π , and we have a square root π and square root this $k t$, and e minus x square over $4 k t$. So, we simplify this is square root π will cancel out and we have a square root $2 k t$ from here, and we have e minus x square $4 k t$.

(Refer Slide Time: 19:32)

Using convolution theorem:

$$u(x,t) = \frac{1}{\sqrt{2\pi}} \cdot [f(x) * g(x)]$$

$$= \frac{1}{\sqrt{2\pi}} \cdot \int_{-\infty}^{\infty} f(\beta) g(x-\beta) d\beta$$

$$= \frac{1}{\sqrt{2\pi}} \cdot \int_{-\infty}^{\infty} f(\beta) \cdot \frac{1}{\sqrt{2kt}} e^{-\frac{(x-\beta)^2}{4kt}} d\beta$$

Subst. $z = \frac{-(x-\beta)}{\sqrt{4kt}} \Rightarrow dz = \frac{d\beta}{\sqrt{4kt}}$

$$h(x,t) = \frac{1}{\sqrt{\pi}} \int_{-\infty}^{\infty} f\left(x + \sqrt{4kt} z\right) \cdot e^{-z^2} dz$$

And now we can use the convolution theorem, because we had this u α t is f hat α , this the inverse transform we know, and now for this also we know that $g(x)$ is the inverse transform of this. So, use the convolution theorem. So, using convolution theorem, we get $u(x,t)$ is 1 over a square root 2π and this $f(x)$ convolution with $g(x)$. So, this is 1 over square root $2, \pi$ and we have this convolution minus infinity to plus infinity f β , and $g(x)$ minus β $d \beta$. So, 1 over square root 2π , and minus infinity to plus infinity, we have f β , and g β is given 1 over $2 k t$ e minus x square.

So, x is now, x minus β over $4kt$, and we have $d\beta$. Again let us simplification we can made, if we take z is equal to β minus x minus β over $4kt$. So, this 1 here. So, that we have this square again. So, this tz will be $d\beta$ over $4kt$. So, square root 2 **square root 2**, we have this square root $4kt$, and this $d\beta$ will be. So, we have 1 over square root π $u(x,t)$, because this square root 2, we can have with this is square root 2. So, that we have exactly this term, and minus infinity to plus infinity f for the β we have to get from here; that will be x plus square root $4kt$ and the z , and we have e minus z square $d\beta$. So, this is a solution of the **of the** problem, and we have in terms of the given function f .

(Refer Slide Time: 21:45)

Ex: $k \frac{\partial^2 u}{\partial x^2} = \frac{\partial u}{\partial t} \quad 0 < x < \infty, \quad t > 0.$

BCs: $u(0,t) = u_0, \quad t \geq 0$

IC: $u(x,0) = 0 \quad 0 < x < \infty.$

u & $\frac{\partial u}{\partial x}$ both tend to zero as $x \rightarrow \infty.$

Sol: Since u is specified at $x=0$ and $0 < x < \infty$, the Fourier sine transform is applicable to this problem.

Take Fourier sine transform

$$\Rightarrow k \alpha \sqrt{\frac{2}{\pi}} u_0 - k \alpha^2 \hat{u}_1(x,t) = \frac{d}{dt} \hat{u}_1(x,t)$$

$$\Rightarrow \frac{d \hat{u}_1}{dt} + k \alpha^2 \hat{u}_1(x,t) = \sqrt{\frac{\pi}{2}} k u_0.$$

So now, we take another problem, where we will apply the Fourier sine **(())** sine transform. So, we have the problem $k \frac{\partial^2 u}{\partial x^2}$ is equal to $\frac{\partial u}{\partial t}$, and x is between 0 and infinity, t is positive. So, the boundary conditions are given $u(0,t)$ is u_0 for t positive greater than equal to 0, and the initial conditions are given $u(x,0)$ is 0. And also that information $\frac{\partial u}{\partial x}$ and u both tend to 0, as x approaches to infinity. So, now the solution, and now we note that this u is specified at x is equal to 0. So, we have this boundary condition u is given at 0, and our range for the x is 0 to infinity. So, we have twice for Cosine or sine transform, but this u is given. So, let us have a look again for these properties. So, if we have the Cosine transform, then we need here $f'(0)$, but this information is not given in the problem for the **for the** sine transform we need $f(0)$.

So, this is given, so we will apply the sine transform the twice is very clear. So, let me just also write, since u is specified at x is equal to 0 and x is between 0 to infinity, the Fourier sine transform is applicable to this problem. So, we take the Fourier transform, taking Fourier transform - Fourier sine transform **sorry** Fourier sine transform, what we get. So, this k is there and for the $\frac{\partial^2 u}{\partial x^2}$ for the second derivative sine transform square root 2π $\alpha f(0)$, and minus α^2 Fourier sine transform. So, we have α , and a square root 2 over π the function value at 0, that is this is $u(0)$, it is given minus we have α^2 , and k is there already; so, $k\alpha^2$, and the Fourier sine transform of u .

So, Fourier sin transform of u , and the right hand side we have with respect to t . So, this will remain as it is this differentiation, and we have the sin transform of u . So, now what we have you a set over dt plus $k\alpha^2 u$ s head αt , and this is square root 2 over π , and we have $k\alpha$ and u naught. Now the, now we need to solve that equations, so far that what we get **we get** the integrating factor.

(Refer Slide Time: 25:58)

Handwritten mathematical derivation on a blue background:

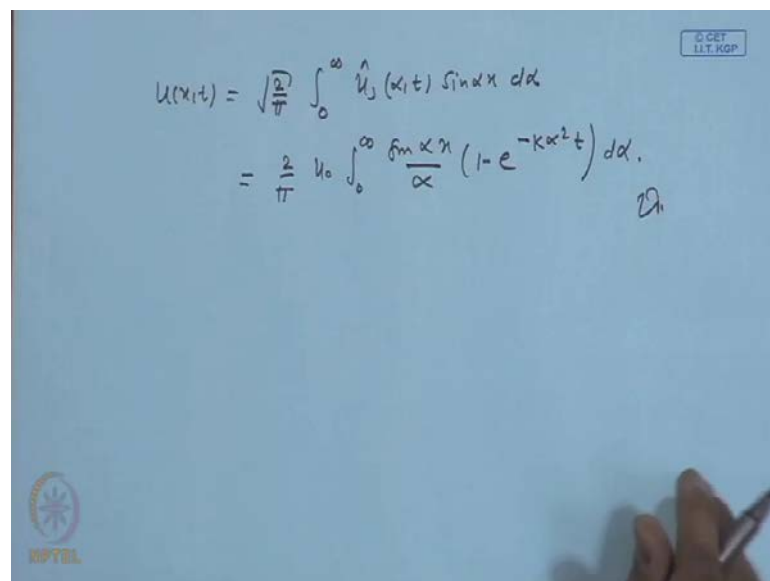
$$\begin{aligned} \text{I.F.} &= e^{k\alpha^2 t} \\ \hat{u}_s \cdot e^{k\alpha^2 t} &= \int \frac{\sqrt{2}}{\pi} k\alpha u_0 \cdot e^{k\alpha^2 t} dt + C \\ \Rightarrow \hat{u}_s &= \frac{\sqrt{2}}{\pi} \cdot \frac{1}{\alpha} \cdot u_0 \cdot e^{k\alpha^2 t} \cdot e^{-k\alpha^2 t} + C e^{-k\alpha^2 t} \\ \hat{u}_s &= \frac{\sqrt{2}}{\pi} \cdot \frac{u_0}{\alpha} + \frac{C \cdot e^{-k\alpha^2 t}}{e^{-k\alpha^2 t}} \\ \text{I.C: } u(x,0) &= 0 \Rightarrow \hat{u}_s(\alpha,0) = 0 \\ \Rightarrow C_0 &= -\frac{\sqrt{2}}{\pi} \frac{u_0}{\alpha} \\ \hat{u}_s(x,t) &= \frac{\sqrt{2}}{\pi} \frac{u_0}{\alpha} (1 - e^{-k\alpha^2 t}) \end{aligned}$$

So, this is the linear equation – linear ordinary differential equation. So, integrating factor will be simply a $k\alpha^2 t$, and then the solution we can get now $e^{k\alpha^2 t}$, and the right hand side we have $\frac{2}{\pi} k\alpha u$ naught, and $e^{k\alpha^2 t} dt$, and plus a constant. Over we have u s hat, this we can take to the right hand side. So, $k\alpha^2 t$, if we integrate this with respect to t this is any way a constant. So, we

will get $e^{-k\alpha^2 t}$ over $k\alpha^2$. So, this $k\alpha^2$ will be cancelled, and we will get $1/\alpha$. So, would we have $2/\pi$, we will get $1/\alpha$ also u_0 , and this $e^{-k\alpha^2 t}$; and from this side, we have with minus $k\alpha^2 t$, and $e^{-k\alpha^2 t}$, and $e^{-k\alpha^2 t}$.

So, this u_0 is $2/\pi$, and we have u_0/α plus $e^{-k\alpha^2 t}$. Now, the initial condition to get these this constant, we have $u(x,0)$ is 0 that is given, if we take the Fourier transform with respect to x , Fourier sine transform. So, we will get $\alpha^2 u_0 = 0$. So, with this condition if we set here this $t = 0$; so, this will be 1, and we have C/α . So, this implies that C is minus $2/\pi$, and u_0/α . So, we have this u_0/α square root $2/\pi$ u_0/α , we take common, because C has also having this factors, we have $1 - e^{-k\alpha^2 t}$.

(Refer Slide Time: 28:30)



$$u(x,t) = \sqrt{\frac{2}{\pi}} \int_0^\infty \hat{u}_s(\alpha,t) \sin \alpha x \, d\alpha$$

$$= \frac{2}{\pi} u_0 \int_0^\infty \frac{\sin \alpha x}{\alpha} (1 - e^{-k\alpha^2 t}) \, d\alpha$$

So, we have this and now we take the inverse - sine inverse transform. So, we will get straightaway this $u(x,t)$, and that is $2/\pi$, and we have 0 to infinity u_0/α sine αx $d\alpha$. So, we can substitute that the $2/\pi$, and also u_0 will come 0 to infinity we have sine αx over α and, we have $1 - e^{-k\alpha^2 t}$ and $d\alpha$, so this is the solution.

(Refer Slide Time: 29:16)

Solve: $\frac{\partial u}{\partial t} = k \frac{\partial^2 u}{\partial x^2}$

$u(x,0) = 0 \quad x \geq 0$

$u_x(0,t) = -\mu; \quad t > 0$

$u \& \frac{\partial u}{\partial x} \rightarrow 0 \text{ as } x \rightarrow \infty.$

Sol: $F_c \left\{ \frac{\partial u}{\partial t} \right\} = k F_c \left\{ \frac{\partial^2 u}{\partial x^2} \right\}$

$\Rightarrow \frac{d}{dt} \hat{u}_c(\alpha, t) = k \left[-\sqrt{\frac{2}{\pi}} u_x(0, t) - \alpha^2 F_c \{ u \} \right]$

$\Rightarrow \frac{d \hat{u}_c}{dt} + k \alpha^2 \hat{u}_c = k \mu \sqrt{\frac{2}{\pi}}.$

IF: $= e^{\int k \alpha^2 dt} = e^{k \alpha^2 t}.$

The next problem, where we will see that we need to apply the Cosine transform. So, next problem solve $\frac{\partial u}{\partial t}$ is $k \frac{\partial^2 u}{\partial x^2}$, and subject to the conditions we have $u(x,0)$ as 0, and for x greater than or equal to 0. And we have u_x the first derivative $\frac{\partial u}{\partial x}$ at $x=0$ as minus μ ; that is given for t positive, and u and again this $\frac{\partial u}{\partial x}$ over $\frac{\partial u}{\partial x}$, this both goes to 0 as, x goes to infinity. So, we have again in the half range. So, we can have possibility to apply Fourier Cosine or Fourier sine transform, but this u_x is given. So, we can apply only Fourier Cosine transform, because we have this first derivative information there.

So, taking this Fourier Cosine transform $\frac{\partial u}{\partial t}$, and k Fourier Cosine transform of $\frac{\partial^2 u}{\partial x^2}$, what we get here, $\frac{d}{dt}$ and this is \hat{u}_c we denote it by this α, t , and we have this k . And now, this Fourier Cosine of this. So, by that formula we have $\frac{2}{\pi}$ and u_x at this 0, t and minus α^2 and Fourier Cosine transform of u . So, we have $\frac{d \hat{u}_c}{dt} + k \alpha^2 \hat{u}_c = k \mu \sqrt{\frac{2}{\pi}}$, and this we will take or this one, because it is a \hat{u}_c only. So, $k \alpha^2$, and u Cosine transform, and here we have k this is given minus μ ; so, minus **minus** will be plus.

So, $k \mu$ and this $\frac{2}{\pi}$, again we solve this linear equation. So, the integrating factor is $e^{k \alpha^2 t}$. So, we have $e^{k \alpha^2 t}$, and then over get the solutions; I am writing the solution directly.

(Refer Slide Time: 31:50)

$$\hat{u}_c = \sqrt{\frac{2}{\pi}} \frac{\mu}{\alpha^2} e^{k\alpha^2 t} + C.$$

IC: $u(x,0) = 0 \Rightarrow \hat{u}_c(\alpha,0) = 0.$

$$\hat{u}_c = \sqrt{\frac{2}{\pi}} \frac{\mu}{\alpha^2} (1 - e^{-k\alpha^2 t}).$$

I.C.T:

$$u(x,t) = \sqrt{\frac{2}{\pi}} \int_0^\infty \hat{u}_c(\alpha,t) \cos \alpha x \, d\alpha$$

$$= \frac{2}{\pi} \mu \int_0^\infty \frac{\cos \alpha x}{\alpha^2} (1 - e^{-k\alpha^2 t}) \, d\alpha.$$

So, that solution will be for that in terms of the u Cosine transform 2 over π μ over α square $e^{k\alpha^2 t}$ plus constant, and again this initial condition is given that $u(x,0)$ is 0 . So, we will get $\hat{u}_c(\alpha,0)$ is 0 , and if we put it here, we will get this c . So, this is 0 , and this will be 1 . So, c will be minus of this. So, then we get this \hat{u}_c is square root 2 over π μ over α square $1 - e^{-k\alpha^2 t}$, and now taking the inverse of sine inverse Cosine transform. So, inverse Cosine transform now to get the $u(x,t)$ that is 2 over π , and 0 to infinity u Cosine αt and $\cos \alpha x \, d\alpha$. So, we substitute here. So, we get 2 over π , we get this μ , and 0 to infinity $\cos \alpha x$ over α square, and we have this $1 - e^{-k\alpha^2 t}$ and $d\alpha$.

So, these were the three applications, where we have use the heat equation in the first was by applying the Fourier transform directly to the problem, in the second we applied the Fourier sine transform depending on the conditions. And in the last example, which we have just an we applied the Cosine transform again depending on the condition.

(Refer Slide Time: 33:44)

Solution of Wave equation:

Ex: $\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2} \quad -\infty < x < \infty$

ICs: $u(x,0) = f(x)$
 $u_t(x,0) = 0$

Bcs: u & $\frac{\partial u}{\partial x}$ both $\rightarrow 0$ as $|x| \rightarrow \infty$

Sol: $\frac{d^2}{dt^2} \hat{u}(\alpha, t) = c^2 (-\alpha^2 \hat{u}(\alpha, t))$

$\Rightarrow \frac{d^2 \hat{u}}{dt^2} + c^2 \alpha^2 \hat{u}(\alpha, t) = 0$

$\Rightarrow \hat{u}(\alpha, t) = C_1 \cos(c\alpha t) + C_2 \sin(c\alpha t)$

So, we go now for the solution of wave equation **solution of wave equation**. So, we take this example that solve the wave **wave** equation $\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$ and x is given minus infinity to plus infinity. So, initial conditions that $u(x,0)$ is $f(x)$ for this x , and $u_t(x,0)$ is 0, and we have boundary conditions that u and $\frac{\partial u}{\partial x}$ both goes to 0 as absolute value of x goes to infinity. So, we take the Fourier transform, because this ranges came from minus infinity to plus infinity. So, we have $\frac{d^2}{dt^2}$ and $\hat{u}(\alpha, t)$ plus $c^2 \alpha^2 \hat{u}(\alpha, t)$ is equal to 0, and for this we have simply minus α^2 , and $\hat{u}(\alpha, t)$.

So, equation is $\frac{d^2 \hat{u}}{dt^2} + c^2 \alpha^2 \hat{u}(\alpha, t) = 0$, and it is general solution one can write $\hat{u}(\alpha, t) = C_1 \cos(c\alpha t) + C_2 \sin(c\alpha t)$ square plus $c^2 \alpha^2$ is equal to 0. So, the roots will be plus **plus** minus this $c\alpha$. So, for that we have the solution $C_1 \cos(c\alpha t) + C_2 \sin(c\alpha t)$.

(Refer Slide Time: 35:50)

$$\begin{aligned} \text{F.T. of I.C.: } u(x,0) &= f(x) \\ \Rightarrow \hat{u}(\alpha,0) &= \hat{f}(\alpha) \\ \text{2. } u_t(x,0) &= 0 \Rightarrow \frac{d\hat{u}(\alpha,t)}{dt}\bigg|_{t=0} = 0 \\ \Rightarrow C_1 &= \hat{f}(\alpha) \\ \frac{d\hat{u}}{dt} &= -C_1 \sin(C\alpha t) (C\alpha) + C_2 \cos(C\alpha t) (C\alpha) \\ 0 &= C_2 \cdot (C\alpha) \Rightarrow C_2 = 0 \\ \hat{u}(\alpha,t) &= \hat{f}(\alpha) \cos(C\alpha t) \\ \text{I.F.T} \end{aligned}$$

So, now, the initial conditions the Fourier transform of initial condition, we have the $u(x,0)$ is $f(x)$ first condition. So, here we get $\hat{u}(\alpha,0)$ is $\hat{f}(\alpha)$, and the second condition we have $u_t(x,0)$ is equal to 0.

So, from here we will get $d\hat{u}(\alpha,t)/dt$ approaches to where this t approaches to 0, and this is again 0. So, from the first condition, so our solution has this $\hat{u}(\alpha,t) = C_1 \cos(C\alpha t) + C_2 \sin(C\alpha t)$. So, when we put t is equal to 0; this is 0, and we have here one. So, we get C_1 straightaway from the first condition, we get C_1 is $\hat{f}(\alpha)$; and from the second condition, so we need to get the derivative first of this. So, $d\hat{u}(\alpha,t)/dt$ is $-C_1 \sin(C\alpha t) (C\alpha) + C_2 \cos(C\alpha t) (C\alpha)$, and we have $C_1 \alpha$ plus C_2 and for sine we have Cosine of $C\alpha t$ and again the derivative of this.

So, C_1 and C_2 α . Now, again if you put this t to 0 in this case this will disappear, this is one. So, this is 0. So, we have C_2 and $C_1 \alpha$ this is 1; so, in any case this C_2 is 0. So, we have $\hat{u}(\alpha,t) = \hat{f}(\alpha) \cos(C\alpha t)$. Now, we take the inverse Fourier transform- **inverse Fourier transform**.

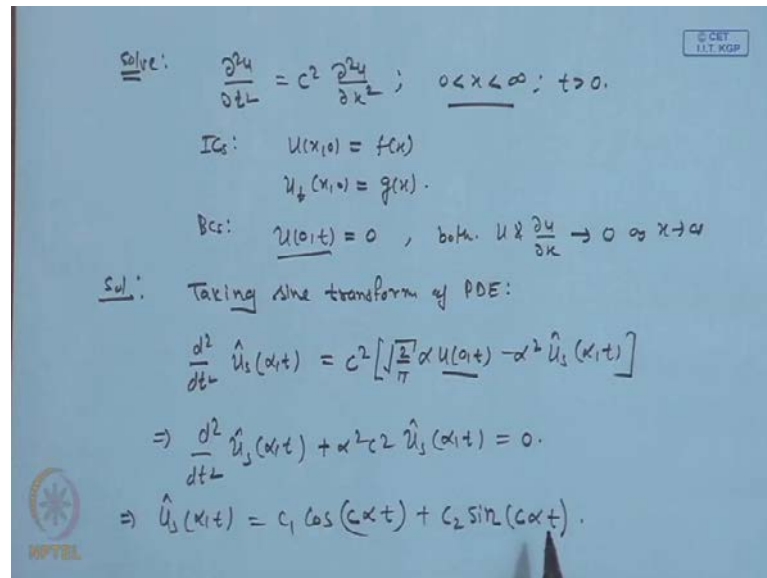
(Refer Slide Time: 37:39)

$$\begin{aligned}
 u(x,t) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(\alpha) \cos(c\alpha t) e^{-i\alpha x} d\alpha \\
 \Rightarrow u(x,t) &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(\alpha) \cdot \frac{(e^{ic\alpha t} + e^{-ic\alpha t})}{2} e^{-i\alpha x} d\alpha \\
 &= \frac{1}{2} \left[\frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(\alpha) e^{-i\alpha(x-ct)} d\alpha + \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} \hat{f}(\alpha) e^{-i\alpha(x+ct)} d\alpha \right] \\
 &= \frac{1}{2} [f(x-ct) + f(x+ct)] \quad \text{Q.E.D.}
 \end{aligned}$$

So, what we get $u(x,t)$ is $\frac{1}{\sqrt{2\pi}}$ times the integral from minus infinity to plus infinity of $\hat{f}(\alpha)$, $\cos c\alpha t$, and then we have $e^{-i\alpha x}$ and $d\alpha$. So, $u(x,t)$ is $\frac{1}{\sqrt{2\pi}}$ times the integral from minus infinity to plus infinity of $\hat{f}(\alpha)$ as it is, this we write in terms of the again exponential function. So, $e^{ic\alpha t} + e^{-ic\alpha t}$ divided by 2, and $e^{-i\alpha x}$ $d\alpha$. So, we have $\frac{1}{2}$, and this $\frac{1}{\sqrt{2\pi}}$, then minus infinity to plus infinity $\hat{f}(\alpha)$, and this **and this** we combine in to 1. So, we have $e^{-i\alpha(x-ct)}$. So, we have $x - ct$ $d\alpha$ plus again this **(())** factor 2π , $\frac{1}{\sqrt{2\pi}}$ times the integral from minus infinity to plus infinity; we have $\hat{f}(\alpha)$, and $e^{-i\alpha(x+ct)}$. So, we have $x + ct$ and $d\alpha$.

So, if we see now that half, and this is the definition of the Fourier inverse. So, we have $e^{-i\alpha(x-ct)}$ instead of x , we have here $x - ct$ here we have $x + ct$. So, we will get $f(x - ct)$ plus $f(x + ct)$, and this is also known as **(())** solution of the wave equation.

(Refer Slide Time: 39:49)



Solve: $\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}; \quad 0 < x < \infty; \quad t > 0.$
 ICs: $u(x, 0) = f(x)$
 $u_t(x, 0) = g(x).$
 Bcs: $u(0, t) = 0$, both u & $\frac{\partial u}{\partial x} \rightarrow 0$ as $x \rightarrow \infty$
 Sol: Taking sine transform of PDE:
 $\frac{d^2}{dt^2} \hat{u}_s(\alpha, t) = c^2 \left[\sqrt{\frac{2}{\pi}} \alpha u(0, t) - \alpha^2 \hat{u}_s(\alpha, t) \right]$
 $\Rightarrow \frac{d^2}{dt^2} \hat{u}_s(\alpha, t) + \alpha^2 c^2 \hat{u}_s(\alpha, t) = 0.$
 $\Rightarrow \hat{u}_s(\alpha, t) = c_1 \cos(c\alpha t) + c_2 \sin(c\alpha t).$

So, just one more example for the wave equation, then we will solve the Laplace equation. So, we have $\frac{\partial^2 u}{\partial t^2} = c^2 \frac{\partial^2 u}{\partial x^2}$, and this is given for t positive. So, in the half range we have this problem. So, initial conditions are $u(x, 0)$, $f(x)$, and $u_t(x, 0)$ is $g(x)$, and boundary conditions are $u(0, t)$. So, u is given as 0, and again the both u and $\frac{\partial u}{\partial x}$, they goes to 0 as x goes to infinity. So, now since the value of u is given as for the formula we will apply the sine transform, because this function value is given at 0.

So, we take the taking sine transform of the PDE. So, we get $\frac{d^2}{dt^2}$ over $\hat{u}_s(\alpha, t)$, and sine transform of this u will denote by $\hat{u}_s(\alpha, t)$, we have c^2 ; and then by the derivative theorem we have $\sqrt{\frac{2}{\pi}} \alpha u(0, t)$; and minus $\alpha^2 \hat{u}_s(\alpha, t)$. So, this is $\frac{d^2}{dt^2} \hat{u}_s(\alpha, t)$, and $u(0, t)$ is given 0. So, this term will disappear, and this will come to the left hand side; this is $\alpha^2 c^2$, and we have $\hat{u}_s(\alpha, t)$ is equal to 0. So, now again we can find general solution, and that will be given by $c_1 \cos(c\alpha t) + c_2 \sin(c\alpha t)$. So, we use this initial conditions now to get these constants.

(Refer Slide Time: 42:31)

At $t=0$: $\hat{u}_s(x,0) = \hat{f}_s(x)$ and $\frac{d}{dt} \hat{u}_s(x,0) = \hat{g}_s(x)$

$\Rightarrow \boxed{c_1 = \hat{f}_s(x)}$ and $\frac{d\hat{u}_s}{dt} = -c_1 \sin(cx) + c_2 \cos(cx)$

$\Rightarrow \hat{g}_s(x) = c_2 \cos(cx)$

$\hat{u}_s(x,t) = \hat{f}_s(x) \cos(cx) + \frac{\hat{g}_s(x)}{cx} \sin(cx)$

Taking I.S.T.

$u(x,t) = \sqrt{\frac{2}{\pi}} \int_0^\infty \hat{u}_s(x,t) \sin \alpha x d\alpha$

So, what we have the first. So, at t is equal to 0, the initial condition we have first u s alpha 0 is \hat{f} alpha, because this was the initial condition. Now, $u(x,0)$ is equal to $f(x)$. So, we take the Fourier sine transform. So, we get u hat s alpha 0, and the Fourier transform of this f . And from the second one we have the derivatives, so d over $d t$ and u s hat alpha 0 is g s hat alpha. So, from the first initial condition, when we put t is equal to 0. In this, so this is 0. So, we have c_1 ; this implies c_1 is this given \hat{f} s hat alpha. So, this is one; for the second one we need to get the derivative of this first, and then we will go to the t is equal to 0.

So, we have $d u$ hat as over $d t$ is minus c_1 , and \cos will be $\sin c$ alpha t , and then we have c alpha plus c_2 and \cos **Cos** c alpha t , and we have C alpha. So, we put t is equal to 0; this **(())** disappear, this is 1 and this is g hat s alpha is equal to c_2 , and we have c alpha. So, c_2 is g s hat alpha over c alpha, and c_1 is this. So, we can get now the solution u s alpha t that is \hat{f} s alpha $\cos c$ alpha t , and plus we have g s hat alpha over the c alpha, and then $\sin c$ alpha t .

Now, we take the inverse taking inverse sine transform we get $u(x,t)$ will be square root 2 over pi, we have 0 to infinity and this function here. So, we have u s hat alpha t and \sin alpha $x d$ alpha. So, this is our u s alpha t , and then multiplied by \sin alpha x . So, we write this now.

(Refer Slide Time: 45:28)

$$\begin{aligned}
 u(x,t) &= \sqrt{\frac{2}{\pi}} \int_0^{\infty} \left[\frac{\hat{f}_s(\alpha)}{2} \cos(c\alpha t) \sin \alpha x + \frac{\hat{g}(\alpha)}{c\alpha} \sin(c\alpha t) \sin \alpha x \right] d\alpha \\
 &= \sqrt{\frac{2}{\pi}} \int_0^{\infty} \frac{\hat{f}_s(\alpha)}{2} \left\{ \sin(\alpha + c\alpha t) + \sin(\alpha - c\alpha t) \right\} d\alpha \\
 &\quad + \sqrt{\frac{2}{\pi}} \int_0^{\infty} \frac{\hat{g}_s(\alpha)}{2c\alpha} \left\{ \cos(\alpha - c\alpha t) - \cos(\alpha + c\alpha t) \right\} d\alpha
 \end{aligned}$$

Consider $g(u) = \sqrt{\frac{2}{\pi}} \int_0^{\infty} \hat{g}_s(\alpha) \sin \alpha u d\alpha$

$$\Rightarrow \int_{x-ct}^{x+ct} g(u) du = \sqrt{\frac{2}{\pi}} \int_0^{\infty} \hat{g}_s(\alpha) \int_{x-ct}^{x+ct} \sin \alpha u d\alpha$$

So, we have $u(x,t)$ as the square root 2 over pi, we have 0 to infinity, we have \hat{f}_s alpha and $\cos c$ alpha t sine alpha x , in one term plus we have \hat{g} alpha over this c alpha, and sine c alpha t and sine alpha x , then d alpha. So, now we expand this, because the two sine alpha, and \cos beta form. So, we take 2 over pi and we have 0 to infinity, and this $\hat{f}(x)$ hat over 2, and this will be now 2 times sine a $\cos b$ will be sine a plus x plus c t alpha plus sine x minus c t alpha d alpha plus; again here 2 over pi with that integral, we have 0 to infinity \hat{g} hat alpha 2 c alpha. So, 2 sine a sine b we have the \cos terms then $\cos a$ minus b our b minus a does not matter. So, x minus c t we write, and then minus $\cos a$ plus b ; so, x plus c t alpha d alpha.

Now for the first round, we can easily write in terms of f taking inverse sine transform, because it is given exactly in that form of a hat alpha, and the sine and again this other term with this sine we will get is straight away \hat{f} alpha; and here also this \hat{f} **sorry** $\hat{f}(x)$ plus c t and this will x minus c t divide by 2, but for this term this alpha is appearing here. So, in this case what we take consider a function $g(u)$. So, $g(u)$ is the Fourier transform of inverse Fourier transform of this \hat{g} hat alpha. So, we have 0 to infinity Fourier sine transform. So, \hat{g} s hat alpha, this is whole Fourier sine transform, and then we have sine alpha u d alpha, and if we integrate this form x minus c t to x plus c t $g(u)$ du 2 over pi, and we also **integrate** change the order of integration.

So, first 0 to infinity will come \hat{g} s alpha, and then this x minus c t to x plus c t will come and sine alpha u d alpha.

(Refer Slide Time: 48:51)

$$\int_{x-ct}^{x+ct} g(u) du = \sqrt{\frac{2}{\pi}} \int_0^{\infty} \frac{\hat{g}(s)}{\alpha} \{ \cos(x-ct)\alpha - \cos(x+ct)\alpha \} d\alpha.$$

$$u(x,t) = \frac{1}{2} (f(x+ct) + f(x-ct)) + \frac{1}{2c} \int_{x-ct}^{x+ct} g(u) du.$$

Ex: Solution of Laplace equation:

Solve $u_{xx} + u_{yy} = 0 \quad -\infty < x < \infty \quad y > 0$

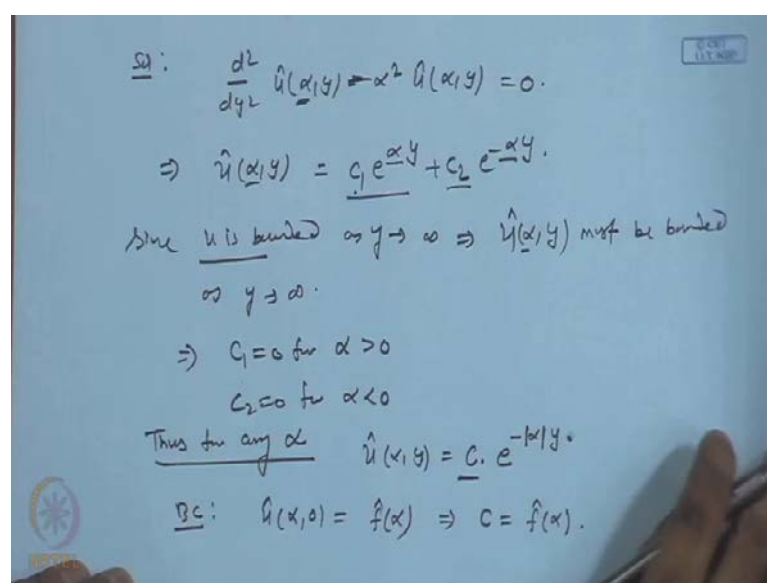
Des $u(x,0) = f(x)$

u is bounded as $y \rightarrow \infty$ $u, \frac{\partial u}{\partial x} \rightarrow 0$ as $|x| \rightarrow \infty$

So, if we integrate this we will get \cos alpha u over alpha, and then put the upper limit, and the lower limits, so what we get in this case that this integral x minus c t to x plus c t $g(u)$ du is 2 over π , we have 0 to infinity \hat{g} s alpha over this alpha, we are getting exactly the term we need there. So, $\cos x$ minus c t alpha minus $\cos x$ plus c t alpha, and then we have d alpha; thus the solution is given by. So, if we just... Now, go back to this, **this** we know what will be the **the** inverse what will be this, and here also we know now from this integral that is equal to that. So, we take with get the solution $u(x,t)$ is half **half** here, and square root 2 over π 0 to infinity f s hat alpha sine x plus c t , we will get the $f(x)$ plus c t .

And then for the second one we have x minus c t $f(x)$ minus c t , and for the second term what we just write it will be just x minus c t over x plus c t , and $g(u)$ du . So, this is the solution. Now, we go for the last example of this lecture, and as a solution of Laplace equation, **Laplace equation**. So, in this case, we solved $u_{xx} + u_{yy}$ is equal to 0 is given from minus infinity to plus infinity by positive, and the boundary conditions are $u(x,0)$ is $f(x)$. Again in the same range, and this is given that u is bounded **is bounded** as y approaches to infinity, and again u and $\frac{\partial u}{\partial x}$ both goes to 0 as $|x|$ goes to infinity.

(Refer Slide Time: 51:22)



$$\text{So: } \frac{d^2}{dy^2} \hat{u}(\alpha, y) - \alpha^2 \hat{u}(\alpha, y) = 0.$$

$$\Rightarrow \hat{u}(\alpha, y) = c_1 e^{\alpha y} + c_2 e^{-\alpha y}.$$

Since u is bounded as $y \rightarrow \infty \Rightarrow \hat{u}(\alpha, y)$ must be bounded as $y \rightarrow \infty$.

$$\Rightarrow c_1 = 0 \text{ for } \alpha > 0$$

$$c_2 = 0 \text{ for } \alpha < 0$$

Thus for any α $\hat{u}(\alpha, y) = c e^{-|\alpha| y}$.

BC: $\hat{u}(\alpha, 0) = \hat{f}(\alpha) \Rightarrow c = \hat{f}(\alpha).$

So, we take the Fourier transform both the side we will get this transformed ordinary differential equation, and then I will directly write the **the** after taking the transform Fourier transform, we will d^2 over dy^2 ; you get αy , and we get minus α square you get αy is equal to 0, and its solution we have $\alpha y c_1$. So, with respect to x we have taking again this Fourier transform: $e^{\alpha y}$ plus $c_2 e^{-\alpha y}$, and the condition is given that u is bounded **u is bounded** as y approaches to infinity, and this will imply straightaway that this it is Fourier transform.

And that we take with respect to x naught with respect to y must be also bounded as y approaches to infinity, and this will tell us because here we have this αy minus αy that c_1 is 0 for α positive, because otherwise this will flow up. So, to have this boundedness, and this c_2 will be 0 for α negative in that case this will beyond bounded. So, for this boundedness. So, in any case thus for any α , what we can write we can eliminate one constant there, because we are anyway getting this αy is some constant times αy .

The **the** absolute value of αy , we cannot get this α , because this c_1 will be 0; if α is positive, and if α is negative this c_2 will be 0. So, what we have some constant, and e power minus αy . Now, with the boundary condition we have u hat $\alpha 0$ is $\hat{f}(\alpha)$, this is given the boundary condition we have taken the Fourier transforms as $y(0)$, we will get this c . So, c will be $\hat{f}(\alpha)$.

(Refer Slide Time: 53:41)

$$\begin{aligned}
 \hat{u}(x, y) &= \hat{f}(\alpha) e^{-|\alpha|y} \\
 \text{Let } g(x) &= \mathcal{F}^{-1}\{e^{-|\alpha|y}\} \\
 &= \frac{1}{\sqrt{2\pi}} \int_{-\infty}^{\infty} e^{-|\alpha|y} \underline{e^{-i\alpha x}} d\alpha \\
 &\quad \text{even} \\
 &= \frac{2}{\sqrt{2\pi}} \int_0^{\infty} e^{-\alpha y} \cos \alpha x d\alpha \\
 &= \sqrt{\frac{2}{\pi}} \int_0^{\infty} e^{-\alpha y} \cos \alpha x d\alpha
 \end{aligned}$$

And now our solution of the transform over the e has $\hat{f}(\alpha)$, and $e^{-|\alpha|y}$. And now, we again apply the same trick to get this inverse this convolution theorem, but we need to get the function before here inverse of **of** this. So, let that $g(x)$ is the Fourier inverse of $e^{-|\alpha|y}$, and then we take as per the definition $\frac{1}{\sqrt{2\pi}}$ from minus infinity to plus infinity $e^{-|\alpha|y} e^{-i\alpha x} d\alpha$. And then, this we write $\cos \alpha x$ and $i \sin \alpha x$. So, since this is the even function, and here \cos is also even function, and with these sine odd function; so, this over the symmetric integral will be 0. So, we have 2 times $\frac{1}{\sqrt{2\pi}}$, and 0 to infinity $e^{-\alpha y}$ and $\cos \alpha x$ will remain.

And now this, we can have $\frac{2}{\sqrt{2\pi}}$, because this will be square 2 here, and we have 0 to infinity $e^{-\alpha y} \cos \alpha x$. And this is a very simple integration, one can just get integrating by parts.

(Refer Slide Time: 55:15)

$$g(x) = \left(\frac{1}{\sqrt{\pi}} \right) \left(\frac{y}{x^2+y^2} \right)$$

$$\Rightarrow u(x,y) = F^{-1} \left\{ \underbrace{\hat{f}(\alpha)}_{\hat{g}(\alpha)} \cdot \underbrace{e^{-|\alpha|y}}_{\hat{g}(\alpha)} \right\}$$

$$= \frac{1}{\sqrt{2\pi}} f * g$$

$$= \frac{1}{\sqrt{2\pi}} \cdot \sqrt{\frac{2}{\pi}} \cdot \int_{-\infty}^{\infty} f(\beta) \cdot \frac{y}{(x-\beta)^2+y^2} d\beta$$

$$\Rightarrow u(x,y) = \frac{1}{\pi} \int_{-\infty}^{\infty} f(\beta) \cdot \frac{y}{(x-\beta)^2+y^2} d\beta$$

So, we will get in this case $g(x)$ as 2 over π , and this will give us y over x square plus y square. So, this integral will give y over x square plus y square, and now we go back to the solution now, and that is the Fourier inverse of this f at α and e minus α y .

So, this was the $\hat{g}(\alpha)$, this is $\hat{f}(\alpha)$. So, by the convolution theorem, we have 1 over square root 2π , and we have f on the convolution of this g . So, 1 over a square root 2π , and this convolution we can write down. So, we have 2 over π , and because this $\hat{g}(\alpha)$ we have this vectors. So, we have taken this, and minus infinity to plus infinity $f(\beta)$ introduce this integrating variable, and then we have y over x , we replace this x by x minus β for the convolution x minus β x square plus y square, and we have this $d\beta$. So, what we get $u(x,y)$ is 1 over π square root, two will cancel this square root π square root π you have 1 over π minus infinity to plus infinity.

We have $f(\beta) y$ over x minus β whole square plus ϕ square $d\beta$, and this solution is also known as the Poisson integral formula. So, that is the here we **we** complete the discussion on Fourier transform. As well as on this rather introductory lectures on transform calculus, and we have mainly discuss the Fourier transform, and the Laplace transform. And each topic was implemented by **by** various well chosen exercises. So, I hope that this lecture which was other introductory of course, will have to understand advance topics in related areas. Thank you.