Advanced Partial Differential Equations Professor Doctor Kaushik Bal Department of Mathematics and Statistics Indian Institute of Technology Kanpur Lecture: 28 Wave Equation in n=2 K+1, K=1, 2,…

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u_{tt} = \Delta u = 0
$$
 in $\mathbb{R}^{N} \times (b, r^{\omega})$
\n14 $u_{tt} = \Delta u = 0$ in $\mathbb{R}^{N} \times (b, r^{\omega})$
\n14 $u_{tt} = \Delta u = 0$ in $\mathbb{R}^{N} \times (b, r^{\omega})$
\n15 $u_{tt} = \frac{1}{2} \times 0$; $u_{tt} = \frac{1}{2} \times 0$
\n16 $u_{tt} = 1$; $\frac{1}{2} \times 0$
\n17 $u_{tt} = 1$; $\frac{1}{2} \times 0$
\n18 $u_{tt} = \frac{1}{2} \times 0$
\n19 $u_{tt} = 1$
\n10 $u_{tt} = 1$
\n11 $u_{tt} = \Delta u = 0$ for $\mathbb{R}^{N} \times (b, r^{\omega})$
\n10 $u_{tt} = 1$
\n11 $u_{tt} = \Delta u = 0$ for $\mathbb{R}^{N} \times (b, r^{\omega})$
\n12 $u_{tt} = \frac{1}{2} \times 0$ for $\mathbb{R}^{N} \times (b, r^{\omega})$
\n13 $u_{tt} = \frac{1}{2} \times 0$ for $\mathbb{R}^{N} \times (b, r^{\omega})$
\n14 $u_{tt} = \Delta u = 0$ for $\mathbb{R}^{N} \times (b, r^{\omega})$
\n15 $u_{tt} = 1$ for $n = 2$ for $n = 3$; Kirchapffs Fr $m = 4$
\n16 $u_{tt} = 1$ for $n = 3$ for $n = 3$

Welcome students, in this week's video, we are going to talk about the norm, the Kirchoff's formula, but for higher dimensions. So, as we have seen that once you can solve the homogeneous problem, you can use the explicit formula to construct a solution for the homogeneous problem via the Duhamel's principle.

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4.4 A
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0 \t{d \t{f(t)} \t{g(t)} \t{
$$

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\left\langle \cdot\right\rangle +\left\langle \cdot\right\rangle
$$

$\begin{array}{cccccccccccccc} \mathbb{Q}^{\dagger} & \mathbb{Q} & \mathbb{Q} & \mathbb{T} & \mathbb{C} & \mathbb{C} & \mathbb{R} \end{array}$

$$
\sigma \leq \sigma
$$

4
$$
u_{t} - \Delta u = 0
$$
 in $\mathbb{R}^{n} \times (0, \infty)$
\n $u(x_{0}) = g(x)$; $u_{\epsilon}(x_{0}) = h(x)$
\n $u(x_{0}) = g(x)$; $u_{\epsilon}(x_{0}) = h(x)$
\n $lim_{t \to 0} g(x) = \frac{h(x)}{h(x)} = \frac{h(x)}{h(x)}$
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5 \text{ b } 0 0 0 T \text{ } 0 0
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6 \text{ c } 1 0
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5 \text{ b } 0 0 0 0 T \text{ } 0 0
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$$
6 \text{ d } 0
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7 \text{ d } 0
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10 \text{ c } 0
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 $\mathbb{R}^{n\times n}$

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V(t) = \lim_{r \to 0} \frac{1}{p_{\epsilon}} \int_{0}^{r} \frac{\tilde{G}(t+v) - \tilde{G}(t+v)}{2r} + \frac{1}{2v} \int_{0}^{r+1} \tilde{H}(v) \, dv \Big|
$$

\n
$$
= \frac{1}{p_{\epsilon}} \int_{0}^{r} \tilde{G}'(t) + \tilde{H}(t) \Big|_{0}^{r}
$$

\n
$$
= \frac{1}{p_{\epsilon}} \int_{0}^{r} \tilde{G}'(t) + \tilde{H}(t) \Big|_{0}^{r}
$$

\n
$$
\therefore \text{ For } n = 2k + 1,
$$

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$$
U(x_{1}b) = \frac{1}{x_{10}} \left[\left(\frac{y}{bt} \right) \left(\frac{1}{6} \frac{3}{3t} \right)^{\frac{n-3}{2}} \left(t^{n-2} \int_{0}^{t} q \, ds \right) + \left(\frac{1}{t} \frac{y}{5t} \right)^{\frac{n-3}{2}} \left(t^{n-2} \int_{0}^{t} h \, ds \right) \Big|_{0}^{r}
$$

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$$
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$$
U(x_{1}b) = \frac{1}{x_{10}} \left[\left(\frac{y}{
$$

 $\begin{array}{ccccccccccccccccc} \mathbb{Q}^{\dagger} & \mathbb{Q} & \mathbb{Q} & \mathbb{Q} & \mathbb{T} & \mathbb{C} & \mathbf{0} & \mathbf{0} \end{array}$

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7. For
$$
0 \leq t \leq \frac{1}{2}
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\n
$$
\int_{0}^{2} (r_{1}b) = \frac{1}{2} \left[\frac{r_{0}}{6}(rt^{2}) - \frac{r_{1}}{6}(t^{2}) \right] + \frac{1}{2} \int_{0}^{2} \frac{kr_{1}}{4}(y) dy \quad (b')^{2}Atmbert \quad \text{(16.1)} \times \frac{1}{2} \text{ and } \frac{1}{2} \text{ (16.1)} \times \frac{1}{2} \text{ (16.1)} \times
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x + 5 = 0
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y = 0
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y = 1
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$$
y = 0
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$$
y = 1
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$$
y = 0
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So, for that the existence theorem for wave equation in odd dimensions:so basically we want to show that the formula which you obtained.

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This is our Kirchoff's formula for n greater than equal 3. And this formula solves our wave equation. So let us end it here.