Advanced Partial Differential Equations Professor Doctor Kaushik Bal Department of Mathematics and Statistics Indian Institute of Technology Kanpur Lecture 23 Maximum Principle for Heat Equation

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Proparhies of Johnston to the heat equation:
(
$$\Omega$$
 is open, smooth or bounded)
 $U_t - (U = 0 \text{ in } D_T = \Omega \times (0, T] - 0)$
The:
(Strong Maximum Principle for the heat equation)
(Ω det $u \in C_T^2(\Omega T) \cap C(\overline{\Omega T})$ Solves 0 . Then max $u = \max u$ where T_t is the parabolic boundary.
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($\overline{\Omega}$ det $u \in C_T^2(\Omega T) \cap C(\overline{\Omega T})$ Solves 0 . Then max $u = \max u$ where T_t is boundary.
($\overline{\Omega}$ the solution in $\overline{\Omega}_{tot} = \Omega \times (\infty t_0) \in \Omega T$ such that
($\overline{\Omega}$ then, u is constant in $\overline{\Omega}_{tot} = \Omega \times (\infty t_0)$
Remarks: for t_T to, the solutions may change?
Since the boundary conditions change?

Welcome, students. In today's class or in this video, essentially, we are going to talk about the consequences of mean value property for the heat operator. So, essentially, what we are going to do is look at some properties of the heat equation based on the mean value theorem.

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Proparties of Jolution to the heat equation:- (2 is open, smooth or bounded)
Ut - QU = Q in
$$\Omega_T = \Omega \times (0, T] - 0$$

Th :- (Strong Maximum Principle for the heat equation)
() dut $u \in C_1^2(\Omega_T) \cap C_1(\overline{\Omega_T})$ solves (). Then max $u = \max u$ where T_t is the parabolic
 $\overline{\Omega_T}$ T_T boundary.
(i) Moreover if Ω is connected and $\exists (X_{01}b_0) \in \Omega_T$ such that
 $u(x_{01}b_0) = \max u$
 $\overline{\Omega_T}$ $then, u is constant in $\overline{\Omega_{00}} = \Omega \times (c_0, t_0)$ t_1 $u \in T_1$ by T_1 boundary.
Remark :- For t_1 to, the solutions may change
since the boundary, conditions change$

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🛡 🖓 🖓 T 🗆 🛞 her $g \in C(\Gamma_{\mp})$ and $f \in C(N_{T})$. Then $\exists \alpha t \mod \sigma ne solution u \in C_{T}^{2}(N_{T}) \cap C(\Lambda_{T})$ Uniqueness of the Heat equation:of the problem UE-QU= of in NF. Z-O U= g on FT. Question: Maximum Principle for the Cauchy Problem in R" T? Proof 5 D. 1.4. Remark 8- Of course u=0 is always a poly of uz-u=0 in $\mathbb{R}^n \times (0, \tau)$ but, one can show all oblive polytions grows rapidly as $(x_1 \rightarrow \infty)$.

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$$(x + 5)$$

$$(y +$$

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Uniquenen of the Heat equation? het ge C(Ft) and fe C(UDT). Then 3 at most one solution $u \in C_{1}^{2}(UT) \cap C(UT)$ of the problem $U_{t} - \Delta u = f$ in N_{T} . 3.0 u = g on T_{T} . Proof 5 D.1.9. Remark of D.1.9. Remark of Of course u = 0 is always a poly of $U_{t} - \Delta u = 0$ in $\mathbb{R}^{n} \times (0, T)$ Remark of Of course u = 0 is always a poly of $U_{t} - \Delta u = 0$ in $\mathbb{R}^{n} \times fe = 0$ but, one can show all other polythisms grows rapidly as $(M \to M)$. (Refer Slide Time: 59:07)

$$(x + 5)$$

$$Proof: Assume, 4at < 1 then \exists e>0 s+ 4a(T+e) < 1.$$
For yells" $x \mid y>0$ and dufine
 $v(x_1e):: u(x_1e) - \frac{M}{(T+e-e)^{n/2}} e^{-\frac{(x-y)^2}{4(T+e)}} (xells", tro)$

$$Clearly, v_{t} - \Delta v = 0 (Using Linearity of the heat operator)$$

$$Fir v = 0 avad eut d = B(y_1v) thin dr = B(y_1v)x(o_1T].$$

$$How, \max v = \max v.$$

$$V_{T} = F_{T}$$

$$Now y, xells", v(x_1o) = u(x_1o) - \frac{M}{(T+e)^{n/2}} e^{-\frac{(x-y)^2}{4(T+e)}} < u(x_1o) = \mathscr{B}(x)$$

$$avd, y (x-y) = v o \le t \le t t then_{T}$$

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$$(x + t) \qquad 0 = T = 0 \qquad 0 < 0 >$$

$$U(x_{t}) = u(x_{t}b) - \frac{M}{(T+b-t)^{n}} e^{-\frac{t^{n}}{t(t+t)}}$$

$$\leq C e^{\alpha|x|^{2}} - \frac{M}{(T+b-t)^{n}} e^{-\frac{t^{n}}{t(t+t)}}$$

$$\leq C e^{\alpha|(u_{t}t)^{2}} - \frac{M}{(T+b)^{n}} e^{-\frac{t^{n}}{u(t+t)}}$$
Now, $V_{d}(T+b) < I = hn(t)$, $I_{d}(T+b) = \alpha_{t} Y$ for some $T > 0$.
Now, $V_{d}(T+b) < I = \alpha_{t} (1 + 1)^{2} e^{(\alpha_{t}+1)^{2}}$

$$\therefore U(x_{t}b) \leq C e^{\alpha_{t}(|U_{t}+1)^{2}} - M [I_{d}(a_{t}x)]^{\frac{n}{2}} e^{(\alpha_{t}+1)x^{2}}$$

$$\leq Sup Q \sim (I_{t} r is setuted safficiently large)$$

$$W$$



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0 < 0 a

Proof: Assume,
$$4dt \leq 1$$
 then $\exists \varepsilon > 0$ st $4d(T+\varepsilon) \leq 1$.
Fix $y \in \mathbb{R}^{n} \setminus [u > 0$ and $define
 $v(x,t): = u(x,t) - \frac{u}{(T+\varepsilon-t)^{n_{2}}} e^{-\frac{|x-y|^{2}}{4(T+\varepsilon+t)}}$ $(x \in \mathbb{R}^{n}, t \ge 0)$
Clearly, $v_{t} - \Delta v = 0$ (Using Linearity of the heat operator)
Fix $v > 0$ and $evt = 0$.
Fix $v > 0$ and $evt = 0$.
Honce, $\max v = \max v$.
 $v = \max v$.$

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Proof : Assume,
$$4ut \leq 1$$
 then $\exists \varepsilon > 0 \ s + 4u(T+\varepsilon) \leq 1$.
Fix $y \in \mathbb{R}^{N}$ $\forall | u > 0$ and define
 $v(x,t) := u(x,t) - \frac{M}{(T+\varepsilon-t)^{N/2}} e^{-\frac{|x-y|^{2}}{4(t+t-t)}} (x \in \mathbb{R}^{n}, t \neq 0)$
Clearly, $v_{t} - \Delta v = 0$ (Using Linearity of the heat operator)
Clearly, $v_{t} - \Delta v = 0$ (Using Linearity of the heat operator)
Fix $v \neq 0$ and eet $d = B(y, v)$ thin $M \neq B(y, v) \times (0, \tau]$.
Hence, $\max v = \max v$. (a)
 $v_{t} = \frac{1}{\tau}$ (b)
Now $v_{t} = x \in \mathbb{R}^{n}$, $v(x; 0) = u(x; 0) - \frac{M}{(T+\varepsilon)^{n}h} e^{-\frac{1}{4}\frac{|x-v|^{2}}{4(t+\varepsilon)}} < u(x; 0) = \mathscr{B}(x)$.
Now $v_{t} = \frac{x \in \mathbb{R}^{n}}{\tau}$, $v(x; 0) = u(x; 0) - \frac{M}{(T+\varepsilon)^{n}h} e^{-\frac{1}{4}\frac{|x-v|^{2}}{4(t+\varepsilon)}} < u(x; 0) = \mathscr{B}(x)$.

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$\begin{aligned} \mathbf{v}(\mathbf{x},\mathbf{t}) &= \mathbf{u}(\mathbf{x},\mathbf{b}) - \frac{\mathbf{A}}{(\mathbf{T}+\mathbf{q}_{*}+\mathbf{b})^{n}\mathbf{h}} \in \frac{\mathbf{v}^{n}}{(\mathbf{T}+\mathbf{q}_{*}+\mathbf{b})^{n}\mathbf{h}} \in \frac{\mathbf{v}^{n}}{(\mathbf{T}+\mathbf{q}_{*}+\mathbf{b})^{n}\mathbf{h}} \\ &\leq Ce^{\alpha|\mathbf{x}|^{2}} - \frac{\mathbf{A}}{(\mathbf{T}+\mathbf{q}_{*}+\mathbf{b})^{n}\mathbf{h}} e^{-\frac{\mathbf{v}^{n}}{\mathbf{u}(\mathbf{T}+\mathbf{q}_{*})}} \\ &\leq Ce^{\alpha|\mathbf{x}|^{2}} - \frac{\mathbf{A}}{(\mathbf{T}+\mathbf{q}_{*}+\mathbf{b})^{n}\mathbf{h}} e^{-\frac{\mathbf{u}^{n}}{\mathbf{u}(\mathbf{t}+\mathbf{q}_{*})}} \\ &\leq Ce^{\alpha|\mathbf{x}|^{2}} - \frac{\mathbf{A}}{(\mathbf{T}+\mathbf{q}_{*})^{n}\mathbf{h}} e^{-\frac{\mathbf{u}^{n}}{\mathbf{u}(\mathbf{t}+\mathbf{q}_{*})}} \\ &\leq Ce^{\alpha|\mathbf{x}|^{2}} - \frac{\mathbf{A}}{(\mathbf{T}+\mathbf{q}_{*})^{n}\mathbf{h}} e^{-\frac{\mathbf{u}^{n}}{\mathbf{u}(\mathbf{u}+\mathbf{q}_{*})}} \\ &\leq Ce^{\alpha|\mathbf{x}|^{2}} - \frac{\mathbf{A}}{(\mathbf{T}+\mathbf{q}_{*})^{n}\mathbf{h}} e^{-\frac{\mathbf{u}^{n}}{\mathbf{u}(\mathbf{u}+\mathbf{q}_{*})}} \\ &\leq Ce^{\alpha|\mathbf{x}|^{2}} - \frac{\mathbf{A}}{(\mathbf{T}+\mathbf{q}_{*})^{n}\mathbf{h}} e^{-\frac{\mathbf{u}^{n}}{\mathbf{u}(\mathbf{u}+\mathbf{q}_{*})}} \\ &\leq Ce^{\alpha|\mathbf{x}|^{2}} - \mathbf{A} \left[\mathbf{u}(\mathbf{u}+\mathbf{x})\right]^{\frac{n}{2}} e^{(\mathbf{u}+\mathbf{x})\mathbf{v}^{2}}. \\ &\leq Ce^{\alpha|\mathbf{u}|^{2}} - \mathbf{A} \left[\mathbf{u}(\mathbf{u}+\mathbf{x})\right]^{\frac{n}{2}} e^{(\mathbf{u}+\mathbf{x})\mathbf{v}^{2}}. \\ &\leq Ce^{\alpha|\mathbf{u}|^{2}} - \mathbf{A} \left[\mathbf{u}(\mathbf{u}+\mathbf{x})\right]^{\frac{n}{2}} e^{(\mathbf{u}+\mathbf{x})\mathbf{v}^{2}}. \\ &\leq Ce^{\alpha|\mathbf{u}|^{2}} + \frac{\mathbf{u}^{n}\mathbf{u}^{2}}{\mathbf{u}^{n}\mathbf{u}^{2}} - \mathbf{A} \left[\mathbf{u}(\mathbf{u}+\mathbf{x})\right]^{\frac{n}{2}} e^{(\mathbf{u}+\mathbf{x})\mathbf{v}^{2}}. \end{aligned}$

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By (1), (11) a (12) one has, V(y,t) < sup & yelknotte[ori]. Let, M-10; sup u = sup g. Rixtorij Rn If 4at \$1 then, to, \$aJ, [\$a, fa].... Apply the above result in each interval P As a corollary, $u_{E-\Delta u} = f$ in $\mathbb{R}^n \times (0, t]$ (Uniqueuen in \mathbb{R}^n) u = q on $\mathbb{R}^n \times 5t = 0$) has atmost one solution, $|u(x,t)| \leq Ce^{\alpha |x|^2} \quad \forall x \in \mathbb{R}^n \lor t \in (0, t]$.

Then we proved a very important, but kind of difficult mean value property, which is not essentially your mean value, but we defined a kind of heat ball and we showed that in that ball you can do all sort of thing that the mean value holds in a different set. And then today we proved that strong maximum principle holds, uniqueness holds and you can use it to even settle the well-poseness problem, just, these ideas are exactly the same as we did for Laplacian. So, I am not really going deep into all that. So, all of this is done. So, with this we are going to finish the heat equation part of this course. Thank you very much.