

Calculus of Variations and Integral Equations

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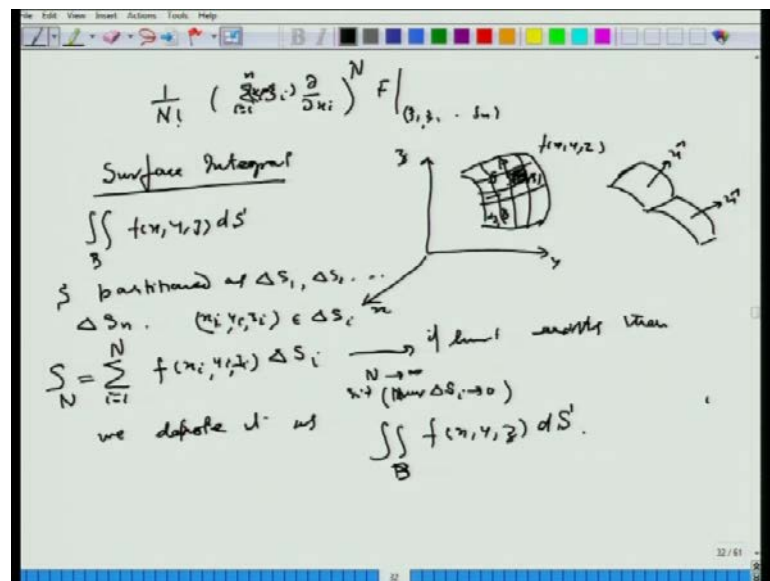
Department of Mathematics and Statics

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Module No. # 01

Lecture No. # 05

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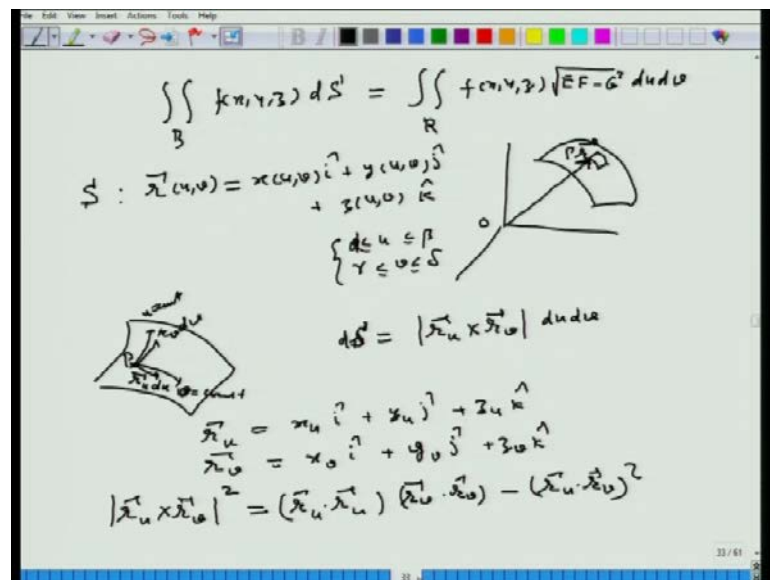
Welcome viewers. The NPTEL lecture series on the calculus of variations; this is the fifth lecture of the series. In the last lecture, we discussed the concept of surface integral where surface b is in three dimension space x, y, z and a function is defined on the surface, that is $f(x, y, z)$ for each point p which is whose coordinates are x, y, z . This point is on the surface, and the values of this function defined on all points on the surface, and it on its boundary b is the surface, and Δb is its boundary; this function is assumed to be continuous on the surface and on its boundary also.

So, then the surface integral double, integral over b $f(x, y, z) dS$, this is capital S and not the arc length is small s ; here dS is the element area on the surface, the typical surface element is like this, shaded darkened curved square or curved rectangle. This surface is partitioned into these element areas dS $\Delta S_1, \Delta S_2$ and so on up to ΔS_n , and x_i, y_i, z_i , is a point in this ΔS_i , then the integral is defined as the limit of the sum

summation i equal to 1 to capital N , when N is the number of these element surface areas; this number of elements partitioning the whole surface, and then we take this summation f of x_i, y_i, z_i , where x_i, y_i, z_i is the point in that element area ΔS_i , times the area of that element that is ΔS_i ; and if this limit exists, we define it as the integral, double integral over B of $f(x, y, z) dS$.

Here, this integral will be defined. It will not be depending on the way we partition it, for any partition when we pass to the limit, it should give you the same value. So, under these sufficient conditions like, f is continuous and surface, this surface is piecewise smooth, and its boundary is piecewise continuously differentiable, then this limit does not depend in the manner we partition it; and so the integral will be defined.

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Here, how we calculate this surface area. We as we know, we have some experience of calculating on the flats planes; the functions which are defined in certain domain like this. So, here x_i in the $x y$ plane, we have this flat surface area d and a function z is defined on this, and we have seen that this surface area is given by this formula. This also is a particular case of the more general one which we are going to consider.

So, here we have for any point p on this surface, like this. Here, this surface is parameterized by two parameters u and v . So, position vector of a point p that is the director vector o, p - starting from o and ending at p . So, here **here**, it will have three components x, y, z . So, each component is a function of u, v ; those parameters and so,

we get x_u, y_u, z_u like that, these are the components. So, the vector r will be given by $x_u i + y_u j + z_u k$. And these u, v are within the range of these bounded intervals α, β and γ, δ .

So, here then this element area dS ; we have seen that this comes out to be the cross, absolute value of the cross product $r_u \times r_v$ $du dv$. Because $r_u du$ is the tangent on the curve v equal to constant, and $r_v dv$ is the tangent element on the curve u equal to constant. So, this curved surface element like this darkened one, will be actually given by, which is dS equal to absolute value of the cross product $r_u \times r_v$; and since $du dv$ is positive it comes out. So, that is the element area dS and. So, here r_u is the partial derivative with respect to the parameter u , and r_v is the partial derivative of r position vector a with respect to v ; so each component gets differentiated partially with respect to those variables, and we get this and so, $r_u \times r_v$ square.

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$$\vec{a} \times \vec{b} = |\vec{a}| |\vec{b}| \sin \theta \hat{u}$$

$$|\vec{a} \times \vec{b}|^2 = |\vec{a}|^2 |\vec{b}|^2 \sin^2 \theta$$

$$= |\vec{a}|^2 |\vec{b}|^2 (1 - \cos^2 \theta)$$

$$= |\vec{a}|^2 |\vec{b}|^2 - |\vec{a}|^2 |\vec{b}|^2 \cos^2 \theta$$

$$= |\vec{a}|^2 |\vec{b}|^2 - (\vec{a} \cdot \vec{b})^2$$
 Taking $\vec{a} = \vec{r}_u$ $\vec{b} = \vec{r}_v$

$$|\vec{r}_u \times \vec{r}_v|^2 = (\vec{r}_u \cdot \vec{r}_u)(\vec{r}_v \cdot \vec{r}_v) - (\vec{r}_u \cdot \vec{r}_v)^2$$

$$= (x_u^2 + y_u^2 + z_u^2)(x_v^2 + y_v^2 + z_v^2) - (x_u x_v + y_u y_v + z_u z_v)^2$$

$$= EF - G^2$$

$$|\vec{r}_u \times \vec{r}_v| = \sqrt{EF - G^2}$$

$$\vec{r}_u \cdot \vec{r}_u = E$$

$$\vec{r}_u \cdot \vec{r}_v = F$$

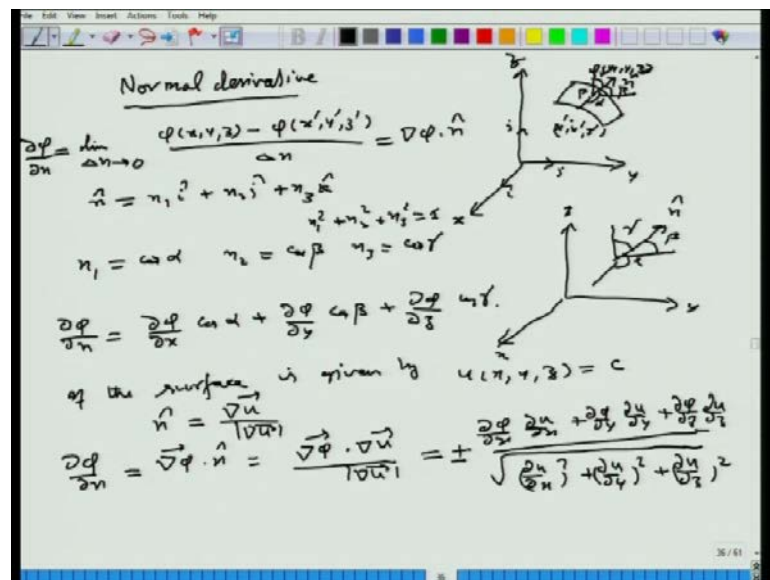
$$\vec{r}_v \cdot \vec{r}_v = G$$

Here we use the formula, which we proved that for any 2 vectors a and b in \mathbb{R}^3 , we see that $a \times b$ is given by the absolute value of $a \times b$ square is given by, absolute value of a square times absolute value of b square minus a dot b; this is the dot product of 2 vectors that is $a_1, b_1, a_2, b_2, a_3, b_3$ where a_i 's and b_i 's are the components of a and b respectively, the square of that.

So, we get this $r_u \times r_v$ absolute value square, like this and so, this $r_u \cdot r_u$ is E , that is how we define that $r_u \cdot r_u$ is E , and $r_v \cdot r_v$; this is F and $r_u \cdot r_v$, this is G .

So, this absolute value of $r_u \times r_v$ square comes out to be $E F$ minus G square. So, therefore, absolute value of $r_u \times r_v$ is square root of $E F$ minus G square; we take here positive square root, and that is what we have here, and so, our formula here; this integral comes out to be in terms of, now these are x, y, z are functions of u, v here. And this we have calculated that dS equal to square root $E F$ minus G square $du dv$. So, that is nothing but this integration over the flat surface $du dv$, which we already know how to calculate.

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So, that is the surface integral formula in terms of a parametric representation u, v , and then also we have Normal general Normal derivative in three dimension. As we have defined the Normal derivative of any function here. So, we have some surface here like this, and some function ϕ is defined on this. So, $\phi(x, y, z)$ which is defined at each point on this surface. So, Normal derivative of this ϕ is. So, Normal derivative of ϕ which was defined like, we defined it into two dimension, we define it limit of Δn ending to 0 of $\phi(x, y, z) - \phi(x', y', z')$. Here is a part of the surface; the close surface like this, where here this is x dash, y dash, z dash inside the domain and the point x, y, z is on the surface.

So, here it is assumed that ϕ is defined at those interior points x , y , z ; so, this limit over ΔN , and we as before we have seen it in two dimensions it turns out to be gradient of ϕ dot \hat{N} . So, where \hat{N} is the unit Normal this \hat{N} , \hat{N} is the unit Normal which is $N_1 i + N_2 j + N_3 k$; i, j, k are unit vectors in x, y, z directions respectively. And we see that this unit vector therefore, $N_1^2 + N_2^2 + N_3^2$; this should be equal to 1. These are direction cosines, they are also given like this; if this Normal makes this angle $\alpha, \beta,$ and γ . So, here is like, I will blow it up in this way; is with x axis, is x, y and z . So, this z the angle is γ with this x axis α , and this β .

So, N_1 is actually $\cos \alpha$, N_2 is $\cos \beta$, N_3 is $\cos \gamma$. Where this \hat{N} is the unit vector; here, this is Normal to the surface and it makes angles $\alpha, \beta,$ γ with the coordinate axis, and so, these are the direction cosines given by $\cos \alpha, \cos \beta, \cos \gamma$. So, we can see that this can be calculated, and so this is also the another notation of this is $\frac{\partial \phi}{\partial n}$. So, we can see that $\frac{\partial \phi}{\partial n}$ is $\frac{\partial \phi}{\partial x}$; gradient vector is $\frac{\partial \phi}{\partial x}$, and then a first component is N_1 . So, that is $\cos \alpha$ plus $\frac{\partial \phi}{\partial y} \cos \beta$ plus $\frac{\partial \phi}{\partial z} \cos \gamma$.

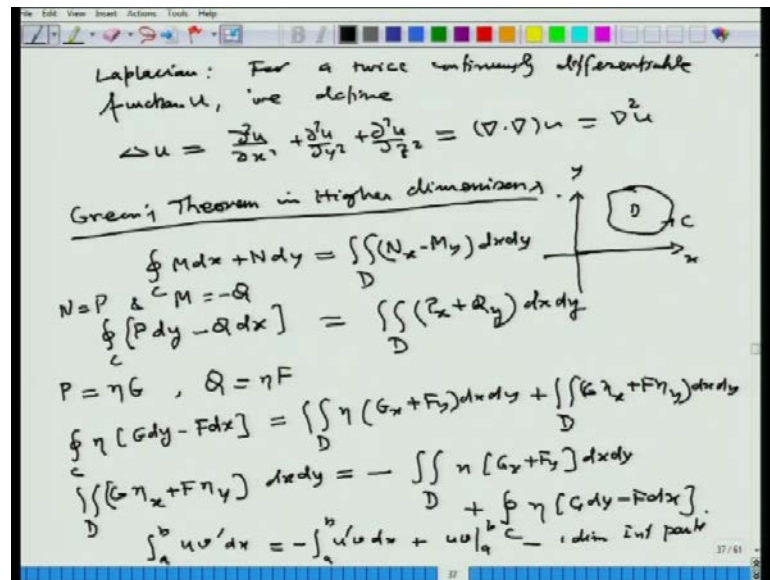
So, this another convenient formula for this. If in terms of direction cosines, we have a certain numbers here, small a , small b , small C then we divide by it is a square root summation $\sqrt{a^2 + b^2 + C^2}$. So, that is also can be seen. Here, if let us say this surface; if this surface is given by z equal to some function z, x, y ; then we know that the normal. So, in that or in terms of u, x, y, z let us say, if you give u, x, y, z equal to some constant c ; if you can solve it for z explicitly, we can express this z as a function of x, y also. So, then we know that the Normal will be gradient of u over. So this, in this case then \hat{N} comes out to be gradient of u over absolute value of the gradient of u , and then we get the directional derivative. So, $\frac{\partial \phi}{\partial n}$ is we have gradient of ϕ dot \hat{N} . So, that will be gradient of ϕ dot gradient of u over gradient of u .

So, this in component form, this is $\frac{\partial \phi}{\partial x} \frac{\partial u}{\partial x} + \frac{\partial \phi}{\partial y} \frac{\partial u}{\partial y} + \frac{\partial \phi}{\partial z} \frac{\partial u}{\partial z}$ over square with plus minus sign, square root here $\sqrt{\frac{\partial u}{\partial x}^2 + \frac{\partial u}{\partial y}^2 + \frac{\partial u}{\partial z}^2}$.

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So, plus sign or minus sign, we have to see that which side we are going to take so that, cosine of the angle should be positive that, and so, that we have to choose either plus sign or minus sign accordingly. If we choose plus sign if $\frac{\partial u}{\partial \phi} \frac{\partial \phi}{\partial N}$, this positive; then we choose plus sign and if this is negative then we choose minus sign.

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Now, here the Laplacian will be defined Laplace operator or Laplacian is actually defined as, the Laplacian of for twice continuously, differentiable function v can be define, we define this Laplacian that of u continuous function u , we define Laplacian of u as $\frac{\partial^2 u}{\partial x^2} + \frac{\partial^2 u}{\partial y^2} + \frac{\partial^2 u}{\partial z^2}$; which is nothing but gradient of gradient $\text{del} \cdot \text{del}$ like this, of u or in short it is written like $\text{del}^2 u$; this another notation for the Laplacian. Now, we go to green's theorem in higher dimension. Firstly, we have seen that green's theorem.

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In higher dimension now we will see. So, recall in two dimension, we had seen that this; if you have this case like in x, y , if you have some certain domain like this. And here, this boundary of this is let us say C , and its positive direction is taken like this anticlockwise, and here we have seen that if M and N are smooth functions here; $u \in C$ that this $M dx + N dy$. This was equal to $\int_D (N_x - M_y) dx dy$. This we had established in the earlier lecture; that were M and r continuously differentiable in D , and the continuous on a piecewise continuous on the boundary, and C is piecewise smooth, then we have

this result which is Green's theorem in two dimension. Now, we can generalize this in higher dimension. So, first let us see certain applications in two dimension itself, so we have, we can write this in a slightly different form like this. So, if we take M equal to like this; so, we can write this as N equal to p, let say N equal to p; and M equal to minus Q. Then we get $p \, dy - Q \, dx$, and this line integral will be then equal to $\oint_C (p \, dy - Q \, dx)$, we get $p \, dx + Q \, dy - dx \, dy$. So, just if moving this minus sign from here, and then that minus sign will come here.

Now, **we can**... So, this is a I mean just same form with different choices of M and N. So, now, we take p equal to some $\eta \, G$, and Q equal to $\eta \, F$; here this $\eta \, G$ and F they are continuously differentiable functions, and its on D and continuous piecewise continuous on the boundary of D, that is on C and C is of course, assume to be piecewise smooth; that means, that each point the tangent is well defined tangent, Normal or defined except at finitely many points. So, **there**... So, putting this P and Q here, we get the following which is on this side we get $\eta \, G$. So, put this P equal to $\eta \, G$; So, we get $\eta \, G \, dy - F \, dx$ equal to here. So, we will have two terms here, this double integral, because there are 2 functions here. So, x derivative will give you $\eta \, G_x$ plus $\eta \, x \, G$.

So, like that we have and collecting the since here, you get $\eta \, G_x + F_y \, dx \, dy$; here this these are all partial derivative, as sub scripted means the partial derivative plus double integral over D $\eta \, G_x + F_y \, dx \, dy$. And so, we take one term on the other side. So, we will get like this, here. So, we have. So, we can write this over D, like this $\eta \, G_x + F_y \, dx \, dy$, we take this on the other side to get minus D $\eta \, G_x + F_y \, dx \, dy$ plus this boundary integral $\eta \, G \, dy - F \, dx$. So, this is the general form of integration by parts; so, Green's theorem gives us the general way of the same formula of which we had in one-dimensional recall, that integration by parts formula was $\int_a^b u \, v' \, dx = u \, v - \int_a^b u' \, v \, dx$ plus a to b and u v evaluated at a to b.

So, this is one-dimensional. So, this is one-dimensional integration by parts; and so, this is two-dimensional integration by parts. So, this is generalization of the earlier integration by parts formula which we discussed. Here, we can see that, here the x derivatives there on η , y derivatives there on η , and G, N, F are not having any derivatives here. So, if you shift these derivatives on G, we get a minus sign here like we got here. So, we got minus sign and η becomes free of derivatives and this x derivative

get shifted on to G; similarly, y derivative get shifted onto F like this, and this eta becomes free of any derivative we gain minus sign here. So, this is the generalization of integration by parts formula of one-dimensional. So, this is two-dimensional integration by parts formula.

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Take $\eta = \psi \Rightarrow G = \phi_x, F = \phi_y$

$$\iint_D \psi \nabla^2 \phi dx dy = - \iint_D [\phi_x \psi_x + \phi_y \psi_y] dx dy + \int_C \psi \frac{\partial \phi}{\partial n} ds$$

$$\iint_D \psi \nabla^2 \phi dx dy = - \iint_D [\nabla \phi \cdot \nabla \psi] dx dy + \int_C \psi (\nabla \phi \cdot \vec{n}) ds$$

$$\iint_D \phi \nabla^2 \psi dx dy = - \iint_D [\nabla \psi \cdot \nabla \phi] dx dy + \int_C \phi (\nabla \psi \cdot \vec{n}) ds$$

$$\iint_D (\phi \nabla^2 \psi - \psi \nabla^2 \phi) dx dy = \int_C (\phi \frac{\partial \psi}{\partial n} - \psi \frac{\partial \phi}{\partial n}) ds$$

$\eta = \psi \quad \iint_D \phi \nabla^2 \phi dx dy = - \iint_D |\nabla \phi|^2 dx dy + \int_C \phi \frac{\partial \phi}{\partial n} ds$

η we take $\phi = 0 \quad P = G \eta_x - \eta G_x$

Similarly, we can have higher dimensional integration by parts formula. So, we can repeat this process like this, if we take. So, take now eta equal to psi, and G equal to phi x, and F equal to phi y. So, here in this we will take eta equal to psi and g equal to phi x and F equal to phi y. So, we get this thing here, and here these x derivative will give you G x x like that; So we get finally, psi del S square phi d x d y equal to minus, this is phi x psi x plus phi y psi y d x d y plus the boundary integral psi del phi by del n d s. Here, this s is the arc length. So, here we will see that, here if you put eta equal to **eta equal to** psi here, and g equal to phi x and F equal to phi y, we get from this; here we get g x will give you phi x x and this will give you phi y y; so that is the Laplacian two-dimensional Laplacian here, each of eta equal to psi, **eta equal to psi** and you get Laplacian there and then in this term here, we get G equal to phi x. So, this is eta x mean psi x and F is phi y and eta is eta y psi y, so we get this term here with minus sign.

This can also be written slightly differently like this, and this is the dot product of the gradient of phi dot gradient of psi. And here this is; obviously, this is we know that this is gradient of phi dot N cap d S. So, this is the in a vector form we have this, and if we

interchange the phi and psi we get phi here, and del square psi d x d y this remains the same, because it is a symmetric. So, we get gradient of psi dot gradient of phi d x d y plus over c.

So, psi here **sorry** this will be phi now, and this will be psi dot N cap d S. So, subtracting this we get the following greens identity, that phi this we can subtract from this. So, phi del square psi minus psi del square phi d x d y; these two terms cancel each other, and so, we get this boundary integral phi, we either write like this or we can write del S phi del n form also. So, this is phi del psi by del n minus psi del phi by del n. So, this is greens; one of the greens identities expressed like this. And if we write phi equal to psi here, we can see that; then we will have here del square this dot product will give you square of these facing. So, if we write phi equal to psi, in this if you write phi equal to psi. So, we will have this phi del square phi d x d y will be minus D mod del phi square plus phi del phi by del n d x.

So, there are various forms of identities obtain like this, which can be by proper choice of M and N or in F and G, we can see various forms of greens identities. We can go to higher order integration by parts also if we take. So, if you take this Q equal to 0, and p equal to G phi x minus **sorry** G eta x eta x minus eta G x. Then we get. So, we substitute here P equal to **P equal to** G eta x minus eta G x and Q equal to 0 here.

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$$\iint_D G \frac{\partial^2 \eta}{\partial x^2} dx dy = \iint_D \eta \frac{\partial^2 G}{\partial x^2} dx dy + \oint_C (G \eta_x - \eta G_x) dy$$

$$\iint_D G \eta_{xx} dx dy = \iint_D \eta G_{xx} dx dy + \oint_C (G \eta_x - \eta G_x) dy$$
 Take $P = \frac{G \eta_x - \eta G_x}{2}$, $Q = \frac{G \eta_x - \eta G_x}{2}$

$$\iint_D G \eta_{xy} dx dy = \iint_D \eta G_{xy} dx dy + \frac{1}{2} \oint_C (G \eta_y - \eta G_y) dx - \frac{1}{2} \oint_C (G \eta_x - \eta G_x) dy$$

Divergence Theorem

$$\vec{F} = F_1 \hat{i} + F_2 \hat{j} + F_3 \hat{k}$$

$$\iiint_V \left[\frac{\partial F_1}{\partial x} + \frac{\partial F_2}{\partial y} + \frac{\partial F_3}{\partial z} \right] dx dy dz = \iint_S \vec{F} \cdot \hat{n} dS$$

$$\hat{n} = n_1 \hat{i} + n_2 \hat{j} + n_3 \hat{k}$$

So, you will get the following thing; integration over D we get $\nabla^2 \eta$ by $\nabla \cdot (\nabla \times \nabla \times \eta)$ and this is over D $\nabla^2 G$ over $\nabla \cdot (\nabla \times \nabla \times G)$ plus the boundary integral $\int_C \nabla \eta \cdot \mathbf{n} - \eta \nabla \cdot \mathbf{n}$. Because we have chosen Q equal to 0 here, Q equal to 0. So, we get only see we have taken Q equal to 0 here. So, this term will not be there this. We will have only $\nabla \cdot \mathbf{p}$ and \mathbf{p} is chosen like this, $\nabla \times \nabla \times \eta$ minus $\nabla \times \nabla \times G$. So, we get at the same thing here, **and we get**... So, this is what is now higher order, I mean higher order derived integration by parts here. Because now, 2 derivatives are being shifted on G; here $\nabla \times \nabla \times$ we can write it like this also in short form $\nabla \times \nabla \times$; this is over D $\nabla \times \nabla \times G$ and plus $\nabla \cdot \mathbf{p}$ minus $\nabla \cdot \mathbf{p}$.

So, here this $\nabla \times \nabla \times$ first we shift one derivatives, we give 1 minus sign; and then again we shift $\nabla \times$ derivative, we get those $\nabla \times \nabla \times$ derivatives on G and η **eta** becomes free. So, this is shifting of derivatives; here, because of 2 derivatives are being shifted, we get plus sign here, and on the boundary we get this term. So that is the integration by parts formula for second order derivatives. Similarly, we can take now, P equal to something and Q equal to 0 So, then we will have P equal to $\nabla \cdot \mathbf{y}$ minus $\nabla \cdot \mathbf{y}$ by 2 and Q equal to $\nabla \times \nabla \times \eta$ minus $\nabla \times \nabla \times G$ by 2. So, then we get for mixed integral, mixed derivatives thus $\nabla \times \nabla \cdot \eta$ $\nabla \times \nabla \cdot G$. So, this is now, these these derivatives are shifted onto this. So, η becomes free and $\nabla \times \nabla \cdot \eta$ over D, and you get those boundary terms $\int_C \nabla \eta \cdot \mathbf{n}$ minus $\eta \nabla \cdot \mathbf{n}$ plus $\int_C \nabla \cdot \mathbf{y}$ minus $\mathbf{y} \cdot \mathbf{n}$; **sorry** minus **minus** sign here, minus $\int_C \nabla \times \nabla \times \eta$ minus $\nabla \times \nabla \times G$.

So, this is for the next derivative. Now, we go to three dimension case and we will have the following result. So, **here we use the divergence theorem which is**...

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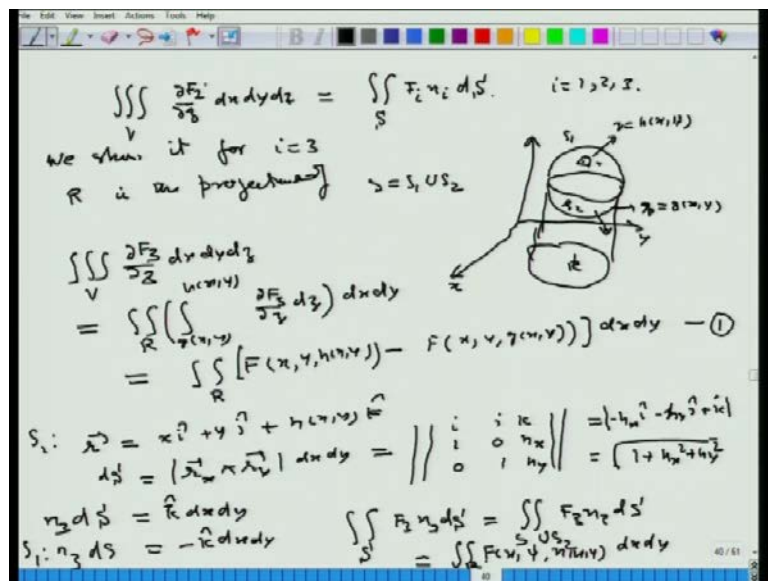
What it says, here this for given any bounded volume v here, and f . So, here divergence means $\nabla \cdot \mathbf{F}$, here $\nabla \cdot \mathbf{F} \, dx \, dy \, dz$ is actually equal to the surface integral or the surface of this v , $\int_S \mathbf{F} \cdot \mathbf{N} \, dS$, here we have the following picture. So, certain bounded volume V is enclosed by the surface S . So, this surfaces is S at any point the outward normal is given by \mathbf{N} as usual \mathbf{N} , the points are there in the domain in V . So, the this is what is the divergence theorem. In component form we get this; so, $\mathbf{F} \cdot \mathbf{N}$ components F_1, F_2, F_3 ; $F_1 \mathbf{i} + F_2 \mathbf{j} + F_3 \mathbf{k}$. So, then we get here $F_1 \frac{\partial F_1}{\partial x} + F_2 \frac{\partial F_2}{\partial y} + F_3 \frac{\partial F_3}{\partial z}$; that is the divergence of \mathbf{F} , because ∇ operator is $\mathbf{i} \frac{\partial}{\partial x} + \mathbf{j} \frac{\partial}{\partial y} + \mathbf{k} \frac{\partial}{\partial z}$. So, dot product gives you this. $\mathbf{F} \cdot \nabla \mathbf{F}$

x, y, z ; here, F these component ϕ_i are assuming to be continuously differentiable in V , and continuous on the surface S - And S is piecewise smooth that is the on those patches, whereas S is smooth and is well defined that is continuous function, at every point except on certain boundary of those patches; like it could be a like this.

So, in these patches let say, S_1, S_2, S_3 and S_4 ; this N is a continuously defined on those patches S_1, S_2 . So, here this side surface integral side, we get $F_1 N_1$ plus $F_2 N_2$ plus $F_3 N_3$ plus $F_4 N_4$ these are scalars components. Here, N cap as usual have components $N_1 i$ plus $N_2 j$ plus $N_3 k$; this is unit normal.

So, those are defined those cosines direction cosines of the normal here. So, how do we see this theorem, here we will component wise we show that, this is integral of this $\text{del } F_1$ $\text{del } x$ is $F_1 N_1 dS$ and so on. So, we show let us do it for the 1 the last term.

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So, last term we will show that this, over V $\text{del } F_3$ over $\text{del } z$ $dx dy dz$ equal to this over S ; this is $F_3 n_3 dS$. So, like that we have. So, are what we are to show that like this, that F_i , this i equal to 1, 2, 3. So, we show it for i equal to 3. So, what we will have, this way again will do it for a simple case like this that here, we will take the surface S parameterized by the x, y here coordinates, and So, this the surface S is projected on to this R here, this is the projection of this, and let say surface is like this. It take the simple case that surface has 2 parts - the upper part, the lower part we write it as z equal to $g(x, y)$, and the upper part z equal to $h(x, y)$.

So, assuming that the surface can be represented like this, and this R is the projection of this surface; here this is the projection of this. And so, R is the projection of projections of, of both the parts let us say this is S_1 , and this is part is S_2 So, S is $S_1 \cup S_2$. So, then we can see that this triple integral over V $\text{del } F^3$ over $\text{del } z \, dx \, dy \, dz$, here will be like this over R , and then here z component will be $g(x, y)$ to $h(x, y)$ of $\text{del } F^3$ over $\text{del } z \, dz$ and $dx \, dy$. And so, it will be over R here, F^3 . So, z differentiation and integration will cancel each other. So, we will get the boundary terms x, y and $h(x, y)$ minus F evaluated at $x, y, g(x, y), dx \, dy$.

So, here we can see that, this is what we will see that this is actually equal to the right hand side of that. So, here we see that if we write this parametric representation like this x, y . So, then the for the upper surface, we see that R comes out for any point P here, position vector R will be given by $x \, i$ plus $y \, j$ plus z component is there x on S_2 ; $h(x, y)$ and so, this element area here, will be so, dS here; will be $r_x \times r_y$, absolute value $dx \, dy$, and we can see that this is nothing but $\text{so, } r_x \dots$ So, absolute value of their determinant like i, j, k and here, you have $1 \, r_x$ is 1 and this is 0 and this is h_x here and $0, 1 \, h_y$ here, absolute value of that. So, we get here minus $h_x \, i$.

So, let us absolute value of that minus $h_x \, i$ minus $g \, h_y \, j$ and plus k . So, we here this square root of 1 plus h_x^2 plus h_y^2 , that is what we will get here. And so, here we see that this is the, here normal this $\cos \gamma$ will be positive, and on the lower part we will see that; here we get. So, this side is nothing but S of $F^3 \cdot N^3 \, dS$, we have seen that, here let us first seen that what is this actually element area. See here, so $N^3 \, dS$; this is the vector form of this element area will be actually equal to, N^3 is $k \, dS$ will come out to be simply $dx \, dy$. So, you will have this $N^3 \, dS$ will be like this, because they component N^3 component is just 1 here. $\text{So, } N^3 \dots$ So, N^3 component is the k here. So, we get this.

Now, on the lower part we get $N^3 \, dS$ on S_1 ; we get this as minus, because here you will have in the negative direction. So, will, but minus $k \, dx \, dy$. So, we can see that, here this is you get minus sign on this. So, we get finally, the same thing here $F \cdot y$. So, this dot product n^3 will be 1 here, n^3 will be... n^3 is a k component. So, it surplus value will be plus 1 here, and here you because of this you get minus. So, this minus sign will be adjusted and so, and so, this is one.

So, we get from this integral surface is $F_3 n_3 dz$, $F_1 dx$ on the lower S as the union S_1, S_2 , so on S_2 we get positive one. So, that is $h \times y$. So, let us first like this; **this** is S_1 union S_2 . So, $F_3 n_3 dz$.

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$$- \iint_S F_3(x, y, z) n_3 dz$$
 In component form

$$\iiint_V [(F_1)_x + (F_2)_y + (F_3)_z] dx dy dz = \iint_S [F_1 \cos \alpha + F_2 \cos \beta + F_3 \cos \gamma] dS$$

$$n = n_1 i + n_2 j + n_3 k, \quad n_1 = \cos \alpha, \quad n_2 = \cos \beta, \quad n_3 = \cos \gamma$$

$$F_1 = \eta F, \quad F_2 = \eta G, \quad F_3 = \eta H$$

$$\iiint_V (F_1 n_x + F_2 n_y + F_3 n_z) dx dy dz = - \iiint_V \eta [F_x + G_y + H_z] dx dy dz + \iint_S \eta [F \cos \alpha + G \cos \beta + H \cos \gamma] dS$$
 Take $\eta = \psi, \quad F = \psi_x, \quad G = \psi_y, \quad H = \psi_z$

$$\iiint_V \psi \nabla^2 \psi dx dy dz = - \iiint_V \nabla \psi \cdot \nabla \psi dx dy dz + \iint_S \psi \frac{\partial \psi}{\partial n} dS$$

$$\iiint_V [\psi \nabla^2 \psi - \nabla \psi \cdot \nabla \psi] dx dy dz = \iint_S [\psi \frac{\partial \psi}{\partial n} - \nabla \psi \cdot \nabla \psi] dS$$

So, first on S_2 that is with positive sign $F \times y$ $h \times y$ and $d \times d y$, because dot product here will have this 1 and so, $d \times d y$ and over this region R minus **minus** over this region R again, and $F_3 \times y$ $g \times y$ $d \times d y$. So, we see that it matches with the quantity here for each i similarly, we can do for other components, and so we get overall by adding all the terms we get the divergence theorem.

Now, So, here in the component form, it can be written like this. As you have seen already them. So, in component form, we have this that over triple integral V , we get F_x plus F_y plus F_z $d \times d y$ $d z$; we get over this S F_1 that is or we can write it in cosine form $\cos \alpha$ plus $F_2 \cos \beta$ plus $F_3 \cos \gamma$ $d S$. Where N_1 is $\cos \alpha$ as usual, N_2 is $\cos \beta$ N_3 is $\cos \gamma$. N has components, $N_1 i$ plus $N_2 j$ plus $N_3 k$.

So, we have this; now, if we choose as before we choose F_1 **sorry**, this was F_1 $F_1 x$, F_2 the components $F_3 y$. So, F_1, F_2, F_3 ; let me write it neatly. This is $F_1 x$ plus $F_2 y$ plus $F_3 z$.

Let the divergence of F here. And so, we get these from the divergence theorem; now, if we choose F_1 equal to let us say η F_2 equal to ηG plus F_3 **sorry** comma F_3 equal

to η ; where again η, F, G, N, H are assumed to be sufficiently smooth; and so, we kept these over V , here $F \eta_x + G \eta_y + H \eta_z dx dy dz$, and we get one more term here, we take it on the other side with minus sign now; and so, we will have over V $F_x + G_y + H_z dx dy dz$ of that boundary term is now surface here. So, we get surface integral over the surface S η , and then you have $F \cos \alpha + G \cos \beta + H \cos \gamma dS$.

So, you can see that, this is now three-dimensional integration by parts; here, derivatives on η . So, here η_x, η_y, η_z ; now, we shift these derivatives on F, G, N, H . We get this minus sign here, and so, we get $F_x + G_y + H_z$ here, and η becomes free of derivatives and here, this is the boundary term; here, boundary is the surface. So, we get this.

And if we put, η equal to ψ , and F equal to ϕ_x, G equal to ϕ_y , and H equal to ϕ_z ; we get this high order integration by parts. So, $\psi \nabla^2 \phi - \nabla \psi \cdot \nabla \phi$; this will have 3 components here $\phi_x, \phi_y, \phi_z \nabla \psi$ and plus this boundary term over S $\psi \nabla \phi \cdot \mathbf{n} dS$, which is now three-dimensional direction normal derivative dS .

If we interchange the ϕ and ψ , and then subtract we get the similar formula which we had earlier. So, this over V of $\phi \nabla^2 \psi - \psi \nabla^2 \phi - \nabla \phi \cdot \nabla \psi + \nabla \psi \cdot \nabla \phi$ will be, this term will cancel and we will have only the boundary term $\phi \nabla \psi \cdot \mathbf{n} - \psi \nabla \phi \cdot \mathbf{n} dS$. So, the more general form of Green's theorem which is actually consequence of the divergence theorem. So, with these preliminaries whatever things are required, we have discussed here in all these lectures, next time on in we will actually start our lectures on the calculus of variation, we introduce certain concepts again; and then we will start the Euler's conditions and various concepts related to the calculus of variations. Thank you very much for viewing the lecture.