

Calculus Of Variations and Integral Equation

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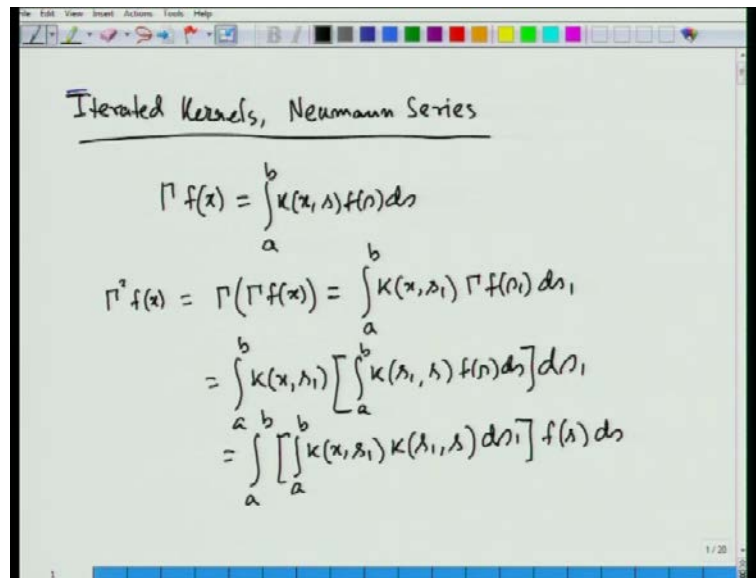
Indian Institute of Technology, Kanpur

Module No. #01

Lecture No. #34

Welcome viewers, once gain to the lecture series of NPTEL on the topic Integral Equation. In the last lecture, we were discussing the successive approximation or eternity method for solving non homogeneous Fredholm integral equation of the second kind. And we have considered one example in the last lecture, to find out a solution using that particular method. Now in these lecture we are again going to address the same eternity method in order to define the resolvent kernel, and in terms of resolvent kernel we are going to describe solution of the non homogeneous Fredholm integral equation of the second kind.

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The image shows a whiteboard with handwritten mathematical equations. The title is "Iterated kernels, Neumann Series". The equations are:

$$\Gamma f(x) = \int_a^b k(x, s) f(s) ds$$
$$\Gamma^2 f(x) = \Gamma(\Gamma f(x)) = \int_a^b k(x, s_1) \Gamma f(s_1) ds_1$$
$$= \int_a^b k(x, s_1) \left[\int_a^b k(s_1, s) f(s) ds \right] ds_1$$
$$= \int_a^b \left[\int_a^b k(x, s_1) k(s_1, s) ds_1 \right] f(s) ds$$

So, in these lecture, we are going to consider the topic that is iterated kernels, which ultimately leads us to Neumann series which is will be used to solve the Fredholm integral equations. In as per the previous discussion you can recall, we have introduced

this notation for integral operator that is capital gamma f x is equal to integral a to b k of x, s f s d s, we have introduced this notation. Now, in order to obtain the iterated kernels that we have done for Volterra integral equations, we can write these gamma 2 f x is nothing but, the integral operator gamma is operating upon gamma f x, so that means, this integral operator gamma is operating upon gamma f x and therefore, we can write this is equal to integral a to b k of x, s 1 gamma f of s 1 d s 1.

So, here this f s is repressed by gamma f s 1 and we have considered this dummy variable as s 1, in order to define the integral operator gamma on gamma f x, and then using the definition for gamma s 1, we can write integral a to b k of x, s 1 then integral a to b k of s 1, s f s d s then d s 1. Now, rearranging the terms that means, interchanging the order of integration we can write this is actually integral a to b k of x, s 1 then k s 1, s d s 1 these result can be integrated from a to b multiplied with f s then d s.

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The image shows a whiteboard with handwritten mathematical derivations. At the top, it defines the first kernel: $k_1(x, s) = k(x, s)$. Below that, it defines the second kernel as an integral of the first kernel over a dummy variable s_1 : $k_2(x, s) = \int_a^b k(x, s_1) k(s_1, s) ds_1$. This is then simplified to $k_2(x, s) = \int_a^b k(x, s_1) k_1(s_1, s) ds_1$. Next, it defines the first integral operator $\Gamma^2 f(x) = \int_a^b k_2(x, s) f(s) ds$. Finally, it defines the second integral operator $\Gamma^3 f(x) = \Gamma(\Gamma^2 f(x)) = \int_a^b k(x, s_1) \Gamma^2 f(s_1) ds_1$, which is further expanded to $\Gamma^3 f(x) = \int_a^b k(x, s_1) \left[\int_a^b k_2(s_1, s) f(s) ds \right] ds_1$. The whiteboard also shows a standard software toolbar at the top and a page number '2' at the bottom.

Now, if we define that $k_1(x, s)$ this stands for $k(x, s)$ same as we have done in case of Volterra integral equations therefore, $k_2(x, s)$ can be defined by integral a to b k of x, s 1 then k s 1, s d s 1, so this actually integral a to b k x, s 1, now repressing this k s 1 s by k 1 s 1, s d s 1 we get the second iterated kernel $k_2(x, s)$.

And therefore, $\Gamma^2 f(x)$ comes out to be integral a to b $k_2(x, s) f(s) ds$ this is the expression for $\Gamma^2 f(x)$, next if we calculate $\Gamma^3 f(x)$ in terms of iterated kernel, then we can find this gamma is operating upon $\Gamma^2 f(x)$ similarly, as previous what

we have done that is integral a to b k of x, s 1 gamma 2 f of s 1 d s 1. Now, from here we can write gamma 2 f s 1 this will be equal to integral a to b k of x, s 1 then integral a to b k 2 s 1, s f s d s this with d s 1.

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$$\Gamma^2 f(x) = \int_a^b \left[\int_a^b k(x, s_1) k_2(s_1, s) ds_1 \right] f(s) ds$$

$$k_3(x, s) = \int_a^b k(x, s_1) k_2(s_1, s) ds_1$$

$$\Gamma^3 f(x) = \int_a^b k_3(x, s) f(s) ds$$

$$k_n(x, s) = \int_a^b k(x, \xi) k_{n-1}(\xi, s) d\xi, \quad n = 2, 3, \dots$$

$$k_1(x, s) = k(x, s)$$

Again interchanging the order of the integration, we can write gamma 3 f x this is equal to integral a to b then integral a to b k of x, s 1 then k 2 s 1, s d s 1 multiplied with f s d s and now, if we define that k 3 x, s is equal to integral a to b k of x, s 1 then k 2 s 1, s d s 1, so therefore, gamma 3 f x will be equal to integral a to b k 3 x, s f s d s.

So, proceeding in this particular way, we can find nth iterated kernel k n x, s that is equal to integral a to b k of x, xi k n minus 1 xi, s d xi in all this definition for k 2 x s k 3 x, s here, these dummy variable s 1 can be replaced by xi, so that means, in general k n x, s equal to integral a to b k x, xi k n minus 1 xi, s d xi and these particular result holds for n equal to 2, 3 and so on. And where k 1 x, s is exactly equal to k of x, s and therefore, you can recall the solution for the Volterra integral equation, what we have considered in the last lecture that was...

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$$\begin{aligned}
 y(x) &= f(x) + \sum_{n=1}^{\infty} \lambda^n \Gamma^n f(x) \\
 &= f(x) + \sum_{n=1}^{\infty} \lambda^n \int_a^b k_n(x,s) f(s) ds \\
 &\quad | \lambda | L_2 (b-a) < 1 \\
 &= f(x) + \int_a^b \left(\sum_{n=1}^{\infty} \lambda^n k_n(x,s) \right) f(s) ds \\
 &= f(x) + \lambda \int_a^b \left(\sum_{n=1}^{\infty} \lambda^{n-1} k_n(x,s) \right) f(s) ds \\
 &= f(x) + \lambda \int_a^b \left(\sum_{n=0}^{\infty} \lambda^n k_{n+1}(x,s) \right) f(s) ds
 \end{aligned}$$

$y(x)$ is equal to $f(x)$ plus sigma in running's from 1 to infinity lambda to the power n gamma n operated upon $f(x)$, this was the result of the **solution** integral equation; that **this** is the solution for the Fredholm integral equation. And now, using these earlier results that is for gamma $3 f(x)$ gamma $2 f(x)$ and in general, you can write also this gamma $n f(x)$ in terms of this n th order iterated kernel, we can write this is equal to $f(x)$ plus sigma in running's from 1 to infinity lambda to the power n in integral a to b $k_n(x, s) f(s) ds$ this is the result (Refer Slide Time: 09:17).

And assuming satisfaction of this condition that is modulus lambda $L_2 (b-a) < 1$, assuming this condition hold where L_2 is actually maximum value of the kernel $k(x, s)$ it is modulus within the interval a, b cross a, b that is within a square therefore, we can interchange this summation and integral sign. Because, in the last lecture we have already proved the uniform convergence of this infinite series and therefore, this is equal to $f(x)$ plus integral a to b sigma n running's from 1 to infinity lambda to the power n $k_n(x, s) f(s) ds$ this entire expression multiplied with $f(s) ds$.

Now, taking one lambda outside the integral sign we can write, this is equal to $f(x)$ plus lambda integral a to b sigma n running's from 1 to infinity lambda to the power $n-1$ $k_n(x, s) f(s) ds$. And now, changing the range of variation for n we can get this is equal to $f(x)$ plus lambda integral a to b sigma n running's from 0 to infinity then it will be lambda to the power n $k_{n+1}(x, s) f(s) ds$, so therefore, this infinite series that is

sigma n running's from 0 to infinity lambda to the power n k n plus 1 x, s, this is actually resolvent kernel.

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$$R(x, s; \lambda) = \sum_{n=0}^{\infty} \lambda^n k_{n+1}(x, s)$$

$$= k_1(x, s) + \lambda k_2(x, s) + \lambda^2 k_3(x, s) + \dots \rightarrow \infty$$

$$y(x) = f(x) + \lambda \int_a^b R(x, s; \lambda) f(s) ds$$

And this resolvent kernel it is denoted by $R(x, s; \lambda)$ and that is equal to sigma n running's from 0 to infinity lambda to the power n k n plus 1 x, s, so that is actually equal to $k_1(x, s) + \lambda k_2(x, s) + \lambda^2 k_3(x, s) + \dots$ up to infinity. And therefore, with this resolvent kernel $R(x, s; \lambda)$ we can write solution of the Fredholm integral equation is $y(x) = f(x) + \lambda \int_a^b R(x, s; \lambda) f(s) ds$ this is actually solution to the given problem; and this series that is $k_1(x, s) + \lambda k_2(x, s) + \lambda^2 k_3(x, s) + \dots$ up to infinity this series actually call the Neumann series. And this is the solution of this Fredholm integral equation, of the second kind which is a non homogeneous equation in terms of the resolvent kernel.

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$$R(x, s; \lambda) = \sum_{n=0}^{\infty} \lambda^n k_{n+1}(x, s)$$

$$= k_1(x, s) + \lambda k_2(x, s) + \lambda^2 k_3(x, s) + \dots + \infty$$

$$y(x) = f(x) + \lambda \int_a^b R(x, s; \lambda) f(s) ds$$

Ex. $y(x) = 1 + \lambda \int_0^1 (1 - 3xs) y(s) ds$

$$k(x, s) = 1 - 3xs = k_1(x, s)$$

$$k_2(x, s) = \int_0^1 k(x, \xi) k_1(\xi, s) d\xi$$

Now, we consider one interesting example, this example you can find in many books for example, the book by Karneval as well as Hildebrand in different books you can find this very famous example, and these example will address again in some later lectures in order to compare the different methods by for the solution of Fredholm integral equation.

Now, here we are considering the problem that is $y(x)$ is equal to 1 plus lambda integral 0 to 1 $(1 - 3xs) y(s) ds$ we have to solve this problem. So, therefore, our kernel $k(x, s)$ this is equal to $1 - 3xs$. Now, first we calculate few initial iterates that is $k_2(x, s)$, $k_3(x, s)$ and so on, and then using the Neumann series we can calculate that resolvent kernel and then in terms of resolvent kernel we write down the solution for the given problem. So, here this $k(x, s)$ is nothing but, your $k_1(x, s)$ next we have to calculate this $k_2(x, s)$, by definition this is integral 0 to 1 $k(x, \xi)$ multiplied with $k_1(\xi, s) d\xi$, so with this definition that is $k(x, s)$ equal to $1 - 3xs$ and $k_1(x, s)$ equal to $1 - 3xs$.

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$$\begin{aligned}
 &= \int_0^1 (1-3x\xi)(1-3\xi s) d\xi \\
 &= \int_0^1 [1 - 3(x+s)\xi + 9x s \xi^2] d\xi \\
 &= 1 - \frac{3}{2}(x+s) + 3xs \\
 k_3(x, s) &= \int_0^1 k(x, \xi) k_2(\xi, s) d\xi \\
 &= \int_0^1 (1-3x\xi) \left(1 - \frac{3}{2}(\xi+s) + 3\xi s\right) d\xi \\
 &= \dots = \frac{1}{4} (1-3xs) = \frac{1}{4} k_1(x, s)
 \end{aligned}$$

We can write this is equal to integral 0 to 1 1 minus 3 x xi this multiplied with 1 minus 3 xi s d xi, so this is equal to integral 0 to 1 1 minus 3 x plus s, this multiplied with xi plus 9 x s xi square d xi this one, and after integration we can find this will be equal to 1 minus 3 by 2 x plus s plus 3 x s this will be the result, so this is actually our second iterated kernel k 2 x, s. Using this definition for k 2 x, s not definition this is actually we have derived k 2 x, s, so this expression we can calculate k 3 x, s.

So, k 3 x, s by definition integral 0 to 1 k x, xi then k 2 xi, s d xi this is equal to integral 0 to 1 1 minus 3 x xi this multiplied with 1 minus 3 by 2 xi plus s plus 3 xi s d xi and after with respect to xi you can arrive at this result, this will be equal to 1 by 4 1 minus 3 x s. So, these result is very much important, because from here you can observe this k 3 x, s is nothing but, 1 by 4 k 1 x, s, so what, we have assumed k 1 x, s and that is actually your given kernel k x, s.

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$$\begin{aligned}
 k_4(x, s) &= \int_0^1 k(x, \xi) k_3(\xi, s) d\xi \\
 &= \int_0^1 k(x, \xi) \frac{1}{4} k_1(\xi, s) d\xi \\
 &= \frac{1}{4} \int_0^1 k(x, \xi) k_1(\xi, s) d\xi = \frac{1}{4} k_2(x, s) \\
 k_5(x, s) &= \int_0^1 k(x, \xi) k_4(\xi, s) d\xi \\
 &= \frac{1}{4} \int_0^1 k(x, \xi) k_2(\xi, s) d\xi \\
 &= \frac{1}{4} k_3(x, s) = \left(\frac{1}{4}\right)^2 k_1(x, s)
 \end{aligned}$$

So, with these result that is $k_3(x, s)$ is equal to $\frac{1}{4}$ multiplied with $k_1(x, s)$, you can calculate $k_4(x, s)$, now $k_4(x, s)$ is equal to $\int_0^1 k(x, \xi) \frac{1}{4} k_1(\xi, s) d\xi$, so this is equal to $\frac{1}{4} \int_0^1 k(x, \xi) k_1(\xi, s) d\xi$, so that is equal to $\frac{1}{4} k_2(x, s)$ because, $\int_0^1 k(x, \xi) k_1(\xi, s) d\xi$ is nothing but, $k_2(x, s)$.

So, similarly, if you calculate $k_5(x, s)$ this will be equal to $\int_0^1 k(x, \xi) k_4(\xi, s) d\xi$, now $k_4(x, s)$ is equal to $\frac{1}{4} k_2(x, s)$, so using this result you can write this is equal to $\frac{1}{4} \int_0^1 k(x, \xi) k_2(\xi, s) d\xi$ this result you can obtain, this will be equal to $\frac{1}{4} k_3(x, s)$ because, $\int_0^1 k(x, \xi) k_2(\xi, s) d\xi$ is nothing but, $k_3(x, s)$ so this is nothing but, $\frac{1}{4} k_3(x, s)$ this will be the result for k_5 . Now, already we have obtained that $k_3(x, s)$ is equal to $\frac{1}{4} k_1(x, s)$, so this is equal to $\frac{1}{4} \left(\frac{1}{4}\right) k_1(x, s)$ so with these few results, we can claim that in general will be having this recursive formula.

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$$k_{n+1}(x, s) = \frac{1}{4} k_n(x, s), \quad n \geq 2$$

$$R(x, s; \lambda) = k_1(x, s) + \lambda k_2(x, s) + \lambda^2 k_3(x, s) + \lambda^3 k_4(x, s) + \lambda^4 k_5(x, s) + \lambda^5 k_6(x, s) + \dots + \infty$$

$$= k_1(x, s) + \lambda k_2(x, s) + \lambda^2 \frac{1}{4} k_1(x, s) + \lambda^3 \frac{1}{4} k_2(x, s) + \lambda^4 \frac{1}{4} k_3(x, s) + \lambda^5 \frac{1}{4} k_4(x, s) + \dots + \infty$$

$$= k_1(x, s) + \lambda k_2(x, s) + \frac{\lambda^2}{4} k_1(x, s) + \frac{\lambda^3}{4} k_2(x, s) + \frac{\lambda^4}{4^2} k_1(x, s) + \frac{\lambda^5}{4^2} k_2(x, s) + \dots + \infty$$

That is $k_{n+1}(x, s)$ this is equal to $\frac{1}{4} k_n(x, s)$ this result is valid for n greater than equal to 2, so this is actually one important step that we have obtained. So, from here, we can write $R(x, s; \lambda)$ that means, with this recursive relation and with some few initial **iter** iterates of the kernel, we can calculate the resolvent kernel $R(x, s; \lambda)$ for the given problem.

So, this is equal to $k_1(x, s) + \lambda k_2(x, s) + \lambda^2 k_3(x, s) + \lambda^3 k_4(x, s) + \lambda^4 k_5(x, s) + \lambda^5 k_6(x, s) + \dots + \infty$.

And now, we can use this result for eternity kernels and some initial results to get this will be equal to $k_1(x, s)$ as usual there is no change, no change for $k_2(x, s)$ then λ^2 it will be $\frac{1}{4} k_1(x, s) + \lambda k_2(x, s)$ and then λ^3 to the power 4 $\frac{1}{4} k_3(x, s) + \lambda^2 k_4(x, s) + \dots + \infty$.

Then using the result in last two terms, that is $k_3(x, s)$ is equal to $\frac{1}{4} k_1(x, s)$ and $k_4(x, s)$ equal to $\frac{1}{4} k_2(x, s)$ we can write, this is equal to $k_1(x, s) + \lambda k_2(x, s) + \lambda^2 \frac{1}{4} k_1(x, s) + \lambda^3 \frac{1}{4} k_2(x, s) + \lambda^4 \frac{1}{4^2} k_1(x, s) + \lambda^5 \frac{1}{4^2} k_2(x, s) + \dots + \infty$, so we have one set of term where $k_1(x, s)$ is there and other set of terms involving $k_2(x, s)$.

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$$\begin{aligned}
 &= \left(1 + \frac{\lambda^2}{4} + \frac{\lambda^4}{4^2} + \dots\right) k_1(x, s) + \lambda k_2(x, s) \left(1 + \frac{\lambda^2}{4} + \frac{\lambda^4}{4^2} + \dots\right) \\
 &= \left(1 + \frac{\lambda^2}{4} + \frac{\lambda^4}{4^2} + \dots\right) (k_1(x, s) + \lambda k_2(x, s)) \\
 &= \frac{k_1(x, s) + \lambda k_2(x, s)}{1 - \frac{\lambda^2}{4}}, \quad |\lambda| < 2 \\
 &= \frac{1 - 3xs + \lambda \left[1 - \frac{3}{2}(x+s) + 3xs\right]}{1 - \frac{\lambda^2}{4}} \\
 &= \frac{1 + \lambda - \frac{3}{2}\lambda(x+s) - 3xs(1-\lambda)}{1 - \frac{\lambda^2}{4}}
 \end{aligned}$$

So, these expression is equal to 1 plus lambda square by 4 plus lambda to the power 4 by 4 square plus dot dot, this multiplied with $k_1(x, s)$ and for the rest of the term, if you take common lambda and $k_2(x, s)$ then this will be multiplied with 1 plus lambda square by 4 plus lambda to the power 4 by 4 square plus dot dot. So, ultimately we are having this expression that is 1 plus lambda square by 4 plus lambda to the power 4 by 4 square plus dot dot up to infinity these multiplied with $k_1(x, s)$ plus lambda $k_2(x, s)$ this pre multiplied infinite series you can easily observe this an geometric series, and this geometric series with first term 1 and common ratio lambda square by 4.

So, this will be equal to $k_1(x, s)$ plus lambda $k_2(x, s)$ these divided by 1 minus lambda square by 4 and criteria for convergence is given by modulus lambda less than 2 and after substituting the expression for $k_1(x, s)$ and $k_2(x, s)$, you can find this is 1 minus 3 x s plus lambda into 1 minus 3 by 2 x plus s plus 3 x s this whole divided by 1 minus lambda square by 4. So, that means, this is equal to 1 plus lambda minus 3 by 2 lambda times x plus s minus 3 x s multiplied with 1 minus lambda divided by 1 minus lambda square by 4, so this is actually the sum for the Neumann series, and also this is the expression for the resolvent kernel.

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$$y(x) = f(x) + \lambda \int_0^1 R(x,s; \lambda) f(s) ds$$

Ex. 1. $y(x) = 1 + \lambda \int_0^{\pi} \sin(x+s) y(s) ds$

2. $y(x) = f(x) + \lambda \int_0^1 e^{x-s} y(s) ds$

3. $y(x) = \sin x - \frac{x}{4} + \frac{1}{4} \int_0^{\frac{\pi}{2}} x s y(s) ds$

4. $y(x) = \frac{3}{2} e^x - \frac{1}{2} x e^x - \frac{1}{2} + \frac{1}{2} \int_0^1 s y(s) ds$

So, with these resolvent kernel, if you substitute into the expression that is $y(x)$ is equal to $f(x)$ plus λ integral 0 to 1 $R(x,s; \lambda) f(s) ds$ then you will be having solution to the given Fredholm integral equation. And now, before going to the next part, I am giving some exercise for your practice, you can solve these problems first one, $y(x)$ is equal to 1 plus λ integral 0 to π $\sin(x+s) y(s) ds$, second problem $y(x)$ is equal to $f(x)$ plus λ integral 0 to 1 $e^{x-s} y(s) ds$.

Number 3, $y(x)$ this is equal to $\sin x - \frac{x}{4} + \frac{1}{4} \int_0^{\frac{\pi}{2}} x s y(s) ds$ and number 4, $y(x)$ this is equal to $\frac{3}{2} e^x - \frac{1}{2} x e^x - \frac{1}{2} + \frac{1}{2} \int_0^1 s y(s) ds$, so all these problems you can solve by the method of resolvent kernels.

Now, before going to the next topic I discuss briefly, an interesting result that is involved with the resolvent kernel, and where we can show that resolvent kernel actually satisfies an integral equation of Fredholm type. But, that will be in terms of two variables x and s , where $f(x)$ can be replaced by the given kernel and deduction is very straightforward.

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$$\begin{aligned}
 R(x, s; \lambda) &= k_1(x, s) + \lambda k_2(x, s) + \lambda^2 k_3(x, s) + \dots + \infty \\
 &= k(x, s) + \sum_{n=1}^{\infty} \lambda^n k_{n+1}(x, s) \\
 &= k(x, s) + \lambda \sum_{n=1}^{\infty} \lambda^{n-1} k_{n+1}(x, s) \\
 &= k(x, s) + \lambda \sum_{n=0}^{\infty} \lambda^n k_{n+2}(x, s) \\
 &= k(x, s) + \lambda \sum_{n=0}^{\infty} \lambda^n \int_a^b k(x, \xi) k_{n+1}(\xi, s) d\xi \\
 &= k(x, s) + \lambda \int_a^b k(x, \xi) \left[\sum_{n=0}^{\infty} \lambda^n k_{n+1}(\xi, s) \right] d\xi \\
 &= k(x, s) + \lambda \int_a^b k(x, \xi) R(\xi, s; \lambda) d\xi
 \end{aligned}$$

We start with a definition, that is $r \times s$ lambda this is equal to $k_1(x, s)$ plus $\lambda k_2(x, s)$ plus $\lambda^2 k_3(x, s)$ plus dot dot up to infinity; and we can write this is equal to since, $k_1(x, s)$ **you know** this is equal to actually $k(x, s)$, then we can write this is plus summation in running's from 1 to infinity $\lambda^n k_{n+1}(x, s)$. So, that means, this is the rest of the part is written under the summation notation, and now if you take one lambda outside the summation notation, this will be $k(x, s)$ plus λ sigma n equal to 1 to infinity $\lambda^{n-1} k_{n+1}(x, s)$.

Now, when you are substituting n equal to 1, so first index of lambda is going to be 0, so changing this limit of the sum, we can write this is $k(x, s)$ plus λ sigma n equal to 0 to infinity then it will be $\lambda^n k_{n+1}$ will be converted into $k_{n+2}(x, s)$. And now, here for $k_{n+2}(x, s)$ we can write the formula for iterated kernel, so that means, this will be equal to $k(x, s)$ plus λ sigma n equal to 0 to infinity λ^n integral a to b $k(x, \xi) k_{n+1}(\xi, s) d\xi$ here, we are just writing the formula for iterated kernel of $k_{n+2}(x, s)$ is equal to integral a to b $k(x, \xi) k_{n+1}(\xi, s) d\xi$.

Now, already we have proved the uniform convergence of these part, so therefore, we can interchange the summation and integral sign, so after interchanging you will have $k(x, s)$ plus λ then integral a to b $k(x, \xi)$ then sigma n running's from 0 to infinity $\lambda^n k_{n+1}(\xi, s) d\xi$. Now, this n running's from 0 to infinity

lambda to the power n k n plus 1 xi, s is nothing but, our resolvent kernel written in terms of xi and s. So, therefore, we can write this is equal to k x, s plus lambda integral a to b k of x, xi then R of xi, s lambda d xi, so if you look at the final expression, so that means, we have obtained R x, s colon lambda is equal to k x, s plus lambda integral a to b k x, xi r xi s lambda d xi.

So, that means, the solution of the Fredholm integral equation given equation was y x equal to f x plus lambda integral a to b k x, s f s d s this was the solution of the Fredholm integral equation.

Now, here this y is replaced by R x s lambda and f is replaced by k x, s, so therefore, you can see this dissolvent kernel satisfies a similar type of integral equation, this is one important observation. Now, we are going to consider an algebraic method where you can see, we have to solve a system of linear equations, and by solving that system of linear equations by some technique.

We can find out the solution of the Fredholm integral equation which is a non homogeneous Fredholm integral equation and with degenerate kernel, so that means, kernel is separable. And in that case, we can see the solvability condition depends upon the solution or uniqueness of the solution for the system of linear equation, so first of all we described this method and then we will consider a simple example.

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The image shows a whiteboard with handwritten mathematical equations. The equations are as follows:

$$y(x) = f(x) + \lambda \int_a^b k(x, s) y(s) ds$$

$$k(x, s) = \sum_{r=1}^n p_r(x) q_r(s)$$

$$y(x) = f(x) + \lambda \int_a^b \left[\sum_{r=1}^n p_r(x) q_r(s) \right] y(s) ds$$

$$= f(x) + \lambda \sum_{r=1}^n p_r(x) \int_a^b y(s) q_r(s) ds$$

$$y_r = \int_a^b y(s) q_r(s) ds, \quad r = 1, 2, 3, \dots, n$$

$$y(x) = f(x) + \lambda \sum_{r=1}^n y_r p_r(x)$$

The whiteboard also features a toolbar at the top with various drawing tools and a page number '12' at the bottom.

So, we are considering equation of the form $y(x) = f(x) + \lambda \int_a^b k(x, s) y(s) ds$ this is the given equation, kernel is separable, so that means, $k(x, s) = \sum_{r=1}^n p_r(x) q_r(s)$, this is a separable equation. And if we substitute these expression $k(x, s)$ into this integral, under this integral sign then you will be having $y(x)$ this is equal to $f(x) + \lambda \int_a^b \sum_{r=1}^n p_r(x) q_r(s) y(s) ds$, this expression multiplied with $y(s)$ as the kernel is separable.

So, we can take this $p_r(x)$ outside the integral sign and therefore, will be having this expression $f(x) + \lambda \sum_{r=1}^n p_r(x) \int_a^b y(s) q_r(s) ds$, now this kernel is separable, so that means, $p_r(x) q_r(s)$ they are known, whenever r ranging from 1 to n , but y is unknown quantity.

So, if we introduce the notation that is y_r this stands for $\int_a^b y(s) q_r(s) ds$ where r equal to 1, 2, 3 dot dot up to n , then these expression $y(x) = f(x) + \lambda \sum_{r=1}^n p_r(x) y_r$ comes out to be $y(x)$, this is equal to $f(x) + \lambda \sum_{r=1}^n p_r(x) y_r$.

So, now you can see by some how we are able to calculate this scalar quantities y_r , where r ranging from 1 to n , then immediately will be having solution to this problem, because $y(x) = f(x) + \lambda \sum_{r=1}^n p_r(x) y_r$. In order to find this solution, we can do one thing q_m where r ranging from 1 to n this is known, we can multiply both side of this equation by q_m where m is taking any value within the range 1 to n .

(Refer Slide Time: 38:51)

$$y(x)q_m(x) = f(x)q_m(x) + \lambda \sum_{r=1}^n y_r p_r(x)q_m(x), \quad 1 \leq m \leq n$$

$$\int_a^b y(x)q_m(x)dx = \int_a^b f(x)q_m(x)dx + \lambda \sum_{r=1}^n y_r \int_a^b p_r(x)q_m(x)dx$$

$$b_m = \int_a^b f(x)q_m(x)dx \quad \alpha_{mr} = \int_a^b p_r(x)q_m(x)dx$$

$$y_m = b_m + \lambda \sum_{r=1}^n \alpha_{mr} y_r, \quad m = 1, 2, 3, \dots, n$$

$$\begin{bmatrix} y_1 \\ y_2 \\ \vdots \\ y_n \end{bmatrix} = \begin{bmatrix} b_1 \\ b_2 \\ \vdots \\ b_n \end{bmatrix} + \lambda \begin{bmatrix} \alpha_{11} & \alpha_{12} & \dots & \alpha_{1n} \\ \alpha_{21} & \alpha_{22} & \dots & \alpha_{2n} \\ \vdots & \vdots & \ddots & \vdots \\ \alpha_{n1} & \alpha_{n2} & \dots & \alpha_{nn} \end{bmatrix} \begin{bmatrix} y_1 \\ y_2 \\ \vdots \\ y_n \end{bmatrix}$$

So, therefore, we can write $y(x)q_m(x)$ this is equal to $f(x)q_m(x)$ plus $\lambda \sum_{r=1}^n y_r p_r(x)q_m(x)$, this is the result we are getting by multiplying q_m ; where $1 \leq m \leq n$. And then integrating from the range a to b , we can find $\int_a^b y(x)q_m(x)dx$ this is equal to $\int_a^b f(x)q_m(x)dx$ plus $\lambda \sum_{r=1}^n y_r \int_a^b p_r(x)q_m(x)dx$, this one.

Now, we need two notations for $\int_a^b f(x)q_m(x)dx$ and $\int_a^b p_r(x)q_m(x)dx$, if we denote by b_m that is the $\int_a^b f(x)q_m(x)dx$ this is the definition for b_n , and this $\int_a^b p_r(x)q_m(x)dx$ this is defined by α_{mr} this one, then this result that is $\int_a^b y(x)q_m(x)dx = \int_a^b f(x)q_m(x)dx + \lambda \sum_{r=1}^n y_r \int_a^b p_r(x)q_m(x)dx$, can be written as y_m is equal to $b_m + \lambda \sum_{r=1}^n \alpha_{mr} y_r$. Now, when you multiplied the expression $y(x)q_m(x) = f(x)q_m(x) + \lambda \sum_{r=1}^n y_r p_r(x)q_m(x)$, then we have mentioned that m is ranging from 1 to n .

So, that means, we can find these type of n equations which have given by $y_m = b_m + \lambda \sum_{r=1}^n \alpha_{mr} y_r$, where $m = 1, 2, 3, \dots, n$ and therefore, we are having a system of equations which can be written as y_1, y_2, \dots, y_n , that is into a matrix form this is equal to b_1, b_2, \dots, b_n plus λ multiplied by $\alpha_{11}, \alpha_{12}, \dots, \alpha_{1n}$, then $\alpha_{21}, \alpha_{22}, \dots, \alpha_{2n}$, up to $\alpha_{n1}, \alpha_{n2}, \dots, \alpha_{nn}$.

2 n proceeding this way, last row will be alpha n 1, alpha n 2, up to alpha n n; this multiplied with y 1, y 2 up to y n, now this matrix equation is nothing but, a system of linear equation.

(Refer Slide Time: 43:49)

The image shows a whiteboard with the following handwritten mathematical content:

$$Y = \begin{bmatrix} y_1 \\ y_2 \\ \vdots \\ y_n \end{bmatrix}, \quad B = \begin{bmatrix} b_1 \\ b_2 \\ \vdots \\ b_n \end{bmatrix}, \quad A_{n \times n} = \begin{bmatrix} a_{11} & a_{12} & \dots & a_{1n} \\ a_{21} & a_{22} & \dots & a_{2n} \\ \vdots & \vdots & \ddots & \vdots \\ a_{n1} & a_{n2} & \dots & a_{nn} \end{bmatrix}$$

$$Y = B + \lambda AY$$

$$\Rightarrow (I_n - \lambda A_{n \times n})Y = B$$

$$|I_n - \lambda A_{n \times n}| \neq 0$$

If we introduce the notations that is capital Y is equal to y 1 y 2 up to y n this one, then capital B equal to this column matrix b 1, b 2 up to b n, and capital A which is an n cross n matrix, this stands for alpha 1 1, alpha 1 2 up to alpha 1 n, in this way alpha 2 1, alpha 2 2, up to alpha 2 n finally, alpha n 1, alpha n 2 up to alpha n n this is a n cross n matrix. And therefore, the matrix equation can be written as Y equal to B plus lambda A Y, now this Y is simply rewritten as I n times Y that is identity matrix, so that means, from here we are having a system of equation I n minus lambda, where a is an n cross n matrix, this matrix multiplied with Y this is equal to capital B.

So, if this matrix I n minus lambda A n cross n is invertible, then we will be having unique solution, so that means, whenever determinant of I n minus lambda A n cross n this is not equal to 0, then we will be having unique solution. And if this is equal to 0, that means, if determinant I n minus lambda A n cross n equal to 0, then we will be having either infinite number of solution or no solution, that we will be discussing the next lecture.

But, the point is that if we are able to find out some hallows of lambda, such that this determinant is non 0, so therefore, we can find unique solution for this system of linear

equations, and once we are able to find out unique solutions y_1, y_2, y_3 up to y_n , these has the unique solutions, so then the expression $y = B + \lambda \int_a^b K(x,s)y(s)ds$ is uniquely determined and this is nothing but, the solution of the given Fredholm integral equation (Refer Slide Time: 46:31).

(Refer Slide Time: 46:57)

The image shows a whiteboard with handwritten mathematical derivations. At the top, it defines a vector $Y = \begin{bmatrix} y_1 \\ y_2 \\ \vdots \\ y_n \end{bmatrix}$, a vector $B = \begin{bmatrix} b_1 \\ b_2 \\ \vdots \\ b_n \end{bmatrix}$, and a matrix $A_{n \times n} = \begin{bmatrix} a_{11} & a_{12} & \dots & a_{1n} \\ a_{21} & a_{22} & \dots & a_{2n} \\ \vdots & \vdots & \ddots & \vdots \\ a_{n1} & a_{n2} & \dots & a_{nn} \end{bmatrix}$. Below this, it shows the equation $Y = B + \lambda AY$, which is rearranged to $(I_n - \lambda A_{n \times n})Y = B$. A condition $|I_n - \lambda A_{n \times n}| \neq 0$ is noted. An example is given: $y(x) = x e^x - x + \int_0^1 x y(s) ds$. For this example, $f(x) = x e^x - x$, $k(x,s) = x = p_1(x) q_1(s)$, $p_1(x) = x$, and $q_1(s) = 1$.

So, now, we consider one example, here we consider the example, we can solve by this method, this is a very simple example, that is $y(x) = x e^x - x + \int_0^1 x y(s) ds$, so just for your understanding this $f(x)$ is equal to as usual $x e^x - x$. Now, kernel $K(x,s)$ this is equal to x , so therefore, this will be equal to as per our notation $p_1(x) q_1(s)$ where $p_1(x)$ this is equal to x and $q_1(s)$ this is equal to 1, so now, if we just solve this equation by the method.

(Refer Slide Time: 47:56)

The image shows a whiteboard with the following handwritten mathematical steps:

$$y(x) = x e^x - x + x \int_0^1 y(s) ds = x e^x - x + x y_1$$

$$y(x) = x e^x - x + x y_1$$

$$\int_0^1 y(x) dx = \int_0^1 (x e^x - x) dx + y_1 \int_0^1 x dx$$

$$\Rightarrow y_1 = [x e^x - e^x]_0^1 - \left[\frac{x^2}{2}\right]_0^1 + y_1 \left[\frac{x^2}{2}\right]_0^1$$

$$\Rightarrow y_1 = [e - e + 1] - \frac{1}{2} + \frac{1}{2} y_1$$

$$\Rightarrow y_1 = 1$$

$$y(x) = x e^x - x + x \cdot 1 = x e^x$$

We have just discussed, this becomes $y(x)$ is equal to $x e^x - x + x y_1$. So, with our notation that we have introduced this is equal to $x e^x - x + x y_1$, so that means, this $x y_1$ is actually continuation from the expression $\int_0^1 y(s) ds$ from 1 to y_1 .

So, in that stage we have multiplied both side by $q(x)$ and then we have integrated, here we have only one q that is $q(x)$, so $q(x)$ is going to be 1, so that means, we have to integrate this result $y(x)$ equal to $x e^x - x + x y_1$ both sides with respect to x .

So, we are multiplying this expression both sides with respect to x from 0 to 1, means we are actually multiplying this equation by $q(x)$ and then integrating from 0 to 1, so therefore, we are having $\int_0^1 y(x) dx$ this is equal to $\int_0^1 (x e^x - x) dx + y_1 \int_0^1 x dx$. So, $\int_0^1 y(x) dx$ is our y_1 , so this y_1 is equal to after integration it will be $x e^x - x$ limit from 0 to 1, then minus $\frac{x^2}{2}$ limit 0 to 1 plus $y_1 \frac{x^2}{2}$ limit 0 to 1.

So, from here we will be having y_1 this is equal to $e - e + 1$, this two things are coming from the upper limit, then minus $x e^x$ at x equal to 0 is 0 and then from here, we will be having this is equal to $-e + 1$ **sorry**, this will be actually the

result of integration will be $x e^{-x}$ (Refer Slide Time: 50:49). So, therefore, $e^{-x} + 1$ then from here you will be having $\frac{1}{2} e^{-x}$, so this e^{-x} cancels with e^{-x} this is half, this will go on the right hand side, so ultimately you will be having $y = \frac{1}{2} e^{-x}$.

So, with $y = \frac{1}{2} e^{-x}$ if you substitute on the first line, then we can find $y' = -\frac{1}{2} e^{-x}$ that is $y' = -\frac{1}{2} e^{-x}$ equal to $x e^{-x} + \frac{1}{2} e^{-x}$, so this is equal to $x e^{-x}$, so by calculating this $y = \frac{1}{2} e^{-x}$ we have obtained $y = \frac{1}{2} e^{-x}$ as a solution.

So, that means, what we have discussed today, that in Fredholm integral equation which are of non homogeneous type, non homogeneous Fredholm integral equation with separable kernel that can be converted into a system of linear equations. And here we have considered a simple example, where we have obtained a unique solution and for a specific value of λ , now in case of this separable kernel this integral equation can be converted into a problem of finding solution for a system of linear equation.

And depending upon uniqueness of the solution of the system of linear equation, which is actually in turns depending up on the magnitude of λ this there may be unique solution, may be no solution, may be infinite number of solution will be having corresponding conclusion for the solution of the Fredholm integral equation.

And in next few lectures, we will try to relate these idea with the concept of resolvent kernel, where this resolvent kernel can be obtained in a unique fashion or not and those theories are actually Fredholm theory for solving integral equation which are known as actually Fredholm integral equation.

And where we will be discussing, three particular theorems of Fredholm and after discussing some other problems of these type where the integral equation of Fredholm integral type with separable kernel can be converted into linear system of linear equation; and by solving those equation will be discussing the rest of the theory for Fredholm integral equation, that is Fredholm theorem one, Fredholm theorem two, Fredholm theorem three and there is actually one **one** correspondence between existence, if unique solution and non existence of the solutions.

So, today I can stop at this point in the next lecture we will be considering few more example of these type, and with help of a particular example, we can try to understand

how this type of situation comes into the picture that these may have unique solution, may not have unique solution, and in case of this problem does not possess unique solution, what will be the to our solution for the Fredholm integral equation, so thank you for your attention.