

**Introduction to Methods of Applied Mathematics**  
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**Module No # 06**  
**Lecture No # 26**  
**Fourier Series and Fourier Transform**

Welcome to all of you in the next class of this course so in yesterday lecture we were discussing about Fourier series so let me continue with the same topic what is Fourier series?

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Lecture 29

Fourier Series:

$L_2(0, 2l) \rightarrow f \in L_2(0, 2l)$   
 $\int_0^{2l} |f(x)|^2 dx < \infty$

$f(x) = f(x+2l) \quad \forall x \in \mathbb{R}$

$f(x) = \sum_{k=-\infty}^{\infty} c_k \cdot e^{ikx}$

$l = \pi$   
 $= \sum_{k=-\infty}^{\infty} c_k e^{ikx}$

$\{e^{ikx}\} \rightarrow L_2(0, 2\pi)$   
 $h \in \mathbb{Z}$   
 $\omega_k(x) = e^{ikx}$

$\langle fg \rangle = \int_0^{2l} f(x) \overline{g(x)} dx$

$\|f\|_2 = \sqrt{\int_0^{2l} |f(x)|^2 dx}$

Fourier series so in last lecture we were taking function which belongs to this space. So let us have some knowledge about this space this is a set of a this is a space of all square integrable function it means if f belongs to L2 0 and dx is less than infinite okay. So more over one more condition which we have written in this space is means this function which belongs to this space is a periodic function of period 2l. All of us know what is the periodic function so it means f of x will be f of x +2l for all x belong to R.

Now in the last lecture I have written this as a Fourier series okay where ck's are called Fourier coefficients. So if you look at just for simplicity I am considering l = Pi in that case this expansion will become ikx okay. So here you look at one thing that e to the power ikx gives you an orthonormal basis for space 0 to 2 Pi for different k belongs to Z. Or here basically we are writing k is equal to 1 infinity sorry minus infinity to infinity okay k is minus infinity to infinity.

So if you look at this space more carefully how we are generating this space if it denote this function by this so basically this is let us say  $w(x)$  so this is a dilation of a simple one element what is that one element that one element is  $e^{i\alpha x}$ . So we are generating whole space by dilation of a single function which is this function okay. So this function is by dilating this functions we are generating the whole space  $L^2(0, 2\pi)$ .

So you should notice that there is a dilation property later on when we will connect this to the this Fourier series to the wavelet series so you will see that dilation will also be in the wavelet series. So at that point of time we will continue with this dilation property and that is the motivation that right now I am mentioning to that property okay. So now as this  $L^2(0, 2\pi)$  is equipped with norm and inner product what is the inner product on this space?

Inner product is given by this definition okay this is a inner product equipped on this space norm is this space. So this is the definition of an norm on this space this is the definition of a inner product on this space. Now the next question comes why we are considering function belongs to this space okay that is the questions I can so now can I write a Fourier series for a function which is not belongs to this space? Of course you can write down Fourier series.

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$f(x) = \sum_{k=-\infty}^{\infty} c_k e^{ikx}$   
 $L_2$  sense of convergence:  
 $S_N(x) = \sum_{k=-N}^N c_k e^{ikx}$   
 $\int |f(x) - S_N(x)|^2 dx \rightarrow 0$  as  $N \rightarrow \infty$

**Gibbs Phenomena**  
 $f(x) = \begin{cases} x & 0 \leq x \leq 0.5 \\ x-1 & 0.5 < x \leq 1 \end{cases}$   
 $N=15$   
 $N=32$   
 $N=64$   
 $x=0.5$

But the questions comes what will be because it is a infinite series  $f$  of  $x$  you know it is a infinite series okay. So it is a infinite series may converge or may not converge so and when we talk

about the converge of infinite series the notions of the convergence are also different they are at different notions of the convergence of a infinite series uniform convergence point wise convergence L2 sense convergence.

So I am now mentioning all of that notion of a convergence in details but what I am doing right now I will tell you what is L2 sense of convergence? So for that reason I am denoting the partial sum okay so this series is set to converge in L2 sense provided this property this tends to 0 as  $N$  tends to infinity. This is the meaning of L2 sense of convergence so basically error is  $f(x) - s_n(x)$  is the error between the function and it is partial sum which is tending to 0 in L2 that is meaning which you could observe from the given formula.

So that is the reason and so if function belongs to this space  $L^2(0, 2\pi)$  only then we can guarantee that Fourier series will converge at least in  $L^2$  sense. This is basically a theorem then proof of that theorem can be seen in any standard partial differential equation text we are Fourier series is explained in detail way so okay. I have given you very explanation this  $L^2$  sense of convergence because if you are writing a infinite series you should be sure in what sense it is converging okay.

So now the one of the very I should say disadvantage of Fourier series I am going to explain and what is that disadvantage is Gibbs phenomena okay. So what is this Gibbs phenomena let me explain you with the help of a one function is given like that okay this is the function if you want to look at the graph this will be  $x$  so 0.5, 0, 1 okay what you see in how if I wanted to approximate this function what will I do? I will approximate this function with the help of this partial sum  $s_n(x)$  which is given by this formula okay for different values of  $n$ .

So for different values of  $n$  means  $n = 16$  I can try for a  $n = 32$   $n = 64$  so that you can also implement in a Matlab by Fourier approximations and once you compute  $s_n(x)$  for  $n = 16, 32, 64$  what you will observe that as roughly I can give a plot for  $n = 16$  this will be like this for  $n = 64$  this will be more closer okay. So what as you increasing the  $n$  what will be the observation that wherever the function is as some smoothness it is okay function this approximation will be close to the approximation of the function.

But if function is non smooth near at  $x = 0.5$  so at this point how large and you take some oscillation will always be there and this will be spreaded throughout the domain also. Means oscillation will be less for let us say  $n = 1000, 2000$  but you will be not be able to remove this oscillation at any cost okay and why this happening? Because this is happening because of the non-smoothness this continuity at  $x = 0.5$  okay.

So this should also be 0.5 or that also less than sin should also be okay because at  $X = 0.5$  because there is a discontinuity in this function at  $x = 0.5$ . So when we are taking a infinite sum of this exponential functions then we are able to resolve this discontinuity that is what this Fourier series says but once I am working with a finite some  $k = -n$  to  $n$  we are not able to resolve this discontinuity proper and because you know  $f(x)$  is a discontinuous function and we are writing that function with a help of a sin and Cosine exponential i.e.  $x$  is sin and Cosine which are continuous functions.

So finite some of it continuous functions cannot given you this continuous function and that is one of the main drawback of the Fourier series which is called in terms of a Gibbs phenomena okay. So this is Gibbs phenomena I have already explained you so now what we do that I can tell you as  $n$  tends to infinity it will be more closer. So this as I am saying this is a disadvantage of a Fourier series and later on I will show you that this disadvantage is not in wavelet series that is why wavelet series is better than Fourier series.

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$$\int_0^{2\pi} |f(x)|^2 dx = 2\pi \sum_{k=-\infty}^{\infty} |c_k|^2$$
 Parseval relation

$l \rightarrow \infty$

Fourier series (FS)

Fourier integral

Complex form of (FI) is called, Fourier Transform (FT)

Fourier transform?

$$\hat{f}(\omega) = \int_{-\infty}^{\infty} e^{-i\omega x} f(x) dx \quad f \in L_1(\mathbb{R})$$

$$|\hat{f}(\omega)| \leq \int |f(x)| dx = \|f\|_1$$

Fourier spectrum

Because this Gibbs phenomena is completely removed in wavelet series okay so that is and now one of the other thing which Fourier coefficient satisfy this relation which is called Parseval relation which is called Parseval relation this is also very famous relations if you okay because if you remember earlier we were also talking some conditions but this is also one of the like if you remember when we were discussing orthonormal basis, Riesz basis we were writing some condition  $c_k$  is basically when  $f$  comma  $f$  inner product of  $f$  with  $v_k$  which  $v_k$ 's are exponential of  $ikx$  okay.

So this is also what is the coefficient of Fourier series are satisfying this relation is called Parseval relation. So now from Fourier series I was writing for a function which has a period to  $l$  but now  $L$  tends to infinity means period tends to infinity Fourier series is converted to a Fourier integral okay. And complex form of Fourier integral is called Fourier transform complex form of Fourier integral in short form is called Fourier transform.

Again in short form Fourier transform can be written with FT and Fourier series can be written with FS. So now let me define what is Fourier transform because first I will define you Fourier series, Fourier transform and then wavelet series, wavelet transform that is the sequence because only then you will be able to appreciate the beauty of wavelet analysis over Fourier analysis. So now what is Fourier transform? Now I am defining a Fourier transform by this formula for  $f$  belongs to  $L^1\mathbb{R}$ .

We have already seen this kind of a spaces  $L^1\mathbb{R}$   $L^2$   $L^p\mathbb{R}$  for  $p$  general greater than or equal to 1 so I do not to redefine again what is this  $L^1\mathbb{R}$  spaces. So this is a Fourier transform a function  $f_x$  which belongs to  $L^1\mathbb{R}$  this is a proper integral so if I am talking about infinite series improper integral I also have to talk about convergence of them. So for that reason I am taking  $f$  belongs to  $L^1\mathbb{R}$  because if  $f$  belongs to  $L^1\mathbb{R}$  this Fourier transform always exist that I can show you also what will be this thing?

This thing will be so this is basically if you look at this is the this [norm](#) so if  $f$  belongs to  $L^1$  this is less than infinitive so it means Fourier transform exists this improper integral converges. So that is why we are working out with  $f$  belongs to  $L^1\mathbb{R}$  now we will study first of all let me give you physical interpretations why we study Fourier transform why we study sometime we want to

analyze functions for mathematicians signal for engineers in a domain with respect to the frequency okay.

Because you know what Fourier transform is doing Fourier transform is decomposing the function into the various waves. So sometime if you want to analyze the function in different domain and that what is the different domain that different domain is called frequency domain and then Fourier transform is also sometime called Fourier spectrum okay Fourier transform is called also Fourier spectrum.

We analyze what is the need of it we analyze the functions in a different domain which is called a frequency domain and Fourier transform is also called after sometime Fourier spectrum. So now we will study different properties of different transform which will be used subsequently.

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$$\begin{aligned} \hat{f}'(\omega) &= \int_{-\infty}^{\infty} e^{-i\omega x} \cdot f'(x) dx \\ &= e^{-i\omega x} \cdot f(x) \Big|_{-\infty}^{\infty} - \int_{-\infty}^{\infty} (-i\omega) e^{-i\omega x} f(x) dx \\ &= (i\omega) \int_{-\infty}^{\infty} e^{-i\omega x} f(x) dx \\ &= i\omega \hat{f}(\omega) \end{aligned}$$

$$\Rightarrow \hat{f}'(\omega) = i\omega \hat{f}(\omega)$$

$$\hat{f}(\omega) \in L_1(\mathbb{R})$$

$$(F^{-1}\hat{f})(x) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{i\omega x} \hat{f}(\omega) d\omega$$

So if I say  $f'(x)$  is also in  $L^1(\mathbb{R})$  then one of the property I can prove this this what is this? e to the power  $-i\omega x$   $f'(x) dx$  that is why chosen this  $f'(x)$  is  $L^1(\mathbb{R})$  because only then I can define this I can talk about the existence of a Fourier transform. So if I use integration by parts so we are using integration by parts where so this is  $i\omega f(x) dx$  okay this is an  $L^1(\mathbb{R})$  functions so  $L^1(\mathbb{R})$  function always tends to 0 at  $x$  tends to plus minus infinity that is the property of a function which belongs to that basis.

So basically this is  $i\omega \hat{f}$  of a okay so in a nutshell this is a famous property of the Fourier transform from this property we can deduce one more thing that effect of  $\omega$  tends to 0 as  $\omega$  tends to plus minus infinite this can be proved directly from here because once you write down this in the left hand side this is very much clear this property we will use later on that is why I have defined here.

And the next property of this Fourier transform is yes this is now if we know from one domain we are going to the frequency domain we should also know how to come back otherwise there is a no point of defining some integral transform. Fourier transform is also integral transform as you must have already studied in earlier. So now we are defining inverse Fourier transform so inverse Fourier transform we are defining by this formula provided  $\hat{f}$  is  $\omega$  is also in  $L^1 \mathbb{R}$ .

Because only then we can again talk the convergence of a this improper integral okay we can talk that is why we can choosing  $\hat{f}$   $\omega$  belongs to  $L^1 \mathbb{R}$ . It is not necessary that if  $f_x$  belongs to  $L^1 \mathbb{R}$   $f$  at  $\omega$  will also be  $L^1 \mathbb{R}$  that I can show by one example.

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$$f(x) = \begin{cases} e^{-x} & x \geq 0 \\ 0 & \text{otherwise} \end{cases}$$

$$\hat{f}(\omega) = \int_{\mathbb{R}} e^{-i\omega x} \cdot e^{-x} dx$$

(integral is not over  $\mathbb{R}$  Its from 0 to infinity)

$$= \int_{\mathbb{R}} e^{-(i\omega+1)x} dx$$

$$= \left[ \frac{e^{-(i\omega+1)x}}{-(i\omega+1)} \right]_{-\infty}^{\infty}$$

$$= \frac{1}{1+i\omega}$$

$$\int \left| \frac{1}{1+i\omega} \right| d\omega$$

$$f(x) \in L^1(\mathbb{R}) \Rightarrow \hat{f}(\omega) \in L^1(\mathbb{R})$$

I am taking one function 0 otherwise okay this is the function I am taking so I am now I am defining what will be Fourier transform for this function with this will be  $1 + 1 + i\omega$  okay. So if you look at now we have define whether this function is  $L^1 \mathbb{R}$  of not this is the Fourier transform of this function so I have to choose whether this function is in not in  $L^1 \mathbb{R}$ . So I will I

can take this function separately. Suppose this is not an L1R function okay so I have given you this by one example that if function is L1R does not necessarily means  $f$  at  $\omega$  may not be L1R okay.

So that is why we are we whenever we were defining a inverse Fourier transform separately we were making this assumption. So now we are looking at the Fourier transform of a one anther functions which all of you must have seen earlier and that function is called as a Gaussian function.

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The image shows a handwritten derivation of the Fourier transform of a Gaussian function. The steps are as follows:

$$f(x) = e^{-ax^2} \quad a > 0 \quad (\text{Gaussian function})$$

$$\hat{f}(\omega) = \int_{-\infty}^{\infty} e^{-ax} \cdot e^{-i\omega x} dx = \int_{-\infty}^{\infty} e^{-a(x + \frac{i\omega}{2a})^2 - \frac{\omega^2}{4a}} dx$$

$$= e^{-\frac{\omega^2}{4a}} \int_{-\infty}^{\infty} e^{-a(x + \frac{i\omega}{2a})^2} dx$$

$$= \frac{e^{-\frac{\omega^2}{4a}}}{\sqrt{a}} \int_{-\infty}^{\infty} e^{-t^2} dt$$

The final result is underlined:  $\hat{f}(\omega) = \frac{\sqrt{\pi}}{a} e^{-\frac{\omega^2}{4a}}$

This is called Gaussian functions we are defining a Fourier transform of this function. So how I can write down  $ax + i\omega$  I can always add and subtract something so using this terminology I have done that. So now I can take this term common okay so what now we have to simplify this integral as well whole thing  $x + i\omega / 2a$  whole square I can take  $t$  so then after substituting this we will get  $e$  to the power  $-t$  square  $dt$  okay.

So this is now we have to calculate this integral we all know this is under root  $\pi$  but if you will ask me how to calculate because we do know the **anti** derivative of this function that is why in that case some special techniques are device to calculate this kind of a integral and we will based on Taylor's series etc., that is just for your knowledge. So  $f$  **hat**  $\omega$  will be under root  $\pi / a$   $e$  to the power minus  $\omega$  square by  $4a$  okay.



So this will be the Fourier transform of a Gaussian function sorry this under root sin will be because this value is under root Pi by under root a. So what do you observe specifically about the Fourier transform of a Gaussian can you think for a while the main point is that Fourier transform of a Gaussian function is also in the form of a Gaussian this is also a form of a Gaussian that is the specific point to observe which we will exploit later on that Fourier and we will work out with this Gaussian function because the Fourier transform of Gaussian is again a Gaussian function.

So this is a important point which all of you should note for a future lectures now I am coming to define a another concept which is used in some of the analysis which is called convolutions okay. So I am defining what is called convolution?

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Convolution:

$$(f * g)(x) = \int_{-\infty}^{\infty} f(t)g(x-t) dt$$

$f, g \in L_1(\mathbb{R}) \quad f * g \in L_1(\mathbb{R})$

$$\int |(f * g)(x)| dx = \int \left| \int f(t)g(x-t) dt \right| dx < \infty$$

$$\|f * g\|_1 \leq \|f\|_1 \|g\|_1$$

$f * g = \mathcal{F}^{-1} \{ \mathcal{F}f \mathcal{F}g \}$

$$(f * g)^\omega = (f^\omega) (g^\omega)$$

$\forall f, g \in L_1(\mathbb{R})$

$$f^\omega = \frac{(f * g)^\omega}{g^\omega}$$

Okay convolution of a two functions is defined in this way okay so what we are doing we are multiplying the function f with the function g not exactly with delayed by parameter tau sorry this should be tau okay and then we are multiplying this function with that delayed by tau to the another function g and then we are integrating with respect to this delay parameter. This is how convolution in this defined convolution is very important it has lot of applications in electrical engineering probabilities [statistics](#) etc., many applications and I assume all of you must have a heard about this already.

So this is how we define convolution but again this is a improper in integral so I have to talk when it exist. So for that reason I am saying if  $f$  and  $g$  belongs to  $L^1\mathbb{R}$  then convolution will be also be a  $L^1\mathbb{R}$  this is just a 1 line proof because this will be equal to  $f$  of  $\tau$   $g$  of  $x - \tau$   $d\tau$   $dx$ . Now using Fubini's theorem you can very well because  $f$  and  $g$  both are in  $L^1\mathbb{R}$  you can very well prove that  $f$  and  $g$  will also be  $L^1\mathbb{R}$  it means  $f$  and  $g$  norm as sorry this is a convolution okay.

So the next property of the convulsion which we can prove that convolution is also commutative that is also this you can take it a exercise and you can prove it by change of variables that convolution is a commutative property for all  $f$  and  $g$  belongs to  $f$  and  $L^1\mathbb{R}$  it is also distributive will be equal to  $f$  star of  $g$  of  $h$  for all  $fgh$  belongs to  $L^1\mathbb{R}$ . Now the next question is identity operation is also exist for a convolution it means  $f$  for all  $f$  belongs to  $L^1$  does identity exist for a convolution operator this is my next question okay.

So can you think for a while for this does identity exist for a convolution operations you think in the mean time I am proving one another theorem which I will be needing to answer to this questions does identity exists for a convolution operation and that theorem is called convolution theorem.

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Convolution theorem:  $f, g \in L^1(\mathbb{R})$  then  $(f \star g)^\wedge(\omega) = \hat{f}(\omega) \hat{g}(\omega)$

$$\begin{aligned} (f \star g)^\wedge(\omega) &= \int e^{-i\omega x} (f \star g)(x) dx \\ &= \int e^{-i\omega x} \left[ \int f(\tau) g(x-\tau) d\tau \right] dx \quad \text{Fubini th.} \\ &= \int f(\tau) \left[ \int e^{-i\omega x} g(x-\tau) dx \right] d\tau \quad \text{change of variables } x-\tau=y \\ &= \int f(\tau) \left[ \int e^{-i\omega(\tau+y)} g(y) dy \right] d\tau \\ &= \int f(\tau) e^{-i\omega\tau} \left[ \int e^{-i\omega y} g(y) dy \right] d\tau \\ &= \hat{f}(\omega) \hat{g}(\omega) \end{aligned}$$

Convolution theorem what this theorem says? This theorem says if  $f$  and  $g$  belongs to  $L^1\mathbb{R}$  then this will always be the case i will prove this statement okay. Now I am using Fubini's theorem

of  $\tau dx$  into  $d\tau$  now I am using a change of variable  $x - \tau = y$  from this step to this step we have using the Fubini's theorem how order of integration basically in a nut shell this theorem says that how order of integration can be changed okay.

So this I can split into 2 terms and then  $f$  of  $\tau e$  to the power  $-i\omega\tau + i\omega y$   $g(y) dy$  and  $d\tau$ . So basically you can see this is equal to  $f$  at  $\omega g$  at  $\omega$  okay. So we have proved convolution theorem it is a very beautiful theorem because it has many applications one of the application is it helps in computing the convolutions with the very reduced complexity because otherwise how will you compute the convolution? You will compute this convolution with the this definition which I have given you here okay but with the help of this convolution theorem what you are doing?

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$$= \int f(\tau) \cdot e^{-i\omega\tau} \left[ \int e^{-i\omega y} g(y) dy \right] d\tau$$

$$= \hat{f}(\omega) \hat{g}(\omega)$$

$f * g(x) = \text{inverse Fourier transform} \left( \hat{f}(\omega) \hat{g}(\omega) \right)$

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$\exists f \in L^1(\mathbb{R}) \quad f * e = f \quad \forall f \in L^1(\mathbb{R})$

$\therefore \hat{f}(\omega) \hat{e}(\omega) = \hat{f}(\omega) \implies \hat{e}(\omega) = 1 \quad e^{\hat{}}(\omega) \rightarrow 0 \text{ as } \omega \rightarrow \pm\infty$

I can compute the convolution with this inverse Fourier transform of  $f$  at  $\omega g$  at  $\omega$  okay. So complexity reduced if I am computing the inverse Fourier transform of multiplication of 2 function  $f$  at  $\omega$  and  $g$  at  $\omega$  and then I am computing the conversions instead of going with the direct definitions. Okay so and convolutions as many application and we will using this convolution theorem to support my answer of my previous question that does identity exist for a convolution operations okay.

So how because now in the left hand side I will be using convolution theorem so it means I am taking the Fourier transform of this both sides so this will give me this by using convolution

theorem which implies  $e$  at  $\omega$  is 1 so is it possible? Can you think is it possible? No, existence of identity means they exist  $e$  belongs to  $L^1\mathbb{R}$  existence of identity means there should exist  $e$  which belongs to  $L^1\mathbb{R}$  which satisfy this property but here we are saying  $e$  at  $\omega$  is 1 which is not true why is it not true?

Because if any function which belongs to  $L^1\mathbb{R}$  is space it should tends to 0 at  $x$  tends to plus minus infinity. So means if  $e$  at  $\omega$  is 1 it means if you remember I have used not I have used I have proved 1 property here  $f$  at  $\omega$  should tends to 0 as  $\omega$  tends to plus minus infinity okay. For a function which belongs to  $L^1\mathbb{R}$  so they are  $e$  belongs to  $L^1\mathbb{R}$  so Fourier transform should and because of that reason I proved this property as I said I will use this property later on.

So Fourier transform of this should tends to 0 as  $\omega$  tends to plus minus infinity but they are it is just a minute but you look here this is not this according to the previous property of a Fourier transform this should tends to 0 as  $\omega$  tends to plus minus infinity. But this is contradictory so it means [their](#) does not exist convolution identity or we have to say something more how we can approximate this convolution identity and with this words.

I am closing my today's lecture and whatever [exercise](#) I have asked you to prove about convolution operations you could look at the proof of them because it really helps and in next lecture what we will do we will tell you how to work out with the approximation of the convolution identity okay thank you very much with this I am closing my lecture.