

Fourier Analysis and its Applications
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11 The Poisson kernel

We still have to show that the solution obtained in the last slide does attain the value $f(\theta)$ on the boundary. Here are some exercises

Exercises:

1. Show that the Poisson kernel is non-negative and

$$\frac{1}{2\pi} \int_{-\pi}^{\pi} \Pi_r(\theta - t) dt = 1. \quad (2.19)$$

2. Show that

$$\lim_{r \rightarrow 1^-} |u(re^{i\theta}) - f(\theta)| = 0. \quad (2.20)$$

Hint: Write $f(\theta)$ as the integral w.r.t over $[-\pi, \pi]$ of $f(\theta)\Pi_r(\theta - t)/(2\pi)$. Now let $\epsilon > 0$ and I be an interval of length ϵ centered at θ . The integral over I is small for one reason and the integral over $[-\pi, \pi] - I$ is small for a different reason. We shall return to these shortly.

Poisson formula for a ball There is a corresponding result for the ball in \mathbb{R}^3 but to derive that we need to spend a little time with associated Legendre equations.

Theorem: Suppose given a continuous function $f(\mathbf{x})$ on the unit ball B centered at the origin in \mathbb{R}^3 then the solution of the boundary value problem

$$\Delta u = 0 \text{ on } B, \quad u|_{\partial B} = f$$

is given by

$$u(\mathbf{y}) = \int_{\partial B} \frac{(1 - \|\mathbf{x}\|^2)f(\mathbf{x})dS(\mathbf{x})}{(1 + \|\mathbf{x}\|^2 - 2\|\mathbf{x}\| \cos \alpha)^{3/2}}, \quad (2.21)$$

where α is the angle between \mathbf{x} and \mathbf{y} .

For the proof see page 180 of *G. B. Folland, Fourier Analysis and its applications, Amer. Math. Soc, Indian Reprint by Universities Press, New Delhi, 2012.*

Solving the heat equation We discuss this as series of solved exercises:

- (i) Let $u(x, t)$ be a smooth solution of the heat equation in the upper half-plane $t \geq 0$ such that $u(x, t) = u(x + 2\pi, t)$ for all $x \in \mathbb{R}$, $t \geq 0$. Show that the energy $\int_{-\pi}^{\pi} (u(x, t))^2 dx$ is a monotone decreasing function of time. Prove the same result if the integral is over any interval of length 2π . Deduce that a smooth 2π -periodic solution of the heat equation with a given initial condition, if it exists, is unique.

Solution: Call the integral $E(t)$ which really represents the *Energy* and the method itself is called the *Energy Method*.

$$E'(t) = 2 \int_{-\pi}^{\pi} u(x, t) u_t(x, t) dx = 2 \int_{-\pi}^{\pi} u(x, t) u_{xx}(x, t) dx$$

Integrate the last equation by parts and we get $E'(t) = -2 \int_{-\pi}^{\pi} (u_x(x, t))^2 dx \leq 0$.

So we get

$$0 \leq E(t) \leq E(0).$$

Now if the initial value $u(x, 0) = 0$ then $E(0) = 0$ and so $E(t) = 0$ for all $t \geq 0$ and we infer that the solution is identically zero. From this the uniqueness follows.

We now turn to existence.

- (ii) Suppose that $f(x)$ is a 2π -continuous periodic function. Further conditions will be imposed as we go along. The solution of $u_t = u_{xx}$ is sought in the form

$$u(x, t) = a_0(t) + \sum_{n=1}^{\infty} (a_n(t) \cos nx + b_n(t) \sin nx) \quad (2.22)$$

Let us substitute (2.22) into the heat equation $u_t = u_{xx}$ and we see that

$$a'_0(t) = 0, \quad a'_n(t) = -n^2 a_n(t), \quad b'_n(t) = -n^2 b_n(t) \quad (2.23)$$

$$a_0(t) = \alpha_0, \quad a_n(t) = \alpha_n e^{-n^2 t}, \quad b_n(t) = \beta_n e^{-n^2 t}.$$

Let us put $t = 0$ in equation (2.22) and we see that

$$f(x) = \alpha_0 + \sum_{n=1}^{\infty} (\alpha_n \cos nx + \beta_n \sin nx) \quad (2.24)$$

which is the Fourier expansion of the function $f(x)$ and can be determined easily since $f(x)$ is already prescribed. So finally the solution $u(x, t)$ assumes the form:

$$u(x, t) = \alpha_0 + \sum_{n=1}^{\infty} \exp(-n^2 t) (\alpha_n \cos nx + \beta_n \sin nx) \quad (2.25)$$

We need to check that as $t \rightarrow 0+$ the series displayed in (2.25) does converge to (2.24). The problem amounts to a passage across the summation sign of a limiting operation namely

$$\lim_{t \rightarrow 0+} \sum_{n=1}^{\infty} (\dots) = \sum_{n=1}^{\infty} \lim_{t \rightarrow 0+} (\dots)$$

We shall not carry out this program in detail but observe that at least when $f(x)$ is continuous and *piecewise smooth*, the Fourier series converges uniformly and so passage to limit across the summation sign is valid. See page 203 of the book of *Bachman-Narici* already cited earlier.

So we have established the existence of the solution for continuous piecewise smooth initial condition.

Exercise: Determine the solution of the heat equation with initial condition $f(x) = \pi^2 - x^2$ on $|x| \leq \pi$ extended as a 2π -periodic function on the entire real line.

Abel summability: Let (a_n) be a sequence of complex numbers. Recall Abel's limit theorem from elementary analysis:

Theorem: If $\sum_{n=0}^{\infty} a_n$ converges then the function $f(t)$ given by

$$f(t) = \sum_{n=0}^{\infty} a_n t^n$$

which is holomorphic on $|t| < 1$ (why?) satisfies

$$\lim_{t \rightarrow 1^-} f(t) = \sum_{n=0}^{\infty} a_n \quad (2.26)$$

where, the limit is taken as t varies over the real axis.

The theorem is non trivial because in the equation (2.26)

$$\text{LHS} = \lim_{t \rightarrow 1^-} \left(\sum_{n=0}^{\infty} a_n t^n \right) = \lim_{t \rightarrow 1^-} \lim_{N \rightarrow \infty} \left(\sum_{n=0}^N a_n t^n \right)$$

whereas

$$\text{RHS} = \lim_{N \rightarrow \infty} \sum_{n=0}^N a_n = \lim_{N \rightarrow \infty} \lim_{t \rightarrow 1^-} \left(\sum_{n=0}^N a_n t^n \right)$$

So the result is about the validity of an exchange of limits ! We say that the series

$$\sum_{n=0}^{\infty} a_n \quad (2.27)$$

is Abel summable if the associated power series

$$\sum_{n=0}^{\infty} a_n t^n \quad (2.28)$$

has radius of convergence 1 and further, if $f(t)$ is the sum function the following limit

$$\lim_{t \rightarrow 1^-} f(t)$$

exists. This limit is then defined to be the Abel sum of the series $\sum a_n$. It is important that the limit is taken as t varies along the real axis. This condition can be relaxed somewhat (*Stolz region*)

Abel summability and Fourier series Let us now consider a *continuous* 2π -periodic function on the real line with Fourier series

$$a_0 + \sum_{n=1}^{\infty} (a_n \cos nt + b_n \sin nt) \quad (2.29)$$

We know that the Fourier series (2.29) may not converge pointwise but a nice analogue of Abel-summability holds. Let us consider the associated series

$$a_0 + \sum_{n=1}^{\infty} r^n (a_n \cos nt + b_n \sin nt) \quad (2.30)$$

The series (2.30) converges for $0 \leq r < 1$ and denoting by $u(re^{it})$ its sum, it is of interest to know whether for a fixed $\theta \in \mathbb{R}$,

$$\lim_{(r,t) \rightarrow (1,\theta)} u(re^{it}) = f(\theta). \quad (2.31)$$

Recall the Poisson kernel:

$$\Pi_r(s) = \frac{(1-r^2)}{1+r^2-2r\cos(s)}$$

We already have the formula for $u(re^{it})$ namely,

$$u(re^{it}) = \frac{1}{2\pi} \int_{-\pi}^{\pi} \Pi_r(s-t) f(s) ds = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(t-s) \Pi_r(s) ds \quad (2.18)$$

Also,

$$f(\theta) = \frac{1}{2\pi} \int_{-\pi}^{\pi} f(\theta) \Pi_r(s) ds.$$

Subtracting from (2.18), we get

$$u(re^{it}) - f(\theta) = \frac{1}{2\pi} \int_{-\pi}^{\pi} (f(t-s) - f(\theta)) \Pi_r(s) ds$$

We exploit the fact that the Poisson kernel is a *positive kernel*:

$$|u(re^{it}) - f(\theta)| \leq \frac{1}{2\pi} \int_{-\pi}^{\pi} |f(t-s) - f(\theta)| \Pi_r(s) ds \quad (2.32)$$

Let $\epsilon > 0$ be arbitrary. Uniform continuity of f gives a δ so that

$$|f(x) - f(y)| < \frac{\epsilon}{2}, \quad \text{for } |x - y| < 2\delta.$$

Break the integral (2.32) in two pieces:

(a) Integral over $[-\delta, \delta]$. On this interval we have for $|t - \theta| < \delta$,

$$\frac{1}{2\pi} \int_{-\delta}^{\delta} |f(t-s) - f(\theta)| \Pi_r(s) ds < \frac{\epsilon}{2} \int_{-\pi}^{\pi} \frac{\Pi_r(s)}{2\pi} ds = \frac{\epsilon}{2}. \quad (2.32)'$$

We have used the fact that the integral of $\Pi_r(s)$ over $[\pi, \pi]$ is 2π .

(b) For the integral over $|s| > \delta$ we argue as follows. Let M be an upper-bound for $|f(s)|$.

$$\frac{1}{2\pi} \int_{|s|>\delta} |f(\theta-s) - f(\theta)| \Pi_r(s) ds < \frac{M}{\pi} \int_{|s|>\delta} \Pi_r(s) dt \quad (2.32)''$$

Now we observe that on $|s| > \delta$,

$$0 \leq \Pi_r(s) \leq \frac{1-r^2}{1+r^2-2r\cos\delta} \longrightarrow 0, \quad \text{as } r \rightarrow 1.$$

So for a suitable $\eta > 0$ and $1 - \eta < r < 1$, we have

$$\Pi_r(s) < \frac{\epsilon}{4M}$$

and the integral $(2.32)''$ over $|s| \geq \delta$ is again less than $\epsilon/2$. Consolidating $(2.32)'$ and $(2.32)''$, we get for $|t - \theta| < \delta$ and $1 - \eta < r < 1$, the estimate

$$|u(re^{it}) - f(\theta)| < \epsilon.$$

and (2.31) follows namely,

$$\lim_{(r,t) \rightarrow (1,\theta)} u(re^{it}) = f(\theta). \quad (2.31)$$

Summarizing,

Theorem: Suppose $f(t)$ is a continuous 2π -periodic function on the real line with Fourier series (2.29) then, for fixed $\theta \in \mathbb{R}$, the function $u(re^{it})$ which is the sum of (2.30) converges to $f(e^{i\theta})$ as $(r, t) \rightarrow (1-, \theta)$.