

Partial Differential Equations
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Lecture – 59
Maximum Principle for Heat Equation

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Problem 1

Show that

$$\int_{\mathbb{R}} \exp(-x^2) dx = \sqrt{\pi}.$$

Note that

$$\int_{\mathbb{R}} \exp(-x^2) dx = 2 \int_0^{\infty} \exp(-x^2) dx$$

as the integrand is an even function.

Welcome to Tutorial on heat equation. In this tutorial, we are going to solve a few problems associated to heat equation. The first problem is a problem in calculus which is integral of e power $-x$ square is root pi. We have used this information in our analysis of heat equation that is why I decided to do this proof. This involves a small trick, I will explain you that. So integral of e power $-x$ square \mathbb{R} is equal to two times the integral on 0 to infinity of the same integral e power $-x$ square because it is an even function

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Solution to Problem 1

Denote $I = \int_0^{\infty} \exp(-x^2) dx$. Then

$$\begin{aligned} I^2 &= \left(\int_0^{\infty} \exp(-x^2) dx \right) \left(\int_0^{\infty} \exp(-y^2) dy \right) \\ &= \int_0^{\infty} \left(\int_0^{\infty} \exp(-x^2) dx \right) \exp(-y^2) dy \\ &= \int_0^{\infty} \int_0^{\infty} \exp(-x^2 - y^2) dx dy \\ &= \int_{\theta=0}^{\theta=\frac{\pi}{2}} \int_{r=0}^{r=\infty} \exp(-r^2) r dr d\theta \end{aligned}$$

using Polar coordinates in the plane.

Therefore, it is enough to compute 0 to infinity e power $-x$ square dx. Then the answer we are interested in will be two times I. So to compute I, we look at I square, I squared is I into I that is this quantity into the same quantity. Because I am writing two times I have used the x in one integral and y in the other integral. Anyway, this is so called a dummy variable, so I can use any variable I want. Now, what I do is that this quantity I take it inside.

Actually, I am assuming that this is finite quantity and that is why everything is fine. So I take this inside and I get this. Now, this is nothing but 0 to infinity 0 to infinity e power $-x$ square into e power $-y$ square is e power $-x$ square $-y$ square dx dy. And that is equal to in the polar coordinates $r = \text{square root of } x \text{ square} + y \text{ square}$. So therefore, $-x$ square $-y$ square becomes $-r$ square and the area element dx dy becomes $r dr d\theta$.

Now we have to find out the limits for the integration variables. Look at this integral, this integral is on 0 infinity cross 0 infinity that is the first quadrant. So the first quadrant, radius goes from 0 to infinity. But the part of the circle which in the first quadrant corresponds to $\theta = 0$ to $\theta = \pi/2$. For example, these are the circles I am integrating on. The angle is 0 to $\pi/2$ here, $\pi/2$ is the angle here.

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Solution to Problem 1 (contd.)



From the last slide, we have

$$\begin{aligned} I^2 &= \int_{\theta=0}^{\theta=\frac{\pi}{2}} \int_{r=0}^{r=\infty} \exp(-r^2) r \, dr \, d\theta \\ &=_{s=r^2} \int_{\theta=0}^{\theta=\frac{\pi}{2}} \left(\frac{1}{2} \int_{s=0}^{s=\infty} \exp(-s) \, ds \right) d\theta \\ &= \int_{\theta=0}^{\theta=\frac{\pi}{2}} \left(\frac{1}{2} \right) d\theta = \frac{\pi}{4} \end{aligned}$$

Therefore $I = \frac{\sqrt{\pi}}{2}$. Consequently, $\int_{\mathbb{R}} \exp(-x^2) \, dx = \sqrt{\pi}$.

So I square equal to this, this is what we have seen on the last slide. Now it is very easy to integrate. Put $s = r$ square, then the inside integral with respect to r becomes with respect to s in this form and this is easily integrable e power minus $-s$, integral is e power $-s$ by -1 limits 0 to infinity. So you get 1 from there, so total is half. So integral with respect to θ from 0 to π by 2 of 1 by 2 that is π by 4 . So I square is π by 4 , therefore I is root π by 2 .

Therefore, this integral which is two times I is root π . This is a very simple trick that we use in calculus. To compute an integral of one variable, we go through computation of integral into variables. Luckily, because of the form of the integrand, we can convert that into polar coordinates and we could easily compute.

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Problem 2



Solve the Cauchy problem

$$u_t - u_{xx} = 0 \quad x \in \mathbb{R}, \quad t > 0$$

$$u(x, 0) = e^{3x} \quad x \in \mathbb{R}.$$

Without worrying about uniqueness issues, we will compute a solution using the formula

$$u(x, t) = \frac{1}{\sqrt{4\pi t}} \int_{\mathbb{R}} \exp\left(-\frac{|x-y|^2}{4t}\right) \varphi(y) \, dy$$

derived for the Cauchy data φ . See **Lecture 7.2**

Now let us solve the Cauchy problem for homogeneous heat equation. Of course, we know a certain formula for the solution of the homogeneous heat equation and the corresponding Cauchy problem with a ϕ of x we have an integral, but that integral is rarely integrable, only in very special cases where we are very lucky. Namely $C\phi$ of x we have taken e^{3x} because the heat kernel is also exponential type, right.

So therefore, this exponential should help us that is how we are able to integrate, otherwise if it is some arbitrary function of x which is continuous and bounded, we do know that it gives rise to a solution but then that is it, we cannot really compute. We have to be really lucky to be able to compute. So without worrying about uniqueness issues because we know that Cauchy problem solution is not unique, it is unique only when you are looking at special classes.

Let us not worry about that and simply compute this integral and see what we get. So you will compute a solution, maybe that is the solution, that is the only solution in some class or maybe there are more solutions, we do not know it because this may not belong to the class for which uniqueness theorem holds. In fact, we are not discussed much about uniqueness theorems for which I have just given you the reference of book by DiBenedetto on PEDs.

You can look at at your leisure time. But as far as this course is concerned, we simply want to compute this integral with ϕ of y equal to e^{3y} , can we do it or not? So, this formula we have derived in lectures 7.2.

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Solution to Problem 2



$$\begin{aligned}
 u(x, t) &= \frac{1}{\sqrt{4\pi t}} \int_{\mathbb{R}} \exp\left(-\frac{|x-y|^2}{4t}\right) \exp(3y) dy \\
 &= \frac{1}{\sqrt{4\pi t}} \int_{\mathbb{R}} \exp\left(-\frac{x^2 + y^2 - 2xy}{4t} + 3y\right) dy \\
 &= \frac{1}{\sqrt{4\pi t}} \int_{\mathbb{R}} \exp\left(\frac{-x^2 - y^2 + 2xy + 12yt}{4t}\right) dy \\
 &= \frac{1}{\sqrt{4\pi t}} \int_{\mathbb{R}} \exp\left(\frac{-(y-x-6t)^2 + 12xt + 36t^2}{4t}\right) dy.
 \end{aligned}$$

So, now this is exponential into exponential, therefore it is exponential of this plus this, so I have written down exactly that. So, this is same as this. So, I am only working in the power of this e, exponential e to the power something, so that power only I am manipulating. Now, I express that like this, so called completion of squares. The advantage now is that I am going to set $y - x - 6t$ by something equal to some variable.

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Solution to Problem 2 (contd.)

$$u(x, t) = \frac{1}{\sqrt{4\pi t}} \int_{\mathbb{R}} \exp\left(\frac{-(y-x-6t)^2 + 12xt + 36t^2}{4t}\right) dy.$$

Do a change of variable $p = \frac{y-x-6t}{\sqrt{4t}}$ and obtain

$$\begin{aligned} u(x, t) &= \frac{1}{\sqrt{\pi}} \int_{\mathbb{R}} \exp(-p^2) \exp(3x + 9t) dp. \\ &= \exp(3x + 9t) \frac{1}{\sqrt{\pi}} \int_{\mathbb{R}} \exp(-p^2) dp. \\ &= \exp(3x + 9t) \quad \square \end{aligned}$$

So, I am going to put $y - x - 6t$ by $\sqrt{4t}$ as P . So, then this term the first term by $\sqrt{4t}$ that simply becomes P^2 and what remains is $12xt$ by $4t$ is $3x$; $36t^2$ by $4t$ is $9t$ and dy by $\sqrt{4t}$ is your dp from here, dp is 1 by $\sqrt{4t}$ dy . So, this is exactly same as this. Now, if you see exponential of $3x + 9t$ it does not depend on P , so it just comes outside.

And what we have is $\frac{1}{\sqrt{\pi}}$ into integral over \mathbb{R} of $e^{-P^2} dp$, we just found out this is a $\sqrt{\pi}$. So $\frac{1}{\sqrt{\pi}}$ by $\sqrt{\pi}$ that is 1 . So, what you have is exponential of $3x + 9t$. So, we are very lucky that we could actually compute the integral.

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Problem 3

For $T > 0$, let \mathcal{R} be defined by

$$\mathcal{R} = \{(x, t) : 0 < x < \pi, 0 < t < T\}.$$

Let u be a solution to the problem

$$\begin{aligned} u_t - u_{xx} &= 0 && \text{on } \mathcal{R}, \\ u(0, t) = u(\pi, t) &= 0 && \text{for } 0 \leq t \leq T, \\ u(x, 0) &= \sin^2 x && \text{for } 0 \leq x \leq \pi. \end{aligned}$$

Show that $0 \leq u(x, t) \leq e^{-t} \sin x$ holds for $(x, t) \in \mathcal{R}$.

Hint: Use Maximum principle.

Now, let us move on to the problem 3. For T positive let R be defined by this rectangle $0, \pi$ cross $0, T$. Let u be a solution to this problem, which is initial boundary value problem, homogeneous heat equation, boundary conditions are 0 , initial condition sine square x . Show that 0 is less than or equal to u of x, t is less than or equal to e power minus $-t$ sine x holds for x, t in the rectangle. Hint is to use maximum principle.

The maximum principle allows us to compare certain things, right. So, therefore the maximum principle using that is the hint here. So, what we have to show is actually two inequalities. We have to show 0 is less than or equal to u, x, t and u, x, t is less than or equal to e power $-t$ into sine x .

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Solution to Problem 3

Proof of $u(x, t) \geq 0$: Follows from minimum principle

- By minimum principle,

$$\min_{\bar{\mathcal{R}}} u(x, t) = \min_{\partial_p \mathcal{R}} u(x, t)$$

- But $\min_{\partial_p \mathcal{R}} u(x, t) = 0$.
- Therefore

$$u(x, t) \geq 0, \text{ for every } (x, t) \in \mathcal{R}.$$



So let us prove the first one that is $u(x, t)$ is greater than or equal to 0. It follows from minimum principle. By minimum principle, we know that minimum of u on R closure is actually minimum of u on the parabolic boundary of R . What is the parabolic boundary of R ? This is R , so parabolic boundary is this, union of these three lines and u is 0 here, u is 0 here, and u is sine square x here that is what is given to us.

So, minimum of u on the parabolic boundary is 0 because of this because minimum of sine square x on this 0 to π is 0. So minimum is 0. Therefore, minimum of u on R closure is 0, which means u is greater than or equal to 0 on R closure, in particular on R . So, this completes the proof of the first inequality.

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Solution to Problem 3 (contd.)



Proof of $u(x, t) \leq e^{-t} \sin x$:

- Observe that $v := v(x, t)$ defined for $(x, t) \in \mathcal{R}$ by

$$v(x, t) := u(x, t) - e^{-t} \sin x$$

satisfies $v_t - v_{xx} = 0$. As u and $e^{-t} \sin x$ solve homogeneous Heat equation.

Now let us look at the second inequality, so we want to prove this. So if you want to apply maximum principle you can only apply that to solution solve heat equation. So, therefore, we naturally ask this question whether this is a solution of heat equation. Yes, we have already seen when we were solving heat equation using separation of variables method that these are the things which are coming as the terms in the series. So, is this a solution? Yes, it is a solution.

So, what we do now is we define a new function v which is $u - e^{-t} \sin x$. This satisfies $v_t - v_{xx} = 0$ because u is already a solution to the homogeneous heat equation, this also solves homogeneous heat equation. Therefore, this equation being linear and homogeneous, sum of two solutions or the difference of two solutions is a solution still. So, v

is a solution to this. And what is that we want to show v of x, t is less than or equal to 0. So, showing this is same as showing v is less than or equal to 0.

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Solution to Problem 3 (contd.)

In fact, $v := v(x, t)$ is a solution to the pr



$$\begin{aligned} v_t - v_{xx} &= 0 && \text{on } \mathcal{R}, \\ v(0, t) = v(\pi, t) &= 0 && \text{for } 0 \leq t \leq T, \\ v(x, 0) &= \sin^2 x - \sin x && \text{for } 0 \leq x \leq \pi. \end{aligned}$$

- By maximum principle applied to v ,

$$\max_{\overline{\mathcal{R}}} v(x, t) = \max_{\partial_p \mathcal{R}} v(x, t)$$

So, in fact, we will look at the problem that v satisfies. So, homogeneous heat equation v is a solution that we have already checked on \mathcal{R} and v of 0, t and v of π, t are also 0, v of $x, 0 = \sin^2 x - \sin x$, please check this, it is very simple. So, I am not going to check here. So, this is a problem that v satisfies. Now, by maximum principle applied to this v , maximum of v on \mathcal{R} closure is nothing but maximum of v on the parabolic boundary.

So, therefore we will look at what is the maximum of v on the parabolic boundary. Parabolic boundary consists of three lines; one is the line $x = 0, \pi$ on which $v = 0$ okay and the other line is the x axis part between 0 to π . So, we need to look at what happens to this $\sin^2 x - \sin x$. $\sin^2 x - \sin x$ is always less than or equal to 0 on this interval because $\sin x$ lies between 0 to 1 that is the reason, we will see the details.

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Solution to Problem 3 (contd.)



- But $\max_{\partial p \mathcal{R}} v(x, t) = 0$.
 - This is because $v(x, 0) = \sin^2 x - \sin x \leq 0$ for $0 \leq x \leq \pi$.
 - $v(0, t) = v(\pi, t) = 0$
- Therefore
$$v(x, t) \leq 0, \text{ for every } (x, t) \in \mathcal{R}.$$
- This implies that $u(x, t) \leq e^{-t} \sin x$.

So, maximum of v on this is zero. This is because v of $x, 0$ is sine square $x - \sin x$ that is always less than or equal to 0 and v of $0, t$ and v of π, t both of them are 0. Therefore, v of x, t is less than or equal to 0 for every x, t in \mathcal{R} that implies u of x, t is less than or equal to $e^{-t} \sin x$. This is what we wanted to show.

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Problem 4

For $T > 0$, let \mathcal{R}_T be defined by

$$\mathcal{R}_T = \{(x, t) : 0 < x < l, 0 < t < T\}$$



Let u be a solution to the problem

$$\begin{aligned} u_t - u_{xx} &= 0 \quad \text{on } \mathbb{R} \times (0, \infty), \\ u(0, t) = u(l, t) &= 0 \quad \text{for } 0 \leq t, \\ u(x, 0) &= \varphi(x) \quad \text{for } 0 \leq x \leq l. \end{aligned}$$

Let $M(T)$ denote the maximum value of u on the rectangle $\overline{\mathcal{R}_T}$.

Show that $M(T)$ is a constant function of T .

So, let us look at problem number 4. For T positive let \mathcal{R}_T be defined as this rectangle x between 0 to 1 and t between 0 to T , $t = T$. So, this is what is called \mathcal{R}_T . So, for example if you take T dash another line, $t = T$ dash, this entire region will be \mathcal{R}_T dash. So, let u be a solution to the homogeneous heat equation in $\mathbb{R} \times (0, \infty)$ and the boundary condition at $x = 0$ and $x = l$ are 0 and u of $x, 0 = \varphi(x)$. Let M of T denote the maximum value of u on the rectangle \mathcal{R}_T and closure, $\overline{\mathcal{R}_T}$. Show that M of T is a constant function of T .

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Solution to Problem 4



- Let $T < T'$.
- Since $\mathcal{R}_T \subset \mathcal{R}_{T'}$, we have $\overline{\mathcal{R}_T} \subset \overline{\mathcal{R}_{T'}}$.
- Therefore, $M(T) \leq M(T')$.
- By Maximum principle,

$$M(T) = \max_{\partial_p \mathcal{R}_T} u(x, t), \quad M(T') = \max_{\partial_p \mathcal{R}_{T'}} u(x, t)$$

So let T be less than T' and \mathcal{R}_T is a subset of $\mathcal{R}_{T'}$, $\mathcal{R}_{T'}$ is a bigger rectangle. Therefore, \mathcal{R}_T closure also is contained in $\mathcal{R}_{T'}$ closure. Therefore, M of T is less than or equal to M of T' . M of T after all is the maximum of u on a smaller set, M of T' is a maximum on a bigger set, therefore this holds. So, by maximum principle what is M of T is a maximum of u on the parabolic boundary of \mathcal{R}_T . Similarly, M of T' is maximum value of u on the parabolic boundary of $\mathcal{R}_{T'}$.

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Solution to Problem 4 (contd.)

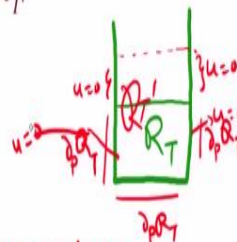


$$M(T) = \max_{\partial_p \mathcal{R}_T} u(x, t), \quad M(T') = \max_{\partial_p \mathcal{R}_{T'}} u(x, t)$$

- Note that

$$\underline{\partial_p \mathcal{R}_T} \subset \underline{\partial_p \mathcal{R}_{T'}}$$

u on both of these sets takes the same values.



Parabolic boundary of \mathcal{R}_T is a subset of parabolic boundary of $\mathcal{R}_{T'}$, u on both these sets takes the same values. Let us just have a look at this picture. This is \mathcal{R}_T , this is $\mathcal{R}_{T'}$. So this part, this part, this part is a parabolic boundary of \mathcal{R}_T . What is the parabolic boundary of $\mathcal{R}_{T'}$? It is bigger than this. What is the new portion? It is only this much and what happens to you the u there? The u is 0 that is the boundary condition and u is already 0 here,

so u does not take any new value so that maximum increases, therefore maximum is the same.

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Solution to Problem 4 (contd.)



$$M(T) = \max_{\partial_p \mathcal{R}_T} u(x, t), \quad M(T') = \max_{\partial_p \mathcal{R}_{T'}} u(x, t)$$

- Note that

$$\partial_p \mathcal{R}_T \subset \partial_p \mathcal{R}_{T'},$$

u on both of these sets takes the same values. Why?

- Thus $M(T) = M(T')$ □

So, M of $T = M$ of T dash. A similar statement can be made for the minimum also. Thank you.