

Nonlinear System Analysis
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Lecture – 20
Introduction to Bifurcation Theory – 1

Hi everyone. My name is Ramkrishna Pasumarthy, I am a faculty member at the Department of Electrical Engineering IIT, Madras. This is a course we are running on Non-linear System Analysis. We are using some lectures from the existing database of NPTEL and then adding up couple of other modules to make the series of lectures a lot more coherent.

I will be handling a couple of modules in this course and so far you would have learnt a lot of both qualitative and quantitative analysis of non-linear systems typical distinction between that of linear and non-linear you would have talked about or learnt about existence of multiple equilibria, existence of limit cycles for example. There were things like finite escape time and some other things like existence and uniqueness of solutions where the uniqueness was not always guaranteed and there are some very straightforward examples even in the one-dimensional case that you would that you would look at.

We also saw a and a stability analysis via a linearization of a non-linear system what we also traditionally call as the Jacobi linearization. And then we saw how the local behavior around that equilibrium point is in some way related to the linearization around that or the linear approximation around that equilibrium point. And we characterized equilibrium points according to we gave them certain names depending on where the eigen values of the linearized system was were lying, especially the second order we characterize them as nodes, focus, center, also talking about stable nodes, unstable nodes, and stable focus, unstable focus and so on.

So, what we will learn today is some very useful aspect of non-linear systems which are not very common in linear systems. I will start with motivating with the linear systems, but we kind of know why we can avoid that kind of phenomena. So, that is called the theory of

Bifurcations, right. So, we will typically look at what happens when a system parameter is changed, ok.

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Things that appear good, may turn bad suddenly!

1. Many a times we see that a dynamical system is behaving nicely in some operating point, suddenly becomes haywire without any apparent warning
2. The forthcoming slides give many examples of such phenomena!
3. How can a differential equation with a continuous dynamics on the right hand side cause such abrupt change in behaviour?
4. Thus we enter the zone of bifurcations!

$\dot{x} = f(x)$

So, what we will eventually learn is that things that appear good may turn bad suddenly and vice versa, what was going well with some small change maybe some environmental condition it can just go pretty bad, something which is not desirable. It also happen the other way. Since we are not going too well and then you can say some site change happens you can see well we are still in some kind of a safe region. We will try to characterize those behavior qualitatively, ok.

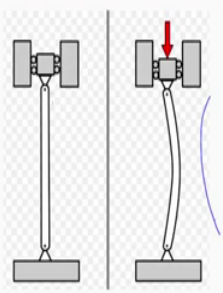
So, just to motivate a bit. So, many times we see that system is behaving nicely around some operating point and suddenly it just behaves you know in a way with which we which we would not anticipate it to without and these things typically come without any apparent

warning, ok. So, we will give examples of such phenomena. We will try to characterize them in terms of the theory that we have learned so far, right and also try to answer this question of how can a differential equation with a continuous dynamics on the right hand side.

So, which essentially means I am talking of \dot{x} equal to f of x , where the right hand side is smooth, right. So, if this is this discontinuous then you may say oh because the right hand side is discontinuous something strange is expected to happen, would actually happens when you talk of say switch systems where you switch between two stable systems you could actually end up in in a in a resulting unstable system, ok. So, how can a differential equation with a continuous dynamics on the right hand side cause such abrupt change in behavior, and this change in behavior is what we term as bifurcations, and we will see how we can categorize this class of bifurcation points.



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Eg 1 - Loaded Beam



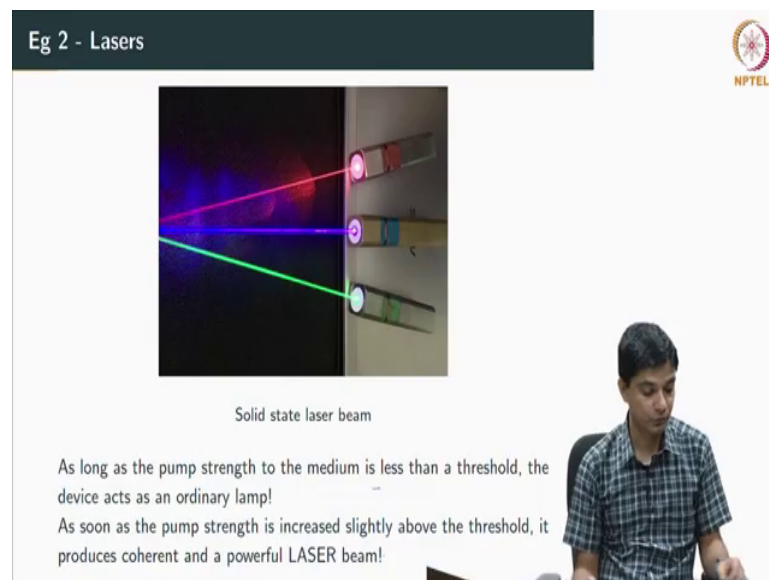
Beam under vertical loading

As long as the vertical loading of a beam is under some threshold L_T (no matter how close), the beam remains in equilibrium at vertical position. But as soon as the load exceeds L_T (even a little bit), suddenly it buckles to one side.



So, one example that we could we would have seen a lot and it is also quite kind of intuitive is you just look at a beam under vertical balancing, right. So, as the load increases where for normal load the beam just behaves as good as it should, but if the load exceeds a certain value it buckles to one side, right. So, you can see it bends either to this side or it also be like bending to this side to the to the opposite side. And we so far have not quantified or even look at looked at it qualitatively why really this happens and what does it actually mean in terms of systems.


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You could actually also look at some other examples of lasers where, as the as the strength is less, it behaves as an ordinary lamp and as the pump strength is increased beyond a certain threshold, it becomes a powerful laser beam.

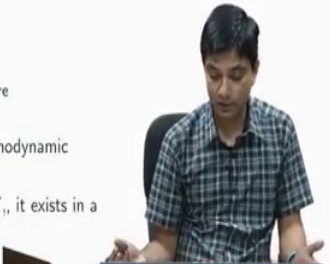
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Eg 3 - Boiling water (Phase Transition)



Abrupt change of phase with temperature

As long as temperature is below 100°C , water in thermodynamic equilibrium with its surroundings exists in liquid state
But as soon as the temperature slightly exceeds 100°C , it exists in a gaseous state



NPTEL

Boiling a water, right. It does not really come with any apparent warning, even though we know that at some level of temperature the state of water changes into a gaseous state, ok.

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Eg 4 - An Insect Outbreak



Spruce budworm is a notorious pest

Spruce budworm is a spurious pest that exhibits a sudden outbreak in the forests of Eastern Canada, attacking and defoliating an entire forest of fir trees in about four years!

Again, this outbreak happens suddenly without any apparent indicators or warning signs




And we also look at insect outbreak, right. So, you start with some small set of insects and all of sudden you see that the entire crop or entire forest is infected with pest. Same thing with even you can call about the recent virus outbreak, where you may have initially to begin with few number of people few 10s or few 100s and you see now close to 100,000 people are infected across the globe and it is just keeping on increasing. And these things do not come with any apparent warning or any prior model, right. So, we will try to qualitatively analyze this kind of behavior, ok.

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Some history

1. It was Poincare who first asked - how does the qualitative behaviour of a system change when a parameter in its dynamics is altered?

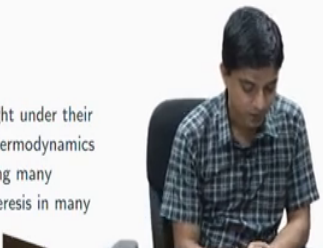



Henri Poincare

Chaos
James Gleick

2. For physicists, bifurcations were already sitting right under their nose as phase transitions (like boiling water) in thermodynamics

3. The theory of bifurcations is important in analyzing many phenomena such as jumps, catastrophes and hysteresis in many natural phenomena




So, where does all this come from? So, this famous French physicist Henri Poincare was first asked how does the qualitative behavior of a system change when a certain parameter is altered, ok.

So, this well this was kind of observed, but never given any formal notion until quite long or until lots of advances in in physics, right, ok. So, this theory is important in analyzing phenomena like jumps and you would have read about catastrophes, hysteresis in in many natural phenomena. And for people who are interested in a lot of this theory there is a book by James Gleick called on chaos theory.

It is a popular read not too much mathematical or academic, but it is a good read. So, so people who are interested beyond what is happening in the textbook or in the course here can

actually have a have a look through that. It is a book titled chaos, I do not remember the title completely by James Gleick assuming I have spell it correct, ok.

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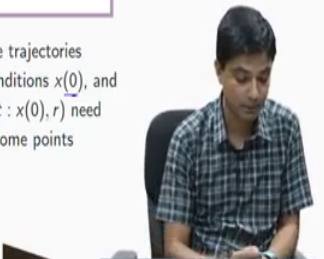
Bifurcation - what is it formally?

Handwritten notes:
 $\dot{x} = f(x, r)$
 F_1, m, a, x

Definition

Consider a dynamical system $\dot{x} = f(x, r)$ where $x \in \mathbb{R}^n$ is the state and $r \in \mathbb{R}^m$ is a parameter of the system. Bifurcation analysis is the study of qualitative change in behaviour of the system trajectories with changes in r .

The secret of the trade of bifurcations is that while the trajectories $x(t; x(0), r)$ indeed depend smoothly on the initial conditions $x(0)$, and on the parameters r , the limiting behaviour $\lim_{t \rightarrow \infty} x(t; x(0), r)$ need not depend smoothly on r and may jump abruptly at some points.

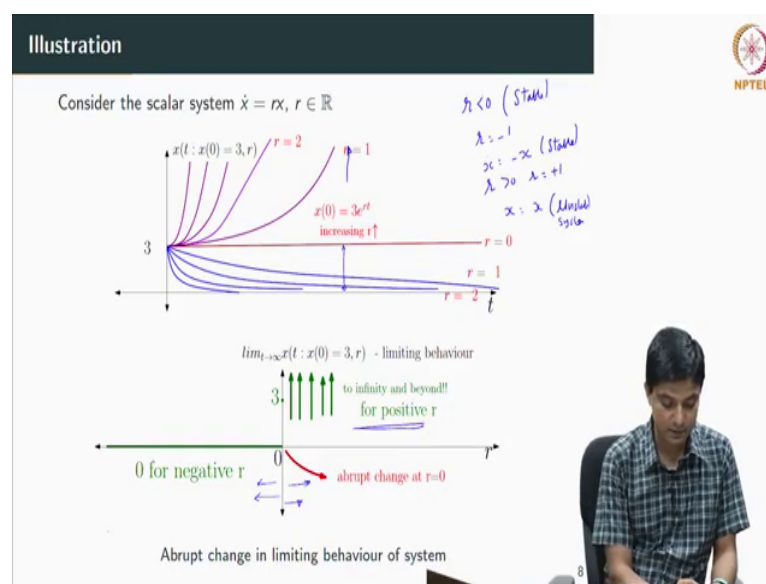


So, what is what is bifurcation? We will first just define it mathematically and then forget about the mathematical definition and then look at systems the way we know. So, we start with a dynamical system $\dot{x} = f(x, r)$ and also r . And what does this actually mean? So, if I write down say the Newton's second law I will have x double dot is mass times acceleration, right or this is m sorry or f equal to $m a$ which is yeah $m \ddot{x}$. And if I look at this typical equation this has to do with the state \ddot{x} whereas, this is a parameter which I usually do not pay much attention to while I am designing systems so far, or what I have whatever I have learned so far, right.

So, we will now pay attention to parameters that occur here, right. So, that is what we call as r , where x is the state and r is a parameter of the system. So, with this study of bifurcation analysis we will see how the qualitative behavior of the system changes or how the solution of the system, how the stability of the system changes with some changes in r , ok.

The secret here is that while these trajectories, right, so you have x at time t starting with some initial condition $x(0)$ and some value of parameter is depending smoothly on the initial conditions and the parameters are the limiting behavior need not smoothly depend. And we will see why this could happen, right. It may abruptly jump it may not necessarily depend smoothly on r and may abruptly jump sometimes, ok.

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So, let us start with a very simple scalar linear system. So, I start with an x system \dot{x} equal to r , r of x , and for I just if I just plot it the vector field. I see that, I can typically look at r


being less than 0, when I call it stable, right \dot{x} let us say r is minus 1. So, \dot{x} is minus of x I know its stable. Then say r going greater than 0, let us say for instance r equal to plus 1. We have \dot{x} equal to x , this is an unstable system, ok

So, r as this value of r changes from negative to positive we have in this green region a stable system, the system behaves as a stable system and all of a sudden as this parameter would increase would change from negative of r would cross 0, you are having an unstable system, right. It goes to infinity and beyond for positive r . For here it is kind of kind of well behaved.

And I can just plot this trajectories x versus time for different values of r and you see here as r increases until this region r equal to 0, your trajectories converge to 0, r equal to 0 your constant and for all r greater than 0 your trajectory just keep going up for arbitrary initial conditions, ok.

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Linear systems - the full story is already known!

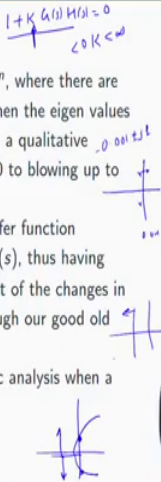
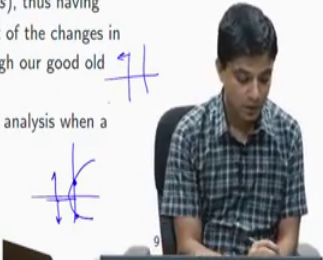


$1 + K G(s) H(s) = 0$
 $-\infty < K < \infty$

For the n dimensional linear system $\dot{x} = Ax$, $A \in \mathbb{R}^{n \times n}$, where there are n^2 parameters through the matrix A , we know that when the eigen values of A cross from left to right half plane (or vice versa), a qualitative change in limiting behaviour occurs [from settling to 0 to blowing up to infinity]

For the SISO system governed by the open loop transfer function $\frac{Y(s)}{U(s)} = G(s)$, subjected to gain feedback, $u(s) = -Ky(s)$, thus having closed loop transfer function $\frac{KG}{1+KG}$, we know the effect of the changes in control gain K on the limiting behaviour of $y(t)$ through our good old friend - the ROOT LOCUS diagram!!

We now want to invent a root locus like diagrammatic analysis when a parameter is varied in a non-linear system!

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So, in the case of linear systems, let us just start with SISO systems, right. Like for single input single output systems, I have a nice notion of transfer function, I have a nice notion of feedback and I know say for example, what does what does the root locus tell you. If I just say $1 + KG(s)H(s) = 0$ is the characteristic equation of the closed loop I am just interested in checking how does the system behave for values of K going from 0 to infinity and I just get some plot sometimes which they just start here, here and they are stable for all K , sometimes they are unstable. So, I can find this value of K for which a system goes from stable to unstable.

Like in this case the for all values of σ less than 0 I was stable, all values of σ greater than 0 the system was unstable. So, I know this this characterization, in terms of root locus. And therefore, while design you typically as do not want to design a stabilizing controller where the poles of the closed loop system are just slightly to the to the left of the complex plane. They say this for example, $-0.001 \pm j1$ may not be a very smart design because a slight change in parameter could or slight change in environmental conditions could change the parameter and push it to an unstable region, and you know that -0.001 is actually unstable even though it is very like close to the to the origin, right.

So, we do we unknowingly did this analysis while we were doing the root locus in the case in our first course on control engineering, what are the how does the system behavior change as the value of K varies from 0 to infinity. You can also draw the negative root locus saying what happens from minus infinity to 0 and we find use the Routh Hurwitz criteria to find out what are the values of K for the system is unstable.

It could turn out in many cases or some cases that the system will be stable for all K , for some cases for K slightly greater than 0 you will be unstable and a and a bunch of bunch of other things too. And therefore, we also learn relative stability in Routh Hurwitz where we say, place all my poles to the left of minus 1, ok.

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Saddle-Node bifurcation

Consider the system

$$\begin{cases} \dot{x}_1 = \mu - x_1^2 \\ \dot{x}_2 = -x_2 \end{cases}$$

Let us analyse the dynamics of the system as the parameter $\mu \in \mathbb{R}$ varies.

1. The system has two equilibrium points:
2. $(\sqrt{\mu}, 0)$ Stable Node, $(-\sqrt{\mu}, 0)$ Saddle point.
3. As μ decreases, the saddle and node approach each other and collide at $\mu = 0$ and
4. For $\mu < 0$ the system has no equilibrium points.


$x_2 = 0$
 $x_1^2 = \mu \Rightarrow x_1 = \pm \sqrt{\mu}$

Jacobi linearization

$$J = \begin{bmatrix} -2x_1 & 0 \\ 0 & -1 \end{bmatrix}$$

$x_1 = \pm \sqrt{\mu}$

Saddle Node



So, let us come to the non-linear case and just start with few examples, ok. So, let us let me start with this second order system $\dot{x}_1 = \mu - x_1^2$, $\dot{x}_2 = -x_2$ and we will check what happens to the behavior of this system as μ varies, ok.

So, first is I will just quickly do it for this example, and the remaining words I will just leave it you to solve. So, what are the equilibrium points? For sure $x_2 = 0$ is the equilibrium point and then another equilibrium point is $x_1^2 = \mu$ which means this is I am looking at $x_1 = \pm \sqrt{\mu}$, ok. So, I do it the standard way. I just look at the Jacobi linearization, resulting in the Jacobian matrix which in this case looks like this, $\begin{bmatrix} -2\sqrt{\mu} & 0 \\ 0 & -1 \end{bmatrix}$, ok.

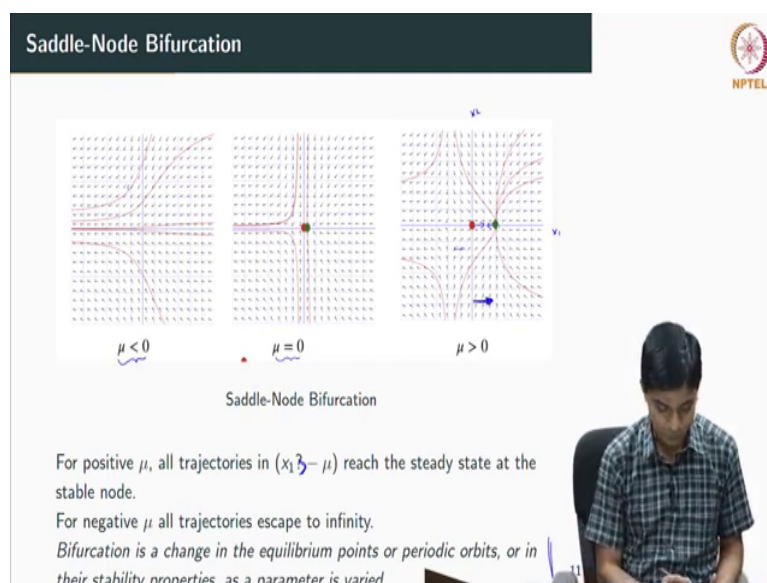
And the other equilibrium, so the system therefore, has two equilibrium points; you have $\mu, 0$ which is from this table from this Jacobian you can find out is a stable node and

minus μ comma 0 which is a minus square root of μ comma 0 which is a saddle point, ok. So, what will be interesting is to check what happens as μ decreases, right.

As μ decreases the saddle and the node approach each other collide at μ equal to 0 and for μ less than 0 the system has no equilibrium points. Yes let us check this, right. So, what happens when μ is less than 0 I am looking for solutions is x_1 is plus minus square root of minus μ , all right, ok.

When μ is less than 0, this the guy inside, this is a negative number and I would definitely not deal with imaginary equilibrium points, right. So, as μ , so these are the two equilibrium points, one is stable node, another is a saddle point. As μ decreases from certain values to 0, the saddle node and the saddle and the node, so you are looking at something like this, right. You have a plus square root of μ you have a minus square root of μ , ok, this is stable this is saddle. And as μ decreases, they will keep moving to this guy will keep moving to the right, this guy will keep moving to the left, ok.

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So, let us see this here from this, ok. So, for μ greater than 0 what we have is, so this is this is a stable node, this is this is a saddle point and as μ decreases these two keep moving towards each other, they meet each other at μ equal to 0, ok. Let us not really worry about what is the stability of μ equal to 0 that we will keep for later discussion, and for μ less than 0 the system does not have an equilibrium, ok. So, for positive μ all trajectories, so if I look at here this diagram, for positive μ all trajectories in this region x_1 , so I am just drawing the diagrams like this x_1 and x_2 in the standard phase portraits.

For all trajectories for which x_1 is greater than minus μ they reach the steady state which is a stable node in in at asymptotically. For negative μ all trajectories, so it is for negative μ all trajectories no matter wherever you start you escape to infinity along these lines along the


arrows. Here at least for μ in which in this region you can converge to this this this green and the green dot is this I just used the green for stability red for unstable, ok.

So, one thing that we observed here is as the μ changes from positive to negative, you have two equilibrium points and for μ less than 0 the equilibrium points disappear totally, right. So, this is this something very strange, right.

In the linear case you had system only going from stable to unstable region for some certain values of K . Here well, the equilibrium points one was stable, one was unstable as they keep coming close to each other with decreasing μ they just disappear, right. There is no equilibrium point, there is no steady state when μ is less than 0 all trajectories escape to infinity ok.

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Transcritical Bifurcation



Consider the system

$$\begin{aligned}\dot{x}_1 &= \mu x_1 - x_1^2 \\ \dot{x}_2 &= -x_2\end{aligned}$$

$\mu(x_1 - x_1^2) < 0$
 $x_1(\mu - x_1) < 0$
 $(0,0) \quad (\mu,0)$

1. The system has two equilibrium points
2. The point $(0,0)$ is
 - Stable Node for $\mu < 0$
 - Saddle for $\mu > 0$
3. The point $(\mu,0)$ is
 - Stable Node for $\mu > 0$
 - Saddle for $\mu < 0$

Handwritten notes and stability analysis:

For $(0,0)$:
 $\mu < 0$: $\begin{bmatrix} \mu & 0 \\ 0 & -1 \end{bmatrix}$ at $(0,0)$
 $\mu > 0$: $\begin{bmatrix} 0 & 0 \\ 0 & -1 \end{bmatrix}$ at $(\mu,0)$

For $(\mu,0)$:
 $\mu < 0$: $\begin{bmatrix} \mu & 0 \\ 0 & -1 \end{bmatrix}$ at $(\mu,0)$
 $\mu > 0$: $\begin{bmatrix} 0 & 0 \\ 0 & -1 \end{bmatrix}$ at $(\mu,0)$

Stability summary table:

	$\mu < 0$	$\mu > 0$
$(0,0)$	Stable	Saddle
$(\mu,0)$	Saddle	Stable

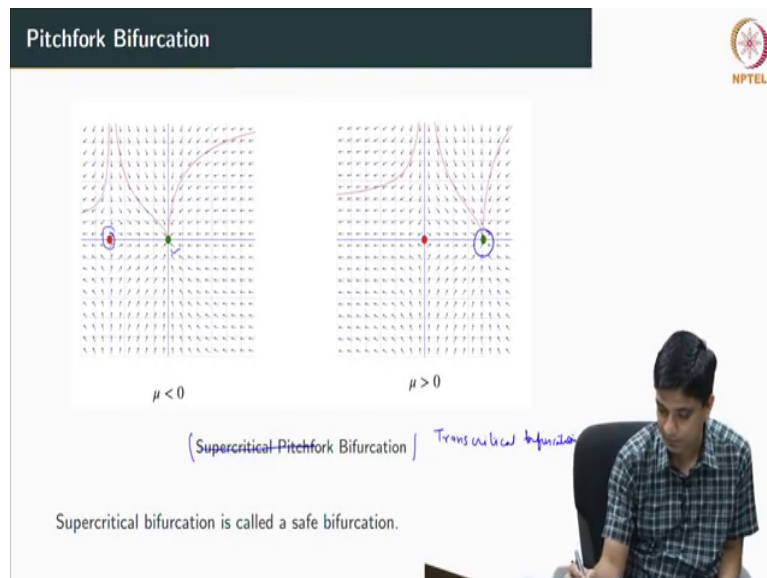
Similarly, I now look at this system where I have $\dot{x}_1 = \mu x_1 - x_1^2$, $\dot{x}_2 = 0$, so this equation will pretty much remain the same. So, this system will have two equilibrium points. So, if I just look at $\mu x_1 - x_1^2 = 0$. So, the equilibrium points would be of course the origin, another equilibrium point will be $\mu, 0$, ok.

And Jacobian would be should be easy to compute the Jacobian at, so the Jacobian at $0, 0$, ok. I now I will skip some of these steps because you would already have know how to compute this you have $0, -1$ and the Jacobian at $\mu, 0$ would look something like this $-\mu, 0, 0, -1$, ok.

Now, look at this at $0, 0$. At $0, 0$ for $\mu < 0$ we get that all the eigen values are negative, right. So, this will be a stable node for $\mu < 0$. And this will become a saddle, so say for μ is equal to -1 my eigen values here would be $-1, -1$ this corresponds to a stable node. For $\mu > 0$ let us say some arbitrary value say μ equal to 1 , my eigen values should be $+\mu, -1$ that corresponds to a saddle, ok.

Now, look at this equilibrium point at $\mu, 0$. $\mu, 0$, for $\mu > 0$ for $\mu > 0$ say μ equal to 1 my eigen values are $-\mu, -1$. It is a stable node, and it is a saddle for $\mu < 0$, ok. What is happening here? Let us just check here, right, ok.

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The some pictures got exchanged in in the slide. But this is essentially corresponding to not a super critical, but this is like a transcritical bifurcation, ok. So, so apologies for the for the typo here or for the for mixing up the images, ok.

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Pitchfork Bifurcation


Consider the system

$$\begin{aligned}\dot{x}_1 &= \mu x_1 - x_1^3 \\ \dot{x}_2 &= -x_2\end{aligned}$$

$\mu < 0$	$\mu > 0$
1 - eq $(0,0)$ stable	3 - eq, 1 point $(0,0)$ Saddle $(-\mu,0), (+\mu,0)$

1. For $\mu < 0$, there is a unique equilibrium point $(0,0)$, which is a stable node.
2. For $\mu > 0$, the system has three equilibrium points saddle at $(0,0)$, and two stable nodes at $(\pm\mu,0)$.
3. As μ crosses the bifurcation point $\mu = 0$, the stable node at the origin bifurcates into a saddle and other two equilibria are stable nodes.





What happens is look at for mu less than 0, ok. And if I look at 0 comma 0 is stable and for mu greater than 0 I will just call it (Refer Time: 21:20) reverse is a saddle, ok. Now, similarly the point mu comma 0 when mu is less than 0 is a saddle and this is stable for mu greater than 0.

So, as we cross the bifurcation points the stable node becomes a saddle and the saddle becomes a stable node. So, they just kind of exchange their properties, so here, right. So, this was for mu less than 0 the mu less than 0, the origin was stable, ok, mu less than 0, origin was stable and then the other equilibrium point well this was this was like a saddle, right. And as you keep increasing mu you cross mu equal to 0, they just change their properties, right. So, what was stable now becomes unstable and then there is another equilibrium point here at this this equilibrium point which becomes a saddle point, for mu less than 0, right.

So, so that is, so two things we have seen, right. In the first case where I have a saddle and the node they meet together for increasing values, for decreasing values of μ and then they just disappear. Now, I have two equilibrium points where the stable node here it becomes a saddle as I cross the bifurcation point and the saddle becomes a stable node, ok. So, ok this is a text with transcritical bifurcation. I just messed up those pictures, ok. I will I will clear I will correct those slides, right.

So, the equilibrium points they exist through all values of μ , for μ less than 0 I have equilibrium points, μ greater than 0 I have equilibrium points, ok. So, only thing that changes contrast to what happened before, right μ less than 0 equilibrium point disappear. Equilibrium points persist only they flip their characteristics, right. The point 0, 0 changes from a stable to a saddle and μ 0 changes from a saddle to a stable node, ok, right, ok.

Now, the third kind of bifurcation which we will look at is what is called as the pitchfork bifurcation. So, again I am looking at a system which looks like this. So, I have a system of $\dot{x} = \mu x - x^3$ and $\dot{y} = 0$ remains the same. For μ less than 0 you can see there is a unique equilibrium point and we can check that this is a stable node, ok.

And for μ greater than 0 the system will have 3 equilibrium points and those will be a saddle at 0, 0 and two stable nodes at $(\pm\sqrt{\mu}, 0)$, right, ok. Let us quickly check what is happening. First this equilibrium and this equilibrium, so you can just say μ less than 0, μ greater than 0 and if I say the number of equilibria I have 1 equilibrium point, here I have 3 equilibrium points, ok. Now, here well 0, 0 was stable and this 0 comma 0 the origin becomes an unstable it becomes a saddle, and then two more equilibrium points appear one at minus $\sqrt{\mu}$ comma 0 and one at plus $\sqrt{\mu}$ comma 0.

Look at how interesting, right. So, we had cases where equilibrium was disappearing here all of a sudden two additional equilibria emerge, ok. So, as μ crosses the bifurcation points, the stable node at the origin it bifurcates into a saddle this the stable node becomes a saddle and two other equilibria are created which are stable nodes, ok. How does the picture for this look like? Ok.

So, this is, called the transcritical, but the super critical pitchfork, ok. So, for μ less than 0, I had one nice stable point all of a sudden what happens is that this equilibrium point when μ goes to slightly higher value than 0 or even larger values it becomes unstable.

Well, this is, so generally this is called a safe bifurcation that is because even though this system the equilibrium the origin changes from a stable to an unstable for very small values of μ you see here the system does not go to infinity, it just settles at this point and or settles at this point. So, for small values of μ , this actually well this is that is what a disaster that is happening. As my system does not jump from being stable all of a sudden to being unstable, it just shifts and sits up at a point which is which is like close by, right.

So, the trajectories do not like go away to infinity, right, ok. So, and then now I have something called if I just flip the sign of x^3 , the second term here I have the reverse behavior, right.


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Pitchfork Bifurcation

Consider the system

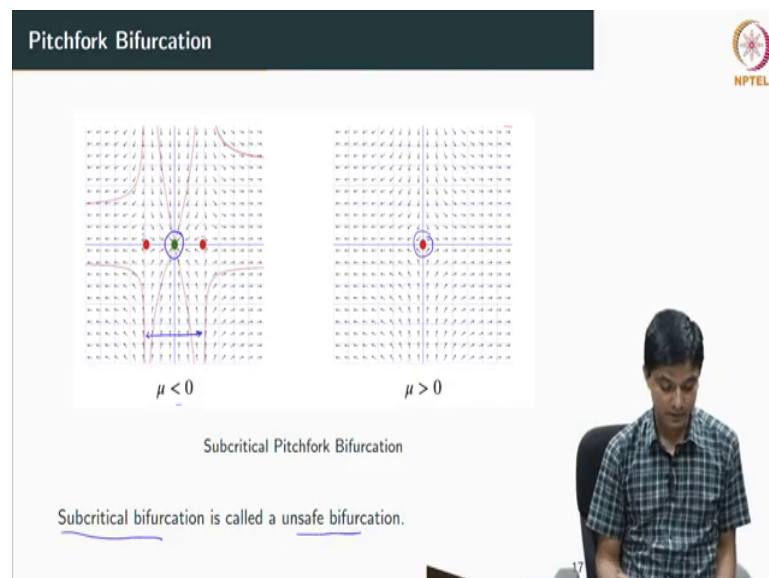
$$\begin{aligned}\dot{x}_1 &= \mu x_1 + x_1^3 \\ \dot{x}_2 &= -x_2\end{aligned}$$

1. For $\mu < 0$, the system has three equilibrium points: stable node at $(0, 0)$, and two saddles at $(\pm\mu, 0)$.
2. For $\mu > 0$, there is a unique equilibrium point $(0, 0)$, which is a saddle.



So, for μ less than 0, I will have 3 equilibrium points, a stable node at 0, 0 and two saddles at plus minus μ comma 0, ok.

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So, I have picture, right, yeah. So, for μ less than 0 this is stable. So, this is all good now these two are unstable, ok. Assume this is my operating point and a sudden change in parameter or sudden change in conditions pushes the parameter from μ less than 0 to μ greater than 0. What you see is that these two red dots which were appearing here they disappear.

So, 3 equilibriums they change to one equilibrium, not only that this one equilibrium is unstable therefore, you would not you would never want this to happen, right because whatever is you are in in the stable region, so even in this case is if you are within this region here you will be pushed to the to the to the green dot, ok.

Whereas, once you bifurcate into this region then all the trajectories become unstable, right. So, all trajectories go to infinity and therefore, a sub critical bifurcation is called an unsafe bifurcation, ok.

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Hopf Bifurcation



Stable focus $-1 \pm j\omega$
unstable focus $\pm 1 \pm j\omega$


When a stable node loses stability at a bifurcation point, an eigen value of the Jacobian passes through zero.

What happens when a stable focus loses stability?


(A pair of complex eigen values could pass through the imaginary axis.

Consider the system

$$\begin{aligned}\dot{x}_1 &= x_1(\mu - x_1^2 - x_2^2) - x_2 \\ \dot{x}_2 &= x_2(\mu - x_1^2 - x_2^2) + x_1\end{aligned}$$



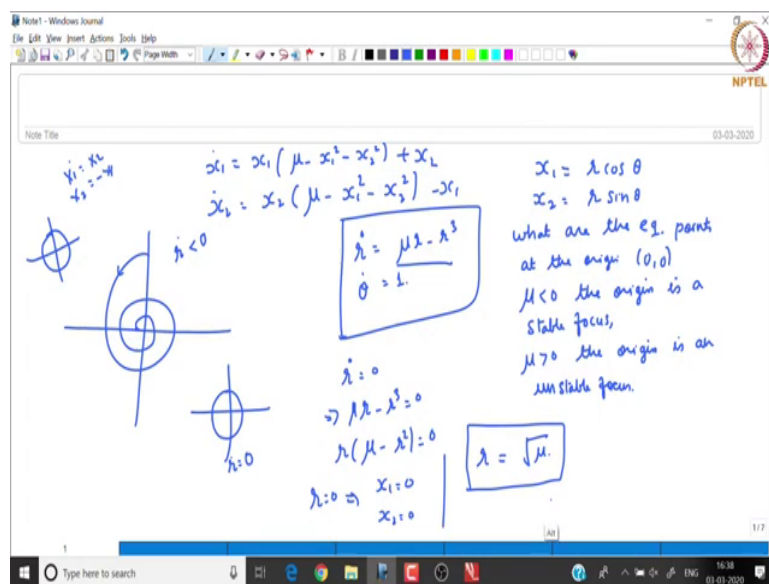
$\mu = 0$



So, to summarize or to look at what really is happening in this in all the all of these examples is that when μ equal to 0 in this case you are encountering a 0 eigen value here, when μ equal to 0, 0 eigen value of the linearization, right and when μ equal to 0 here you still encounter a 0 eigen value, right. Similar things will happen also here. When μ equal to 0 you encounter again a 0 eigen value, right, ok. So, whenever a stable node loses its stability at bifurcation point an eigen value of the Jacobian passes through 0, ok.

So, this, we were just looking at a stable node so far and stability is not just or all stable equilibrium are not just nodes, it could always be or they can always be as a stable focus. So,

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So, I have $\dot{x} = x^2 - y^2$ and $\dot{y} = 2xy$. This may look a little complicated to solve, but some of our earlier courses taught us coordinate transformation, right. So, where you just go from rectangular to polar, and in these coordinates I have something nice: $\dot{r} = r$ and $\dot{\theta} = 1$. So, I will just analyze this in the polar coordinates now, ok.

So, what are the equilibrium points? Ok. So, this simple inspection would show that the system has a unique equilibrium point at the origin. Even from here you can find, right that system has a unique equilibrium point at origin and that is the only thing that happens.

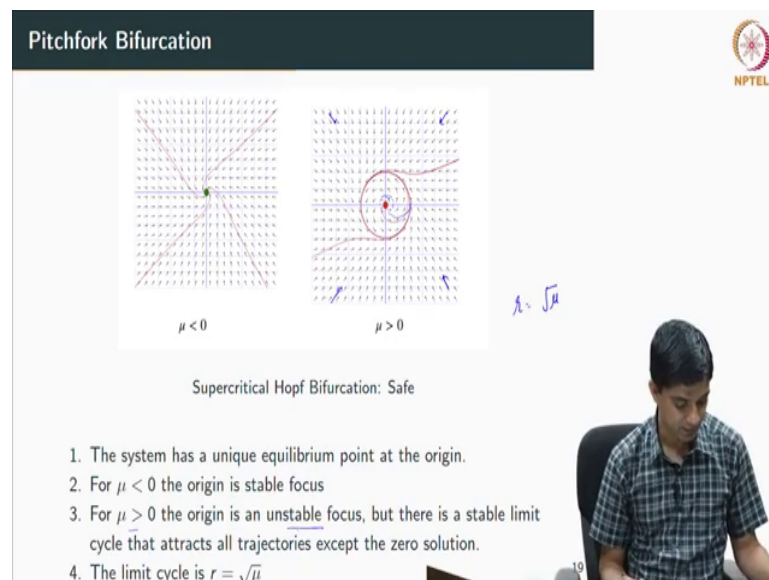
So, let us check what happens, it is for and the standard analysis you can find that for μ less than 0 the origin is a stable focus and for μ greater than 0 the origin is an unstable focus, ok. So, let us look at, ok. So, how does a focus look like? So, it is a stable focus would just spiral this way to the origin which means a certain radius \dot{r} is negative is less than 0 and some $\cos 1 \theta$. For a periodic orbit you will have that \dot{r} is equal to 0, that the radius does not change, right.

So, if you look at system which is a the simple harmonic oscillator $\ddot{x} = -x$, it will have a phase portrait which is not just a circle, right, and if \dot{r} is greater than 0 then you will have an unstable trajectory. So, this is a stable one and then \dot{r} will be the opposite sign, ok.

So, just check, let us check there is something interesting happen here when \dot{r} equal to 0. \dot{r} equal to 0 implies $\mu r - r^3$ equal to 0 or $r(\mu - r^2)$ equal to 0, r equal to 0. Well, not very interesting that is r equal to 0 is the origin x_1 equal to 0 x_2 equal to 0, r equal to 0 I have x_1 equal to 0, x_2 equal to 0 whereas, the other equilibrium point r will be square root of μ , ok.

So, something we can infer from here, right, that there is a possibility of a periodic orbit and possibly this could also be a limit cycle, right of radius r equal to square root of μ , ok. So, let us just check what happens with the phase portraits of this.

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When μ less than 0 what I have is that the origin is a stable focus all trajectory spiral to the origin when μ greater than 0 this becomes an unstable equilibrium. But what happens is all the trajectories go and converge to a limit cycle whose radius is given by r equal to square root of μ , ok. For μ greater than 0 origin is unstable, but there is a stable limit cycle that attracts all trajectories except the 0 solution, right.

So, what happened at the 0 solution? 0 solution is r equal to 0, nothing changes here, right. So, in this case we have a limit cycle and you see all trajectories, trajectories starting from here, from here, from here or from here they will converge to this limit cycle, ok. So, in this case what is essentially happening is when you change from a stable focus, so where eigen values could be $-1 \pm j\omega$ to an unstable focus.

Say, with eigen values of $1 \pm j\omega$, right. You a pair of complex eigen values actually pass through the imaginary axis. So, you had things here, things here and the bifurcation point at μ equal to 0. Well, your eigen values actually pass through the imaginary axis, right, ok. So, this is a bit of qualitative behavior of how equilibrium points can change with changes in the system parameter.

So, I will end this lecture here. And next time when we meet, we will do a little more analysis on when do these kind of equilibriums occur or we will try and derive certain necessary and sufficient conditions where I can guarantee that a system undergoes a certain kind of equilibrium. And again, we will revisit all the examples that we motivated our lecture with.

Thanks for listening.