

Introduction to Nonlinear Dynamics
Prof. Gaurav Raina
Department of Electrical Engineering
Indian Institute of Technology, Madras

Module -06
Lecture-29

2-Dimensional Flows, Bifurcations, Lecture 3

(Refer Slide Time: 00:01)

2 Dimensional Flows Bifurcations

Half bifurcation

Note that saddle-node, pitchfork, and transcritical bifurcations always involve the collision of two or more fixed points.

We now consider a new bifurcation that has no counterpart in one-dimensional systems – and it provides a way for a fixed point to lose stability without colliding with any other fixed points of the system.

Half bifurcation

Suppose that a 2-dimensional system has a stable fixed point. What are all the possible ways it could lose stability as a parameter μ varies?

Clearly, the eigenvalues of the Jacobian are key!

If the fixed point is stable, the eigenvalues λ_1, λ_2 must both lie in the left half plane ($\text{Re } \lambda < 0$).

Since the λ 's satisfy a quadratic equation, with real coefficients, there are two possible options:

- (i) both the eigenvalues are real and negative
- (ii) they are complex conjugates

The diagrams show the complex plane with the imaginary axis (Im λ) and real axis (Re λ). In (i), two real eigenvalues λ_1, λ_2 are shown on the negative real axis. In (ii), a pair of complex conjugate eigenvalues λ_1, λ_2 is shown in the left half plane.

We now deal with the half bifurcation. Note that the Saddle node, the Pitchfork and the Transcritical Bifurcations always involve the collision of two or more fixed points, we now consider a new bifurcation that has no counter parts in one dimensional systems and it provides way for a fixed point to lose stability without colliding with any other fixed points of the system. So we now get on to the half bifurcation.

Suppose that a two-dimensional system has a stable fixed point, what are all the possible ways it could lose stability as a parameter μ varies. Now clearly the Eigen values of the Jacobian are key, if the fixed point is stable. The Eigen values λ_1 and λ_2 must both lie in the left half of plane, where $\text{Re } \lambda < 0$. Since the λ 's satisfy a quadratic equation with real coefficient.

There are two possible options one, both the Eigen values are real and negative. So, we plot the imaginary of λ versus the real of λ , you have two Eigen values which are real and

negative. Number two, they are complex conjugates, so we plot again imaginary lambda and real of lambda and we have complex conjugate Eigen values.

(Refer Slide Time: 02:32)

To destabilize the fixed point, one or both the eigenvalues need to cross into the right half plane as μ varies.

With the saddle-node, transcritical and pitchfork bifurcations we explored the cases in which a real eigenvalue passes through $\lambda = 0$.

One now considers the other possible scenario, in which two complex conjugate eigenvalues simultaneously cross the imaginary axis into the right half plane.

The Hopf bifurcation comes in two flavours

- Supercritical Hopf
- Subcritical Hopf

To destabilise the fixed point one or both of the Eigen values need to cross into the right half plane as μ varies. With the saddle nodes, transcritical and pitchfork bifurcations, we explored the cases in which the real Eigen value passes through $\lambda = 0$. One now considers the other possible scenario in which two complex conjugate Eigen values simultaneously cross the imaginary axis into the right half plane. The half bifurcation comes in two flavours a supercritical half and a subcritical half.

(Refer Slide Time: 03:50)

Supercritical Hopf Bifurcation

Suppose that we have a physical system, which has a control parameter μ

And for $\mu < \mu_c$ (some critical value of μ) the system settles to some equilibrium through exponentially damped oscillations

If μ keeps increasing, at some stage the equilibrium state will lose stability. In many cases the resulting motion is a small amplitude, sinusoidal, limit-cycle oscillation about the original equilibrium state

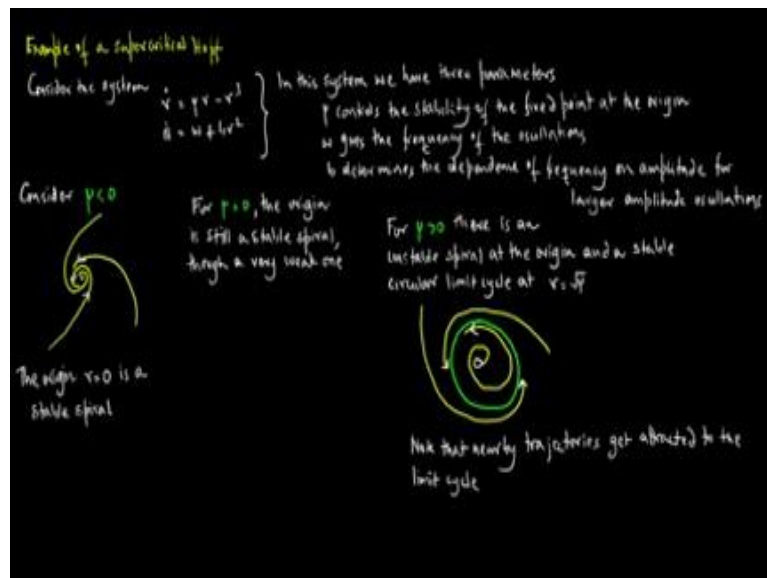
In terms of the flow in phase space, a supercritical Hopf bifurcation occurs when a stable spiral changes into an unstable spiral which is surrounded by a small amplitude limit cycle

Hopf bifurcations can occur in phase spaces of any dimension $n \geq 2$

First, we discuss the supercritical half bifurcation. Suppose that we have a physical system which has a control parameter μ and for $\mu < \mu_{\text{critical}}$. For some critical value of μ , the system settles to some equilibrium through exponentially damped oscillations. So that is an example of damped oscillations when μ is less than μ_{critical} . So if μ keeps increasing at some stage the equilibrium state will actually lose stability.

In many cases the resulting motion is a small amplitude sinusoidal limit cycle oscillation about the original equilibrium state. So that is the example of the limit cycle, where $\mu > \mu_{\text{critical}}$. In terms of the flow in phase space a supercritical half bifurcation occurs when a stable spiral changes into an unstable spiral, which is surrounded by small amplitude limit cycle. Half bifurcations can occur in phase spaces of any dimension n greater than or equal to 2.

(Refer Slide Time: 05:55)



Now let us look at the example of a supercritical half. So, consider the system $\dot{r} = \mu r - r^3$, $\dot{\theta} = \omega + br^2$. In this particular system we have three parameters, μ controls the stability of the fixed point at the origin. ω gives the frequency of the oscillations and b determines the dependence of frequency on amplitude for larger amplitude oscillations.

So, consider μ less than 0, the origin $r = 0$ is a stable spiral. So, $\mu = 0$, the origin is still a stable spiral though a very weak one. So, for μ greater than 0, there is an unstable spiral at the origin and a stable circular limit cycle at the square root of μ . Note that nearby trajectories get attracted to the limit cycle.

(Refer Slide Time: 07:54)

To get a sense of how the eigenvalues behave during the bifurcation, we rewrite the system in Cartesian coordinates.

With $\begin{cases} x = r \cos \theta \\ y = r \sin \theta \end{cases}$ Then $\begin{aligned} \dot{x} &= \dot{r} \cos \theta - r \dot{\theta} \sin \theta \\ &= (\dot{r} - \omega r) \cos \theta - r(u + br^2) \sin \theta \\ &= (r - (x^2 + y^2))x - (u + b(x^2 + y^2))y \\ &= rx - uy + \text{cubic terms} \end{aligned}$

Similarly, we get $\dot{y} = \omega x + ry + \text{cubic terms}$

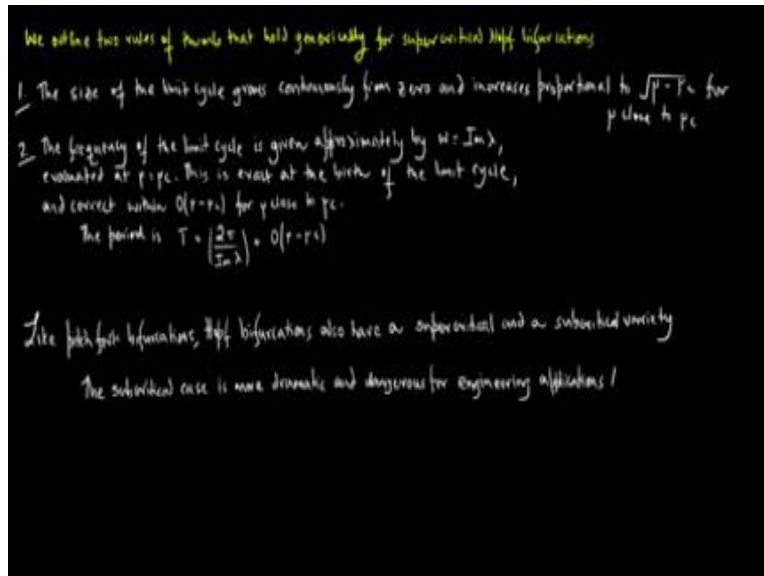
So the Jacobian at the origin is $A = \begin{pmatrix} r - u & -\omega \\ \omega & r \end{pmatrix}$ which has eigenvalues $\lambda = r \pm i\omega$

So the eigenvalues cross the imaginary axis from left to right as r increases from negative to positive values.

To get a sense of how the Eigen values behave during the bifurcation. We rewrite the system in Cartesian coordinates. So, we write $x = r \cos \theta$ and $y = r \sin \theta$, then $\dot{x} = \dot{r} \cos \theta - r \dot{\theta} \sin \theta$, which $= \mu r - r^3 \cos \theta - r \omega \sin \theta$, which is $= \mu x - x^3 - y^2 x - \omega y$ and the whole term times $x - \omega x + b \text{ times } x^3 + y^2 x$ and the whole term times y , which is $= \mu x - \omega y + \text{cubic terms}$.

And similarly, we get $\dot{y} = \omega x + \mu y + \text{cubic terms}$. So, the Jacobian at the origin is $A = \begin{pmatrix} \mu & -\omega \\ \omega & \mu \end{pmatrix}$, which has Eigen values $\lambda = \mu \pm i\omega$. So, the Eigen values actually cross the imaginary axis from left to right as μ increase from negative to positive values.

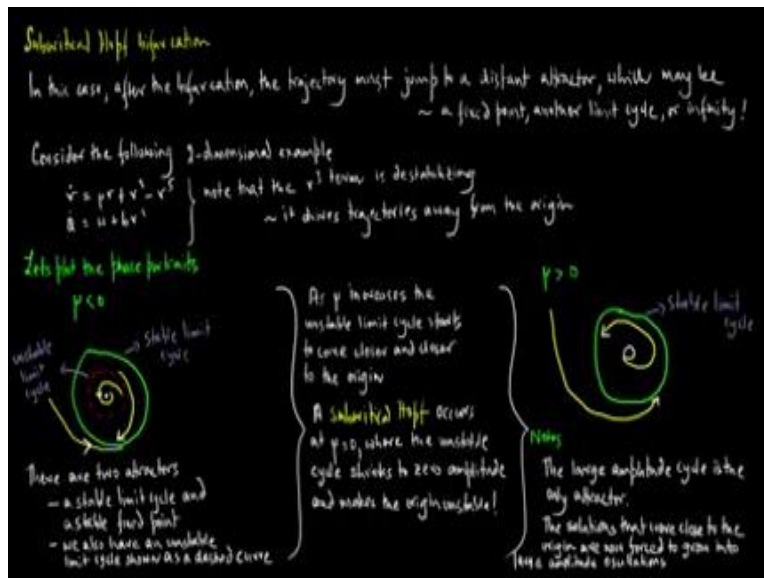
(Refer Slide Time: 10:03)



We now outline two rules of thumb that hold generically for supercritical half bifurcations. The first is that the size of the limit cycle grows continuously from zero and increases proportional to $\mu - \mu_{\text{critical}}$ square root, for μ close to μ_{critical} . The second is that the frequency of the limit cycle is given approximately by $\omega = -\text{Im}(\lambda)$ evaluated at $\mu = \mu_{\text{critical}}$.

This is exact at the birth of the limit cycle and correct within order $\mu - \mu_{\text{critical}}$, for μ close to μ_{critical} . The period is capital $T = 2\pi / \text{Im}(\lambda) + \text{order } \mu - \mu_{\text{critical}}$ terms. Now like pitchfork bifurcations, half bifurcations also have a supercritical and a subcritical variety. The subcritical case is much more dramatic and dangerous for engineering applications.

(Refer Slide Time: 11:50)

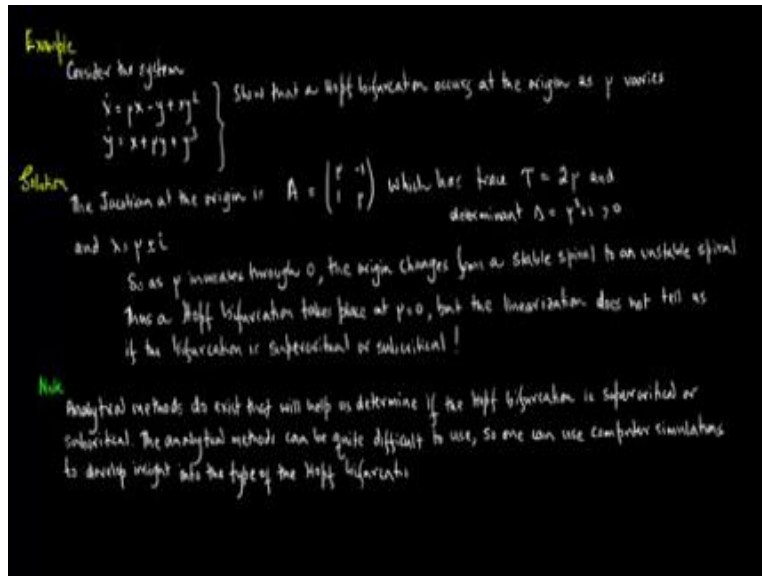


We now consider a subcritical half bifurcation. In this case after the bifurcation the trajectory must jump to a distance attractor which may be a fixed point, another limit cycle or of to infinity. So consider the following two dimensional examples, $\dot{r} = \mu r + r^3 - r^5$ and $\dot{\theta} = \omega + br^2$. Note that the r^3 term is destabilising ie it drives trajectories away from the origin.

Now let us plot the phase portrait with μ less than 0. We have a stable limit cycle and we have a unstable limit cycle. In this system, there are two attractors a stable limit cycle and a stable fixed point. We also have an unstable limit cycle shown as a dashed curve. As μ increases the unstable limit cycle starts to come closer and closer to the origin. A subcritical half occurs at $\mu = 0$, where the unstable cycle shrinks to zero amplitude and makes the origin unstable.

Now we consider μ greater than 0, we are left with the stable limit cycle. Here is some notes the large amplitude limit cycle is the only attractor in the system. The solutions that were close to the origin are now force to grow into large amplitude oscillations.

(Refer Slide Time: 14:56)



So, let us look at an example, consider the system $\dot{x} = \mu x - y + xy^2$ and $\dot{y} = x + \mu y + y^3$, show that a half bifurcation occurs at the origin as μ varies. So, let us look at the solution, the Jacobian at the origin is $A = \begin{pmatrix} \mu & -1 \\ 1 & \mu \end{pmatrix}$, which has trace 2μ and determinant $\mu^2 + 1$ is greater than 0 and $\lambda = \mu \pm i$.

So, as μ increases through zero the origin changes from a stable spiral to an unstable spiral. Thus, a half bifurcation takes place at $\mu = 0$, but the linearization actually does not tell us if the bifurcation is supercritical or subcritical. Now we note that the analytical methods do exist that will help us to determine, if the half bifurcation is supercritical or subcritical. But the analytical methods can be quite difficult to use. And so, one can first use computer simulations to develop insight into the type of the half of bifurcation.

(Refer Slide Time: 17:17)



Now the topic of discussion in this lecture was the half bifurcation. The half bifurcation is a very important one in both science and in engineering applications. Essentially it is bifurcation that leads to the birth of a limit cycle as a parameter in the system varies. It is worth mentioning that this bifurcation only occurs in dimensions two or higher. So, it does not have a counter part in one dimensional flows. This bifurcation comes with two variants a supercritical half and the subcritical half.

In a supercritical half, we get the birth of a small amplitude stable limit cycle as a parameter in the system varies. In the subcritical half is more dramatic bifurcation, where after the bifurcation trajectory may actually go off to a completely different fixed point, may go off to a distant limit cycle or may actually go off to the infinity. So, between these two variants the subcritical half is the one that should be avoided specially in engineering applications.