

**Introduction to Nonlinear Dynamics**  
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**Module -06**  
**Lecture-19**  
**2-Dimensional Flows, Linear Systems, Lecture 3**

(Refer Slide Time: 00:01)

**2-DIMENSIONAL FLOWS**  
**LINEAR SYSTEMS**

We wish to study the general case of an arbitrary  $2 \times 2$  matrix, and classify the possible phase portraits that can occur.

A 2 dimensional linear system is

$$\begin{cases} \dot{x} = ax + by \\ \dot{y} = cx + dy \end{cases} \quad \left. \begin{matrix} a, b, c, d \text{ are} \\ \text{parameters} \end{matrix} \right\}$$

Written compactly in vector form

$$\dot{\mathbf{x}} = A\mathbf{x},$$

where  $A = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$

$$\mathbf{x} = \begin{pmatrix} x \\ y \end{pmatrix}$$

Note: We will use green to denote vectors

**Classification of Linear Systems**

We seek trajectories of the form

$$\mathbf{x}(t) = e^{\lambda t} \mathbf{v}, \text{ where } \mathbf{v} \neq 0$$

is some fixed vector to be determined and  $\lambda$  is a growth rate also to be determined.

To find conditions on  $\mathbf{v}$  and  $\lambda$ , we substitute  $\mathbf{x}(t) = e^{\lambda t} \mathbf{v}$  into  $\dot{\mathbf{x}} = A\mathbf{x}$  to obtain

$$\lambda e^{\lambda t} \mathbf{v} = e^{\lambda t} A \mathbf{v}$$

Cancelling  $e^{\lambda t}$  gives  $A\mathbf{v} = \lambda \mathbf{v}$ , so the solutions exist if  $\mathbf{v}$  is an eigenvector of  $A$  with corresponding eigenvalue  $\lambda$ .

So the solution  $\mathbf{x}(t) = e^{\lambda t} \mathbf{v}$  is an eigen solution.

In this lecture we deal with classification of linear systems. We wish to study the general case of an arbitrary 2 by 2 matrix and classify the possible phase portraits that can occur. A two dimensional linear system is  $\dot{x} = ax + by$  and  $\dot{y} = cx + dy$  where  $a, b, c, d$  are parameters. So written compactly in vector form, we get  $\dot{\mathbf{x}} = A\mathbf{x}$ , where  $A$  is a  $2 \times 2$  matrix and  $\mathbf{x}$  is  $\begin{pmatrix} x \\ y \end{pmatrix}$ . Note that we will be using green to denote vectors.

So we seek trajectories of the form  $\mathbf{x}(t) = e^{\lambda t} \mathbf{v}$ , where  $\mathbf{v}$  is not equal to zero. Is some fixed vector to be determined and  $\lambda$  is a growth rate, which is also to be determined. To find conditions on  $\mathbf{v}$  and  $\lambda$ , we substitute  $\mathbf{x}(t) = e^{\lambda t} \mathbf{v}$  into  $\dot{\mathbf{x}} = A\mathbf{x}$  to obtain  $\lambda e^{\lambda t} \mathbf{v} = e^{\lambda t} A \mathbf{v}$ .

So cancelling  $e$  to the  $\lambda t$  gives  $A \text{ times } v = \lambda \text{ times } v$ . So the solution exists, if  $v$  is an Eigen vector of  $A$  with corresponding Eigen value  $\lambda$ . So the solution  $x$  of  $t = e$  to the  $\lambda t$  times  $v$  is an Eigen solution.

(Refer Slide Time: 02:51)

The eigenvalues of a matrix  $A$  are given by the characteristic equation  $\det(A - \lambda I) = 0$  where  $I$  is the identity matrix.

For a  $2 \times 2$  matrix  $A = \begin{pmatrix} a & b \\ c & d \end{pmatrix}$ , the characteristic equation becomes

$$\det \begin{bmatrix} a - \lambda & b \\ c & d - \lambda \end{bmatrix} = 0$$

Expanding the determinant yields

$$\lambda^2 - T\lambda + \Delta = 0$$

where

$$T = \text{trace}(A) = a + d$$

$$\Delta = \det(A) = ad - bc$$

then

$$\lambda_1 = \frac{T + \sqrt{T^2 - 4\Delta}}{2}, \quad \lambda_2 = \frac{T - \sqrt{T^2 - 4\Delta}}{2}$$

are the solutions of  $\lambda^2 - T\lambda + \Delta = 0$

Note, that the eigenvalues depend only on the trace and the determinant of  $A$ .

Normally, the eigenvalues are distinct:  $\lambda_1 \neq \lambda_2$ . In this case the corresponding eigenvectors  $v_1$  and  $v_2$  are linearly independent and span the entire plane.

In fact, any initial condition  $x_0$  can be written as a linear combination of eigenvectors, for example

$$x_0 = c_1 v_1 + c_2 v_2$$

So the general solution is

$$x(t) = c_1 e^{\lambda_1 t} v_1 + c_2 e^{\lambda_2 t} v_2$$

The Eigen values of a matrix  $A$  are given by the characteristic equation determinant of  $A - \lambda I = 0$ , where  $I$  is the identity matrix. So for a  $2 \times 2$  matrix  $A$  with entries  $a, b, c$  and  $d$ , the associated characteristic equation becomes the determinant of  $a - \lambda, b, c, d - \lambda = 0$ . Expanding the determinants yields  $\lambda^2 - \tau \lambda + \Delta = 0$ . Where  $\tau$  is the trace of  $A$  which is equal to  $a + d$  and  $\Delta$  is the determinants of  $A = ad - bc$ .

Then  $\lambda_1 = \frac{\tau + \sqrt{\tau^2 - 4\Delta}}{2}$  and  $\lambda_2 = \frac{\tau - \sqrt{\tau^2 - 4\Delta}}{2}$ . Then  $\lambda_1$  and  $\lambda_2$  are the solutions of  $\lambda^2 - \tau \lambda + \Delta = 0$ . Note that the Eigen values depend only on the trace and the determinant of  $A$ , normally the Eigen values are distinct, so  $\lambda_1$  is not equal to  $\lambda_2$ . In this case the corresponding Eigen vectors  $v_1$  and  $v_2$  are linearly independent and span the entire plane.

In fact any initial condition  $x$  not can be written as a linear combination of Eigen vectors. For example  $x$  of not  $= c_1 v_1 + c_2 v_2$ , so the general solution is  $x$  of  $t = c_1 e^{\lambda_1 t} v_1 + c_2 e^{\lambda_2 t} v_2$ .

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**Example** Solve the initial value problem

$$\begin{cases} \dot{x} = x + y \\ \dot{y} = 4x - 2y \end{cases} \text{ Subject to the initial condition } (x_0, y_0) = (2, -3)$$

Writing in matrix form we get

$$\begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix} = \begin{pmatrix} 1 & 1 \\ 4 & -2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}$$

We first find the eigenvalues of the matrix  $A$

For the matrix  $T = \text{trace}(A) = -1$   
 $\Delta = \det(A) = -6$

The characteristic equation is  $\lambda^2 + \lambda - 6 = 0$   
 which gives  $\lambda_1 = 2, \lambda_2 = -3$

Now we need to find the eigenvectors

Given an eigenvalue  $\lambda$ , the corresponding eigenvector  $v = (v_1, v_2)$  satisfies

$$\begin{pmatrix} 1-\lambda & 1 \\ 4 & -2-\lambda \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix}$$

So for  $\lambda_1 = 2$ , we get

$$\begin{pmatrix} -1 & 1 \\ 4 & -4 \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \text{ which gives } (v_1, v_2) = (1, 1) \text{ or any scalar multiple thereof}$$

For  $\lambda_2 = -3$ , we get

$$\begin{pmatrix} 4 & 1 \\ 4 & 1 \end{pmatrix} \begin{pmatrix} v_1 \\ v_2 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \end{pmatrix} \text{ which gives } (v_1, v_2) = (1, -4)$$

So let us consider an example, solve the initial value problem  $\dot{x} = x + y$  and  $\dot{y} = 4x - 2y$ . Subject to the initial condition  $x(0) = 2, y(0) = -3$ . So writing in matrix form we get  $\begin{pmatrix} \dot{x} \\ \dot{y} \end{pmatrix} = \begin{pmatrix} 1 & 1 \\ 4 & -2 \end{pmatrix} \begin{pmatrix} x \\ y \end{pmatrix}$ . So we first find the Eigen values of the matrix  $A$ , for the matrix the trace of  $A = -1$  and the determinant of  $A = -6$ .

The characteristic equation is  $\lambda^2 + \lambda - 6 = 0$ , which gives  $\lambda_1 = 2, \lambda_2 = -3$ , as solutions. Now we need to find the Eigen vectors, so given an Eigen value  $\lambda$  the corresponding Eigen vector  $v$ , suppose to be  $v_1$  and  $v_2$  satisfies  $(1 - \lambda) v_1 + 1 v_2 = 0$ .

So for  $\lambda_1 = 2$ , we get  $-1 v_1 + 1 v_2 = 0$  which gives  $v_1 = v_2$  or in fact any scalar multiple thereof. For  $\lambda_2 = -3$ , we get  $4 v_1 + 1 v_2 = 0$ , which gives us  $v_1 = -\frac{1}{4} v_2$ . (Refer Slide Time: 08:36)

In Summary

Write the general solution as a linear combination of the eigensolutions

The general solution is  $x(t) = c_1 \begin{pmatrix} 1 \\ 1 \end{pmatrix} e^{2t} + c_2 \begin{pmatrix} 1 \\ -4 \end{pmatrix} e^{-3t}$

We now need to compute  $c_1$  and  $c_2$  to satisfy the initial condition  $(x_0, y_0) = (2, -3)$

At  $t=0$ , the general solution becomes  $\begin{pmatrix} 2 \\ -3 \end{pmatrix} = c_1 \begin{pmatrix} 1 \\ 1 \end{pmatrix} + c_2 \begin{pmatrix} 1 \\ -4 \end{pmatrix}$

Which gives  $\begin{aligned} 2 &= c_1 + c_2 \\ -3 &= c_1 - 4c_2 \end{aligned}$

Which yields  $c_1 = 1, c_2 = 1$

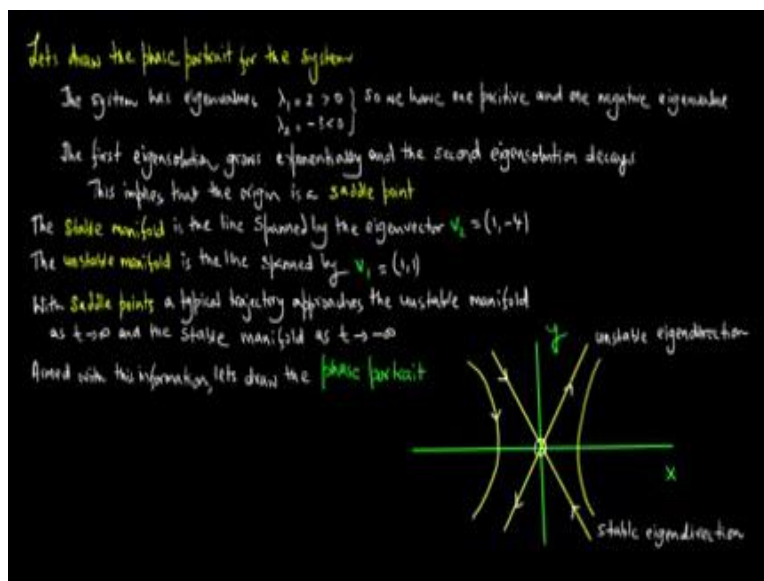
So the solution to the initial value problem is:

$$x(t) = e^{2t} + e^{-3t}$$

$$y(t) = e^{2t} - 4e^{-3t}$$

In summary  $v_1 = \begin{pmatrix} 1 \\ 1 \end{pmatrix}$  and  $v_2 = \begin{pmatrix} 1 \\ -4 \end{pmatrix}$ . So we can write the general solution as a linear combination of the Eigen solutions. The general solution is  $x(t) = c_1 \begin{pmatrix} 1 \\ 1 \end{pmatrix} e^{2t} + c_2 \begin{pmatrix} 1 \\ -4 \end{pmatrix} e^{-3t}$ . We now also need to compute  $c_1$  and  $c_2$  to satisfy the initial condition  $x(0) = 2, y(0) = -3$ . At  $t = 0$ , the general solution becomes  $\begin{pmatrix} 2 \\ -3 \end{pmatrix} = c_1 \begin{pmatrix} 1 \\ 1 \end{pmatrix} + c_2 \begin{pmatrix} 1 \\ -4 \end{pmatrix}$ . Which ends up giving  $\begin{aligned} 2 &= c_1 + c_2 \\ -3 &= c_1 - 4c_2 \end{aligned}$ , which yields  $c_1 = 1$  and  $c_2 = 1$ . So finally we have the solution to the original initial value problem, which is  $x(t) = e^{2t} + e^{-3t}$ ,  $y(t) = e^{2t} - 4e^{-3t}$ .

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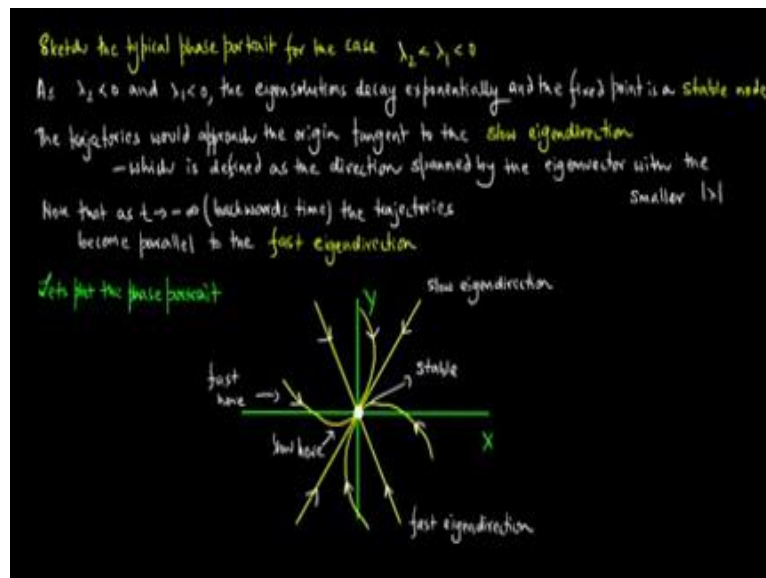


Now let us draw the phase portrait for the system. The system has two Eigen values 2 and -3, so we have one positive and one negative Eigen value. The first Eigen solution grows

exponentially and second Eigen solution actually decays. So this implies that the origin is a saddle point. The stable manifold is the line spanned by the Eigen vector  $v_2 = 1 \ -4$  and the unstable manifold is the line spanned by the Eigen vector  $v_1 = 1 \ 1$ .

With saddle points a typical trajectory approaches the unstable manifold as  $t$  tend to infinity and the stable manifold as  $t$  tends to minus infinity. So now armed with all of this information, let us draw the phase portrait. So we plot  $y$  versus  $x$ , we first highlight the stable Eigen direction and then highlight the unstable Eigen direction. And we now go ahead and complete the phase portrait and that is the completed phase portrait.

(Refer Slide Time: 12:47)



Now let us sketch the typically phase portrait for the case  $\lambda_2$  less than  $\lambda_1$  is less than zero. As  $\lambda_2$  is less than zero and  $\lambda_1$  is less than zero. The Eigen solutions decay exponentially and the fixed point is a stable node. The trajectories would approach the origin, tangent to the slow Eigen direction, which is defined as the direction spanned by the Eigen vector with the smaller absolute value of  $\lambda$ .


Note that as  $t$  tends to minus infinity that is going backwards in time the trajectories become parallel to the fast Eigen direction. So now let us plot the typical phase portrait. So we plot  $y$  versus  $x$  we highlight the slow Eigen direction. Then we highlight the faster Eigen direction that is the stable node, that is were it could be fast, that is were it would be slow. And now we go ahead and complete the phase portrait.

(Refer Slide Time: 14:53)


What happens as the eigenvalues are complex numbers?

In this case the fixed point is either a center or a spiral

**Center**



**Spiral**



Note that in the center the origin is surrounded by a family of closed orbits. The centers are neutrally stable as nearby trajectories are neither attracted nor repelled from the fixed point.

Let's look at the eigenvalues in some detail

$$\lambda_{1,2} = \frac{1}{2} \left( T \pm \sqrt{T^2 - 4\Delta} \right)$$

So we get complex eigenvalues when  $T^2 - 4\Delta < 0$

Let's write the eigenvalues as  $\lambda_{1,2} = \alpha \pm i\omega$ , where  $\alpha = T/2$ ,  $\omega = \frac{1}{2} \sqrt{4\Delta - T^2}$

Now  $\omega \neq 0$ . The eigenvalues are distinct, and general solution is  $x(t) = c_1 e^{\lambda_1 t} v_1 + c_2 e^{\lambda_2 t} v_2$

The  $\lambda$ 's are complex so the  $c$ 's and the  $v$ 's are complex.

Now what happens in the case if the Eigen values are complex numbers? In this case the fixed point are either a center or a spiral. Now we go ahead and plot the center, so that is what typical centers would look like. Now we plot the spiral, essentially the trajectories spiral inwards. Note that in the center, the origin is surrounded by family of closed orbits. The centers are neutrally stable as nearby trajectories are neither attracted nor repelled from the fixed point.

Now let us look at the Eigen values in some detail. So  $\lambda_{1,2} = \frac{1}{2} \text{trace} \pm \frac{1}{2} \sqrt{\text{trace}^2 - 4\Delta}$ . So we get complex Eigen values when  $\text{trace}^2 - 4\Delta < 0$ . So let us write the Eigen values as  $\lambda_{1,2} = \alpha \pm i\omega$ , where  $\alpha$  is equal to  $\text{trace}/2$  and  $\omega = \frac{1}{2} \sqrt{4\Delta - \text{trace}^2}$ .

Now  $\omega$  is not equal to zero, so the Eigen value are distinct and the general solution is  $x(t) = c_1 e^{\lambda_1 t} v_1 + c_2 e^{\lambda_2 t} v_2$ . Now note that the  $\lambda$ 's are complex and so  $C$ 's and the  $V$ 's are also complex.

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What this essentially means is that  $x(t)$  involves linear combinations of  $e^{(\alpha \pm i\omega)t}$

By Euler's formula  $e^{i\omega t} = \cos \omega t + i \sin \omega t$

So  $x(t)$  is a combination of terms involving  $e^{\alpha t} \cos \omega t$  and  $e^{\alpha t} \sin \omega t$

Hence we have exponentially decaying solutions if  $\alpha = \text{Re}(\lambda) < 0$  and growing solutions if  $\alpha > 0$

the fixed points are stable spirals and unstable spirals

If  $\alpha = 0$  (pure imaginary eigenvalues), then the solutions are periodic with period  $T = \frac{2\pi}{\omega}$

The oscillations have fixed amplitude and the fixed point is a center!

Now what this essentially means is that  $x$  of  $t$  involves linear combinations of  $e$  to the  $\alpha + i$   $\omega$  times  $t$ . By Euler's formula  $e$  to the  $i \omega t + \omega t + i \sin \omega t$ , so  $x$  of  $t$  is a combination of terms involving  $e$  to the  $\alpha t \cos \omega t$  and  $e$  to the  $\alpha t \sin \omega t$ . Hence we have exponentially decaying solutions, if  $\alpha$  which is the real part of  $\lambda$  is less than zero and growing solutions, if  $\alpha$  is greater than zero.

So the fixed points are stable spirals, if  $\alpha$  is less than zero and unstable spirals, if  $\alpha$  is greater than zero. If  $\alpha = 0$ , ie we have pure imaginary Eigen values. Then the solutions are periodic with period capital  $T = 2\pi$  by  $\omega$ . The oscillations have fixed amplitude and the fixed point is the center.

(Refer Slide Time: 19:19)

What happens as the eigenvalues are equal?

So far we have been assuming that the eigenvalues are distinct

Let's assume that  $\lambda_1 = \lambda_2 = \lambda$ , and consider two cases

Case (i) There are two independent eigenvectors corresponding to  $\lambda$ , or

Case (ii) There is only one

Case (i)

In this case, the eigenvectors span the plane and every vector is an eigenvector with the same eigenvalue

Let's get a better understanding of this.

Write an arbitrary vector  $x_0$  as a linear combination of the two eigenvectors  $x_0 = c_1 v_1 + c_2 v_2$


Then  $A x_0 = A(c_1 v_1 + c_2 v_2) = c_1 \lambda v_1 + c_2 \lambda v_2 = \lambda x_0$

So  $x_0$  is an eigenvector with eigenvalue  $\lambda$

Then if  $\lambda \neq 0$ , all trajectories are straight lines through the origin ( $x(t) = e^{\lambda t} x_0$ ) and the fixed point is a star node

If  $\lambda = 0$ , the whole plane is just filled with fixed points - the system is  $\dot{x} = 0$

for  $\lambda < 0$ , we get

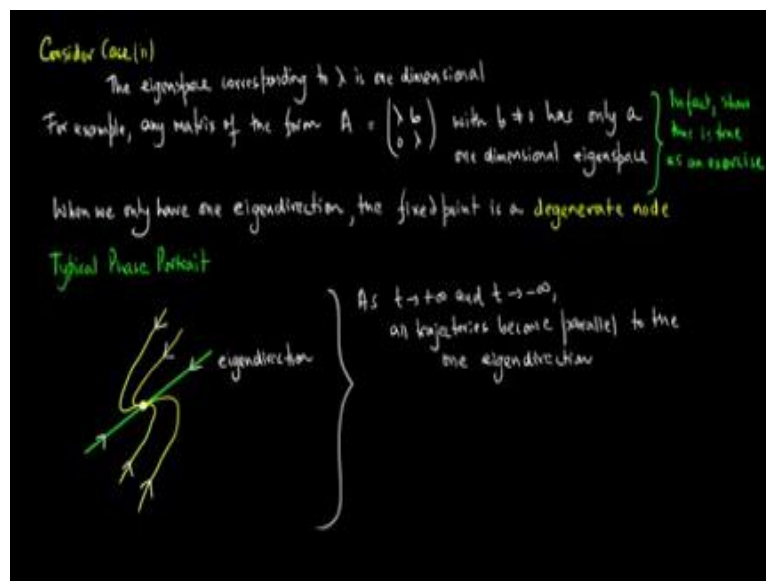


Now what happens as the Eigen values are equal, so far we have been assuming that the Eigen values are distinct? So let us assume that  $\lambda_1 = \lambda_2 = \lambda$  and consider two cases. Case1, there are two independent Eigen vectors corresponding to  $\lambda$  or case2 there is only one. So now let us consider case1 in some detail. In this case the Eigen vectors spanned the plane and every vector is an Eigen vector with the same Eigen value  $\lambda$ .

So let us get a better understanding of this, write an arbitrary vector  $x$  not as a linear combination of the two Eigen vectors. So  $x \neq c_1 v_1 + c_2 v_2$ , then  $A \text{ times } x \neq A \text{ times } c_1 v_1 + c_2 v_2 = c_1 \lambda v_1 + c_2 \lambda v_2 = \lambda \text{ times } x$ . And so  $x$  is an Eigen vector with Eigen value  $\lambda$ . Then if  $\lambda$  is not equal to zero, all trajectories are straight lines through origin  $x$  of  $t = e$  to the  $\lambda t \text{ times } x$  and the fixed point is a star mode.

If  $\lambda = 0$ , the whole plane is simply filled with fixed points. The system is  $\dot{x} = 0$ , for  $\lambda$  is not equal to zero. We get the following phase portrait, where we essentially have straight lines that are following through the origin and point is the star mode.

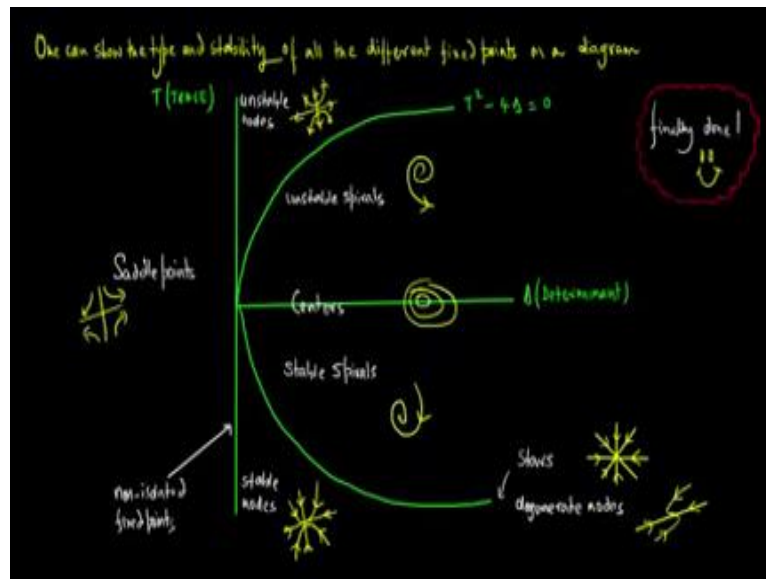
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Now and let us consider case2 the Eigen space corresponding to  $\lambda$  is one dimensional. So for example any matrix of the form  $A = \begin{pmatrix} \lambda & 0 \\ 0 & \lambda \end{pmatrix}$ , with  $b$  not equal to zero has only a one dimensional Eigen space. In fact we request that, to show this is true as an exercise. When we have only one Eigen direction the fixed point is a degenerate mode.

So now let's plot the typical phase portrait that is the only Eigen direction that completes the typical phase portrait as  $t$  tends to positive infinity and  $t$  tends to negative infinity all the trajectories become parallel to the one Eigen direction.

(Refer Slide Time: 24:13)



One can now show the type and stability of all the different fixed points on a single diagram. So we plot the trace versus the determinant  $\delta$  and we also plot the curve  $t^2 - 4\delta = 0$ . So that is where the saddle points live, that is where the unstable nodes are, that is region for unstable spirals, centers, stable spirals, stable nodes and along with curve you will have stars and degenerate nodes. And finally this is the point for non isolated fixed points.

So now we will plot the simple caricatures, for the typical phase portraits. We have unstable spiral, a center, stable spiral, stable nodes, stars and finally degenerate nodes. So we are finally, finally done that is hard work I think we deserve smile face at the end of it all.

(Refer Slide Time: 26:04)



The intention of this lecture was to look at the classification of two dimensional linear systems in terms of the qualitatively different phase portraits that could have come from such systems. So when you are looking at a two dimensional system you have equations of the form  $\dot{x} = ax + by$  and  $\dot{y} = cx + dy$  where  $a$ ,  $b$ ,  $c$  and  $d$  are parameters. And in the last question of the form, what happens when you have negative Eigen values that they are distinct.

What happens when you have complex Eigen values? What happens when the Eigen values are in fact the same? So essentially what we did in the lecture was we systematically went through all the numerous cases and then eventually put it all together in a plot, while you plotted the trace versus determinant, highlighting qualitatively different phase portraits that could occur depending on where you are in the trace versus determinant plot.