

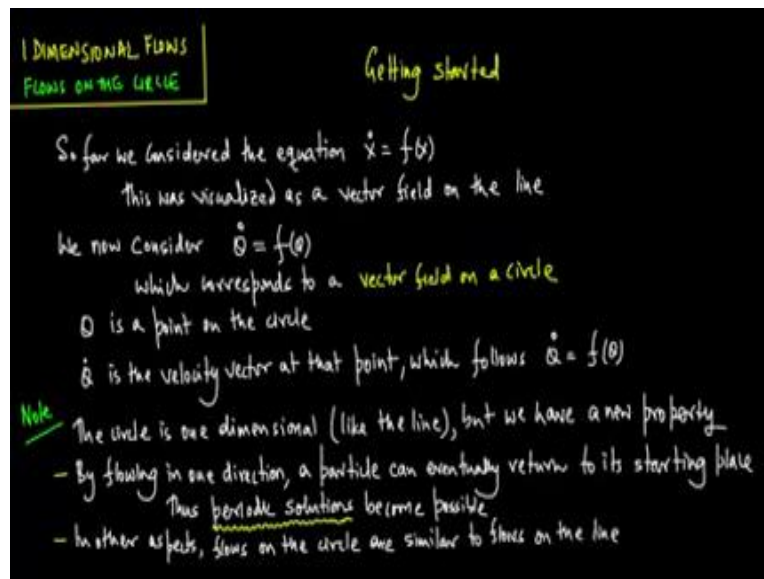
Introduction to Nonlinear Dynamics
Prof. Gaurav Raina
Department of Electrical Engineering
Indian Institute of Technology, Madras

Module -05

Lecture-15

1-Dimensional Flows, Flows on the Circle, Lecture 1

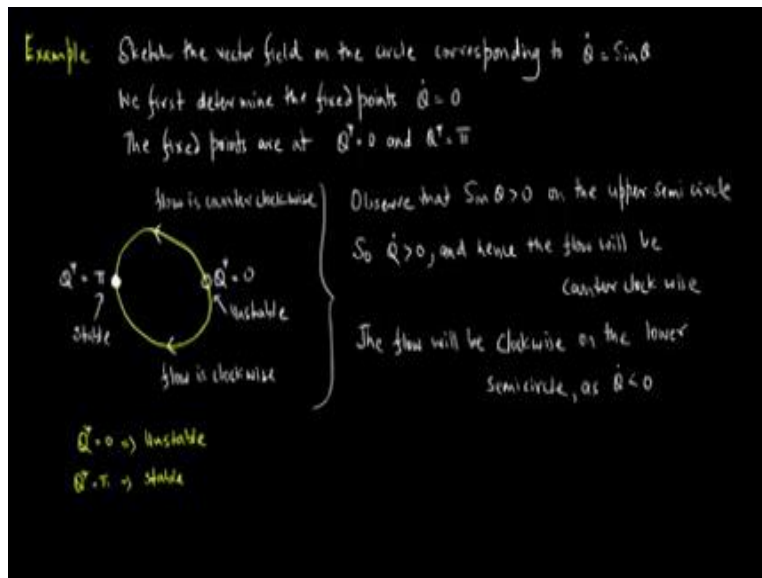
(Refer Slide Time: 00:01)



We are still dealing with one dimensional flows except that now our focus is on the circle. So, we will not be dealing with flows on the circle rather than on flows on the line. So, let us get started. So far, we have considered the equations $\dot{x} = f(x)$ and this equation was visualised as a vector field on the line. We now consider $\dot{\theta} = f(\theta)$, which corresponds to a vector field on a circle.

So θ is a point on the circle and $\dot{\theta}$ is the velocity vector at that point, which follows $\dot{\theta} = f(\theta)$. Here are some notes. The circle is one dimensional just like the line, but we have a new property by flowing in one direction a particle can eventually return to its starting place. Thus, in effect periodic solutions actually become possible. In other aspects, flows on the circle are actually quite similar to flows on the line.

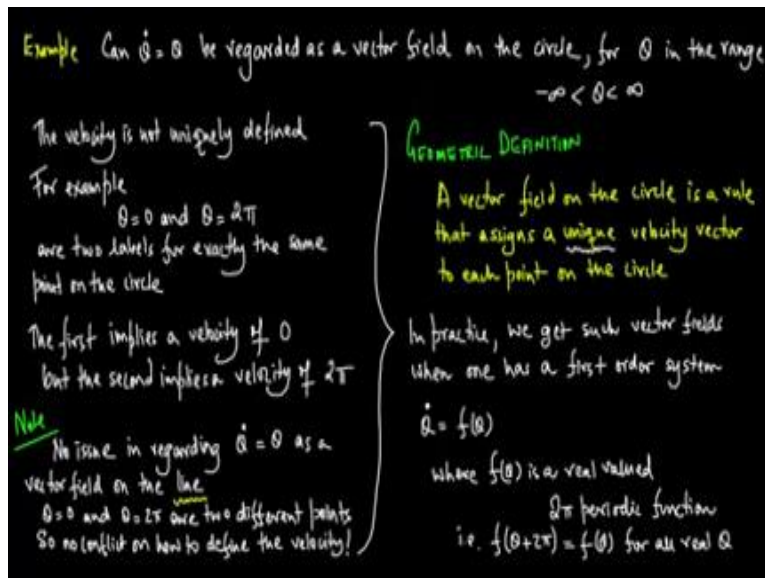
(Refer Slide Time: 01:46)



Now let us consider an example. Sketch the vector field on the circle corresponding to $\dot{\theta} = \sin \theta$. We first determine the fixed points by putting $\dot{\theta} = 0$ the fixed points are at $\theta^* = 0$ and $\theta^* = \pi$. So we plot the circle and highlight $\theta^* = 0$ and $\theta^* = \pi$. Now observe that $\sin \theta$ is greater than zero, on the upper semi circle, so $\dot{\theta}$ is greater than zero and hence the flow will be counter clockwise.

The flow will be clockwise on the lower semicircle as $\dot{\theta}$ is less than zero. So, the flow is counter clockwise on the upper semicircle and the flow is clockwise on the lower semicircle, so highlighting the fixed points one is unstable and the other is stable. So $\theta^* = 0$ is unstable and $\theta^* = \pi$ is stable.

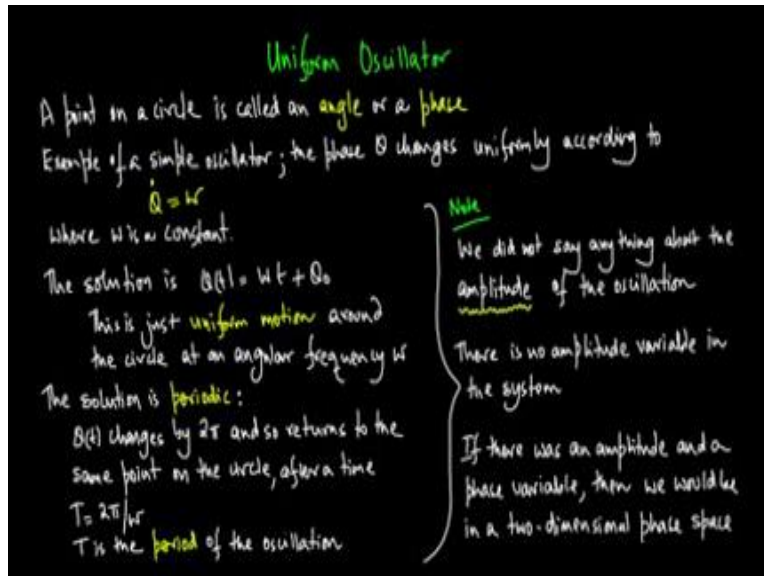
(Refer Slide Time: 03:17)



Now let us consider another example. Can $\dot{\theta} = \theta$ be regarded as a vector field on the circle for θ in the range, where θ is less than infinite, greater than minus infinity. The problem is that the velocity is not uniquely defined. For example, $\theta = 0$ and $\theta = 2\pi$ are two labels, for exactly the same point on the circle. The first implies a velocity of zero, the second implies a velocity of 2π . Now note that there is actually no issue in regarding $\dot{\theta} = \theta$ as the vector field on the line, $\theta = 0$ and $\theta = 2\pi$ are two different points.

So, there is actually no conflict on how to define the velocity. So, we now go ahead and offer a geometric definition. A vector field on the circle is a rule that assigns a unique velocity vector to each point on the circle. In practice, we get such vector fields when one has a first order system $\dot{\theta} = f(\theta)$. Where $f(\theta)$ is a real valued 2π periodic function i.e. $f(\theta + 2\pi) = f(\theta)$ for all real θ .

(Refer Slide Time: 05:23)



Let us consider a uniform oscillator; a point on a circle is called an angle or a phase. So, example of a simple oscillator is as follows, the phase θ changes uniformly according to $\dot{\theta} = \omega$ where ω is a constant. The solution is $\theta = \omega t + \theta_0$. Now this is just uniform motion around the circle at an angular frequency of ω . Note that the solution is periodic θ changes by 2π and so returns to the same point on the circle after a time $T = 2\pi/\omega$. T is the period of the oscillation.

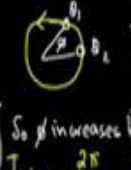
Note that we actually did not say anything about the amplitude of the oscillation. There is actually no amplitude variable in the system, if there were amplitude and a phase variable then we would be in a two-dimensional phase space.

(Refer Slide Time: 07:18)

Example Two runners A and B run at a steady pace around a circular track. It takes A T_1 seconds to run once around the track. It takes B T_2 seconds to run once around the track. } $T_2 > T_1$
 Note that A will periodically overtake B.
Q How long does it take for A to lap B once, on the assumption that they both start together?

Solution
 Let $\dot{\theta}_1$ be the speed of A. This says: A runs at a steady pace and completes a circuit every T_1 seconds.
 Then $\dot{\theta}_1 = \omega_1$, where $\omega_1 = 2\pi/T_1$. Similarly: we suppose that $\dot{\theta}_2 = \omega_2$, $\omega_2 = 2\pi/T_2$ for B.
 The condition for A to lap B is that the angle between them has increased by 2π . We get $\dot{\phi} = \dot{\theta}_1 - \dot{\theta}_2 = \omega_1 - \omega_2$.
 So if we define the phase difference $\phi = \theta_1 - \theta_2$, we need to find how long it takes for ϕ to increase by 2π .
 So ϕ increases by 2π after a time

$$T_{\text{lap}} = \frac{2\pi}{\omega_1 - \omega_2} = \left[\frac{1}{T_1} - \frac{1}{T_2} \right]^{-1}$$



Now let us consider an example two runners A and B run at a steady pace around a circular track. It takes A, T_1 seconds to run once around the track and it takes B, T_2 seconds to run once around the same track, assuming that T_2 is greater than T_1 . So, note that A will in fact periodically overtake B. So, the question is how long does it take for A to lap B once, on the assumption that they actually both start together?

So here is the solutions let $\dot{\theta}_1$ be the speed of A, when $\dot{\theta}_1 = \omega_1$ where $\omega_1 = 2\pi / T_1$. This says A runs at a steady pace and completes a circuit every T_1 seconds. Similarly, we suppose that $\dot{\theta}_2 = \omega_2$, where $\omega_2 = 2\pi / T_2$ for B. The condition for A to lap B is that the angle between them has to increase by 2π . So, if you define the phase angle, the phase difference $\phi = \theta_1 - \theta_2$.

We need to find how long it takes for ϕ to increase by 2π . So that is the simple schematic of the circular track. We note down $\dot{\theta}_1$, $\dot{\theta}_2$ and phase difference ϕ . So, we get $\dot{\phi} = \dot{\theta}_1 - \dot{\theta}_2$ which is equal to $\omega_1 - \omega_2$. So, the phase difference ϕ actually increases by 2π after a time $T_{\text{lap}} = 2\pi / \omega_1 - \omega_2$, which is equal to $1 / (1/T_1 - 1/T_2)$ to the -1.

(Refer Slide Time: 10:36)



This was a short lecture, which got a started-on flows on the circle, we are still dealing with one dimensional flows. However, we have moved away from flows of the line to flows on the circle. So, we now deal with the equation of the form $\dot{\theta} = f(\theta)$ where θ is point on the circle and $\dot{\theta}$ is the velocity vector at that point. And there is one key difference between flows on the line and flows on the circle and that as follows, say imagine that you are on a circle rather than a line and if you keep going around.

Essentially what will happen is that after a point a time you come back to the same point, so what will happen is that in flows on the circle periodic solutions actually becomes possible and that is the key differentiator between the flows on the line and flows on the circle. And in this lecture, we started off by introducing very simple uniform oscillator of the form $\dot{\theta} = \omega$.