

**Networks and Systems**  
**Prof V G K Murti**  
**Department of Electrical Engineering**  
**Indian Institute of Technology, Madras**

**Lecture - 19**  
**Fourier Transforms (7)**

In the last lecture, we saw how the Fourier transform methods can be employed to evaluate the transients in a network. A few comments about these methods: essentially, when we used to find out the transient behavior of a network, it is generally after a switching operation. Let us assume that some excitation function is introduced into the circuit at  $t$  equals to 0. And we wish to find out the response of the network following the switching operation. We should make 2 assumptions first of all let us assume that the network is initially relaxed.

That means, it has no initial energy is stored in the network which other words means the capacitors are initially uncharged and the inductors carry no current. This assumption of the network being initially relaxed does not entail any loss of generality. Because, if you indeed there are initial charges on the capacitors and initial current in the inductors. They can be replaced by equivalent sources and these equivalent sources can be considered to be additional excitations.

The principle of superposition is brought into play to evaluate the responses to the various excitations and adding them up to find the total solution. Therefore, we shall assume that the network is initially relaxed and we are interested in finding over the response of the network following this switching of the excitation at  $t$  equal to 0. A second assumption we will make is that the natural response of the network is such is 1 such that, it dies down the negligible proportion in course of time.

This is common to most network that we deal with particular deceptive network and therefore, this is once again does not seriously disturb the generality of this equation. However, the fourier methods can also be applied where the natural frequency the network have 0 real part in other words there is sustained natural response oscillation, sinusoid oscillations or a dc component. But the analysis become little more complicated, but we will avoid situation of this type in this course.

That means, we will assume that the natural response decays with time under these assumptions when we have an excitation switched in at  $t = 0$ . We expect the response also to be 0 for  $t < 0$ . Because the network is assumed to be causal and therefore, any excitation applied at  $t = 0$  cannot produce a response prior to that. So, we expect the response to be 0 for  $t < 0$  and then become nonzero from  $t = 0$  onwards for positive  $t$ .

Now, the Fourier method tells us that any such excitation function can be split up as the sum of tiny sinusoids starting from  $t = -\infty$  onwards and going up to  $t = \infty$ . Now, each of the sinusoids produces a response which can be evaluated using the steady state sinusoid circuit theory methods, that is using the frequency response techniques.

Then, all these responses can be added up to find the total response. So, the sum of the steady state responses is what we obtained for the Fourier integral approach. And it turns out that, the sum of the various steady state responses for the types of networks that we have talked about is also will become 0 up to  $t = 0$ .

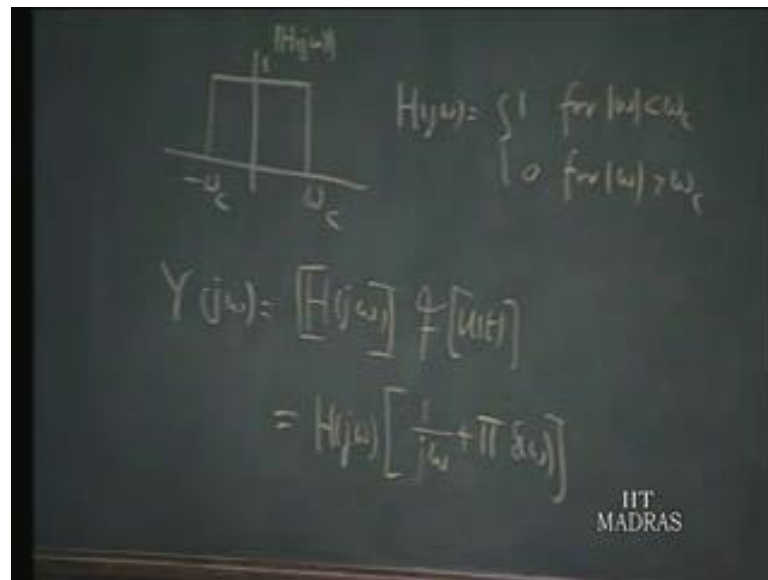
So, this is as should be because no response can occur before the application of the excitation and we further observe that, if there are any natural response terms produced at  $t = -\infty$  at the time the excitation starts all the sinusoids start from  $t = -\infty$ .

All these natural responses will decay to negligible proportions by the time  $t = 0$  arrives and this is the period which we are interested in. We are interested in finding out the behavior of the network from  $t = 0$  onwards. That means, the natural response terms which might have emanated at  $t = -\infty$  would have decayed to negligible proportions, when we come to  $t = 0$  onwards.

The sum of the steady state response that we obtain will include transient terms natural responses starting from  $t = 0$  that is what we are really interested in. So, the steady state responses will add up to 0 up to  $t = 0$ . And from  $t = 0$  onwards the response that we get the total response that we get include not only the forced response, but also the transient response or the natural response of the network starting from  $t = 0$ . That is how it goes.

Now, let us take a second example to illustrate the application of the Fourier theory to a network or system which is characterized by the frequency response. So, let me take filter network and see how the Fourier theory network, Fourier theory can be applied to this network.

(Refer Slide Time: 06:27).

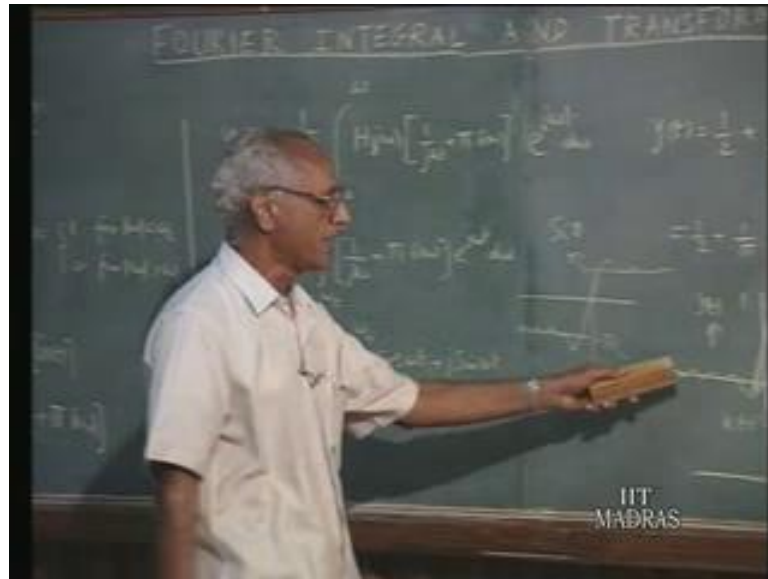


Let us consider an ideal low pass filter and we apply a unit step function as its excitation and we like to find out what is the response is. Now, the ideal low pass filter has a frequency response function like this. So, the frequency response function of the ideal low pass filter can be considered to be 1 for omega magnitude less than omega c for omega magnitude greater than omega c.

That is this is kind of response function that we have now, under these conditions this is the system function of the low pass filter. Therefore,  $y(j\omega)$  the Fourier transform of the output is  $h(j\omega)$  the system response function frequency response function times the Fourier transform. This is  $h(j\omega)$  times the Fourier transforms of  $u(t)$  which is the input function. Therefore, this will be  $h(j\omega)$  times the Fourier transform of the step function is  $1/j\omega + \pi\delta(\omega)$ .

So, this is  $y(j\omega)$ , if you find the inverse fourier transform of this you would naturally get the  $y(t)$  this is what we are interested in.

(Refer Slide Time: 08:27)



So,  $y$  of  $t$  is obtained as  $\frac{1}{2\pi}$  from minus infinity to plus infinity of  $h$  of  $j\omega$  times  $\frac{1}{j\omega} + \pi\delta(\omega)$   $e^{j\omega t}$   $d\omega$ . This is the inverse Fourier transform function. Now, we are integrating with respect to  $\omega$  and we know that  $h$  of  $j\omega$  vanishes when  $\omega$  magnitude exceeds  $\omega_c$ .

Therefore, we need to carry out this integration from between the limits minus  $\omega_c$  to plus  $\omega_c$  because outside the integral vanishes. Therefore, this minus  $\omega_c$  to plus  $\omega_c$   $\frac{1}{j\omega} + \pi\delta(\omega)$   $e^{j\omega t}$   $d\omega$ . Now, we observe as far the second term is concerned  $\pi\delta(\omega)$  is multiplied by  $e^{j\omega t}$   $d\omega$ . Now, the multiplying factor  $\delta(\omega)$  we can substitute  $\omega$  equal to 0 in that.

Therefore, this becomes 1 and we are integrating a  $\delta(\omega)$  over the range which covers  $\omega = 0$ . Therefore, that results in  $\pi$ . That means, the result of this portion of the integration will result in  $\pi$  and therefore, the  $\pi$  multiplied by  $\frac{1}{2\pi}$  and become half plus we have minus  $\omega_c$  to plus  $\omega_c$   $e^{j\omega t}$  can be written as,  $\cos(\omega t) + j\sin(\omega t)$  divided by  $j\omega$ .

Now, this portion  $\cos(\omega t)$  divided by  $j\omega$  is an odd function of  $\omega$  and we are integrating between the symmetrical limits from minus  $\omega_c$  to plus  $\omega_c$ .

Consequently, this is being the odd function and you are integrating between symmetrical limits minus  $\omega_c$  to plus  $\omega_c$  the contribution of this portion become 0. Therefore, the contribution arises only from  $\sin \omega_c t$  by  $\omega_c$  term and  $\sin \omega_c t$  by  $\omega_c t$  is an even function of  $\omega_c$  therefore, I can write this as  $2 \int_0^{\omega_c} \sin \omega_c t \text{ by } \omega_c d\omega_c$ .

That means, this will become  $y$  of  $t$  becomes half plus 2 times we must have also  $\frac{1}{2\pi}$  from this because  $\frac{1}{2\pi}$  times this  $\frac{1}{2\pi}$  you must have. So,  $\frac{1}{2\pi}$  you must have. So,  $2 \int_0^{\omega_c} \sin \omega_c t \text{ by } \omega_c \sin \omega_c t \text{ by } \omega_c t \text{ by } \omega_c d\omega_c$ . Or this is half of  $\frac{1}{\pi}$  this can be put as  $\sin$  integral of  $\omega_c$ .

Then, after all I can write this I can introduce  $t$  here and  $d\omega_c t$  here and this can be thought of a  $\sin \theta$  by  $\theta d\theta$ . And once you have this is taken as  $\theta$  this will become  $\omega_c t$ . So, this can be put in this form. Now, if you evaluate this you have this is the constant dc term as far  $\sin$  integral is concerned. You recall the  $\sin$  integral we have a characteristic like this it goes to  $\pi$  up on 2 and minus  $\pi$  up on 2 this is  $\sin$  integral of  $\theta$ , which we have thought of discussed earlier.

This is multiplied by  $\pi$ ; that means, this quantity will reach asymptotically minus half of negative values of time and plus half for positive values of time. So, if you add to that this half eventually  $y$  of  $t$  will be like this will be some oscillations like this and eventually, this reaches 1 and this  $t$ . So, this will be the response of the ideal low pass filter for unit step input. So, if you have a unit step input here 1 the output will be like this.

And depending up on the cut off frequency  $\omega_c$  for the low pass filter if you make the  $\omega_c$  larger and larger this will become more closer, closer to the step function. That means, the overshoots the oscillation will reduce and this portion of this curve will be steeper and steeper.

So, the more the cut off frequency of the low pass filter; that means, the greater bandwidth of the low pass filter become closer and closer to this. But, interesting thing is you are applying a  $\sin$  wave like this actually the excitation function is 0 up to this point and you are applying this is your  $u$  of  $t$ . But the response has is non 0 for negative values of time.

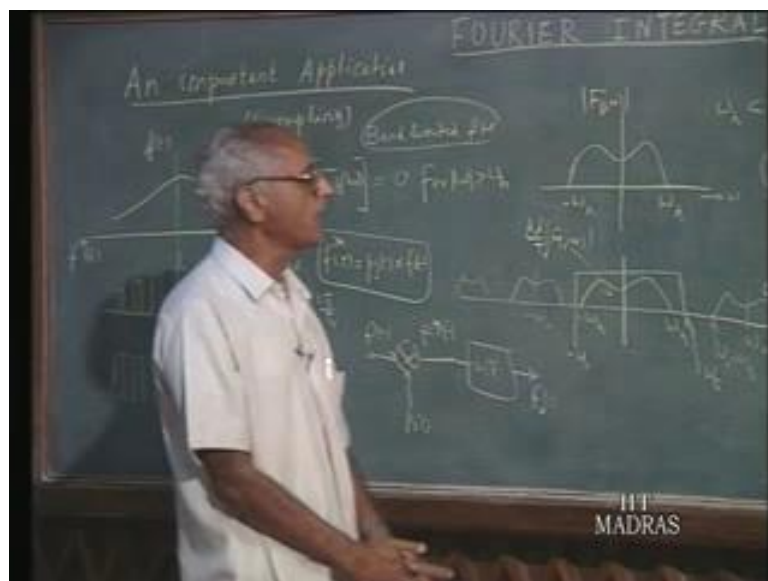
So, even before the excitation is applied there is a response. So, this clearly shows that this is a non causal system. And which means physically this is something which we cannot expect to achieve in the physically and therefore, the ideal low pass filter that we are talking about here this type of characteristic; can only be approximated it can never be realized in practice. Because, we believe that no physical system can give response even before the excitation is applied.

So, this is non causal situation and this is 1 way of showing that an ideal low pass filter of that particular characteristic is not achievable in practice. All I head can do is modify this characteristic in such way that, this portion of the response is avoided. That means, all low pass filters that we can build in practice cannot give this type of response. It can only give a response suppose this is the input then only starts giving response on this point onwards may be something like this and this can be avoided.

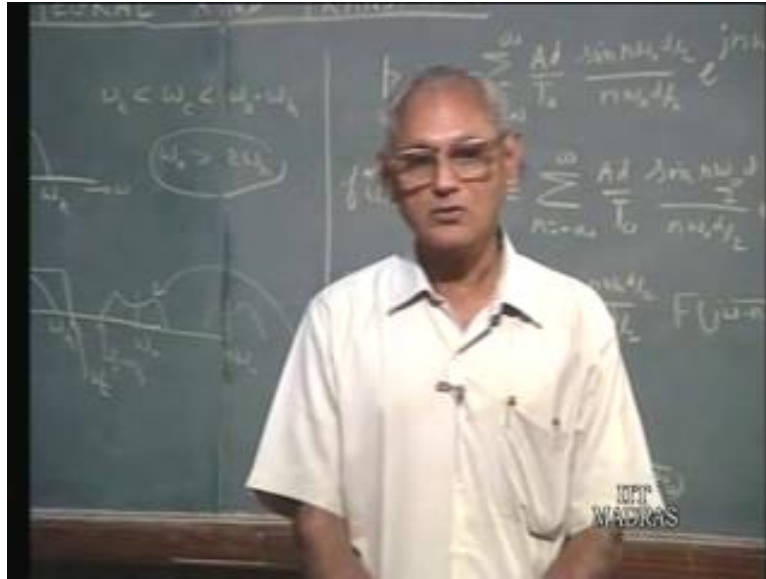
That means the ideal low pass filter characteristic cannot be obtained in practice. However, my point in showing this giving this example is that Fourier transform methods can be applied to network characterized in terms of frequency response as here and this is very convenient tool in such situations.

As an illustration of the application of the Fourier transform method, let us consider 1 application which is referred to as the sampling theorem.

(Refer Slide Time: 16:02).



(Refer Slide Time: 16:45).



Let us consider the function of time like this which is band limited in that the Fourier transform for that is confined to may be like this. So, the Fourier transform is 0 for values of  $\omega$  beyond  $\omega_h$  or we can say that the highest frequency component that is present in  $f$  of  $t$  is  $\omega_h$ .

No component of frequency beyond  $\omega_h$  and the positive side minus  $\omega_h$  on the negative side exists in this particular  $f$  of  $t$ . Such functions are called band limited functions band limited  $f$  of  $t$  that means, the frequency band of the component that exists in  $f$  of  $t$  are limited to this band nothing beyond that. Suppose, we do not consider  $f$  of  $t$  for the entire time duration, but generate from  $f$  of  $t$  a sample version of that. That means, this is  $f$  of  $t$  we have variation; you do not consider the entire variation by take samples from that tiny interval like this.

So, we consider  $f$  of  $t$  for small intervals of duration  $d$  and the period between 2 successive samples is  $T$ . This we may call sampled version of  $f$  of  $t$  we will call that  $f^*$  of  $t$ ; this is sampled version. So, notice  $f$  of  $t$  is the continuous function of time we are considering a sampled version, we are taking samples each time for duration  $d$  and taking the intervals between the successive sample is  $T$ .

Now, it turns out that the Fourier spectrum of this sampled version will be something like this, if this is the original Fourier spectrum. The  $f^*$  of  $t$  the fourier transform for that will be the same spectrum is repeated endlessly in both directions with a

reduced scale factor. If this is  $\omega h$  this is  $\omega h$  and the sampling period is  $T$  not and let  $\omega$  be the corresponding angular frequency. Then the next section is centered around  $\omega$ , this is  $\omega$  and this is the other section is repeated at  $2\omega$  and so on.

So, endlessly it goes in both directions; that means, the sample version incorporates in it the Fourier transform of the original signal to a reduced scale. Besides, it contains several other reproduction of the same original spectrum at different scale. Now, imagine that the sample version is put through the low pass filter so that; we select only this portion of the spectrum. That means, we lift this spectrum out of assembly of the spectra section of the spectra and consider only this portion and avoid all others by putting this samples versions through a low pass filter.

Then the spectrum of this which is the resulting the output of the low pass filter is the same thing as the spectrum except for a scale factor. In other words the low pass filter output would be a replica of  $f$  of  $t$  reduced by a certain scale factor. So, it tells us that, if we have band limited function  $f$  of  $t$  if you sample it in this manner. This sample, from this sample you can extract the original function  $f$  of  $t$ .

Even though we have missed out on the variation in these intervals in between, does not really matter as long as this is band limited function from this sampled version you can recover the original signal. And this is very important in communications particularly because, if you take these samples wide apart. Then in the blank periods you can send the other messages in the same communication channel.

So, communication can be employed to send several messages in succession. So, you sample 1 particular message and then send it at regular intervals. And in the blank periods you can send another messages, a second message, a third message so on and so forth. This interleaving of these messages over the same communication channel is referred to as the time division time multiplexing. And this is very common occurrence to affect utilize the communication channel to send several messages which are inherently band in character.

Now, you can see in order for us to be able to lift this section of the spectrum and remove the others you must have the low pass filter characteristic something like this. So, if the cutoff of the low pass filter is  $\omega_c$ . Then  $\omega_c$  must be larger than



$\omega_h$ , but at the same time you do not want to pick up any portion of this spectrum. So, if this is  $\omega_c$  you have  $\omega_c$  this must be  $\omega_c \leq \omega_c - \omega_h$  and  $\omega_c \geq \omega_h$ .

So, you must be able to choose a  $\omega_c$  which lies in between  $\omega_h$  and  $\omega_c - \omega_h$ . Which in other words means  $\omega_c$  is the cut off frequency must be less than  $\omega_c - \omega_h$  and must be greater than  $\omega_h$ . And this is possible only if  $\omega_c - \omega_h > 2\omega_h$ . So, only if  $\omega_c - \omega_h > 2\omega_h$  you will be able to switch in on a suitable cut off frequency of the low pass filter.

Then  $\omega_c - \omega_h$  must be larger than  $\omega_h$ . In other words what it says is your sampling frequency  $\omega_c$ , must be more than twice the highest frequency component that is present in your original signal;  $\omega_c$  must be larger than  $\omega_h$ . If it so, then you can choose a low pass filter and allow this sample version to go through it and out comes a version a continuous function of time which is a faithful replica of your original  $f(t)$ .

I have not gone to the mathematics; just i am trying to explain how it works out to complete the mathematics. Let us see, what it means is to get the sample version from  $f(t)$  all you have to do is multiply  $f(t)$  by a pulse strain. So, if you have a pulse strain of duration  $d$  we will call this  $p_s(t)$ . Let us say for the sake of generality an amplitude  $a$  duration  $d$  and the period  $p$ . So, if you have such a pulse strain you multiply  $f(t)$  by this pulse strain you get  $f^*(t)$ .

So,  $f^*(t)$  is  $p_s(t) \times f(t)$ . Then we can find out the Fourier spectrum of the  $f^*(t)$  we will go to the mathematics now and show that this is indeed the type of spectrum that you get for  $f^*(t)$ .

Let us see.  $p_s(t)$  is the periodic pulse strain therefore, it can be extended by means of Fourier series and we have done this type of expansions several times in the past. So, let me straight away put down  $A_d \times \frac{1}{p} \sum_{n=-\infty}^{\infty} \sin(n\omega_c \frac{d}{2}) \frac{1}{2} \frac{1}{\omega_c \frac{d}{2}}$  is the  $c_n$  coefficient multiplied by  $e^{jn\omega_c t}$ . That is the time variation  $n$  ranging from minus infinity to plus infinity is your Fourier series expansion for  $p_s(t)$  the periodic pulse strain.

Now,  $f^*(t)$  is  $\sum_n \sin(n\omega_0 t)$  multiplied by  $f(t)$ . So, which means that it can be written in this fashion as the sum of several terms  $\sin(n\omega_0 t)$  by  $f(t)$ . So,  $f^*(t)$  can be thought of as a infinite sum of several terms and this is the portion which depends on the time the rest of it is purely a coefficient.

Now, to find out the Fourier transform of this we can find out the Fourier transform of each individual term and then sum them up. And since, this is a purely a constant as far as time is concerned. So, we can write this as  $\int_{-\infty}^{\infty} \sin(n\omega_0 t) f(t) e^{-j\omega t} dt$ . If the Fourier transform of  $f(t)$  is  $F(j\omega)$  we have seen that the Fourier transform of  $f(t) e^{-jn\omega_0 t}$  is simply  $F(j\omega - n\omega_0)$ . That is multiplication by exponential in the time domain; is equivalent to translation in the frequency domain this is something which we have done.

Now, let us see what we have got  $f^*(j\omega)$  the fourier spectrum of the sample version consists of several terms like this:  $F(j\omega - n\omega_0)$  when  $n = 0$  it is  $F(j\omega)$ ; when  $n = 1$  it is  $F(j\omega - \omega_0)$ . And for each value of  $n$  you have a particular coefficient. So, if you take  $n = 0$  the portion of this corresponding to  $n = 0$  will be  $\int_{-\infty}^{\infty} f(t) dt$  this is equal to  $\int_{-\infty}^{\infty} f(t) dt$  when,  $\omega = 0$  is  $F(j\omega)$ . If  $n = 1$ , this becomes  $\int_{-\infty}^{\infty} f(t) e^{-jn\omega_0 t} dt$  this is  $\int_{-\infty}^{\infty} f(t) e^{-j\omega_0 t} dt$  not all the ways this is  $\int_{-\infty}^{\infty} f(t) e^{-jn\omega_0 t} dt$  by  $\omega_0$ .

So, this portion which corresponds to  $n = 0$  is  $\int_{-\infty}^{\infty} f(t) dt$  times  $F(j\omega)$ . And this is something else  $\int_{-\infty}^{\infty} f(t) dt$  times this quantity corresponds to this.

$$F^*(j\omega) = \sum_{n=-\infty}^{\infty} \frac{Ad}{T_0} \frac{\sin n\omega_0 d/2}{n\omega_0 d/2} F(j\omega)$$

$n=0 \rightarrow \frac{Ad}{T_0} F(j\omega)$

$n=1 \rightarrow \frac{Ad}{T_0} \frac{\sin \omega_0 d/2}{\omega_0 d/2} F(j\omega - \omega_0)$

IIT  
MADRAS

So, we have seen now corresponding the central portion of the spectrum is ad up on t not times f of j omega that is it is reduced version of this. So, the implementation of this is straight forward you have f of t you can have multiplier p s of t. And this is your f star of t you put this a low pass filter and you get f not which is a replica of your original f of t.

This is the general what is called the sampling theorem which referred to a literature of sampling theorem. That is, band limited signal can be recovered from it is sampled version, when the samples are taken at regular intervals with a frequency f not which should be larger than, 2 times the highest frequency component present here.

In practice, because of the limitations of the low pass filter that we have the design of the low pass filters. In practice, you may need to have a sampling rate employ sampling rate which is higher than 2 omega h. Even though in theory any sampling rate greater than 2 omega h will suffice. But, in practice we have to take 3 or 4 times the highest frequency component because the limitation of your design of the low pass filters. To conclude our discussion of the Fourier transform methods let me offer a few remarks at this stage.

Ideally, the Fourier transform is obtained by integrating the f of t over a infinite interval of time from minus infinity to plus infinity. In practice, when you have a signal, it is normally of finite duration and therefore, the integration is to be carried over for the duration of the signal. And this poses no problem. But, on the other hand

if you have a type of signal which is continuous infinitely long for infinitely long signal and you like to find out its Fourier transform in practice. You may not be able to have an analytical description of the signal.

Then such cases what we can do is we can statistically process this information that is, if you take a finite length of the signal like for example, a peak signal or so. Take a finite length and find out its frequency characteristic and where the energy rises on which portion of the spectrum and so on. And it turns out that the statistical processing of the signal is quite useful; that means, even if you take other samples of the signals the essential character of the spectrum does not change.

Therefore, this way it is possible to handle signals even though they are need not be known for an infinitely long time. A second point which I would like to emphasize once again is, that the Fourier transform method enables us to find out the transient behavior of network and system, essentially, using steady state method which we have emphasized once again earlier when we are working out the problems. That is what we are trying to do is we are considering the transient excitation as the sum of infinite number of tiny sinusoids and we find the response to each sinusoid using essentially steady state methods frequency response methods.

Then add all the responses together to find out the total response. When I say summation it is actually the integration that is being done because, these are all a continuous band of frequency that you are having. But, the principle is the same summation is essentially an integration; integration is essentially the summation. Now, it is this fact that transient behavior of networks and systems can be analyzed on the basis of steady state sinusoidal theory. That places the Fourier transform method and Fourier theory and important gives it an important place in communication engineering.

In communication engineering essentially, we have to deal with transient signals and transient processing of transient signals. But, the communication network can be characterized in terms of their steady state behavior under sinusoid excitation condition. And using that information we are able to process the behavior of the network under transient signals.

And some people assert that this is 1 of the most important idea in communication engineering is the way, in which we tackle this transient signals by the Fourier transform method.

Now, we also seen that the Fourier transform methods we have taken networks and found out the transient in the network using the Fourier transform methods. And we mentioned that we have certain restrictions earlier; not only, on the type of excitation signals. Which can be handled by the Fourier transform method and we also said we have placed some restriction on the network also the natural response of the network should die down with time for the Fourier transform method to work in a convenient fashion.

Now, there are several such restrictions when we apply the Fourier transform of method to the transient solution network and systems. A competitor to this is the Laplace method which we will take up later. For the Laplace transform method offers greater conveniences for the analytical solutions of the transient in network and system. Mainly because, it enlarge the class of time functions it can be handled; that means, certain types of functions like  $e^{-2t}$ .

For example, cannot be handled by the Fourier transform method, but Laplace transform method can handle this. Secondly, the several delta functions which may get in the Fourier transforms are avoided when use the Laplace transform method because of these reasons Laplace transform is a more convenient tool for analytical evaluation of transient and linear systems.

The same time where the networks are characterized in terms of their frequency behavior; in terms of their frequency response conditions. Like communication network, like filters for example, the Fourier transform method is very convenient tool. Because, these frequency response characterize of systems and networks can experimentally determine  $h(j\omega)$  can be experimentally determine.

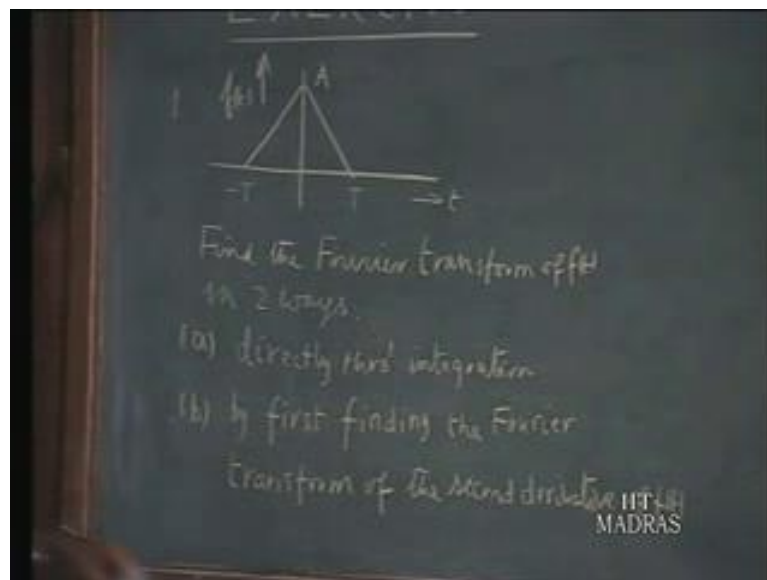
$H(s)$  may not be so easy experimentally determine. So, where the networks are characterized by the frequency response behavior then Fourier transform method trans out to be very convenient tool to employ. For example, the behavior of the low pass filter that we have discussed little earlier.

Also, in terms in certain branches like a control engineering and so on it is the frequency response of a system that we are interested in. To assess the stability of the control system sometime we like to know the frequency response characteristic of this. So, in such cases of course, the Fourier transform method and handling of the various components to the frequency response approach turns out to be very convenient.

They are also well defined methods where once you know the frequency response you can calculate the transient response graphically. So, because of all these reasons Fourier transform method plays the vital role in analysis of network and systems. But as i mentioned for the transient analysis of network perhaps the Laplace transform methods is the more useful tool and that will take up later.

But nevertheless as i mentioned wherever the networks or systems are characterized by the frequency response Fourier transform method turns out to be more convenient tool.

(Refer Slide Time: 36:04)

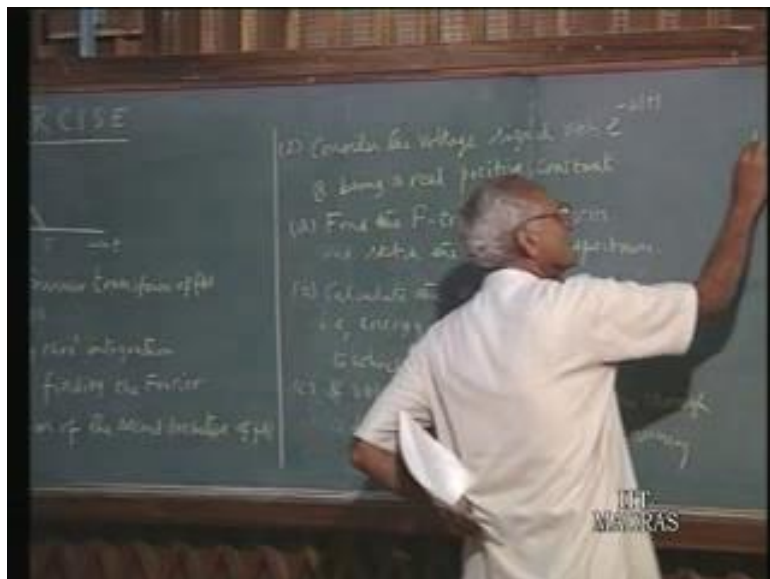


And this stage i will give you set of problems as an exercise for you to work out the Fourier analysis transform methods. Now, let me write down a set of problems for you as an exercise under the topic Fourier transform and the Fourier integral. Let us consider,  $f$  of  $t$  which has got this type of variation  $t$  minus  $t$  this is a this is  $f$  of  $t$ .

Find the Fourier transform of  $f$  of  $t$  in 2 ways; a: directly through integration applying the standard formula for the fourier transform. That is, minus infinity to plus infinity  $\int_{-\infty}^{\infty} f(t) e^{-j\omega t} dt$  that type of integration.

b: By first finding the fourier transform of the second derivative of  $f$  of  $t$ . The motivation for this is if you take the derivative of this you get a pulse here and another pulse here and if you take the second derivative you have impulses 3 impulses. To find out the Fourier transform of the impulses quite easy. So, you find the Fourier transform of the second derivative of  $f$  of  $t$  and from that deduce the Fourier transform of the original function  $f$  of  $t$ .

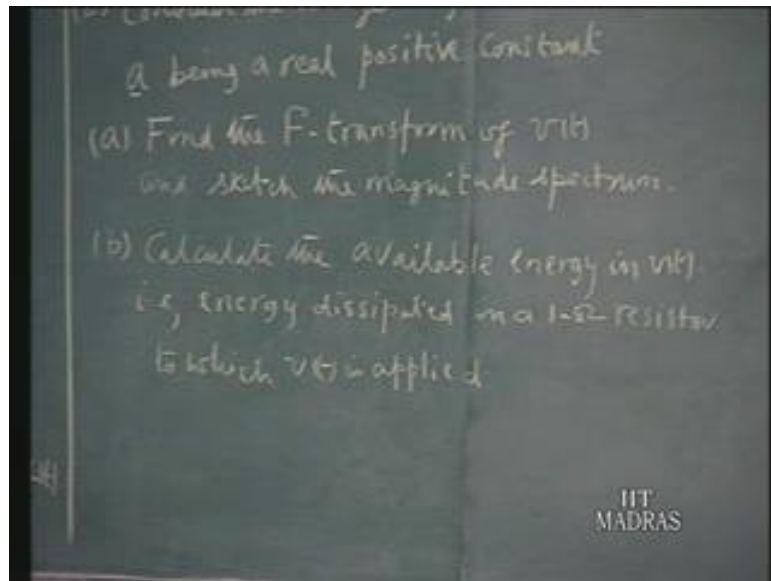
(Refer Slide Time: 38:20).



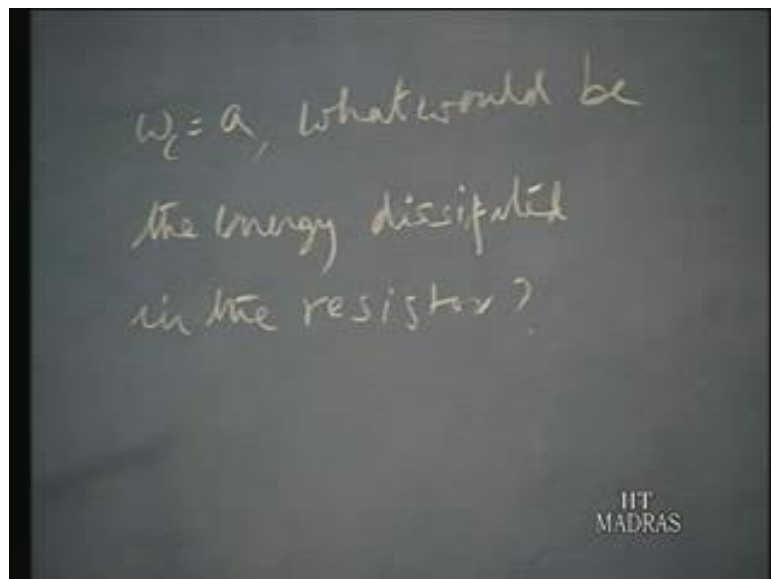
Problem number 2: consider the voltage signal  $v$  of  $t$   $v(t) = e^{-at} u(t)$  where  $a$  is being a real positive constant; a: find the Fourier transform of  $v$  of  $t$  this is voltage signal. Let me say, this is  $v$  of  $t$  find the Fourier transform of  $v$  of  $t$  and sketch the magnitude spectrum.

b: calculate the available energy of  $v$  of  $t$  the signal  $v$  of  $t$ . What you mean by the available energy? If this  $v$  of  $t$  was applied to a 1 ohm resistor what is the energy dissipated in the resistor  $v$  of  $t$  in the resistor  $r$ . That is what is meant by available energy  $v$  of  $t$  that is, energy dissipated in 1 ohm resistor to which  $v$  of  $t$  is applied.

(Refer Slide Time: 41:10)



(Refer Slide Time: 41:53).

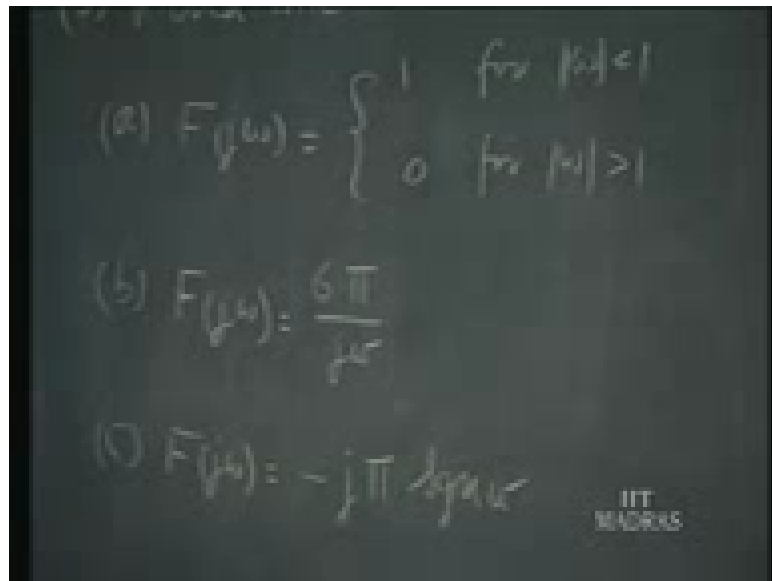


c: if  $v$  of  $t$  were applied to a  $1$  ohm resistor through a low pass filter, with cut off frequency  $\omega_c$  is being equal to  $a$ . What would be the energy dissipated in the resistor?

That is you have a low pass filter which cuts off all frequencies beyond  $\omega_c$   $\omega_c$  greater than  $a$  then, what is the energy dissipated in the resistor.

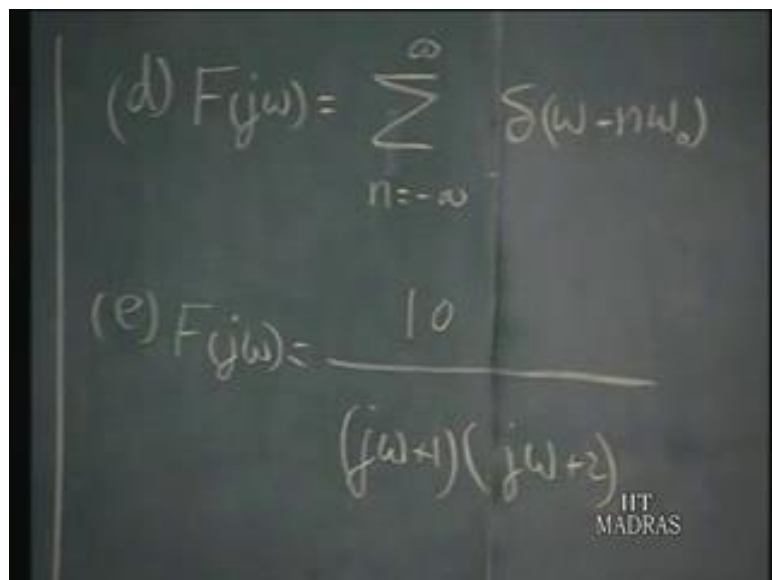


(Refer Slide Time: 42:40).


$$(a) F(j\omega) = \begin{cases} 1 & \text{for } |\omega| < 1 \\ 0 & \text{for } |\omega| > 1 \end{cases}$$
$$(b) F(j\omega) = \frac{6\pi}{\omega}$$
$$(c) F(j\omega) = -j\pi \text{sgn}(\omega)$$

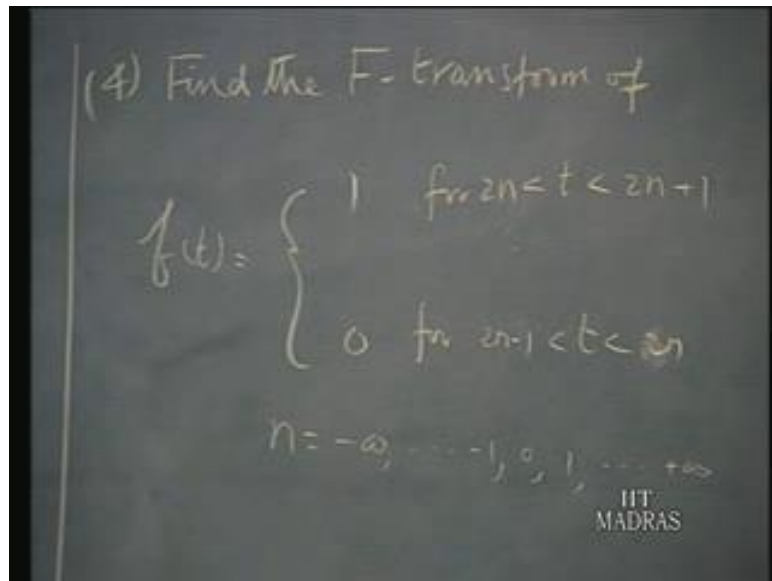
Find the inverse Fourier transforms of the following functions; a:  $f$  of  $j$   $\omega$  equals 1 for  $\omega$  magnitude less than 1 0 magnitude greater than 1, b:  $f$  of  $j$   $\omega$  equals  $6\pi$  up on  $j$   $\omega$ , c:  $f$  of  $j$   $\omega$  equals minus  $j$   $\pi$  signum function of  $\omega$ .

(Refer Slide Time: 43:45).


$$(d) F(j\omega) = \sum_{n=-\infty}^{\infty} \delta(\omega - n\omega_0)$$
$$(e) F(j\omega) = \frac{10}{(j\omega+1)(j\omega+2)}$$

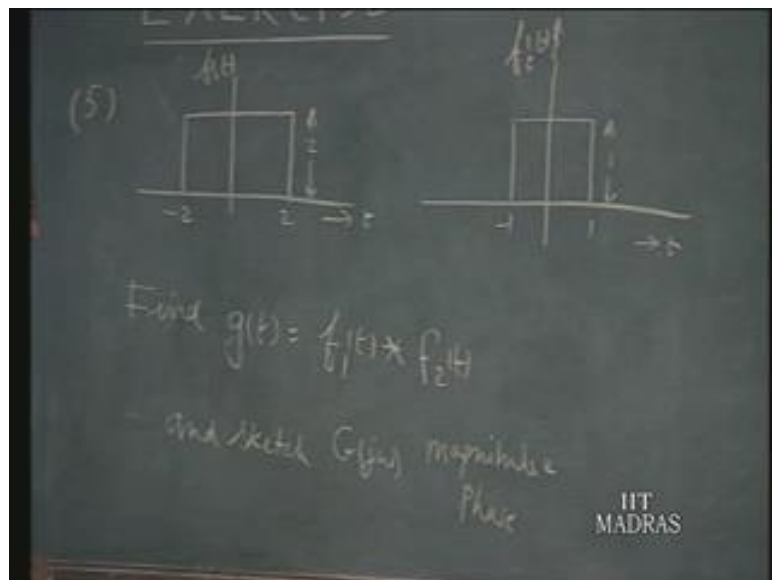
d:  $f$  of  $j$   $\omega$  equals a set of impulse functions periodically occurring intervals of  $\omega$  not, e:  $f$  of  $j$   $\omega$  equals ten up on  $j$   $\omega$  plus 1 times  $j$   $\omega$  plus 2. That is the third problem.

(Refer Slide Time: 44:30).



Fourth problem find the Fourier transform of a function  $f$  of  $t$  defined as 1 for  $t$  between  $2n$  and  $2n+1$  and 0 for  $t$  between  $2n-1$  and  $2n$ . Where,  $n$  ranges from minus infinity minus 1 0 1 etcetera to plus infinity for all integral values of  $n$  positive and negative. That is the fourth problem.

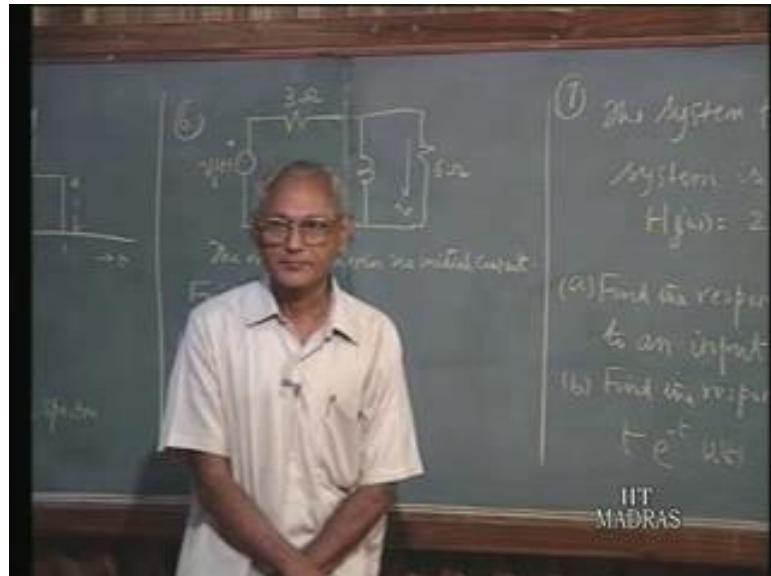
(Refer Slide Time: 45:26)



Let us consider 2 functions  $f_1(t)$  a pulse of amplitude 2 lasting from minus 2 seconds to plus 2 seconds, and another  $f_2(t)$  lasting from minus 1 to 1 second of height 1 unit. Find  $g(t)$  which is obtained by the convolution of  $f_1(t)$  and  $f_2(t)$  and the sketch  $g(t)$

omega which is the Fourier transform of  $g$  of  $t$  sketch  $g$  of  $t$ ; that means, magnitude and phase.

(Refer Slide Time: 47:04).



We will change this  $g$  of  $j\omega$  and sketch the spectra. Find  $g$  of  $j\omega$  the spectra of  $g$  of  $\omega$ . Sixth we have circuit in which there is a voltage source  $v_i$  of  $t$  impressed at cross a circuit like this. This 1 henry inductor 6 ohm resistor and this current is  $i$  not. The inductor carries 0 initial current no initial current that is before the excitation is applied the inductor is carrying no current.

We are asked to find the current  $i$  not of  $t$  for the following excitation  $v_i$  of  $t$  equals  $9\delta(t)$ , b:  $v_i$  of  $t$  equals  $9u(t)$ , c:  $v_i$  of  $t$  equals  $9e^{-t}u(t)$ . So, this is the 6 problem this is the transient analysis of this problem. 7 the system function of a linear system is  $h$  of  $j\omega$  equals  $2 + \frac{1}{j\pi\omega}$ . Find the response of the system to an input  $\delta(t)$  that is impulse response of the system is asked to found out, b: find the response to a input which is equal to  $t$  multiplied by  $e^{-t}u(t)$ .

So, in both these cases use make use the system function find out the input fourier transform multiply them out you get the output fourier transform. Find the inverse Fourier transform you get the response we have not found out we have not discuss the Fourier transform of functions this type. But, we should able to derive them using the

from fundamentals  $t e$  to the power of minus  $t u t$  you have a Fourier transform which is,  $1 \text{ over } j \text{ omega plus } 1 \text{ whole squared}$  that we should be able to derive.