

Analysis and Design Principles of Microwave Antennas
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Lecture - 27
Paraboloid Reflector Antenna (Contd.)

Welcome to this lecture in NPTEL course, we are continuing the discussion on reflect Paraboloid Reflector system. We have seen the aperture field distribution from the feeds radiation pattern, now we will see the directivity of the this whole system; that means, aperture directivity.

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APERTURE DIRECTIVITY

$$\vec{E}(r) = j k_0 \frac{e^{-jk_0 r}}{2\pi r} \left[\hat{a}_x (f_x \cos \phi + f_y \sin \phi) + \hat{a}_y \cos \theta (f_y \cos \phi - f_x \sin \phi) \right]$$

$$\vec{E}_a = E_0 \hat{a}_y \quad ; \quad x^2 + y^2 \leq a^2$$

$$= 0 \quad ; \quad \text{otherwise}$$

$$U_\theta = E_0 \hat{a}_y \iint_{S_a} e^{jk_x x} e^{jk_y y} dx dy$$

ρ, ϕ

So, we know actually in few lectures back we have proved by Fourier transform method that for a uniform constant phase aperture field with linear polarisation, the radiated field is given by the expression $E_r j k \cos \theta e^{-jk r} / 2\pi r$; a $\theta \cos \phi + f_y \sin \phi$ plus a $\phi \cos \theta f_y \cos \phi - f_x \sin \phi$, this we have already proved.

Now, our aperture is circular with radius a in z is equal to 0 plane. So, we can write our aperture field is E_y linearly polarised uniform electric field, we are assuming it. Now, we have seen that it is it has got problems, but first let us see the directivity; that means, what maximum we can achieve, if we can create this uniform field let us arbitrarily I have taken y polarised field in any polarised single polarisation. So, this is

valid for $x^2 + y^2 \leq a^2$ is the circle, circular disc and 0 otherwise.

So, with this aim what will be my Fourier transform f_t , f_t will be $E_0 \hat{a}_y$ sorry $E_0 \hat{a}_y$ that a y then on the aperture e to the power so, this is 1, 1 into e to the power $j k_x x + j k_y y$ $dx dy$. Now, to evaluate this integral this integral evaluation is not so easy. So, I will help you. So, since this is a circular disc we introduce the polar coordinates, but since it is the aperture is here source we say the polar coordinate is ρ instead of ϕ I call it ϕ' . I should have call the ρ dash also, but it would not get confused because the other feeds thing is θ . So, since it does not have any θ so, there is no confusion. But, since there is also ϕ so, the observation field ϕ and source field ϕ' that is why I am showing ϕ' .

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$$\rho = \sqrt{x^2 + y^2}$$

$$x = \rho \cos \phi'$$

$$y = \rho \sin \phi'$$

$$k_x = k_0 \sin \theta \cos \phi$$

$$k_y = k_0 \sin \theta \sin \phi$$

$$f_t = E_0 \hat{a}_y \int_0^a \int_0^{2\pi} e^{j k_0 \rho \sin \theta \cos(\phi - \phi')} \rho d\phi' d\rho$$

$$e^{j \omega \cos(\phi - \phi')} = J_0(\omega) - 2 [J_2(\omega) \cos 2(\phi - \phi') - J_4(\omega) \cos 4(\phi - \phi') + \dots]$$

$$+ 2j [J_1(\omega) \cos(\phi - \phi') - J_3(\omega) \cos 3(\phi - \phi') + \dots]$$

$J_n(\omega) \rightarrow$ Bessel funcn 1st kind of order n .

Now, what is my ρ ? ρ is nothing, but root over $x^2 + y^2$. So, x is equal to $\rho \cos \phi'$, y is $\rho \sin \phi'$ and also we know k_x is $k_0 \sin \theta \cos \phi$ and k_y is $k_0 \sin \theta \sin \phi$. So, I can put it on that f_t expression so, it becomes $E_0 \hat{a}_y$, then 0 to a 0 to 2π e to the power $j k_0 \rho \sin \theta \cos(\phi - \phi')$ $\rho d\phi' d\rho$ ok. Now, this thing e to the power j something \cos that can be, that has a series expansion, anyone knows that this expansion is in terms of Bessel functions. So, we can write if I have $e^{j \text{something} \cos(\phi - \phi')}$, the Bessel expansion is J_0 that ω you can say or ω whatever this is a

constant; that means, j into these.

So, J_2 ω minus $2 J_2 \omega \cos 2 \phi$ minus $J_4 \omega \cos 4 \phi$ minus ϕ dashed etcetera etcetera plus $2 j J_1 \omega \cos \phi$ minus ϕ dash minus $J_3 \omega \cos 3 \phi$ minus ϕ dash plus etcetera. So, big but so, what is the, if I write $J_n \omega$ for various subscript. So, this is the Bessel function which comes in all engineering science problems. So, Bessel function this is J means it is the first kind and n means of order n ok. Now, the if I put it here fortunately the integration of this $\cos 2 \phi$ minus ϕ dash where, ϕ dash this varying from 0 to 2π $\cos 4$. So, everything from 0 to 2π integration that becomes 0 ; only the thing is who does not become 0 is $J_0 \omega$.

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$$f_t = E_0 \hat{a}_y \int_0^a 2\pi J_0(k_0 \rho s \phi) \rho d\rho$$

$$\int_0^z u^{2n} J_{2n-1}(u) du = z^{2n} J_{2n}(z)$$

$$= 2\pi E_0 \hat{a}_y \int_0^{k_0 a s \phi} J_0(k_0 \rho s \phi) \frac{k_0 \rho s \phi}{k_0 s \phi} \frac{d(k_0 s \phi \rho)}{k_0 s \phi}$$

$$= 2\pi a^2 E_0 \hat{a}_y \frac{J_1(k_0 a s \phi)}{(k_0 a s \phi)} \cdot \frac{1}{2}$$

$$= \pi a^2 E_0 \hat{a}_y$$

So, f_t comes out as $E_0 \hat{a}_y$ 0 to a $2 \pi J_0$ naught, we will see that $k_0 \rho \sin \theta \rho d\rho$. So, it is an integration of Bessel function of 0 order. So, there the there is another formula which will help you. So, with this formula you can find out that this will become $2 \pi E_0 \hat{a}_y$ 0 to $k_0 a \sin \theta J_1$ naught $k_0 \rho \sin \theta$ $k_0 \rho \sin \theta$ by $k_0 \rho \sin \theta$ d of $k_0 \rho \sin \theta$ sorry $k_0 \rho \sin \theta \rho$ by $k_0 \rho \sin \theta$.

So, if you do that it finally, becomes $2 \pi a^2 E_0 \hat{a}_y J_1$ $k_0 a \sin \theta$ by $k_0 a \sin \theta$. Now, on the z axis on the z axis θ is 0 so, $k_0 \rho \sin \theta$ fully 0 . So, this J_1 naught 0 by 0 this term will give you J_1 naught 0 by 0 . Now so, that is a (Refer Time: 11:09) rule so, we will have to differentiate these and then put the limit. So,

this part gives you half so, it becomes pi a square E naught a y.

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$\theta = 0, \phi = \frac{\pi}{2}$
 $f_y = \pi a^2 E_0$
 $f_x = 0$
 $\rightarrow E_y(r) = j k_0 \frac{e^{-jk_0 r}}{2\pi r} \theta f_y \hat{a}_\theta$
 $= \frac{j k_0 E_0}{2\pi r} e^{-jk_0 r} \pi a^2 \hat{a}_\theta$
 Radiate Prop intensity along z axis
 $\pi^2 \frac{1}{2} Y_0 |E_0|^2 = \frac{1}{8} k_0^2 E_0^2 Y_0 a^4$
 $P_r = \iint_{S_a} (\text{Poy. vect}) dA$

On the z axis the e plane pattern is theta is equal to 0 is equal to pi by 2. So, we get f y will be pi a square E naught and f x is equal to 0. So, we have f y f x I have got so, we can say that E radiated field that will be j k naught to the power minus j k naught r by 2 pi r, f y you can put or I will write f y a theta. So, that will be j k naught E naught by 2 pi r e to the power minus j k naught r pi a square a theta. So, this will be the power field.

Now, the power density per unit solid angular radiation intensity in this direction will be r square half Y naught E theta square. So, I will say power radiation intensity let us say radiation intensity along z axis that will be this. So, if you do it will become 1 by 8 k naught square E naught square Y naught a to the power 4. Another thing we need the total radiated power P r so, for that will be over the aperture, the pointing vector into d A.

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Pointing vector over aperture $\approx \frac{1}{2} Y_0 E_0^2$

$$P_r = \frac{1}{2} Y_0 E_0^2 \int_0^{2\pi} \int_0^a P \, d\rho \, d\phi$$

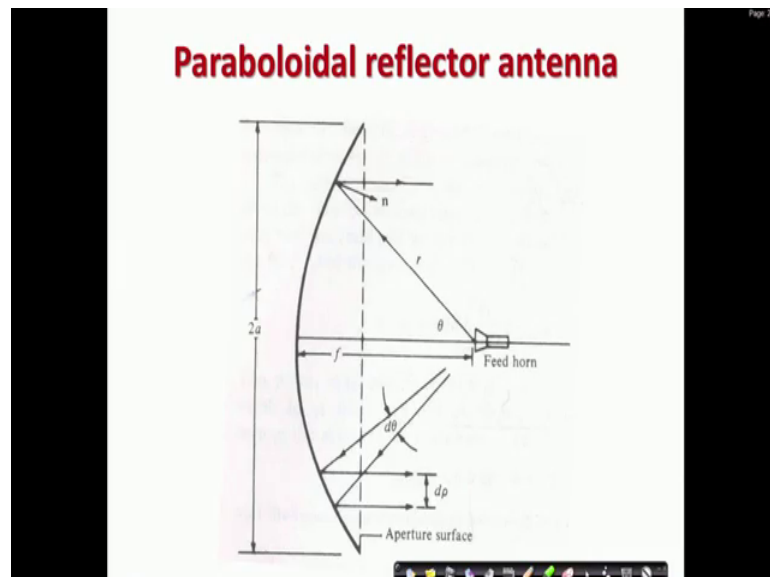
$$= \frac{1}{2} \pi a^2 Y_0 E_0^2$$

$$D = 4\pi \frac{\frac{1}{4} k_0^2 E_0^2 Y_0 a^2}{\frac{1}{2} \pi a^2 Y_0 E_0^2}$$

$$D = \frac{4\pi}{\lambda_0^2} (\pi a^2)$$

Now, pointing vector over aperture is what? Pointing vector over aperture that actually here we can say that though aperture is not very far away, but we are having the already a parallel beams. And, it has been it can be proved that they are the magnetic field is also like the per field; that means, that n cross that thing.

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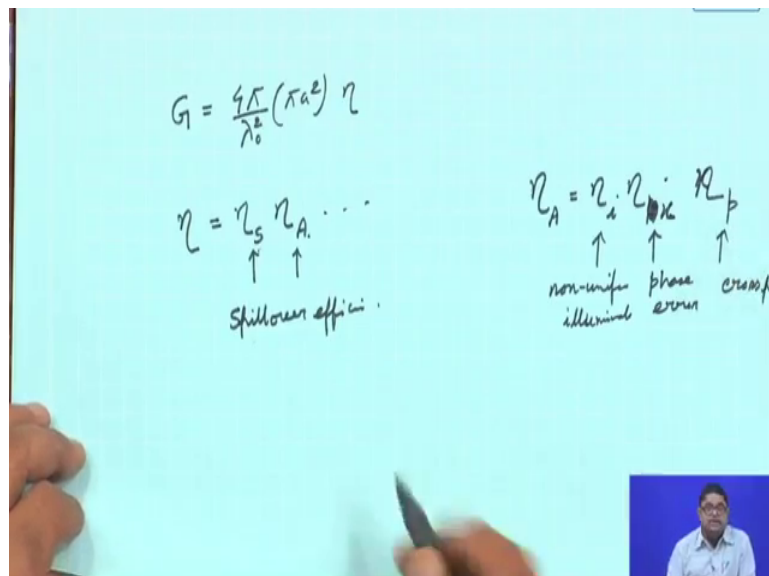
So, we can almost take it that it will be half $Y_0 E_0^2$; so, just like far field it becomes there. So, if this is there the total radiated power P_r will be half $Y_0 E_0^2$ then the area $\int_0^{2\pi} \int_0^a \rho \, d\rho \, d\phi$. If you do it

becomes half pi a square Y naught E 0 square. So, now you are in a position to do directivity 4π into that radiation intensity 1 by $8k$ naught square E naught square Y naught a to the power 4 by total radiated power half pi a square Y naught E naught square. These gives you and I will write it like this, directivity is 4π by lambda naught square such a simple formula, such a useful formula.

Now, what is it pi a square means what? The aperture physical area. So, aperture physical area is getting multiplied by 4π by lambda naught square. So, if you have a dish 1 meter radius then pi. So, 4π square gives you 40 roughly and if you are at 10 gigahertz 3 centimetre means 0.03 into 0.03 so, 0.009. So, I can say 0.01 so that means, 4000 is the directivity; 1 meter radius can give you directivity of 4000, quite remarkable.

But obviously, we cannot when this is directivity; that means, there are no losses. Now, actually there will be losses etcetera. So, the maximum directivity obtainable from an aperture antenna with a constant phase uniform field is physical area multiplied by 4π by lambda naught square; very simple formula easy to remember. But, what will be the gain of the parabolic antenna reflector antenna system that will be G is equal to 4π by lambda naught square into pi a square into efficiencies.

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Now, these efficiencies the first one actually we will see a long list of efficiencies at the end. Initially I am writing eta, that eta to start with I make that there is a two things: eta S eta A later others we will come so, I will say these. Now, this I will see first this then,

this, then others. Now, this η_s this is called spillover efficiency, η_a is aperture efficiency. Now, aperture efficiency I need to talk a bit because whatever we have done that if I have an uniform illumination, if I have a constant phase illumination, if I have a single polarisation aperture field then I get the directivity I have found.

Now, aperture efficiency is violation of all these. So, aperture efficiency will comprise of illumination efficiency because, we have seen that there is a taper and it is desirable from the application point of view. So, it will be multiplied by η_i then it will be that phase. So, uniform phase we have we know that there will be a phase error. So, that is phase and I think that third parties now sorry this is polarisation. So, the next one non-uniform thing let me call it taper η_x .

So, this is due to this is due to non-uniform illumination, this is due to phase error, due to phase non-uniform phase and then η_p it is the cross pole presence of non-linear thing. So, spillover efficiency what is it? Spillover efficiency is that some radiation from the feed is not intercepted by the reflector because, reflector is finite it has an angular aperture. So, it is not intercepting all the rays. So, power is lost there and definitely then we are assuming total power radiated, but if that power is not reflected so, ultimately aperture is not getting that power. So, gain will definitely reduce so, that is called spillover loss and efficiency factor due to that is called spillover efficiency.

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The image shows a handwritten derivation on a light blue background. At the top, it is titled "Spillover Efficiency". The main equation is:

$$\eta_s = \frac{\text{Power intercepted by the reflector}}{\text{Total Power radiated by the feed}}$$

$$= \frac{\int_0^{2\pi} \int_0^{\psi/2} g(\alpha, \phi) \sin \alpha \, d\alpha \, d\phi}{\int_0^{2\pi} \int_0^{\pi} g(\alpha, \phi) \sin \alpha \, d\alpha \, d\phi}$$

Below this, there is another equation for the Directivity of feed, D_f :

$$D_f = \frac{4\pi g(0,0)}{\text{Total radiated power by feed}}$$

There is a small logo in the top right corner that says "L.T.KGP".

So, what is spillover efficiency now? So, I am now talking of spillover efficiency, η_s is

power intercepted by the reflector by total power radiated by the feed. So, what is it? Power intercepted means I have $g(\theta, \phi)$ is the power radiation pattern, $\sin \theta d\theta d\phi$. And so, ϕ 0 to 2π and θ is from 0 to ψ by $\frac{\psi}{2}$ half of the angular aperture because, 0 to 2π I am giving. So, here 0 to ψ by $\frac{\psi}{2}$ this is the power intercepted by the reflector and here 0 to 2π 0 to π $g(\theta, \phi) \sin \theta d\theta d\phi$ that is all, that is spillover efficiency ok.

Now, ultimately we will see that I am not interested to find simply this spillover efficiency or it is we will later see that jointly what happens here. Actually, the others can be separately done, but these two are as physically I explained that if I want to have an uniform illumination then spillover increases, that I said earlier that the of best possible illumination was that sec to the power $4 \sin^2 \theta$, but that gives lot of a. So, these two are coupled and one effects the other. So, that is why we will jointly maximise this instead of trying to do that. Now, to that in instead of evaluating it further what I do, I want to put here the directivity of the feed; that means, this is radiation pattern here I want to put the directivity of the feed so, D_f .

Now, what is directivity of the feed? D_f let me write directivity of feed D_f . What will be its expression? $4\pi g_{00}$ by total radiated power by feed. So, you see that this is nothing, but total radiated power by feed this denominator. So, can I write that η_s is D_f then 0 to 2π 0 to ψ by $\frac{\psi}{2}$ $g(\theta, \phi)$ by actually, this will come out $4\pi g_{00} \sin \theta d\theta d\phi$ ok, I leave it here ok.

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$$\eta_s = D_f \int_0^{2\pi} \int_0^{\psi/2} \frac{g(\theta, \phi)}{4\pi g(\theta, 0)} \sin\theta d\theta d\phi$$

Aperture Efficiency (η_A).

loss in gain due to

- tapered electric field in the aperture
- nonconstant phase for aperture field
- presence of unwanted cross polarized field.

$$\eta_A = \eta_i \eta_p \eta_x$$

illumination eff. phase error cross polarized eff.

Now, next I start I will use that as I said that since I will jointly do. So, I will take this after see finding an expression for aperture efficiency η_A . So, who gives this aperture efficiency? The loss in gain due to first I have a tapered electric field in the aperture, then non-constant phase for aperture field and then presence of unwanted cross polarised field ok. So, actually this expression is available or it can be derived, but I would not do that because that will be more mathematical in nature. So, I can write $\eta_A = \eta_i \eta_p \eta_x$ this is called illumination efficiency, then non-constant phase is phase error efficiency and cross polarised thing is η_x .

So, this is illumination efficiency, this is phase error efficiency, this is cross polarised efficiency. So, I am not going there are some long list of things. So, we assume while deriving the, this illumination efficiency that η_p and η_x they are one, later we will put their values. So, under that assumption it can be derived, if anyone interested we see in the may be in the course page we can upload this the derivation. But finally, what we need is that η_i the illumination efficiency that becomes an expression something like this. This is the expression that is obtained for the illumination efficiency.

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$$\eta_i = \frac{4f^2}{\pi a^2} \frac{\left| \int_0^{2\pi} \int_0^{\psi/2} [g(\theta, \phi)]^{1/2} \tan \frac{\theta}{2} d\theta d\phi \right|^2}{\int_0^{2\pi} \int_0^{\psi/2} g(\theta, \phi) \sin \theta d\theta d\phi}$$

$$g(\theta, \phi) = \sec^4 \frac{\theta}{2}$$

$$\eta_i = \frac{4f^2}{a^2} \tan^2 \frac{\psi}{4} = 1$$

$$a = 2f \tan \frac{\psi}{4} \uparrow$$

Now, if you do not believe this expression, I can give you a check. For uniform aperture field we require that g theta phi should be proportional to \sec^4 theta by 2. Put that in this expression g theta phi values and η_i should turn out to be 1, then that will be indirect way of saying this expression. So, you can do that, that if you put it there you will see this will become \sec^2 theta by 2 \tan theta by 2 d theta. This whole thing will come out because, g theta phi is independent of phi. So, the phi integration will give you 2π and here also you can integrate. So, you will see that under that assumption that g theta phi is \sec^4 theta by 2, this η_i expression becomes $4f^2$ by a^2 \tan^2 pi by 4.

Already, we have found the value of a when we derived that ρ is equal to a is $2f \tan$ pi by 4. So, put that value of a so, η_i becomes 1 putting it here it becomes 1. So, this expression is correct. So, now we will see that how to manipulate these, that we will see that this jointly if we put this two expressions η_A and η_S how to manipulate that, but that we will see in the next class. So, we have now found out the illumination efficiency, found out means this is I have given this formula and the spillover efficiency we have derived. So, now we will see how to play with that as an engineer in the next class.

Thank you.