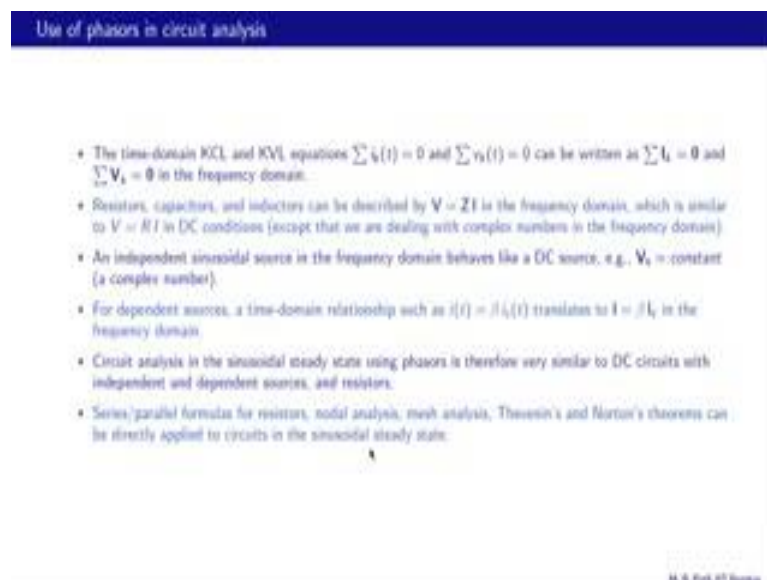


Basic Electronics
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Lecture - 06
Phasors (continued)

Welcome back to Basic Electronics. In the last lecture we have introduced phasors and seen how to interpret them in the sinusoidal steady state. In this lecture we want to apply phasors to RLC circuits, to figure out the steady state currents and voltages. In addition we will look at the maximum power transfer theorem for RLC circuits in the sinusoidal steady state. So, let us start.

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Use of phasors in circuit analysis

- The time-domain KCL and KVL equations $\sum i_k(t) = 0$ and $\sum v_k(t) = 0$ can be written as $\sum I_k = 0$ and $\sum V_k = 0$ in the frequency domain.
- Resistors, capacitors, and inductors can be described by $V = ZI$ in the frequency domain, which is similar to $V = RI$ in DC conditions (except that we are dealing with complex numbers in the frequency domain).
- An independent sinusoidal source in the frequency domain behaves like a DC source, e.g., $V_k = \text{constant}$ (a complex number).
- For dependent sources, a time-domain relationship such as $i(t) = \beta i_k(t)$ translates to $I = \beta I_k$ in the frequency domain.
- Circuit analysis in the sinusoidal steady state using phasors is therefore very similar to DC circuits with independent and dependent sources, and resistors.
- Series/parallel formulas for resistors, nodal analysis, mesh analysis, Thevenin's and Norton's theorems can be directly applied to circuits in the sinusoidal steady state.

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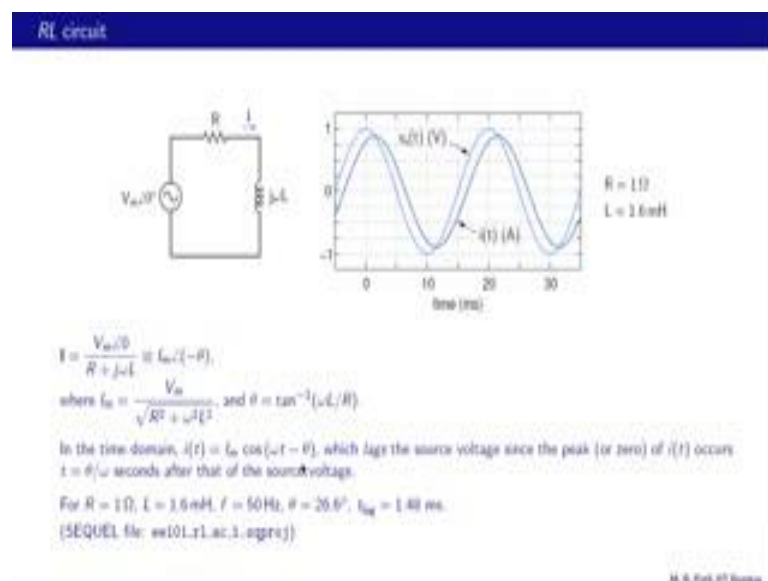
Let us now look at how we can use phasors in circuit analysis; the time domain KCL and KVL equations $\sum I_k = 0$ and $\sum V_k = 0$ can be written as phasor equations, $\sum I_k = 0$ where I_k are current phasors now and $\sum V_k = 0$ where V_k are voltage phasors in the frequency domain.

Resistors, capacitors and inductors can be described by $V = ZI$ in the frequency domain, which is similar to $V = RI$ in DC conditions. The only difference is that we are now dealing with complex numbers, when we talk about this equation. An independent sinusoidal source in the frequency domain behaves like a DC source for example, for a voltage source a sinusoidal voltage source we can say that the

phasor V_s is the constant, which is the complex number. For dependent sources a time domain relationship such as $i(t) = \beta i_c(t)$ translates to the phasor relationship, phasor $I = \beta$ times phasor I_c in the frequency domain. So, the equation looks very similar except we have complex numbers here.

So, from all of these remarks we conclude that circuit analysis in the sinusoidal steady state using Phasors, is very similar to DC circuits with independent and dependent sources and registers; and therefore, all the results that we derived for DC circuits are valid also for sinusoidal steady state analysis therefore, series parallel formulas for resistors, nodal analysis, mesh analysis, Thevenin and Norton's theorems can be directly applied to circuits in the sinusoidal steady state; totally difference of course, is we are now dealing with complex numbers.

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Let us now consider an R L circuit in the sinusoidal steady state; where R and L in series the impedance of the register is R as we have seen before and the impedance of the inductor is $j \omega L$.

This is our sinusoidal source voltage $V_m \cos \omega t$, that is $V_m \cos \omega t$ in the time domain and we are interested in this current. When we use Phasors, this becomes an extremely simple calculation as we will see; let us imagine that we have a DC circuit with a DC source here V_s , a resistor R_1 here and a register R_2 here, what would the current I in that case? The DC current it would be simply V_s divided by $R_1 + R_2$;

so we can use exactly the same equation except instead of V_s we have a Phasors source now that is $V_m \angle 0$ and instead of R_1 and R_2 they have R and $j\omega L$, the impedances of the resistor and the inductor. So, then we can write I as $V_m \angle 0$, this source divided by $R + j\omega L$ as simple as that.

We can write this phasor current as $I_m \angle \text{minus theta}$, where I_m is the magnitude of this expression and that is simply V_m divided by the magnitude of the denominator which is square root $R^2 + \omega^2 L^2$. The angle of the denominator is $\tan^{-1}(\omega L / R)$ and the angle of the numerator is 0. So, the net angle of I is $0 - \tan^{-1}(\omega L / R)$ that is all we get this minus sign over here. So, θ here is $\tan^{-1}(\omega L / R)$.

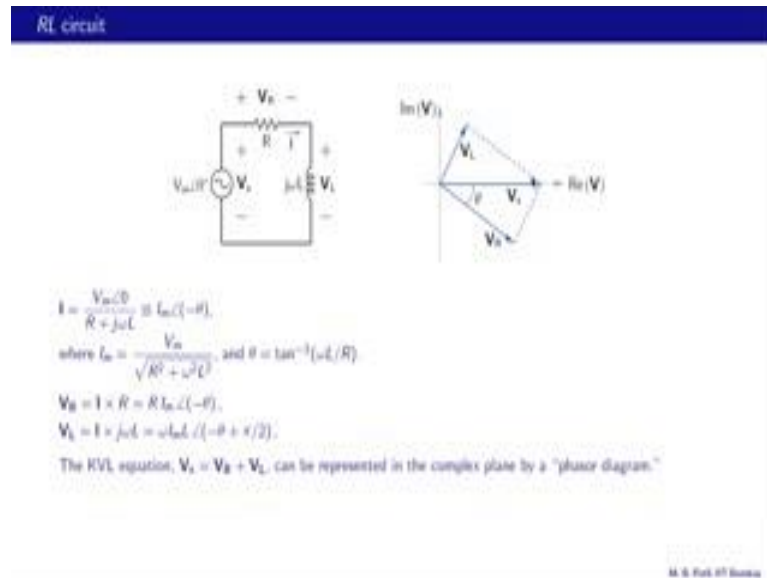
In the time domain we have $i(t)$ equal to $I_m \cos(\omega t - \theta)$, I_m is given by this expression and θ is this one. Now let us take some component values say R equal to 1 ohm, L equal to 1.6 millihenry and f equal to 50 hertz then θ turns out to be 26.6 degrees; and here are the plots for the source voltage and the current this is our source voltage V_m is 1. So, it is one times $\cos(\omega t)$, and this is our current the dark one and let us now try to understand this current plot in terms of this equation.

Our I_m value turns out to be about 0.9 ampere and therefore, we see that the current goes from minus 0.9 amperes to plus 0.9 amperes. And as we have said earlier θ is 26.6 degrees, now how do we relate this value to this plot? This is an angle and this is time over here. So, what we can do is to convert this angle to time; we know that 360 degrees corresponds to one period and that in this case is 20 milliseconds, because $1/f$ is 20 milliseconds and therefore, we can use that fact to find the time that corresponds to this angle that turns out to be 1.48 milliseconds.

Now let us get back to this equation $i(t)$ equal to $I_m \cos(\omega t - \theta)$, and ask the question when does this $i(t)$ go through is speak? That is when $\omega t - \theta$ is equal to 0 and the answer is that happens when t is equal to θ / ω , and that turns out to be exactly 1.48 milliseconds. So, what it means is that the current does not go through the peak here when the voltage goes through each speak, but a little later that is at 1.48 or nearly 1.5 milliseconds, and that is why we say that it lags the source voltage.

Notice how easy this entire calculation was, we did not write down any differential equation, we simply used this expression which is like 2 resistors in series and then we managed to get all the information that we required. So, that is how phasors really help in analyzing circuits in the sinusoidal steady state. So, here is the circuit file you can simulate the circuit and maybe change some component values and look at the results.

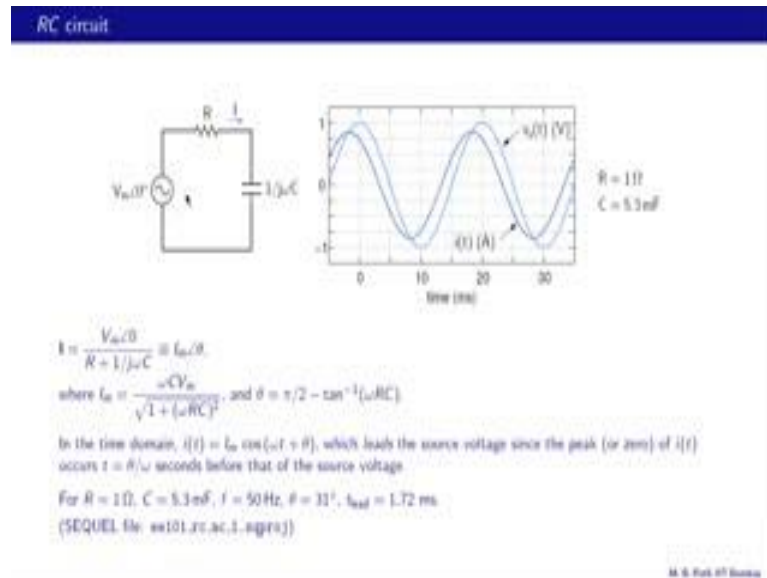
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Let us now represent the KVL equation that is V_s equal to V_R plus V_L in a graphical form in the complex plane, using a phasor diagram; and to do that we require V_s which is V_m angle 0, and V_R , V_L . Let us look at V_R first; what is V_R ? V_R is I times the impedance of the resistor which is R . So, that is R times I_m angle minus theta, because our I is I_m angle minus theta. What about V_L ? V_L is I times the impedance of the inductor, which is $j\omega L$. Now this j is nothing, but angle pi by 2 and I is I_m angle minus theta. So, therefore, this quantity is $\omega L I_m$ times L angle minus theta plus pi by 2.

So, this is our Phasor diagram, this vector is V_s and V_s as angle 0 therefore, it is a longer X axis that is the real part of V . This is our V_R and V_R has angle minus theta; that means, theta in the clockwise direction and this is our V_L , what is the angle of V_L ? It is minus theta plus pi by 2; that means, we go clockwise by theta and then we go anti clockwise by pi by 2, that brings us to this angle all right and now this equation V_s equal to V_R plus V_L is essentially a vector equation, if we add V_L and V_R then we get V_s .

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Next let us take an RC circuit a series RC circuit, and this can be analyzed in exactly the same manner as the RL circuit which we just saw; what do we do in this case? We replace the resistor with its impedance which is R, the capacitor with its impedance which is 1 over j omega C and then we can obtain this current as V m angle 0 divided by R plus 1 over j omega C, that we write as I m angle theta, where I m its given by this quantity you should really verify this; and theta is pi by 2 minus tan inverse omega R C. Now this tan inverse omega R C can vary from 0 to pi by 2, and therefore this angle is basically a positive angle.

In the time domain we write the current as I m cos omega t plus theta, and now let us calculate theta for some component values say R equal to 1 ohm, C equal to 5.3 millifarad and f equal to 50 hertz. For this combination theta turns out to be 31 degrees and that corresponds to a time of 1.72 milliseconds. Now let us ask this question when does i of t go through its peak and that happens when omega t plus theta becomes equal to 0; that means, t is equal to minus theta divided by omega and that time is exactly 1.72 milliseconds.

Let us look at the plots law; this light curve is the source voltage 1 angle 0 and it goes through its peak at t equal to 0. Because there is the cos function, this dark curve is the current and that goes through its peak when t is equal to minus 1.72 milliseconds as we just discussed. In other words the current goes through its peak before the source voltage

goes through its peak, and that is why we say that the current leads the source voltage since the peak of i occurs θ by ω seconds before that of the source voltage.

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RC circuit

$$I = \frac{V_m \cos(\omega t)}{R + 1/j\omega C} = I_m \cos(\omega t - \theta)$$

where $I_m = \frac{\omega C V_m}{\sqrt{1 + (\omega RC)^2}}$, and $\theta = \pi/2 - \tan^{-1}(\omega RC)$

$$V_R = I \times R = R I_m \cos(\omega t - \theta)$$

$$V_C = I \times (1/j\omega C) = (I_m/\omega C) \cos(\omega t - \theta - \pi/2)$$

The KVL equation, $V_s = V_R + V_C$, can be represented in the complex plane by a "phasor diagram."

M. S. Park of Samsung

Let us draw a phasor diagram to represent the KVL equation in this case, that is V is equal to V_R plus V_C . What is V_R ? I times R where I is I_m angle θ . So, therefore, V_R is R times I_m angle θ . What about V_C ? V_C is I times the impedance of the capacitor, which is 1 over $j\omega C$. Now 1 over j is the same as $-j$ that is an angle of $-\pi/2$ and therefore, we get V_C equal to I_m by ωC that is the magnitude, and for the angle we get θ which comes from I and then we have this $-\pi/2$ coming from this 1 over j .

And as we mentioned earlier our θ is a positive angle between 0 and $\pi/2$ all right; with that information we can now draw the phasor diagram, this is our V_s and that is along the X axis, because it has an angle of 0 . What about V_R ? V_R has got an angle θ which is positive, and between 0 and $\pi/2$. So, that is what V_R looks like; what about V_C ? V_C has an angle of $\theta - \pi/2$. So, we go anti clockwise by θ and then become clockwise by $\pi/2$ that bring us to this angle. So, that is our V_C and now we can see that the vector equation $V_s = V_R + V_C$; each satisfied this is our V_s . So, V_R plus V_C brings us to V_s . So, this is the phasor diagram corresponding to this KVL equation in this case.

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Circuit example

$$Z_3 = \frac{1}{j \times 2\pi \times 50 \times 2 \times 10^{-3}} = -j1.6 \Omega$$

$$Z_4 = j2\pi \times 50 \times 15 \times 10^{-3} = j4.7 \Omega$$

$$Z_{EQ} = Z_1 + Z_3 \parallel (Z_2 + Z_4)$$

$$= 2 + (-j1.6) \parallel (10 + j4.7) = 2 + \frac{(-j1.6) \times (10 + j4.7)}{-j1.6 + 10 + j4.7}$$

$$= 2 + \frac{1.6 \angle (-90^\circ) \times 11.05 \angle (25.2^\circ)}{10.47 \angle (17.2^\circ)} = 2 + \frac{17.7 \angle (-64.8^\circ)}{10.47 \angle (17.2^\circ)}$$

$$= 2 + 1.69 \angle (-82^\circ) = 2 + (0.235 - j1.87)$$

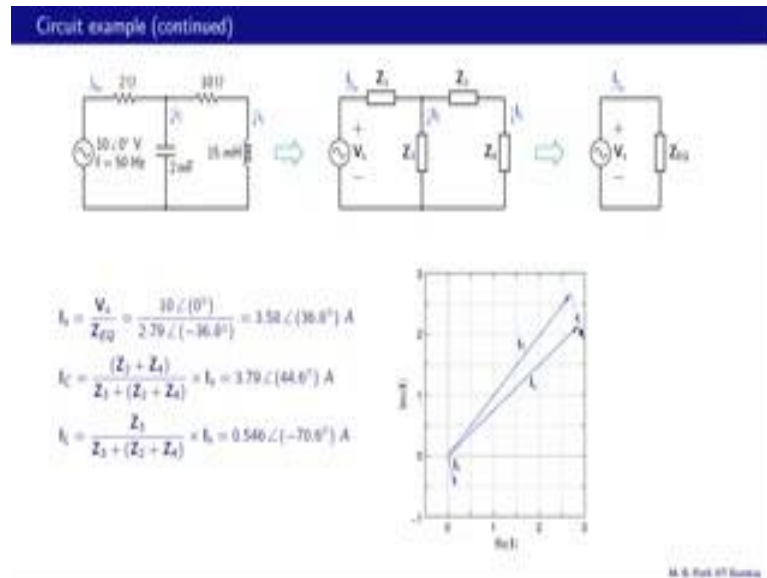
$$= 2.235 - j1.87 = 2.79 \angle (-46.8^\circ) \Omega$$

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Let us consider a little more complex circuit now the one shown here. So, we have a sinusoidal voltage source 10 angle 0 frequency 50 hertz, and we want to find these currents I S, I C and I L. Step number one we convert all the components to their impedances, so this 2 ohms of course remains 2 ohms; 10 ohms remains 10 ohms to millifarads becomes Z 3, where it Z 3 is 1 over j omega Z that turns out to be minus j 1.06 ohms, 15 milli entry becomes Z 4 where is Z 4 is j omega L omega is 2 pi times 50 that is 1. So, that turns out to be j 4.07 ohms. And now next step is we can calculate the equivalent impedance of this combination and this is exactly like series parallel register combination, we have Z 2 and Z 4 in series that combination in parallel with Z 3 and the whole thing in series with Z 1. So, that is what Z equivalent is Z 1 plus Z 3 parallel Z 2 plus Z 4.

You are definitely encouraged to go through all of these steps and arrive at this final result for Z equivalent, but also look up your calculator and it is possible that you will be able to do this calculation in a smaller number of steps, depending on what your calculator allows.

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Now, let us go ahead. So, we have come up to this step, you have found said equivalent we already have V s. So, now, we can calculate I s. So, I s is V s divided by Z equivalent, V s is 10 angle 0 and Z equivalent from the last slide is this number here, so that turns out to be 3.58 angle 36.8 degrees ampere.


Once we get I s we can now get I c by using the current division formula, that is I c equal to Z 2 plus Z 4 divided by Z 2 plus Z 4 plus Z 3 times I s like that and that turns out to be this number, what about I L? You can get I L in 2 ways: one I s minus I c gives us I L by KCL or we can use the current division formula I L is equal to Z 3 by Z 3 plus Z 2 plus Z 4 times I s either way you should get this number for I L. And now let us draw the phasor diagram, which describes the KCL equation at this node. Here is the phasor diagram and notice that we have used the same scale for the X and Y axis that is this distance which represents 1 unit on the X axis also represents 1 unit on the Y axis. And we do that so as to represent angles correctly; that means, a 45 degree angle would you indeed look like a 45 degree angle if we follow this practice.

All right let us now verify whether the KCL equation at this node is satisfied. What is the equation we have I s equal to I c plus I L. Our I s is 3.58 angle 36.08 degrees, that is this vector this magnitude is 3.58, and this angle is 36.08 degrees. I c has a magnitude of 3.79. So, little bit larger than I s and an angle of 44.06 degrees, that is I c this angle is 44.06 degrees. I L it is much smaller in magnitude 0.546 and it has a negative angle,

minus 70.06 degrees. So, that is our I_L ; and now we see that I_c plus I_L is indeed equal to I_s so; that means, KCL is verified.

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Maximum power transfer (sinusoidal steady state)



Let $Z_L = R_L + jX_L$, $Z_{Th} = R_{Th} + jX_{Th}$, and $I = I_m \angle \phi$.

The power absorbed by Z_L is,

$$P = \frac{1}{2} I_m^2 R_L$$

$$= \frac{1}{2} \left| \frac{V_{Th}}{Z_{Th} + Z_L} \right|^2 R_L$$

$$= \frac{1}{2} \frac{|V_{Th}|^2}{(R_{Th} + R_L)^2 + (X_{Th} + X_L)^2} R_L$$

For P to be maximum, $(X_{Th} + X_L)$ must be zero. $\Rightarrow X_L = -X_{Th}$.

With $X_L = -X_{Th}$, we have,

$$P = \frac{1}{2} \frac{|V_{Th}|^2}{(R_{Th} + R_L)^2} R_L$$

which is maximum for $R_L = R_{Th}$.

Therefore, for maximum power transfer to the load Z_L , we need,

$$R_L = R_{Th}, X_L = -X_{Th}, \text{ i.e., } \boxed{Z_L = Z_{Th}^*}$$

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We have looked at the maximum power transfer theorem for linear DC circuits; now let us look at maximum power transfer in the sinusoidal steady state. So, let us consider circuit whose Thevenin equivalent is given by this circuit here namely a voltage source V_{th} in series with an impedance Z_{th} ; both of these of course, are complex numbers. Now we connect a load impedance to the circuit and as a result a current is going to flow and that current is denoted by this phasor I .

We want to find the condition for which the power transfer from this circuit to Z_L is maximum. Let us begin with Z_L equal to $R_L + jX_L$, where R_L is the real part of Z_L and that is the imaginary part of Z_L and similarly let Z_{th} be $R_{th} + jX_{th}$. Let $I = I_m \angle \phi$, I_m is the magnitude of this phasor I here and ϕ is its angle.

Now, the power absorbed by Z_L is given by $p = \frac{1}{2} I_m^2 R_L$, where I_m is the magnitude of this phasor I and R_L is the real part of Z_L . What is I_m ? It is the magnitude of I , and what is I ? It is simply V_{th} divided by $Z_{th} + Z_L$. So, we get $p = \frac{1}{2} \frac{|V_{th}|^2 R_L}{(R_{th} + R_L)^2 + (X_{th} + X_L)^2}$. Let us rewrite this expression as $\frac{1}{2} \frac{|V_{th}|^2 R_L}{(R_{th} + R_L)^2 + (X_{th} + X_L)^2}$ and that multiplied by R_L . Now this $|V_{th}|^2$ is independent of Z_L and as far as we are

concerned that is just a constant. So, what we need to do now is to find conditions on Z_L for which this whole expression is maximum.

Now, for P to be maximum, clearly this denominator must be minimum and that will happen if X_{th} plus X_L squared is 0, because there is a square here the smallest value that this term can take is 0 and therefore, that gives us X_L equal to minus X_{th} and with X_L equal to minus X_{th} this second term disappears, and we get P equal to half mod V_{th} squared divided by R_{th} plus R_L squared times R_L . So, for maximum power transfer we need to maximize this expression now.

How do we do it? We differentiate P with respect to R_L and we equate dP/dR_L to 0 and then find that P is going to be maximum when R_L is equal to R_{th} . So, we have 2 conditions: one the imaginary part of Z_L that is X_L must be equal to negative of the imaginary part of Z_{th} that is X_{th} and the real part of Z_L must be equal to the real part of Z_{th} .

To summarize for maximum power transfer to the load Z_L , we need R_L equal to R_{th} and X_L equal to minus X_{th} ; that means, over load impedance must be the complex conjugate of the Thevenin equivalent impedance that is Z_{th} . So, Z_L must be equal to Z_{th}^* .

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Impedance matching

Calculate the turns ratio to provide maximum power transfer of the audio signal.

$$Z_L = Z_{th}^* \rightarrow \left(\frac{N_1}{N_2}\right)^2 \times 8\Omega = 1k\Omega \rightarrow \frac{N_1}{N_2} = \sqrt{\frac{1000}{8}} = 11.2$$

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Let us now look at an application of the maximum power transfer theorem for sinusoidal steady state here it is an audio amplifier driven by an audio input signal. So, the frequencies here would be in the range 20 hertz to say 16 kilo hertz or so. This audio amplifier is followed by transformer and we will soon comment on why this transformer is required and then finally, we have this speaker. This speaker has a complex impedance which varies with frequency, but in the audio range its resistance is more or less constant typically 8 ohms, and its imaginary part can be ignored. So, the equivalent circuit represent take this entire situation is given by this circuit here, this source here represents the input signal amplified by the gain of the audio amplifier and then this one k resistance here is the output resistance of the audio amplifier, and then we have the transformer with a turns ratio N_1 , N_2 and finally, the speaker which is represented by this a 8 ohms resistance here.

Our objective of course, is to maximize the power transfer from this circuit to the speaker that is how we will hear the loudest sound that is possible with the given input signal. So, let us simplify the circuit, we can transfer this resistance to the other side of the transformer and then it becomes N_1 by N_2 squared times 8 ohms.

Here is our actual problem statement, we look at this circuit and calculate the turns ratio to provide maximum power transfer of the audio signal; what is the maximum power transfer theorem? It says that Z_L must be equal to Z_{th}^* , and in this case since the imaginary parts of Z_{th} and Z_L are 0 all it means is that the real parts must be equal; that means, we must have N_1 by N_2 squared times 8 ohms equal to 1 k like that.

And we can now solve this equation for N_1 by N_2 and that gives us N_1 by N_2 of about 11. So, if we pick a transformer with this translation then it is current it that the audio signal will deliver maximum power to the speaker, and we will hear the loudest possible sound that is possible with this input signal.

In conclusion we have seen how to use phasors to analyze RLC circuits in the sinusoidal steady state; this background is going to be very useful when look at filter circuits be a little later. We also looked at the maximum power transfer theorem for RLC circuits in the sinusoidal steady state. We considered an example which is very important in practice namely how to obtain maximum audio power from a speaker, that is all for now.

See you next time.