

Power System Dynamics and Control
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Lecture No # 22
Synchronization of a Synchronous Machine (Contd.)

We have been studying the synchronization transient of a synchronous machine connected to an infinite bus. What we saw last time was that, you can have a bump less synchronization of a synchronous machine to an infinite bus, which is essentially a fixed voltage source, the machine locks on to the frequency of the **synchronization** of the infinite bus. In case, it is synchronized well, now the thing we noticed last time was that, if the speed of the machine initially is almost the same as that of the infinite bus. And, at the instant of you know the interconnection, the voltage, open circuit voltage of the generator, the voltage phasor you can say is practically the same as the infinite bus voltage phasor, then you get a kind of a bump less transfer.

Of course, you cannot get an perfect transfer without transients, it's practically not possible. There will be some transient, because you cannot might do this match, which I mentioned sometime back exactly. So, what you really see is that, when you connect a synchronous machine to an infinite bus, you do get a small synchronization transient. Most notably, you notice a kind of a what is known as a swing mode or a low frequency oscillation, which usually damps down and that really after that, synchronous machine is in synchronism with the voltage source. Thereafter, if you increase the mechanical torque to the synchronous machine, you find that the machine transfers power from **from** itself to the voltage source or the infinite bus.

So, what you notice there of course, is that as you go on increasing the mechanical power, you will find at that the δ angle increases and in fact there comes a point after which if you try to increase the mechanical power the synchronous machine in fact loses synchronism. There is it kind of fails to reach a steady state, in case you go on increasing the mechanical power beyond the point.

Of course, the important thing is that we had in fact seen precisely such a transient in one of our lecture. In fact, the first lecture, I had shown you a small demonstration clip, it will be good to revisit that clip, you can go back to the first lecture and see that clip that

is precisely what we have tried to stimulate, tried to stimulate in the 21 lecture. That is a synchronous machine connected to a infinite bus and there in thereafter, we go on increasing the mechanical power to point at which it loses synchronism

One of the things which we did not do, when we increase the mechanical power was the mechanical torque to the synchronous machine was that, we did not increase the field voltage simultaneously. In fact, a synchronous machine tends to lose synchronism very easily, in case you try to load it without a concurrent increase in the field voltage.

So, what we will do **this**, in this lecture is? We will first (()) revisit this earlier transient, in which we first give a small step, we first synchronize the machine then, give a small step in the mechanical power of 0.25 per unit and then we try to increase the mechanical power right up to its rated value of 1 per unit, and there we do not increase the field voltage, we will see that it loses synchronism. Thereafter, we will redo this simulation with, by increasing the field voltage concurrently, that is at the same time as we rather at this we increase the field voltage simultaneously, traps simultaneously the better word. We increase the field voltage simultaneously along with the mechanical torque. You will find that under such circumstances the **the** generators able to supply the rated power to the infinite bus.

Now, one small clarification of course, I hope you did not miss it last time was that, when we are simulating the equations of a synchronous machine, we are neglecting the stator transient or the stator flux transients. That is we have replaced $d\psi_d$ by $d\psi_d/dt = 0$ and $d\psi_q$ by $d\psi_q/dt = 0$. The reason is of course, that the transients associated with the these states ψ_d and ψ_q are very fast. So, what we have done essentially is, replace the differential equations which relate the ψ_d and ψ_q fluxes by algebraic equations. One thing you should note here is that, an implicit thing is that, when I am making this assumption, I am really interested in the slower electro mechanical transient. That is why; I can do without the ψ_d and ψ_q differential equation, replace them in fact by the algebraic equations and still get a reasonably correct result.

Now, you may ask that well, why do I do it anyway you did can just as well keep the differential equations corresponding to ψ_d and ψ_q , the reason why of course, I have done that is if I retain the differential equation **in $d\psi$** in ψ_d ψ_q , you get what is known as stiff system, that is a mixture of fast and slow transients are exciting in that

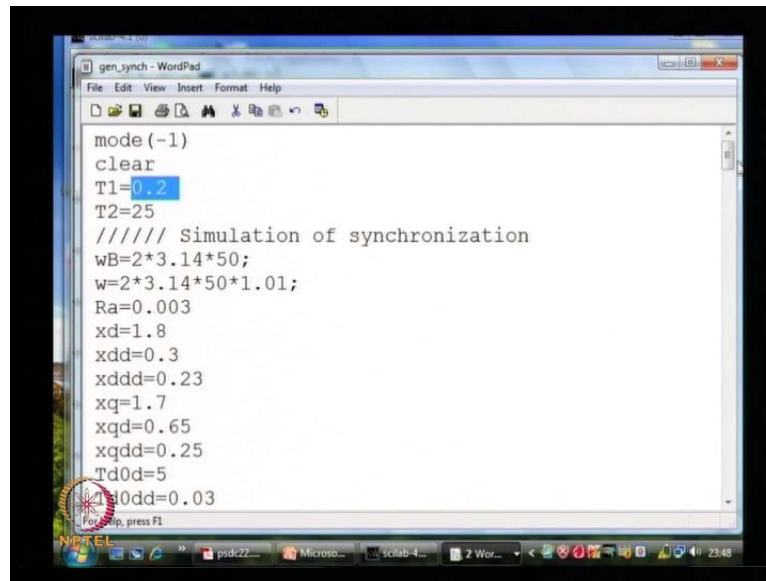
system and as a result of which, simple numerical integration scheme are likely to misbehave. So, what I have really done is, I have made this approximation of neglecting the fast transients, removed in some sense the stiffness of **of** the system and then used Euler method, which is a simplest method to simulate the system.

So, that is why I have done this. If I if of course, I had used I had retained the stiffness in the system, that is I had retained the $\frac{d\psi}{dt}$ and $\frac{d\psi_q}{dt}$ terms in the equations. In such a case, you would find that you would need to use method like trapezoidal rule or backward Euler method. Now the problem in doing that is, once you discretize the differential equations with trapezoidal rule or backward Euler method, remember these are non-linear differential equations. Then, what you really get is are non-linear algebraic equations. Once you discretize the system by **by** these numerical methods, you get non-linear algebraic equations and for every time step, I would need to solve numerically **solve** for every time step or numerically you have to solve the algebraic equations.

On the other hand, Euler method is simple and the per time step computation is very very straight forward. So, that is what we have done so far. Today, I will try to simultaneously increase the field voltage along with a torque increase. And, we can see that in such a case, you can in fact run the synchronous machine at rated power. So, let us just redo the transient we did last time and go head and simulate the transient with field voltage increase simultaneously.

So, let us go on to that. So today's lecture of course, we will be revisiting the transient associated with a synchronizations of a synchronous machine. Now, what we will do is of course, like in the previous lectures, we will do a simulation of this system.

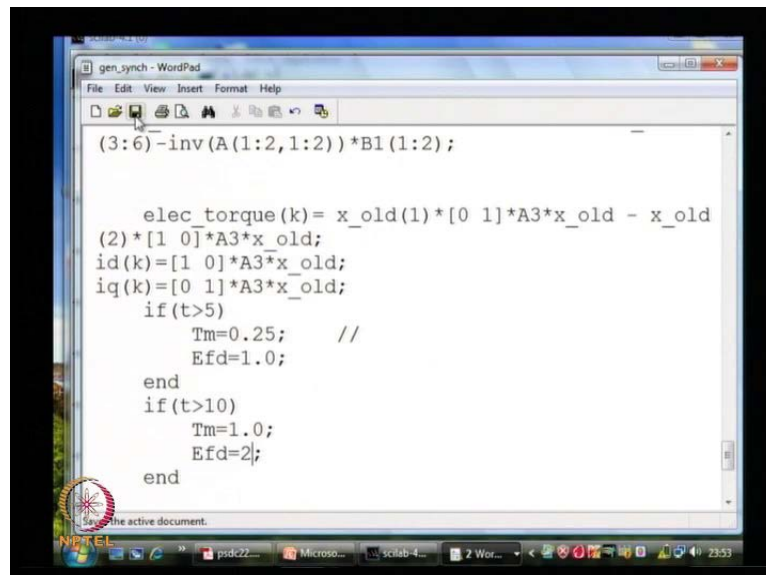
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```
gen_synch - WordPad
File Edit View Insert Format Help
mode(-1)
clear
T1=0.2
T2=25
///// Simulation of synchronization
wB=2*3.14*50;
w=2*3.14*50*1.01;
Ra=0.003
xd=1.8
xdd=0.3
xddd=0.23
xq=1.7
xqd=0.65
xqdd=0.25
Td0=5
Td0dd=0.03
```

So, let me just show you the simulation file here. This is gen underscore sink is the file, I will be doing a 25 second simulation. I will be synchronizing the machine at the point two seconds. This is of course, **the** you see the data of the machine, this something you have done a quiet a few time before in this **in these** lecture.

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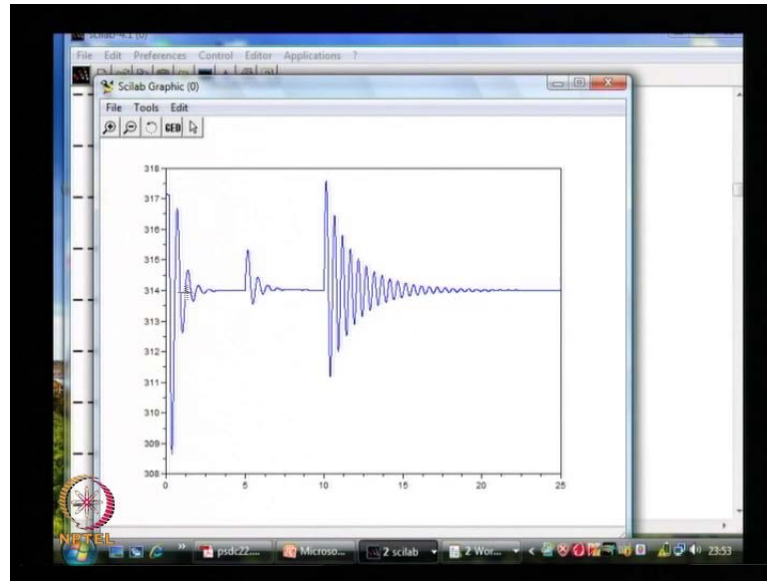


```
gen_synch - WordPad
File Edit View Insert Format Help
(3:6)-inv(A(1:2,1:2))*B1(1:2);

elec_torque(k)= x_old(1)*[0 1]*A3*x_old - x_old
(2)*[1 0]*A3*x_old;
id(k)=[1 0]*A3*x_old;
iq(k)=[0 1]*A3*x_old;
if(t>5)
    Tm=0.25; //
    Efd=1.0;
end
if(t>10)
    Tm=1.0;
    Efd=2;
end
```

So, I will just quickly came through this. As the mechanical power is changed to one per unit, we also double the field voltage. So, simultaneously we double the field voltage and will redo this transient, redo this simulation here.

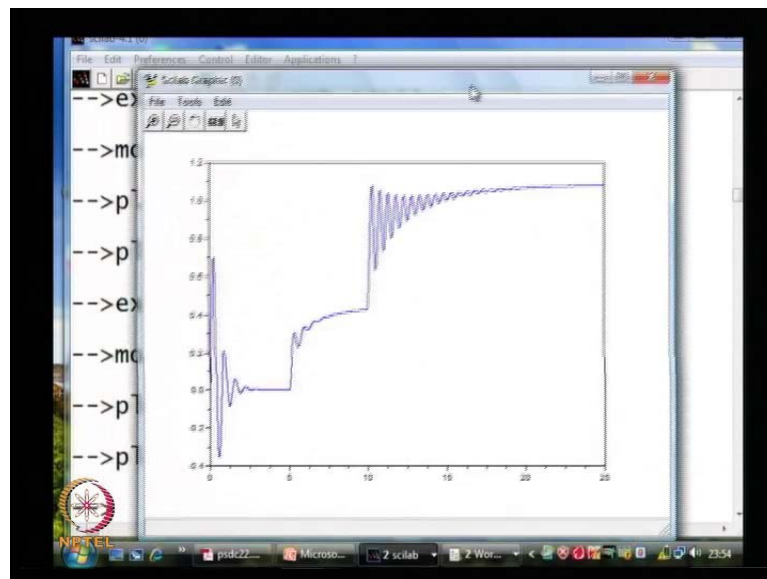
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And I plot, now the speed. I notice that, in this particular case, the machine regains rather is still in synchronism.

So, by increasing the field voltage in some sense have insured that the machine is able to deliver the rated power. And the remain in synchronism and operates stably of course, these there is a electro mechanical oscillation or a swing, which is seen in the speed, whenever we give any such disturbs. Now, we can of course, have a look at how delta looks is well **yeah**.

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So, for the first disturbance of course, delta goes from 0 to 0.4 roughly, this is radians. And the second transient also it is stable, there is a swing and there is a gradual raise as well, rather kind of a oscillate decaying oscillation, which we see here and also decaying exponent, I meant there is also this kind of raise because of, which it settles finely near about one radiant.

So, this is basically how the system behaves in case we increase the mechanical power along with a change in the field voltage. If we do not change the field voltage along with the mechanical power, the system can lose synchronism. So, that is one important thing you should remember. In the next ah next few lectures, we will go on to talking about excitation systems in the sense, that we will try to understand how the field voltage is changed ah by by a control system or an a control system and an excitation system. Field voltage in a synchronous machine changes along with the loading. In fact, we generate in such a way that whenever the loading of a synchronous machine changes, the field voltage is also changed. It is in fact every necessary to do. So, this is something you should keep in mind.


Now, what we will do next is look at another form of analysis, what we will do is? Do a kind of small signal analysis of a synchronous machine. So far, what we have been doing here is, in fact looking at the a numerical simulation of the system. The reason why we do a numerical simulation is that the system is non-linear. Now, we can in fact do a linearized analysis around an operating point using Eigen analysis, the Eigen analysis tool. As we have done before, the only reference here is that (()) this is non-linear we have to first convert the system into a linearized system around an operating point.

So, what we will be really doing is (()) are non-linear system, we cannot inferred the non-linear behavior by Eigen analysis. That is not possible but, what we can do is, if you are at an operating point, that is an at an equilibrium point. If you give small disturbances, you can rewrite the equations in the linear form, which has valid only for small disturbances and inferred the behavior from the Eigen values of the resulting linear differential equations.

So, this is what we will do next at this point, we should look it, where are actually the nonlinearities in our equations. Let us have a look.


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Compact Form of Flux Equations

$$\frac{d}{dt} \begin{bmatrix} \psi_d \\ \psi_q \\ \psi_F \\ \psi_H \\ \psi_G \\ \psi_K \end{bmatrix} = A_1 \begin{bmatrix} \psi_d \\ \psi_q \\ \psi_F \\ \psi_H \\ \psi_G \\ \psi_K \end{bmatrix} + A_2 \begin{bmatrix} i_d \\ i_q \end{bmatrix} + B_1 \begin{bmatrix} v_d \\ v_q \end{bmatrix} + B_2 E_{fd}$$


Now, if you look at the equations of a synchronous machine, **we had** we have been using these compact form of equations in our analysis, that is first of all we have written the flux equations, $\frac{d}{dt}$ of the flux equations is equals to A_1 into the flux. So, this is the states space form of the flux equations. Remember that, A_1 is a function of speed. So, actually all though this is written, it looks almost as if it is linear. A_1 is actually a function of the state, of the states. Speed is the state of the system. Remember that, this is written in a kind of a composite form i_d i_q are in fact related to all the fluxes by another set of equations.

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$$A_3 \begin{bmatrix} \psi_d \\ \psi_q \\ \psi_F \\ \psi_H \\ \psi_G \\ \psi_K \end{bmatrix} = \begin{bmatrix} i_d \\ i_q \end{bmatrix}$$


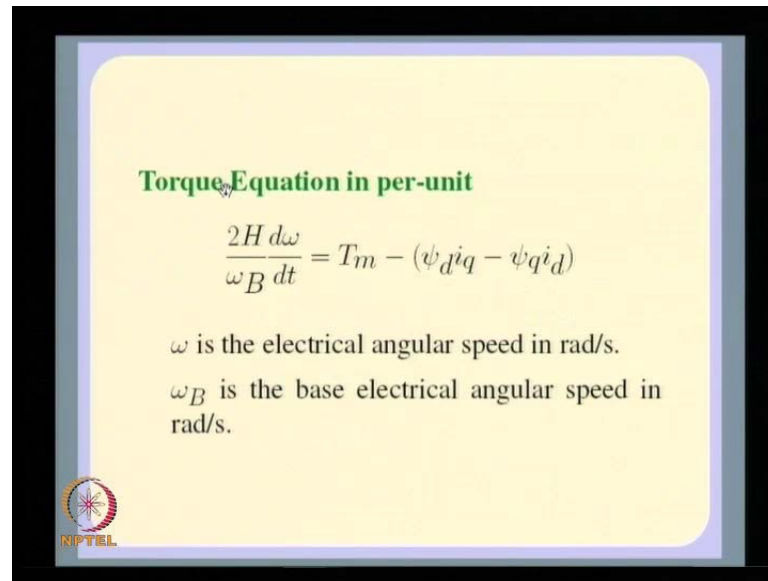
The in fact, the inductance effectively the inductance matrixes or the reactants matrixes in per unit.

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$$A_1 = \begin{bmatrix} 0 & -\omega & 0 & 0 & 0 & 0 \\ \omega & 0 & 0 & 0 & 0 & 0 \\ \frac{1}{T_d'} & 0 & -\frac{1}{T_d'} & 0 & 0 & 0 \\ \frac{1}{T_d''} & 0 & 0 & -\frac{1}{T_d''} & 0 & 0 \\ 0 & \frac{1}{T_q'} & 0 & 0 & -\frac{1}{T_q'} & 0 \\ 0 & \frac{1}{T_q''} & 0 & 0 & 0 & -\frac{1}{T_q''} \end{bmatrix}$$

These equations are of course, in per unit form. A 1 looks like this, it is in fact **linear** non-linear because, A 1 itself is a function of the states. So, A 1 into psi involves effectively product terms of speed and flux. So, that is why it is basically a set of non-linear equations. So, there is nonlinearity here, in fact these are of course, the definitions of **A 1 the** A 1 B 1 B 2 A 3. So, one of the things you will notice here is of course, I think I have **I have I have** effectively written the torque equation, if you look at it in per unit, this also involves speed product terms.


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Torque Equation in per-unit

$$\frac{2H}{\omega_B} \frac{d\omega}{dt} = T_m - (\psi_d i_q - \psi_q i_d)$$

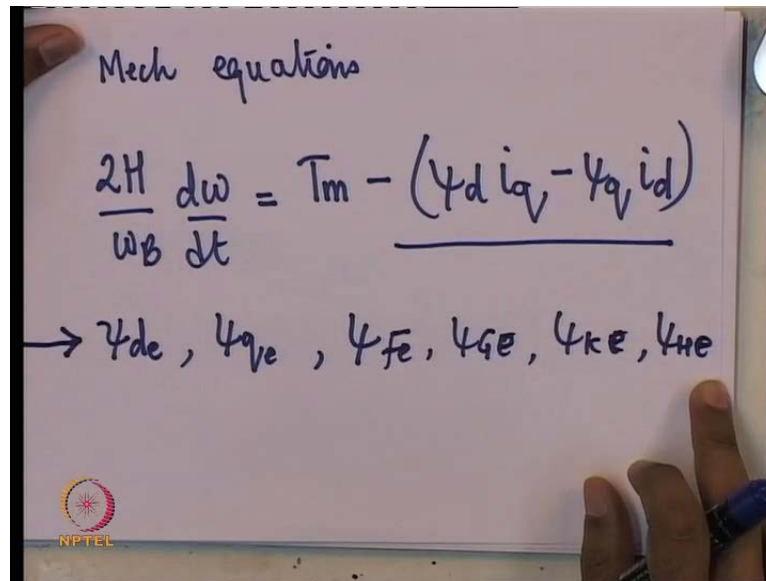
ω is the electrical angular speed in rad/s.
 ω_B is the base electrical angular speed in rad/s.



So, if you look at the basic compact equations of a synchronous machine. This is a product kind of term, even contains is the function of the states and also the torque equation in per unit is a non-linear equation. So, when we are trying to use Eigen analysis tools, remember that the mechanical and electrical equations are in fact coupled and you see the flux and current equation in the mechanical equations and in the flux equations, there is a omega dependants.

So, because of this nonlinearity, we cannot directly apply linear analysis. We will have to do a linearization around an equilibrium point. So, the basic idea is, suppose you are operating at a particular equilibrium point.

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Mech equations

$$\frac{2H}{\omega_B} \frac{d\omega}{dt} = T_m - (\psi_d i_q - \psi_q i_d)$$

→ $\psi_{de}, \psi_{qe}, \psi_{fe}, \psi_{ge}, \psi_{ke}, \psi_{he}$

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So for example, if you look at the mechanical equations, $\frac{2H}{\omega_B} \frac{d\omega}{dt}$ is equal to T_m minus $\psi_d i_q$ minus $\psi_q i_d$. Now, if you look at this, this is a non-linear equation.

Now, if you are operating at a certain equilibrium point and the equilibrium values of ψ_d is ψ_{de} similarly, the equilibrium value of ψ_q is ψ_{qe} and of course, i_q and i_d are dependent on the equilibrium values of ψ_{fe} in addition to these ψ_{fe} , ψ_{ge} , ψ_{ke} and ψ_{he} . So, these are the equilibrium point. So, this e here is denoting equilibria. So, this the equilibrium value of these states. So, if you look at you **rewrite** rewrite these equations again. If your **speed** equilibrium speed of a synchronous machine connected to an infinite bus is well if you look at $\frac{d\delta}{dt}$ it is nothing but, $\omega - \omega_0$, this is the frequency of the infinite bus.

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$$\frac{d\delta}{dt} = \omega - \omega_0 \leftarrow \text{freq, infinite bus.}$$
$$\omega_e = \omega_0 \checkmark$$

$\psi_{de}, \psi_{qe}, \psi_{fe} \dots ?$
 $\delta_e \dots ?$

So, if look at what is the equilibrium speed of the machine? The equilibrium speed of the machine is ω_e and by definition of equilibria, equilibria **are** the value of the states at which all the d by e t is become equal to 0. So, if you look ω_e it should be equal to ω_{naught} . So, the value of ω at equilibrium is ω_e , it is equal to ω_{naught} . So, these are essentially the equilibrium values of the states. So, equilibria are actually obtained by setting $d\delta$ by $d t$ $d\omega$ by $d t$ and all the $d\psi$ by $d t$ is equal to 0. Now, what are the equilibrium values of, what are the actual values of ψ_d , ψ_q , ψ_f and so on. What is the equilibrium value of δ_e for example, what is the equilibrium value of δ ? So, all these are all the states, so one of the equilibrium value we have of course got, then they are 6 flux states for which we have to get the equilibrium values and δ_e equilibrium value has to be obtained.

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The image shows a piece of paper with handwritten mathematical equations. At the top, it says $\psi_{d i_q} - \psi_{q i_d}$. Below that, the electrical torque is given as $T_e = \psi_{d e} i_{q e} - \psi_{q e} i_{d e}$. A large box is drawn around the following expression:
$$= \frac{E_{fd} \cdot V_{LLrms} \sin \delta_e}{X_d} + \left(\frac{1}{X_q} - \frac{1}{X_d} \right) \frac{V_{LLrms}^2}{2} \sin 2\delta_e$$
 In the bottom left corner of the paper, there is a small circular logo with a star and the text 'NPTEL'.

Now remember that, we had obtained the expression for the electrical torque, if you recall what is the electrical torque expression? The electrical torque expression was, in steady state of a synchronous machine was, electrical torque under equilibrium conditions is of course, it is $\psi_{d e} i_{q e}$ minus $\psi_{q e} i_{d e}$. But, we can write it down as, we have derived this in some of our previous lectures, $\sin \delta_e$ upon X_d plus $\left(\frac{1}{X_q} - \frac{1}{X_d} \right) \frac{V_{LLrms}^2}{2} \sin 2\delta_e$. So, the equilibrium value of the torque is nothing but E_{fd} , this is depending on the field voltage, V_{LLrms} is the voltage line to line rms voltage of the voltage source to which the synchronous machine is connected to δ_e of course, is the equilibrium angle.

Now, the important thing is that, **this is the this** if you look at $\psi_{d e} \psi_{d i_q}$ minus $\psi_{q e} i_{d e}$, this is essentially an equilibrium rather the transient, this expression for torque is valid also in transient conditions. But, (()) remember this expression here which I have written down is only valid during steady state. So, please do not use it for transient conditions. So, under equilibrium conditions, this is true.

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$$\frac{d\delta}{dt} = 0 \Rightarrow \omega_e = \omega_0$$
$$\frac{d\omega}{dt} = 0 \Rightarrow T_m = T_e$$

δ_e

So, if I set $d\delta/dt$ is equal to 0, we **we** get the equilibrium value of ω should be equal to ω_0 . Also, by setting $d\omega/dt$ equal to 0, we get T_m should be equal to the equilibrium value of this. Now, if we know what E_{fd} is and what V_{line} to V_{rms} is, from this we can easily back calculate what the value of δ_e is? So, **you** I know what T_e is T_m , so if I given you what the mechanical torque is and what the E_{fd} is. In such a case, you should be able to compute δ_e . So, first state, second state is computed. This is a equilibrium values of ω_e and δ_e .

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$$\omega = \omega_e$$
$$V_d = -V_{um} \sin \delta_e$$
$$V_q = V_{um} \cos \delta_e$$
$$\theta = \omega_0 t + \delta$$

Thereafter, once if you know the equilibrium value of ω is ω_e , also you have got the equations of the **the** differential equations corresponding to the fluxes. So, if you look at them here, I have written them down in compact form here.

So, you set this equal to 0. A_1 is a function of the time constants as well as ω but, you know that the equilibrium value of ω is ω_e . So, you get an **equation** so, you can evaluate what A_1 is, this is set to 0. A_2 into i_d i_q , i_d i_q itself is a function of ψ_d and ψ_q , it is a linear function. So, A_2 into A_3 into ψ basically, because i_d i_q itself is related by this expression. Thereafter, if you look at d_1 , so b_1 is also a constant matrix. V_d and V_q in fact, if you look at what I am writing here, V_d we have seen before by the definition of the voltage source, we have we wrote down the time, the **voltage** voltage of the infinite bus in terms of V_a V_b V_c , as per that definition V_d will come out to be $\sqrt{3} V_{line} \sin \delta_e$ and V_q comes out to be $\sqrt{3} V_{line} \cos \delta_e$.

So, this is of course, obtained we have obtained this before in our course, remember that the rotor position is nothing but, $\omega t + \delta_e$. So, this is the definition of the rotor position, ω is the speed of the infinite bus and because of that and the definition of v_a , v_b and v_c the phase to neutral voltages of the infinite bus, we obtain v_d and v_q in this form. So, **what I wish to say** what I wish to say here is that, first of all we have got the equilibrium value of the speed, then we have got the equilibrium value of δ_e . From equilibrium value of speed and δ_e , we can from the flux equations which are shown here, by setting the right hand side sorry the left hand side here equal to 0. That is put setting d by d_t equal to 0, obtain the equilibrium values of ψ_d , ψ_q , ψ_F , ψ_H , ψ_G and ψ_K .

So, this is basically what is involved in getting first the equilibrium values. Once you have got the equilibrium values in order to do a small signal or a small disturbance **analysis** analysis around an equilibrium point. What we really need to do is? Reformulate the equations or linearized the equations, which is suitable for such small disturbance analysis.

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The image shows a whiteboard with two equations written in black marker. The first equation is $\frac{d\psi}{dt} = A_1 \psi + \dots$. The second equation is $\psi = \underline{\psi_{de}} + \underline{\Delta\psi}$. A hand holding a blue pen is visible on the right side of the board. In the bottom left corner, there is a logo for NIPTEEL.

So for example, we have got $\frac{d}{dt}$ of the flux is equal to A_1 into the flux and so on. Now, what we do is we assume that the fluxes, we are considering only small disturbances are from the equilibrium. So, we will assume that ψ_{de} plus $\Delta\psi$. So, this is the equilibrium value of the flux and this is the deviation from it.

So, we can reformulate the differential equations in terms of only the deviations around the equilibrium. Now, if these deviations are small, one can do a procedure called linearization.

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The image shows a whiteboard with three equations written in black marker. The first equation is $\frac{2H}{\omega_B} \cdot \frac{d\omega}{dt} = T_m - (\psi_d i_{q'} - \psi_q i_{d'})$. The second equation is $\omega = \omega_e + \Delta\omega$. The third equation is $\frac{2H}{\omega_B} \frac{d\Delta\omega}{dt} = T_m - (\psi_d i_{q'} - \psi_q i_{d'})$. A hand holding a blue pen is visible on the right side of the board. In the bottom left corner, there is a logo for NIPTEEL.

So, I will just show you the linearization procedure for a simple case, 2H we just look at this equation, you can actually linearize every equation equal T_m minus $\psi_d i_q$ minus $\psi_q i_d$. This is of course, a non-linear term, now we assume ω is equal to $\omega_{\text{equilibrium}}$ plus the deviation from the equilibrium. This is a known quantity. So, we can reformulate these equations, you can write this as $2H B d$ of ω_e plus $\Delta \omega$, ω_e is a constant, it is the equilibrium value of this speed. So, what we will do is, is equal to T_m , we will assume that T_m is a constant minus $\psi_d i_q$ minus $\psi_q i_d$. Now, $\psi_d i_q$ minus $\psi_q i_d$ in fact can be written down as.

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$$\begin{aligned}
 & (\psi_{de} + \Delta \psi_d) (i_{qe} + \Delta i_q) \\
 & - (\psi_{qe} + \Delta \psi_q) (i_{de} + \Delta i_d) \\
 & = (\psi_{de} i_{qe} - \psi_{qe} i_{de}) \\
 & \quad + \psi_{de} \Delta i_q - \psi_{qe} \Delta i_d \\
 & \quad + \Delta \psi_d i_{qe} - \Delta \psi_q i_{de} \\
 & \quad + \text{2nd order terms}
 \end{aligned}$$

So, it is $\psi_e \psi_d$ plus $\Delta \psi_d$ into i_q plus Δi_q minus $\psi_q \psi_e$ plus $\Delta \psi_q$ into i_d plus Δi_d .

So, this becomes if you neglect, well first let us write it down. So, what we can write it down is $\psi_d i_q$ minus $\psi_q i_d$ plus $\psi_d \Delta i_q$ plus $\psi_q \Delta i_d$ **sorry** this should be a minus, plus $\Delta \psi_d i_q$ **sorry I am sorry** $\Delta \psi_d$ into i_q minus $\Delta \psi_q$ into i_d plus, what are second order terms? that is product of two delta kind of terms. Note that $\psi_d i_q$ minus $\psi_q i_d$ is in fact equal to the equilibrium value of the torque, electrical torque which is of course, going to be equal to the mechanical torque T_m . So, $\psi_d i_q$ minus $\psi_q i_d$ is in fact T_m .

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$$\frac{2H}{\omega_B} \cdot \frac{d\Delta\omega}{dt} = \psi_{de} \Delta i_d - \psi_{qe} \Delta i_q + \Delta\psi_d i_{de} - \Delta\psi_q i_{qe}.$$
$$\dot{\Delta x} = A \Delta x$$
$$\Delta T_m = 0$$
$$\Delta E_{fd} = 0.$$

So, what we have essentially is the linearized equation $\frac{2H}{\omega_B} \frac{d\Delta\omega}{dt}$, if you assume that T_m is a constant in fact, you can write this as $\psi_{de} \Delta i_d - \psi_{qe} \Delta i_q + \Delta\psi_d i_{de} - \Delta\psi_q i_{qe}$.

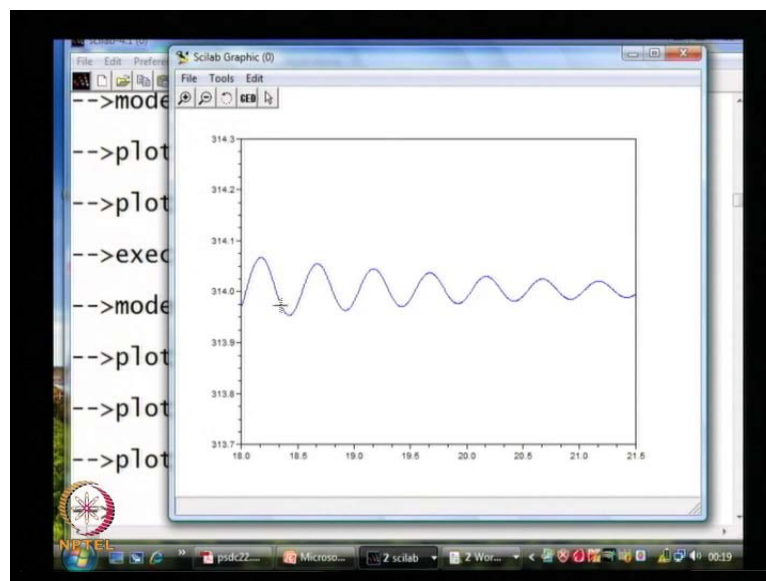
So, this is just an example of how you can do linearization. So, you can in fact linearize the $\frac{d\psi_d}{dt}$ equation and $\frac{d\psi_q}{dt}$ equation in a similar fashion. So, what you eventually get is, you know you have to couple of course; the mechanical and electrical equations the flux equations are coupled. So finally, what you will get is going to be Δx in the linearized states assuming of course, $\Delta T_m = 0$ and $\Delta E_{fd} = 0$. So, we will be assuming the inputs are constants, will be only giving changes in Δx that is giving the initial condition or initial disturbances to the states.

So, around the equilibrium point, the behavior of the equations for small disturbances can be studied using this. So, I have just shows you how you can do a linearization. So, I request you to go **go** back and just try to do the full system linearization and get it in the form $\dot{\Delta x} = A \Delta x$. The properties of the transients for small disturbances around the equilibrium can be got just by analyzing the Eigen properties. That is the Eigen values and Eigen vectors of the resulting state matrix (A) of course, A now is dependent on the equilibrium point because, it really has terms like ψ_{de} and so on. But, is a constant for that equilibrium point, it is the constant. But, it is a function of the value of the states at the equilibrium point.

So, one of the things you should remember is that a non-linear system, which is linearized and then we if we do non-linearized rather linearized analysis on it, Eigen analysis on **on** it, what we will see that the Eigen values which, we get for different equilibrium points are going to be different. So, the transient behavior around different equilibrium points in facts, is going to be different. It is not going to be the same. Did we actually notice this in our simulations, the answer is yes.

So, let us look at the simulations again carefully. So, back to the simulations, if you notice even one of the plots, which I had shown you, **yeah** maybe we can have a look at the speed transient. Well, you may say well, there is no difference in the kind of behavior around an equilibrium point, you have these swings. But, one of the simple things you can have actually look at look for rather is whether there is any change in the frequency of the swings. For example, the frequency of the **swings** swings, we will just do you do this again one second, we just get this back yeah. If you look at the frequency of swings around this equilibrium point **yeah**, I will just **yeah**.

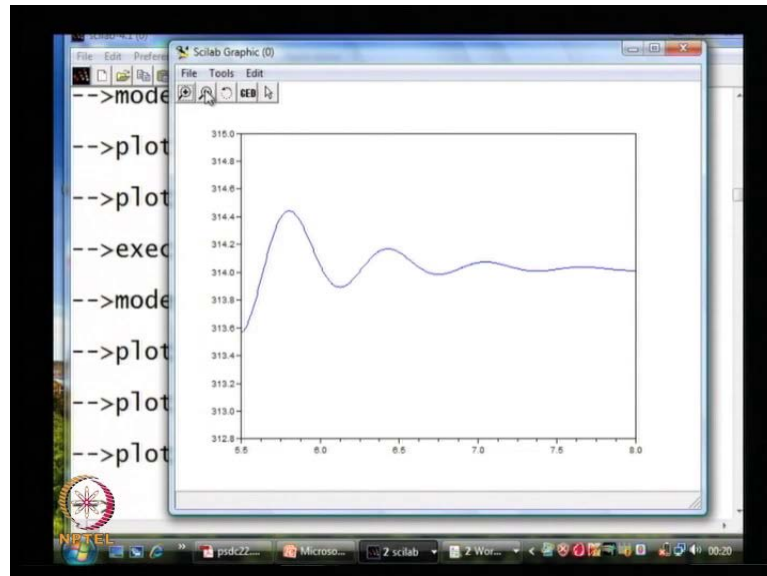
(Refer Slide Time: 32:39)



You look at this frequency of the swings around this equilibrium point, you will find it is well almost half a second. So, this is almost a 2 hertz oscillation, around the equilibrium point corresponding to **I am sorry yeah**, around the equilibrium **corresponding** point corresponding to T_m is equal to 1 and $E_f d$ is equal to 2. The frequency of oscillations

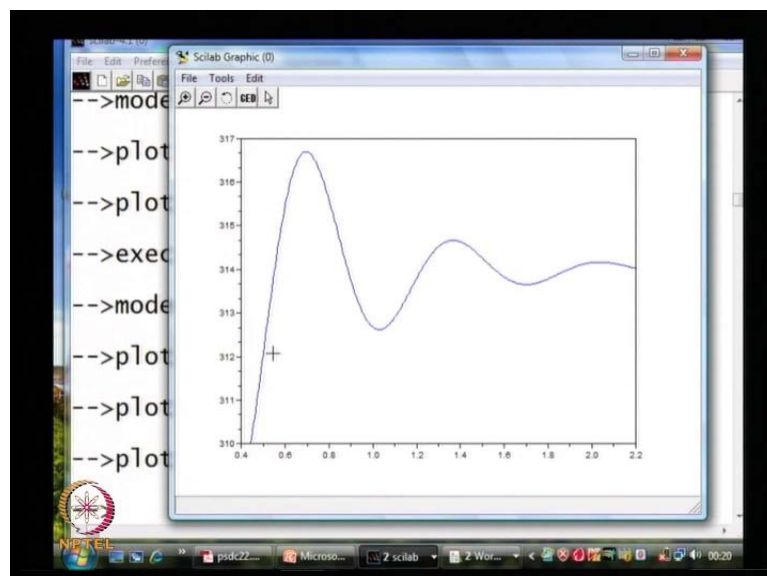
is in fact frequency of oscillations is in fact around 2 hertz. What about the frequency near about the equilibrium corresponding to T_m is equal to 0.25 and $E_f d$ is equal to 1.

(Refer Slide Time: 33:29)



So, let us just have a look at that here. So, if you look at that, this is a roughly 6.1 seconds and this is roughly well, this is near about slightly more than the frequency, the period is more than 0.5. It in fact the frequency seems to be lesser near about T_m is equal to 0.25 per unit and if you look at the behavior just of the synchronization, when mechanical power in fact 0.

(Refer Slide Time: 34:13)



The frequency of oscillation is **is** again roughly 0.7 **yeah**. So, with T_m is equal to 0 the frequency slightly lesser, in fact the period is larger and with larger values of T_m with of course, a simultaneous increase of $E_f d$, we see that the frequency in fact slightly increases, in fact it is higher here compared to here. So, as the operating point changes, we get a change in frequency.

That is not suppressing because, the Eigen values change because, the A matrix whose Eigen values, we are computing is the function of the equilibrium values of the states, which are obtain after linearization. So, that is basically why this happens, in fact one more striking thing, which you can observe from this plot. In fact all the plots revile a great deal of information. You will notice that the rate of decay of this oscillation is much much slower compared to the decay of oscillation here. In no time, you will find that this reaches equilibrium whereas; this takes several seconds to reach an equilibrium.

So, what one can expect after doing a linear analysis, an Eigen value analysis we will do that shortly of course, is that we can infer that whether, this frequency is going to change whether, the lamping is going to change and so on from that analysis itself.

So, let us just verify that, by actually writing down an **analysis** Eigen value analysis program. So, as I mentioned sometime back of course, it involve some **some** deal of effort in the sense that, you will have to linearized the set equations, form the A matrixes, the A matrix of the system, which is the function of the equilibrium point.

So, this is something you **you** will have to do. Now, one of the **one of the** points, which I need to of course, clarify this point the A matrix, which are going to get here, would be initially we will include the $d\psi/dt$ and $d\psi_q/dt$ transient. We will not convert $d\psi/dt$ and $d\psi_q/dt$ differential equation into algebraic equation. We will first consider all of them together.

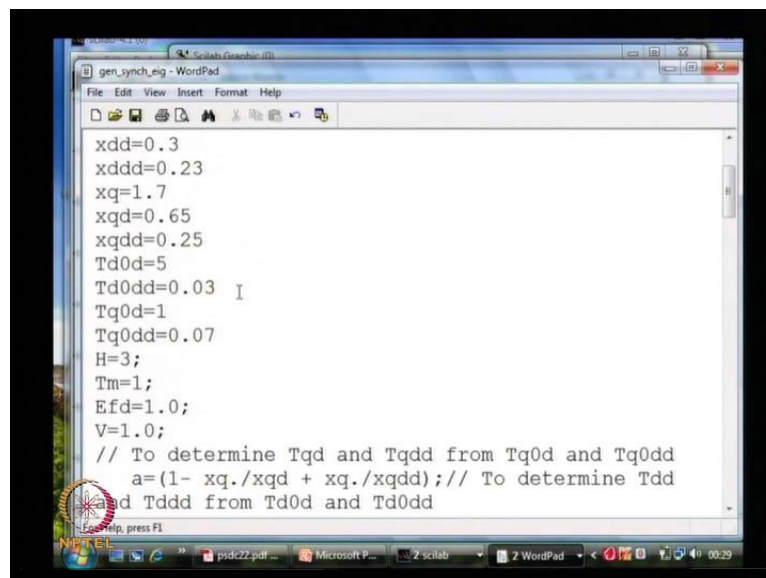
So, what we are really looking at is a model, in which you are going to have this set of equations. So, this set of equations is what we are going to consider with the differential equations in ψ and ψ_q retain, we are not neglecting or replacing $d\psi/dt$ and $d\psi_q/dt$ equal to 0. But, we can always do that, in case we do that we will be converting two of these differential equations to algebraic equations, allowing us to eliminate two of these variables. So, two algebraic equations when we get, we can write ψ and ψ_q in terms

of other states and get rid of them and reduce the number of differential **equation** differential equations.

We in fact did this, when we were doing the numerical simulation, in order to remove the stiffness. When you doing Eigen analysis, there is no compelling reason to remove the stiffness of the system, removing the **the** situation, where they are fast and slow transient that is what I mean, when I say removing stiffness. When you are doing Eigen analysis, there is no compelling reason whereas, for simulation if you want to use simple numerical integration methods, then you have to remove the numeric the stiffness in the equations.

So, let us now do an Eigen analysis of the system, I have written down a program again in (()) lab for doing the Eigen analysis.

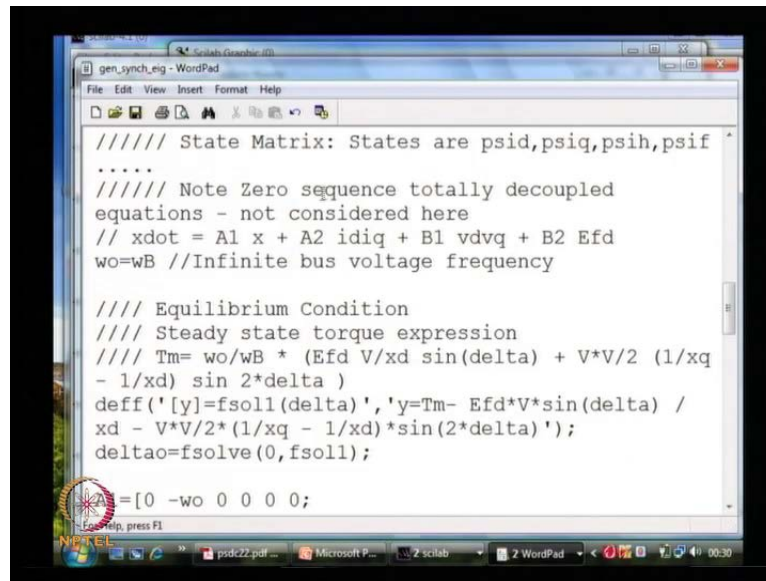
(Refer Slide Time: 38:09)



```
gen_synch_eig - WordPad
File Edit View Insert Format Help
xdd=0.3
xddd=0.23
xq=1.7
xqd=0.65
xqdd=0.25
Td0d=5
Td0dd=0.03 I
Tq0d=1
Tq0dd=0.07
H=3;
Tm=1;
Efd=1.0;
V=1.0;
// To determine Tqd and Tqdd from Tq0d and Tq0dd
a=(1- xq./xqd + xq./xqdd); // To determine Tdd
d Tddd from Td0d and Td0dd
Help, press F1
```

We will just run through it. The important differences between a Eigen value analysis program and a simulation program is that of course, we are talking in terms of a system operating initially at an equilibrium. So for example, if we **look** look at the equilibrium T m is equal to 0 and E f d is equal to one. From this, so I have saved this. From this, we could we would need to calculate the equilibrium value of the state corresponding to speed.

(Refer Slide Time: 38:40)



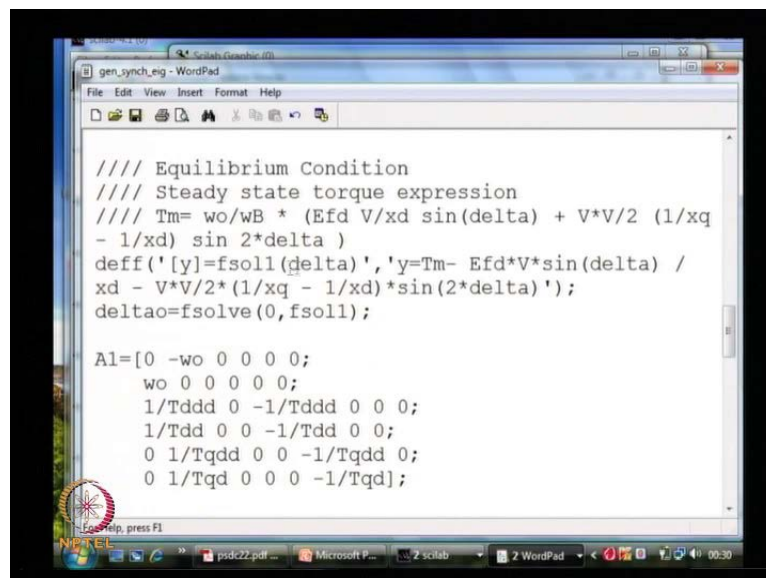
```
gen_synch_eig - WordPad
File Edit View Insert Format Help
///// State Matrix: States are psid,psiq,psih,psif
.....
///// Note Zero sequence totally decoupled
equations - not considered here
// xdot = A1 x + A2 idiq + B1 vdvq + B2 Efd
wo=wB //Infinite bus voltage frequency

///// Equilibrium Condition
///// Steady state torque expression
///// Tm= wo/wB * (Efd V/xd sin(delta) + V*V/2 (1/xq
- 1/xd) sin 2*delta )
deff('y)=fsoll(delta)', 'y=Tm- Efd*V*sin(delta) /
xd - V*V/2*(1/xq - 1/xd)*sin(2*delta)');
deltao=fsolve(0,fsoll);

A1=[0 -wo 0 0 0 0;
wo 0 0 0 0 0;
1/Tddd 0 -1/Tddd 0 0 0;
1/Tdd 0 0 -1/Tdd 0 0;
0 1/Tqdd 0 0 -1/Tqdd 0;
0 1/Tqd 0 0 0 -1/Tqd];
```

But, we know that ω_e is equal to ω_{naught} , which is equal to the infinite bus frequency. That is of course, the equilibrium value of the speed.

(Refer Slide Time: 38:53)



```
gen_synch_eig - WordPad
File Edit View Insert Format Help

///// Equilibrium Condition
///// Steady state torque expression
///// Tm= wo/wB * (Efd V/xd sin(delta) + V*V/2 (1/xq
- 1/xd) sin 2*delta )
deff('y)=fsoll(delta)', 'y=Tm- Efd*V*sin(delta) /
xd - V*V/2*(1/xq - 1/xd)*sin(2*delta)');
deltao=fsolve(0,fsoll);

A1=[0 -wo 0 0 0 0;
wo 0 0 0 0 0;
1/Tddd 0 -1/Tddd 0 0 0;
1/Tdd 0 0 -1/Tdd 0 0;
0 1/Tqdd 0 0 -1/Tqdd 0;
0 1/Tqd 0 0 0 -1/Tqd];
```

The equilibrium value of δ has to be obtained by solving this equation. ω_{naught} and ω_B will of course, assume to be the same, the speed of the infinite bus we will assume to be equal to the base value of the frequency.

So, we can solve for δ . \Once we solve for δ we can also using the algebraic equations **using the algebraic equations** solve for fluxes and once we solve for fluxes, we

can get the equilibrium values. Once we get the equilibrium values, we have to form the A matrixes from an linearized equations.

(Refer Slide Time: 39:37)

```

gen_synch_eig - WordPad
File Edit View Insert Format Help
Am=[0 1;0 0];
Bmpsi=-wB/2/H*[0 0 0 0 0 0;iss(2) -iss(1) 0 0 0 0];
Bmi=-wB/2/H*[ 0 0; -xss(2) xss(1)];
Bm=Bmpsi+Bmi*A3;

Bpsi=[wB*V*cos(deltao) -xss(2);wB*V*sin(deltao) xss
(1); 0 0; 0 0; 0 0; 0 0];

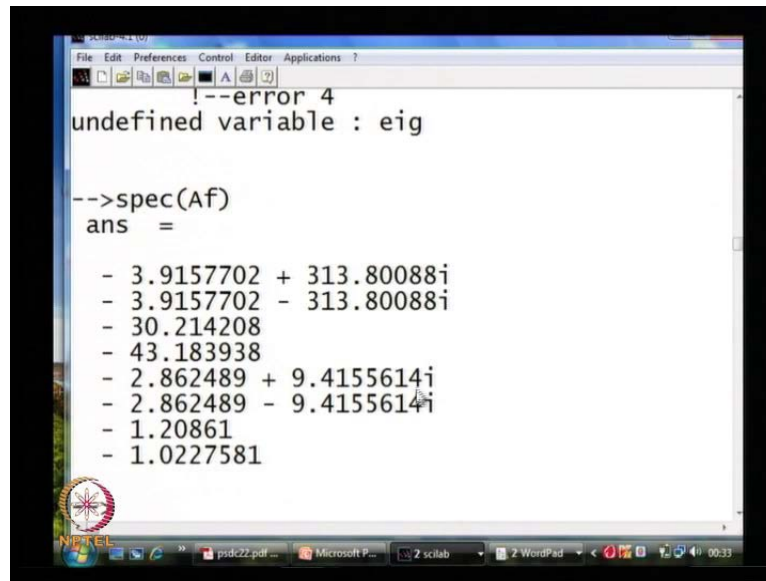
Af=[A Bpsi;Bm Am];

//// Stator transients neglected
Afmod=Af(3:8,3:8)-Af(3:8,1:2)*(Af(1:2,1:2)\Af
(1:2,3:8));
  
```

So for example, if you look at this equation, let me just in this equation, you realize that the components of the A matrix eventually, what I call is A f is consisting of the equilibrium values δ_0 . So, remember that linearized matrixes are a function of the states, the equilibrium values of the states. So, what we will do is, now run this Eigen analysis program. I encourage you to try it out yourself, you try to write a program to do the linearized analysis. So, I will just run it for you here, (no audio 40:21 to 40:39). just one second. Remember that, we are doing the Eigen analysis for system around T_m is equal to 0 and E_{fd} is equal to 1.

So, one thing what we will do is, what is the equilibrium value of δ_0 , the answer is 0. Of course, if mechanical power is 0 T_m is equal to 0, you can directly infer from the steady state torque equation that δ_0 is equal to 0. Now, if I find out the Eigen values of the system, the Eigen values of a system linearized around the equilibrium point corresponding to T_m is equal to 0.

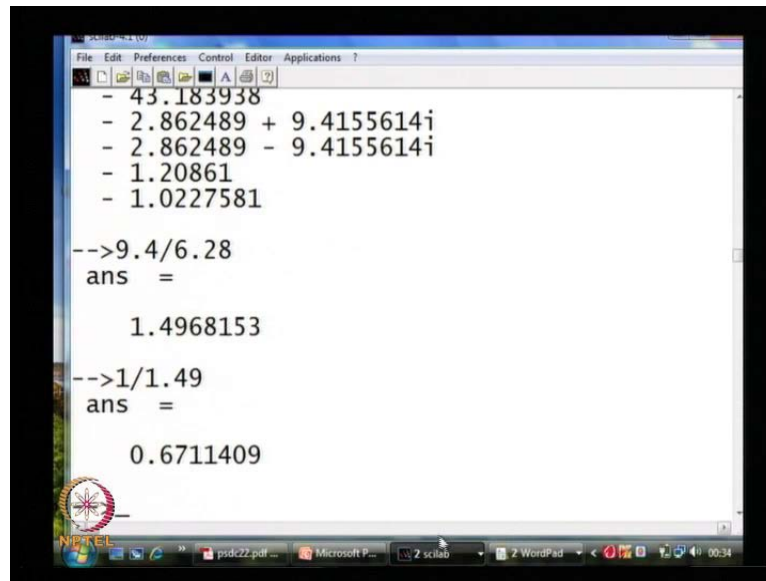
(Refer Slide Time: 41:46)

A screenshot of a MATLAB command window. The window title is 'MATLAB'. The command prompt shows an error: '!--error 4' and 'undefined variable : eig'. Below this, the user enters '-->spec(Af)' and the output is 'ans ='. The output consists of eight eigenvalues: two complex conjugate pairs with large imaginary parts (313.80088i and -313.80088i), and six real eigenvalues. The real eigenvalues are 30.214208, 43.183938, 2.862489, 1.20861, and 1.0227581. The complex eigenvalues are -3.9157702 + 313.80088i and -3.9157702 - 313.80088i. The MATLAB logo and 'NETEL' watermark are visible in the bottom left corner. The taskbar at the bottom shows several open applications including 'psd22.pdf...', 'Microsoft P...', '2 sclab', and '2 WordPad'. The system clock shows '00:33'.

What we find is **what we find is**, the Eigen values are as seen here. What we see is? You have got these two Eigen values, which are very large in magnitude. They almost 314 radians per second, as discussed in the short circuit analysis of **of** a synchronous machine, these are actually corresponding to the d and q states. You will find that these two sort of new Eigen values are coming, because of the inclusion of the electro mechanical equations. The mechanical equations contribute the equations corresponding to ω and δ .

So, basically you get these two extra Eigen values compared to the Eigen value analysis of a short circuit synchronous machine. These two extra Eigen values in fact, indicate that you should be having a damped oscillation of frequency 9.4 radians per second. So, 9.4 radians per second corresponds to divided by 2π , corresponds to 1.49 hertz.

(Refer Slide Time: 43:04)



```
File Edit Preferences Control Editor Applications ?
- 43.183938
- 2.862489 + 9.4155614i
- 2.862489 - 9.4155614i
- 1.20861
- 1.0227581

-->9.4/6.28
ans =
1.4968153

-->1/1.49
ans =
0.6711409
```

So, one upon 1.49 hertz is in fact 0.67 seconds, which is the period of the oscillation corresponding to T_m is equal to 0. So, if you look at this again, this frequency out here should be around 0.7, that is what Eigen value analysis predicts. So, is it actually true, look at it again **yeah**, this is around 1.56 and this is roughly 2.4. So, that is around near about 0.7 to 0.8 second.

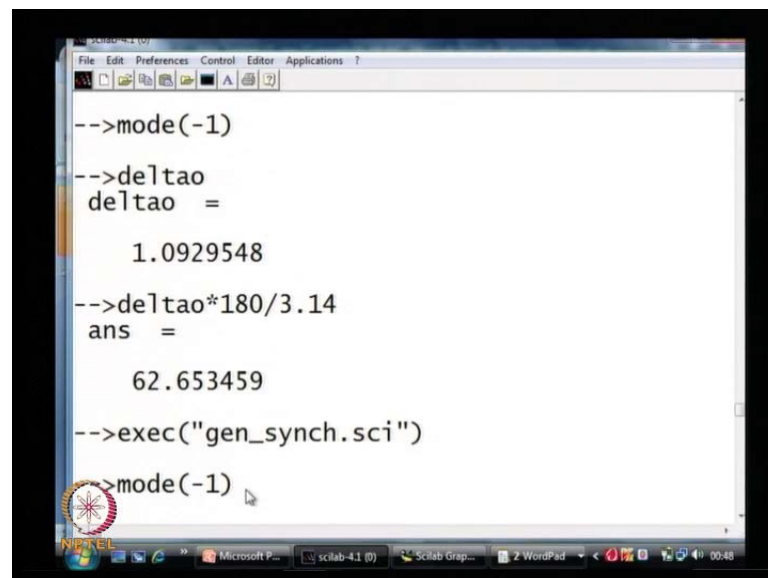
So, actually what we are getting through Eigen analysis is in fact, quite accurate. Also, you can see the rate of change of the peaks, rather you can see that the peak value of this oscillation keeps decaying with time. I leave this as an exercise to you to verify that the real part of this Eigen value corresponds to the rate of decay, which is seen in this figure. We redo the Eigen analysis with Efd is equal to 2 and T_m is equal to 1, so we save this and **rewrite** redo it. **Oops** So, what we need to do is of course, **yeah**. So, if you take out the Eigen values, you will find that the Eigen values you the most striking thing you will notice is that the frequency of oscillations has increased, the damping has come down. This real part, negative real part magnitude has come down. So, what one **one** can get from the simulation essentially is what one should see in the simulation is that the damping is this is of course, the plot of delta, you see that the damping is much faster near the T_m is equal to 0 case, as compared to the damping near T_m is equal to 1.

So, this is the plot of delta remember. The frequency of oscillations here is higher than the frequency of oscillations here. So of course, we can also see this in omega, the **the**

simulation of omega so, I just redo this and plot **and plot** omega. So, I just remove this and plot omega yeah. So, the same thing is of course, seen omega, remember the same modes are been seen in different all **all** almost all the states. So, what you see in delta are the same model characteristics are seen in omega. So, you see that the frequencies at different equilibria are different and the damping also is **is** different. I mean this takes longer to damp out compared to the damping here.

So, the real part of the Eigen value near T_m is equal to 0, can be expected to be larger than the real part, negative real part of course, the magnitude is lower in this. This is actually true. So, this is basically a good correlation between Eigen analysis and the simulation.

(Refer Slide Time: 46:36)



```
-->mode(-1)
-->deltao
deltao =
    1.0929548
-->deltao*180/3.14
ans =
    62.653459
-->exec("gen_synch.sci")
>mode(-1)
```

So, this is what we see here, in fact if you look at the value of delta 0 here, from the Eigen analysis **sorry** the computation of I have to redo the Eigen analysis program. So, the value of delta, equilibrium value of delta, when T_m is equal to 0 and E_{fd} is equal to 1, is 1.09 radiant. So, that actually we can convert this radiant value to a degree value, in fact it is 62 degrees. So, in fact if you look at the simulation, we will redo the simulation and plot the value of delta.

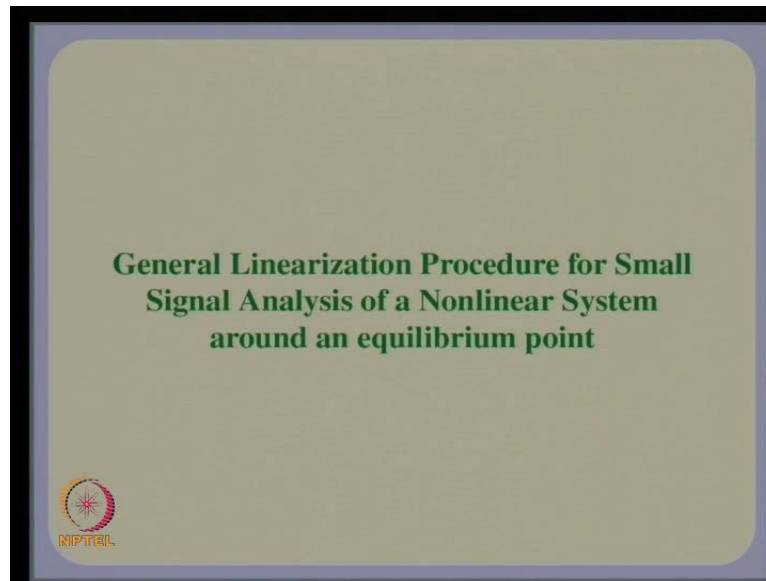
So, what we will do is, this of course, omega we will plot the value of delta with T_m is equal to, so if you look at the simulation, the simulation also if you look at the value of delta which, to with this settles down is also near 1.09. So, what you are really seeing is a

good correlation between the numerical simulation as well as the Eigen analysis of a system linearized around an equilibrium point. So, this is an interesting study in which, we have actually done the linear analysis also of the synchronization transients.

So, I hope you got the hang of **what was** what we are trying to do. We could really, it is very interesting, you know when you do a simulation, you get a certain time response plot but, you feel a great deal of joy, when you are able to correlated it with Eigen value analysis, which one many of us feel that is the more analytically and it gives better insight. So, it is important whenever you do a numerical simulation, whenever you study a transient, it is very important to interpret the results, which are coming correctly and you should be able to correlate it with the kind of inferences you get from other tools.

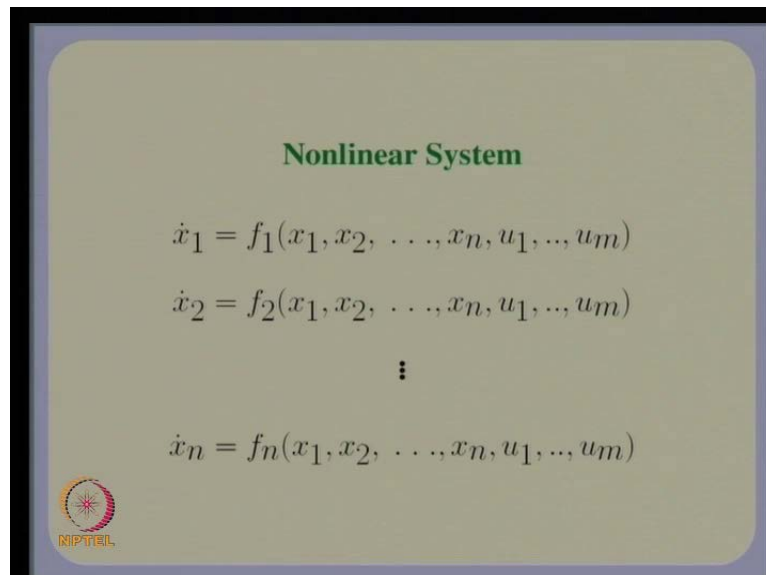
So, I hope you got a flavor of this, in this lecture. In the next lecture, we will try to do a simplified, I will **I will** try introduce you to some simplified models of a synchronous machine. In fact, we have used what is probably a full blown model in all our analysis so far. We will try to do the reverse thing, what we will do is try to make more and more simplifications and come down to a bear bound model or rather the classical model with which, we had actually done some very simple studies right at the beginning of our course. So, we will really know by looking at what kind of simplifications we have done. We can arrive at the toy model, which we considered right at the beginning of the course, that itself is a nice exercise, where we come to know what all we have actually brushed under the carpet in order to get a simple toy model.

(Refer Slide Time: 49:37)



Before we move on, we will discuss again the general linearization procedure for small signal analysis of a non-linear system around an equilibrium point. So, although I have done the linearization of this system, it would be good for you to be familiar and **and** at e is with the general linearization procedure.

(Refer Slide Time: 50:02)



So, if you **if you** have say a non-linear system of this kind, \dot{x}_1 is nothing but, $\frac{dx_1}{dt}$ and you have got n variables of this kind, then the general non-linear equations are given as shown in this screen on the screen.

(Refer Slide Time: 50:27)

Equilibrium Point(s)

$$0 = f_1(x_{1e}, x_{2e}, \dots, x_{ne}, u_{10}, \dots, u_{m0})$$
$$0 = f_2(x_{1e}, x_{2e}, \dots, x_{ne}, u_{10}, \dots, u_{m0})$$
$$\vdots$$
$$0 = f_n(x_{1e}, x_{2e}, \dots, x_{ne}, u_{10}, \dots, u_{m0})$$

u_{10}, \dots, u_{m0} are specified.
More than one set of $x_{1e}, x_{2e}, \dots, x_{ne}$ may satisfy the above equations. Choose the set which is of interest.

NIPTE

Now, the equilibrium points for this system are obtained by setting x_1 dot x_2 dot and so on equal to 0. Of course, in this non-linear system, u denotes the inputs. So, if I given the inputs that is $u_1 0 u_n 0$ are specified, those are the nominal inputs which are given for this equilibrium point, then the equilibrium values of the states may be obtained from these algebraic equation, these again are non-linear algebraic equations. The equilibrium points are obtained by solving these non-linear algebraic equations. Remember that, more than one set of equilibrium states may satisfy the above the equations. So, you have may have more than one set.


Now, in the general linearization procedure, we will try to or we will choose the set, which is of interest. I mean, you are really going to do small signal analysis around an equilibrium point. So, you have to decide the equilibrium points. So, choose the set which is of interest.

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Small Signal Model Development

$$\Delta \dot{x}_1 = \frac{\partial f_1}{\partial x_1} \Delta x_1 + \frac{\partial f_1}{\partial x_2} \Delta x_2 \dots + \frac{\partial f_1}{\partial x_n} \Delta x_n + \frac{\partial f_1}{\partial u_1} \Delta u_1 \dots + \frac{\partial f_1}{\partial u_m} \Delta u_m$$

Partial derivatives **are evaluated** using the equilibrium values of states, $x_{1e}, x_{2e}, \dots, x_{ne}$
 $\Delta x_1 = x_1 - x_{1e}$ denotes the deviation from the equilibrium.
 $\Delta u_1 = u_1 - u_{10}$.




After that, you need to take each differential equation, each non-linear differential equation and linearized it. So, what this involves I mean, what **what** eventually you need to do is, obtain the partial derivatives of f_1 with respect to the states x_1 to x_n as well as the inputs u_1 to u_m . Now, remember that these partial derivatives are evaluated using equilibrium values of the states. So, these partial derivatives are essentially constant.

Remember of course, that Δx_1 here, is a small disturbance from equilibrium. So, Δx_1 is equal to $x_1 - x_{1e}$ or the deviation from the equilibrium. Similarly, Δu_1 is equal to $u_1 - u_{10}$ where u_{10} is a nominal input. So, this is what you get as a linearized equation. Now, in my previous development, we **we** just for example, linearized $\dot{\psi} - \psi$. You may say that well this looks different from that. No, actually if you take out a partial derivatives **of the** function $\dot{\psi} - \psi$ and plug in the equilibrium values you are going to get exactly this same linear model as we got some time back.

So, in fact this is the more general and direct representation of what we will get, if we apply this linearization with Δx and Δu very small. So, this is how you would normally linearize a set of non-linear differential equations around an equilibrium point.

(Refer Slide Time: 53:33)

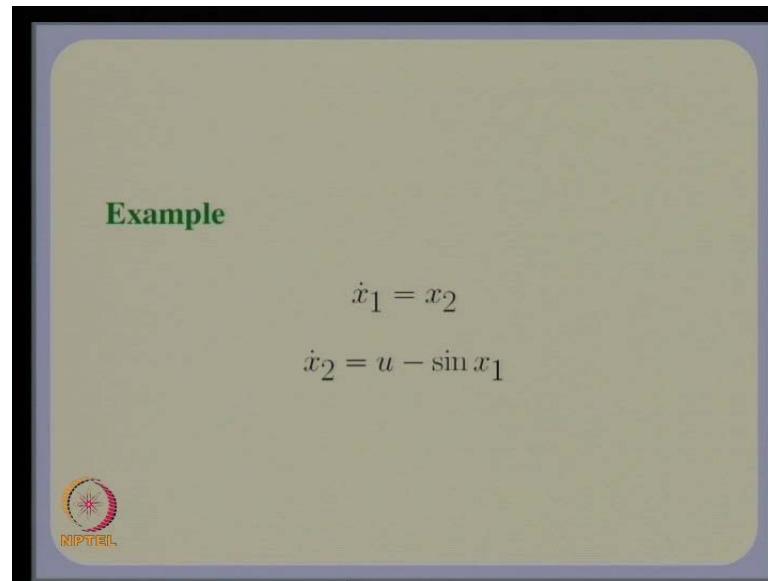
Overall Model

$$\begin{bmatrix} \Delta \dot{x}_1 \\ \vdots \\ \Delta \dot{x}_n \end{bmatrix} = \begin{bmatrix} \frac{\partial f_1}{\partial x_1} & \cdots & \frac{\partial f_1}{\partial x_n} \\ \vdots & \ddots & \vdots \\ \frac{\partial f_n}{\partial x_1} & \cdots & \frac{\partial f_n}{\partial x_n} \end{bmatrix} \begin{bmatrix} \Delta x_1 \\ \vdots \\ \Delta x_n \end{bmatrix} + \begin{bmatrix} \frac{\partial f_1}{\partial u_1} & \cdots & \frac{\partial f_1}{\partial u_m} \\ \vdots & \ddots & \vdots \\ \frac{\partial f_n}{\partial u_1} & \cdots & \frac{\partial f_n}{\partial u_m} \end{bmatrix} \begin{bmatrix} \Delta u_1 \\ \vdots \\ \Delta u_m \end{bmatrix}$$


So, this fact will become cleared shortly remember of course, that you have to linearized each differential equation right from x_1 to x_n and therefore, eventually you will get in fact a Jacobean matrix. So, you know you do this for x_1 but, you can also do it for the other x_2 x_3 to x_n and eventually you will get this Δx is equal to A into Δx plus B into Δu . This A and B of course, are these matrixes.

Note that, these are effectively constant matrixes because; these partial derivatives are evaluated at the equilibrium points. So, equilibrium point essentially is **the** denotes the values of the states. So, you have to plug in the values of the states. Now, this will become clear.

(Refer Slide Time: 54:31)



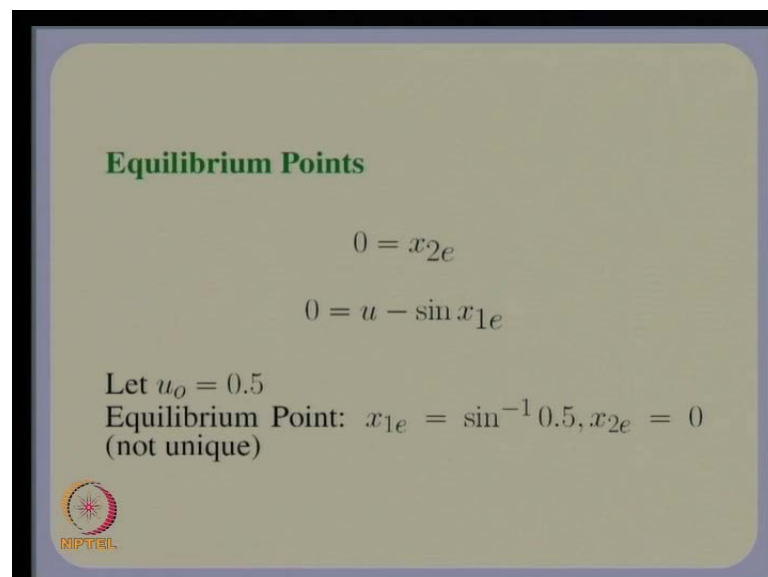
Example

$$\dot{x}_1 = x_2$$
$$\dot{x}_2 = u - \sin x_1$$

NIPTTEL

If you consider this example, \dot{x}_1 is equal to x_2 and \dot{x}_2 is equal to $u - \sin x_1$. So, this is a non-linear set of equations, the nonlinearity obviously is because of the \sin term.

(Refer Slide Time: 54:49)



Equilibrium Points

$$0 = x_{2e}$$
$$0 = u - \sin x_{1e}$$

Let $u_0 = 0.5$
Equilibrium Point: $x_{1e} = \sin^{-1} 0.5, x_{2e} = 0$
(not unique)

NIPTTEL

So, the first step in doing this is, find the equilibrium points for the nominal input 0.5, the equilibrium point is $\sin^{-1} 0.5$ or \sin^{-1} half, which is 30 degrees and x_{2e} is equal to 0. So, x_{1e} is in fact 30 degrees this is of course, the equilibrium points are not

unique. For example, $x_1 e$ could also be pi minus that is 180 minus 30 degrees. That is also in equilibrium point.

So, remember that there is in the non-linear systems, they may be more than one equilibrium point. Let us say the equilibrium point corresponding to **sin inverse** sin inverse 0.5 is equal to 30 degrees and $x_2 e$ is equal to 0, is the equilibrium point of our interest. So, equilibrium point of our interest is 30 degrees and 0. Let us say, suppose this is our equilibrium point of interest.

(Refer Slide Time: 55:44)

Small Signal Model around
 $(x_{1e} = 30^\circ, x_{2e} = 0, u_0 = 0.5)$

$$\begin{bmatrix} \Delta \dot{x}_1 \\ \Delta \dot{x}_2 \end{bmatrix} = \begin{bmatrix} 0 & 1 \\ -\cos 30^\circ & 0 \end{bmatrix} \begin{bmatrix} \Delta x_1 \\ \Delta x_2 \end{bmatrix} + \begin{bmatrix} 0 \\ 1 \end{bmatrix} \Delta u$$

So, you can easily obtain the small signal model, this small signal model is essentially a linear model around this equilibrium point. So, this is eventually the small signal model of the system. One could also in fact, try to find out the equilibrium point; you know the linear model around another equilibrium point. So for, in this case for example, the other equilibrium point is $x_1 e$ is equal to 150 degrees, $x_2 e$ is equal to 0 and u_0 of course, is the nominal value 0.5. And interestingly, the properties of the **the** small signal stability properties of different equilibrium points can be different. So, that is the very interesting and important thing, which you should keep in mind.

So, this is the general procedure, when you want to linearized a non-linear system around an equilibrium point. So, you can directly take partial derivatives and from the Jacobean matrix. Now, I wish to point out again that, in the example we have considered and where I have shown you how to linearized particular differential equation or this you

know the **the** torque equation effectively. Have linearized the electrical torque but, it may not be directly evident that you are in fact you are doing exactly this. In fact, I have in fact taken out the partial derivative of the expression of t_e with respect to ψ_d , ψ_q , i_d and i_q and plugged in the equilibrium value.

So this is but, a more **general way of be** general way of putting it. That is you take out the partial derivatives or evaluate this Jacobean at the equilibrium values of the states. So, eventually what you get is a linear model of a system with a constant coefficient matrixes and another point to be noted is of course, that is the multiple equilibrium points, you could have different small signal models around different equilibrium points and of course, **the properties** the stability properties of the A matrix **a matrixes** are around different operating points are equilibrium points could be different. So, this is something you should keep in mind. Now, we go back to our example again of a synchronous machine.