

Finite Element Method and Computational Structural Dynamics
Prof. Manish Shrikhande
Department of Earthquake Engineering
Indian Institute of Technology, Roorkee

Lecture - 30
Finite Elements of C⁰ Continuity in 2-D and 3-D- XII

Hello friends. So, we were discussing about solution of a bracket problem, plane stress problem, steel bracket problem using Finite Elements. And we choose to discretize the domain by using 4 elements of 3 elements of 4 node rectangle and 1 element of 3 node triangle.

And we saw how the mesh was discretized, and the nodes, and what is the mapping between the local node numbering and global node numbering within the element, and how the displacements are interpolated once the coordinates are defined the of the nodes. The interpolation functions can be derived. And once the interpolation functions are available, they can be related to the derivatives.

(Refer Slide Time: 01:19)

Finite Elements of Two and Three Dimensions
○○○○○○○
Approximations for 2-D Elastic Continuum
○○○○○○○
Finite Elements for Two-Dimensional Domains
○○○○○○○○○○○○○○○○○○○○○○○○○○○○○○●○○○○○○○○○○

Example Problem-5

- ▶ Approximation of the displacement field within the rectangular element:
 $u^{(e)} = Nu_e$
- ▶ The strain field within element: $\epsilon^{(e)} = \mathcal{L}Nu_e = Bu_e$:

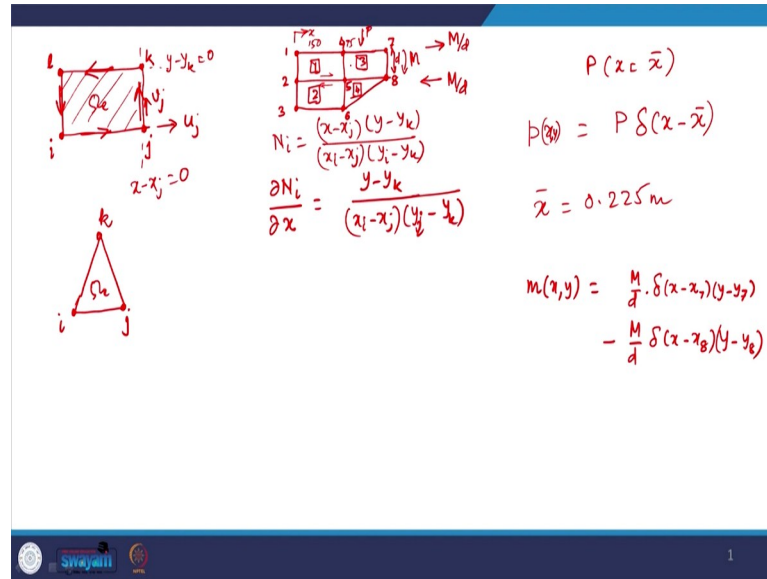
$$\begin{pmatrix} \epsilon_{xx} \\ \epsilon_{yy} \\ \gamma_{xy} \end{pmatrix}^{(e)} = \begin{bmatrix} \frac{\partial}{\partial x} & 0 \\ 0 & \frac{\partial}{\partial y} \\ \frac{\partial}{\partial y} & \frac{\partial}{\partial x} \end{bmatrix} \begin{bmatrix} N_i & 0 & N_j & 0 & N_k & 0 & N_l & 0 \\ 0 & N_i & 0 & N_j & 0 & N_k & 0 & N_l \end{bmatrix} u_e$$

$$= \begin{bmatrix} \frac{\partial N_i}{\partial x} & 0 & \frac{\partial N_j}{\partial x} & 0 & \frac{\partial N_k}{\partial x} & 0 & \frac{\partial N_l}{\partial x} & 0 \\ 0 & \frac{\partial N_i}{\partial y} & 0 & \frac{\partial N_j}{\partial y} & 0 & \frac{\partial N_k}{\partial y} & 0 & \frac{\partial N_l}{\partial y} \\ \frac{\partial N_i}{\partial y} & \frac{\partial N_i}{\partial x} & \frac{\partial N_j}{\partial y} & \frac{\partial N_j}{\partial x} & \frac{\partial N_k}{\partial y} & \frac{\partial N_k}{\partial x} & \frac{\partial N_l}{\partial y} & \frac{\partial N_l}{\partial x} \end{bmatrix} u_e$$

Manish Shrikhande manish.shrikhande@eeq.iitr.ac.in
Finite Element Method and Computational Structural Dynamics
Department of Earthquake Engineering Indian Institute of Technology Roorkee

So, the strain field and those derivatives the strain derivatives that these are essentially the derivatives of partial derivatives of interpolation functions. Now, these partial derivatives are completely defined when we specify the coordinates of the nodes. So, for example, N_i is defined by as $x - x_j$ multiplied by $y - y_k$ divided by $x_i - x_j$ by $y_i - y_k$.

(Refer Slide Time: 02:04)



Now, the first derivative is of course, a function of y and the denominator. So, just to tell you, so $\frac{\partial N_i}{\partial x} = \frac{y - y_k}{(x_i - x_j)(y_i - y_k)}$. So, it is linear in y and similarly all derivative functions. So, N_i is essentially $(x - x_j)(y - y_k)$ divided by numerator normalized at the node $(x_i - x_j), (y_i - y_k)$. So, derivative is simply this remaining expression.

This is completely defined the moment we have the coordinates of all these nodes in the mesh. So, once we map these local node class identifiers i, j, k , and l , with the physical actual node in the global mesh, then these derivatives can be evaluated readily by using these coordinates.

(Refer Slide Time: 03:26)

[illegible]

And that is how the terms, respective terms, derivative terms in the domain integral these are evaluated. So, $\frac{\partial u}{\partial x}$ will be defined by similar shape function derivatives, and once, because those are already defined in terms of nodal coordinates for each node, for each element the. And this is of course, the waiting function that will keep on changing W can be either N, Ni, Nj, Nk, Nl in succession.

So, again the derivative is defined in terms of the coordinate functions, coordinate variables, nodal coordinates. And this expression then can be evaluated with respect to x and y . So, domain integral and the this entire expression will be known as the stiffness matrix of the problem.

So, this is the approximation within the element of any rectangular element, or for that matter even for triangular element, the only difference is the number of variables nodal variables and the size of this matrix N . Other than that, there is no difference between the this representation of displacement field. Of course, this expression of interpolation function will be different for in the case of rectangle and for from the a triangular element.

(Refer Slide Time: 05:05)

Finite Elements of Two and Three Dimensions
○○○○○○
Approximations for 2-D Elastic Continuum
○○○○○○○
Finite Elements for Two-Dimensional Domains
○○○○○○○○○○○○○○○○○○○○○○○○○○○○○○●○○○○○○○○○○

Example Problem-6

- ▶ The interpolation functions for a triangular domain may be given as:

$$N_i = \frac{1}{2\Delta} [(x_j y_k - x_k y_j) + (y_j - y_k)x + (x_k - x_j)y]$$

$$N_j = \frac{1}{2\Delta} [(x_k y_i - x_i y_k) + (y_k - y_i)x + (x_i - x_k)y]$$

$$N_k = \frac{1}{2\Delta} [(x_i y_j - x_j y_i) + (y_i - y_j)x + (x_j - x_i)y]$$
- where, Δ denotes the area of the triangular domain.
- ▶ Approximation of the displacement field within the triangular element: $u^{(e)} = \mathbb{N}u_e$
- ▶ The components of strain field can be obtained as $\epsilon^{(e)} = \mathcal{L}\mathbb{N}u_e = \mathbb{B}u_e$, which are constant within the element—hence the name: “Constant Strain Triangle”.

Manish Shrikhande manish.shrikhande@eeq.iitr.ac.in
Finite Element Method and Computational Structural Dynamics
Department of Earthquake Engineering Indian Institute of Technology Roorkee

So, for a triangular domain the interpolation functions can again be given in terms of nodal coordinates, as we had seen during this during our discussion of triangular family. So, N_i the interpolation function for i th node that we have, this is the i th node location in the local node numbering for element.

So, this is defined in terms of coordinates of node j and node k , as $(x_j y_k) - (x_k y_j)$, $+$, $(y_j - y_k) x$. So, this is the coefficient of x and $x_k - x_j$ multiplied by y . So, that is the coefficient of y and the entire thing normalized by twice the area of triangle. So, this is the twice delta is the determinant of the coordinates $1 \ x_i \ y_i, 1 \ x_j \ y_j, 1 \ x_k \ y_k$. So, that determinant of 3 by 3 matrix is the; is equal to twice the area of the triangle enclosed by these 3 nodes.

So, similarly we have by cyclic permutation, we can had derive the interpolation for node j and node k , and you can see wherever we have j for the i th node interpolation that is replaced by k for j 's node interpolation and so on. So, cyclic symmetry holds in this case.

So, the approximation of the displacement field within triangular element is again similarly given by shape function matrix multiplied by the displacement, nodal displacement. The interpolation, the strain field is again same differential operator, operating on the interpolation function. And because this is all of them are first degree polynomials.

So, the derivatives are constants. So, this strain field is constant over the entire element. So, ϵ_{xx} is constant over the entire element, ϵ_{yy} is constant over the entire element, and similarly γ_{xy} is also constant over the entire element. And hence the name constant strain triangle or CST that is often used for in place of 3 node triangle element.

(Refer Slide Time: 07:39)

Finite Elements of Two and Three Dimensions
○○○○○
Approximations for 2-D Elastic Continuum
○○○○○
Finite Elements for Two-Dimensional Domains
○○○○○○○○○○○○○○○○○○○○○○○○○○○○○○●○○○○○○○○

Example Problem-7

- ▶ Let us now establish the element level equilibrium equations for each element of the finite element mesh. Since the material properties are uniform for all elements, the constitutive relation matrix D for the plane stress condition may be given as:

$$D = \frac{E}{1-\nu^2} \begin{bmatrix} 1 & \nu & 0 \\ \nu & 1 & 0 \\ 0 & 0 & \frac{1-\nu}{2} \end{bmatrix} = 2 \times 10^{11} \begin{bmatrix} 1.0989 & 0.3297 & 0 \\ 0.3297 & 1.0989 & 0 \\ 0 & 0 & 0.3846 \end{bmatrix} \text{ N/m}^2$$
- ▶ Let us disregard the inertia term involving mass and acceleration as the applied loads are static and therefore dynamic effects can be neglected.

Manish Shrikhande manish.shrikhande@eeq.iitr.ac.in
Finite Element Method and Computational Structural Dynamics
Department of Earthquake Engineering Indian Institute of Technology Roorkee

So, in order to establish the element level equilibrium equations, we need the constitutive matrix D for plane stress conditions, what are referred to as coefficients d_{11} , d_{12} , d_{33} , etcetera. So, these are the coefficients of those constitutive matrix.

So, these coefficients are of course, going to be different for plane stress conditions and for plane strain conditions. But for the once we idealize, once we consider, once we reassure ourselves that the problem that is being dealt with is consistent with the assumptions of plane stress conditions, then the constitutive matrix for plane stress condition which is given by this is can be derived in terms of the element properties that is Young's modulus and Poisson's ratio. And this is the constitutive matrix for plane stress conditions.

We will also in this particular case because the loading conditions, they are constant, they are not changing with time. So, there are inertia terms are not significant and we will drop these acceleration terms, right. So, integral of W times $\rho \ddot{u}$. So, this is the acceleration term. So, these are the inertia terms that define the govern the dynamics of

this system, if the loading, if the system is subjected to rapidly changing loads or time varying loads.

Since, in this particular case we are not looking at time varying loads, so this calculation will not be required, otherwise this can be computed in the same way as other domain, other domain integral. The only thing that is being implicit, that will be implicitly assumed is the where, while in this part in this particular expression, we are looking at approximation of the displacement.

In this particular term, requires consideration of the variation of acceleration, how are we going, how are we approximating variation of acceleration within the element domain, right. So, we will revert to this return to this particular point after discussion of this static problem.

So, we ignore the dynamic effects. So, since the loads are static. So, now let us look at the formulation of element level equation.

(Refer Slide Time: 10:35)

Finite Elements of Two and Three Dimensions
○○○○○○○
Approximations for 2-D Elastic Continuum
○○○○○○○
Finite Elements for Two-Dimensional Domains
○○○○○○○○○○○○○○○○○○○○○○○○○○○○○○●○○○○○○○

Example Problem-8
 The element (sub-domain) # 1 is defined by the nodes (2, 5, 4, 1). The element stiffness matrix is defined as:

$$K_1 = t \iint_{\Omega_1} B^T D B \, d\Omega = t \int_{x_i}^{x_j} \int_{y_i}^{y_k} B^T D B \, dx \, dy$$
 Some representative elements of the 8×8 symmetric stiffness matrix of 4-node rectangular element are:

$$k_{11}^{(1)} = \frac{Et}{1-\nu^2} \iint_{\Omega_1} \left[\left(\frac{\partial N_i}{\partial x} \right)^2 + \left(\frac{1-\nu}{2} \right) \left(\frac{\partial N_i}{\partial y} \right)^2 \right] d\Omega$$

$$k_{12}^{(1)} = \frac{Et}{1-\nu^2} \iint_{\Omega_1} \left[\nu \frac{\partial N_i}{\partial x} \frac{\partial N_i}{\partial y} + \left(\frac{1-\nu}{2} \right) \frac{\partial N_i}{\partial y} \frac{\partial N_i}{\partial x} \right] d\Omega$$

$$k_{13}^{(1)} = \frac{Et}{1-\nu^2} \iint_{\Omega_1} \left[\frac{\partial N_i}{\partial x} \frac{\partial N_j}{\partial x} + \left(\frac{1-\nu}{2} \right) \frac{\partial N_i}{\partial y} \frac{\partial N_j}{\partial y} \right] d\Omega$$

Manish Shrikhande manish.shrikhande@eeq.iitr.ac.in
Finite Element Method and Computational Structural Dynamics
Department of Earthquake Engineering Indian Institute of Technology Roorkee

We take element number 1 defined by 2, 5, 4, 1; 2, 5, 4, 1 as the element of our discussion. Process is exactly identical for all the elements. So, we will only refer to this in this discussion. Step by step process for all the elements and in going through the entire every single step, you can see in the solved example in the book that we have my book on finite element and computational structural dynamics.

So, I will not repeat the entire, I will not go through the entire series of calculations for all the elements in this discussion for paucity of time, because of paucity of time. So, the element stiffness matrix can be defined, I mean the all the equations that we have here, if I collect all these terms arrange it in matrix form, I will end up with equation of this type.

So, B^T transpose that is again derivative of the interpolation function. If you recall, strain displacement matrix involves derivative of the interpolation function and since we need to multiply, our weighting function is δv shape function by derivative of the weighting function multiplies with all of these terms. So, that is related to the strain displacement matrix again for Galerkin formulation and that is what we use.

So, essentially it boils down to expression of this type evaluation of expression of this type $B^T D B$. So, B^T transpose is the strain displacement matrix in the case of borrowing the terminology from solid mechanics, D is the constitutive relation matrix, and B is the again strain displacement matrix, and this is to be integrated over the element domain.

So, element domain is this entire element domain which is bounded by range x ranging from x_i to x_j and y ranging from y_j to y_k , right. So, these are the domains, and this is what Ω ranges from x_i to x_j , y_j to y_k , or y_i to y_k , does not matter. So, y_i and y_j they are identical. So, it really does not matter whether I use y_i or y_j . It is the same coordinate so y_i to y_k .

And within this integral, within this domain these are to be evaluated. In this particular case, the element is the directions are oriented along the are parallel, the element edges are parallel to the coordinate direction. So, these integrals can be evaluated very easily or straight forward manner and eventually the 8 by 8 symmetric matrix terms can be evaluated individually as this. So, these are some of the representative evaluations that come out.

So, this is of course, $E/(1-\nu^2)$ is from the constitutive matrix D , and rest of the terms they are coming from the product of $B^T D B$ individual elements that they come. And this is the superscript, 1 refers to element number 1, and subscripts they refer to the dimension. So, there are total 8 by 8 matrix, so 64 elements. So, the k_{11} refers to element in first row and first column for the element stiffness matrix.

So, that is what is evaluated it simply $\int_{\Omega_i} N^T f d\Omega$, x component of the body force and evaluated over the in domain integral domain interval.

(Refer Slide Time: 17:18)

Finite Elements of Two and Three Dimensions
○○○○○○○
Approximations for 2-D Elastic Continuum
○○○○○○○
Finite Elements for Two-Dimensional Domains
○○○○○○○○○○○○○○○○○○○○○○○○○○○○○○●○○○○○

Example Problem-10

The equivalent nodal loads due to body forces:

$$f_1 = t \int_{\Omega_i} N^T f d\Omega$$

$f = [0, -\rho g]^T$: the vector of body forces per unit volume in each coordinate direction, and ρg is the weight density.

$$\begin{aligned}
 f_1 &= t \int_{\Omega_i} \begin{bmatrix} N_i & 0 & N_j & 0 & N_k & 0 & N_l & 0 \\ 0 & N_i & 0 & N_j & 0 & N_k & 0 & N_l \end{bmatrix}^T \begin{pmatrix} 0 \\ -\rho g \end{pmatrix} d\Omega \\
 &= (x_j - x_i)(y_k - y_i) \frac{t}{4} [0 \quad -\rho g \quad 0 \quad -\rho g \quad 0 \quad -\rho g \quad 0 \quad -\rho g]^T \\
 &= [0 \quad -1.7658 \quad 0 \quad -1.7658 \quad 0 \quad -1.7658 \quad 0 \quad -1.7658]^T \text{ N}
 \end{aligned}$$

Manish Shrikhande manish.shrikhande@eeq.iitr.ac.in
Finite Element Method and Computational Structural Dynamics
Department of Earthquake Engineering Indian Institute of Technology Roorkee

So, this is what the nodal forces will look like, evaluation of nodal forces. So, waiting function multiplication is represented by $N^T f$. So, f is the vector of nodal forces. In this case, there is no body force in the along the x direction, but there is a body force component along y direction that is self-weight, and that is referred to here as ρ times g per unit volume.

So, ρ is the density and $-$ sign because weight acts vertically downwards and the positive sign is assumed upward as the positive y direction. So, this is the body force component along x and y direction, and when we multiply this with the weighting weight function matrix or the interpolation matrix to represent the weight function multiplication, domain integral.

And then, evaluate it along the domain x_i to x_j and y_i to y_k , we end up with this expressions, that you can check and straightforward manner. And once we substitute the values for all of these coordinates and for element number 1, and thickness and values of ρ and acceleration due to gravity g , we get this vector for the element forces as this in terms of I mean Newtons, so many Newtons.

So, essentially equally distributed at all the nodes. So, 1.76 Newton downward at all the 4 nodes. So, that is the gravity load self-weight. So, that takes care of the boundary domain integrals in this case. So, this particular integral, both of these in both of these equations they are taken care of, except this inertia term which we are dropping because acceleration is deemed to be insignificant or negligible.

(Refer Slide Time: 19:51)

Finite Elements of Two and Three Dimensions
○○○○○○
Approximations for 2-D Elastic Continuum
○○○○○○○
Finite Elements for Two-Dimensional Domains
○○○○○○○○○○○○○○○○○○○○○○○○○○○○○○●○○○○

Example Problem-11

The equivalent nodal loads due to applied tractions (natural (or, Neumann) boundary condition) on the edges of the element:

$$Q_1 = Q_{1\Gamma_{ij}} + Q_{1\Gamma_{jk}} + Q_{1\Gamma_{kl}} + Q_{1\Gamma_{li}}$$

$$= t \int_{\Gamma_{ij}} N^T T^{(1)} d\Gamma + t \int_{\Gamma_{jk}} N^T T^{(1)} d\Gamma + t \int_{\Gamma_{kl}} N^T T^{(1)} d\Gamma + t \int_{\Gamma_{li}} N^T T^{(1)} d\Gamma$$

$T^{(1)} = [T_x \ T_y]^T = [\hat{N}^T \sigma^{(1)}]^T$: vector of tractions along the element boundary, \hat{N} : the matrix of direction cosines of outward unit normal on the boundary. A representative boundary integral evaluation:

$$Q_{1\Gamma_{ij}} = Q_{1\Gamma_{25}} = t \int_{\Gamma_{25}} N^T T^{(1)} d\Gamma$$

$$= \frac{t(x_5 - x_2)}{2} [T_x^{\Gamma_{25}} \ T_y^{\Gamma_{25}} \ T_x^{\Gamma_{25}} \ T_y^{\Gamma_{25}} \ 0 \ 0 \ 0 \ 0]^T$$

Manish Shrikhande manish.shrikhande@eeq.iitr.ac.in
Finite Element Method and Computational Structural Dynamics
Department of Earthquake Engineering Indian Institute of Technology Roorkee

What remains now is this boundaryintegral. How to evaluate these boundaryintegrals? And these boundaryintegrals are again evaluated by segregating the boundaries, looking at each boundary at one time. So, this is gamma ij, the boundaryintegral is evaluated in counter clockwise direction, and in this particular case this boundaryij is common with the boundary for kl for element two. So, the those boundary terms would get cancelled.

Similarly, if this boundary term along this edge will get cancelled with the boundary term, similar boundary term coming from element number 3, and 1, 2, this boundaryit is of course, the essential boundary specified. So, Neumann boundary condition does not apply here. So, there will be reaction. So, though those reactions are unknown. And this is of course traction free.

So, boundary 4, 1, which corresponds to kl which coincides with the problem boundary the original boundary of the original domain. So, we only need to evaluate the boundaryin along this edge 4, 1, right. All other boundaries conditions are not really required to be evaluated, but nevertheless we will explain the process, how it is done.

So, together the rectangular element, there are any rectangular element is defined by 4 boundaries. So, y_{ij} , y_{jk} , y_{kl} , and y_{lk} , and again these are the order I mean the node numbering is the order in which we are moving. So, counter clock wise, i to j, j to k, k to l and from l to i, ok, fine. So, this one should be li not lk. So, that is a typo of here. Pardon for that error.

And this subscript one refers to element number 1. So, for element number 1 the boundary along the edge ij, for element number 1 boundary along jk edge, for element number 1 boundary along kl edge, and so on. So, that would be given by integral. Again, thickness is of course, constant. So, integral along this edge weighting function multiplied by the traction, tractions for node element number 1. Along this edge and similarly for all these 4 terms, right.

And this traction, there are two components, I mean its two-dimensional problem. So, there can be tractions components along x direction and components along y direction. So, T_x or T_x and T_y and they are given by direction cosine, product of the direction cosine with the normal stress and tangential stress along the bound, along the bound element boundary.

So, this is what we will get, vector of tractions along the element boundary, \hat{N} , the matrix of direction cosines of the outward unit normal. In this, for this particular geometry it is simply 1 and 0, and a representative boundary evaluation is what looks like this. So, $\hat{N} T_1$, so this is what T_x in along this edge T_y along edge 25 and again T_x along 25, T_y along 25. So, both of them divided equally between the two nodes 2 and 5, and then multiplied by what is the length and thickness of course remains.

And that completes the evaluation of the one boundary, one segment of the boundary, and similarly we can go for next segment, and third segment, fourth segment, and that is how we do that. And during the assembly, as we saw in the case of one-dimensional element these for the edge of 25 there would be a corresponding matching contribution coming from the element number 2 which would be from element 5 to 2. For element 1, we counter clockwise direction was from 2 to 5.

For element number 2, the correspond this edge would be evaluated from 5 to 2. And accordingly the direction cosines would of course, be outward normal, so they would be

in opposite directions, and the corresponding terms would be of opposite signs and they would cancel out during the process of assembly, right. At the individual element level of course they do not, they appear, but during the process of assembly they would of course cancel out.

And therefore, generally we generally do not bother about calculation going through this calculation, unless this boundary coincides with the problem boundary. For example, this boundary of 4, 1 corresponds with the or boundary of 4, 1 corresponds with the actual problem boundary. So, we will evaluate it at this point. So, that brings us to the problem.

In this case, we have element number 3. We have this point load at the midpoint of this edge P and then there is a moment M . So, how do we model these? How do we account for these? So, P point load can be easily modeled as a Dirac delta function that is a discontinuity.

So, point load P at x is equal to \bar{x} can be represented as function of P x as $P \delta x - \bar{x}$, and then this function P x can be treated as a continuous function of x and substituted in the integral. Integrant will be evaluated as the by using the property of Dirac delta function, and that would result in point evaluation sampling at \bar{x} , and appropriately the calculations will be done.

So, all that we need to do is substitute for T x , the traction as instead of continuous function, this variation of Dirac delta function at \bar{x} whatever the coordinates of \bar{x} are. So, in this particular case it would be 225 because this is where the x is ranging from and this is at $75 + 150$. So, this is 150 and this is 75, so together its 225. So, it would be x is equal to 0.225 meter. So, \bar{x} is equal to 0.225 meter, for the problem that we are dealing with.

And that leaves us with another issue, how do we model the moment because there are no rotational degrees of freedom here, so how do we incorporate all the forces that we have, the allowance is only for the rectilinear forces, force along x direction, force along y direction, how do we model the moment, applied moment in this case, in this problem. So, the way to model moment is that this moment can be represented can be replaced by appropriate couple and this couple is equal to M by d , where d is this distance, right.

So, this constant couple separated over a distance d that defines the moment. So, while analyzing for moment we need to convert it into an equivalent couple that would be of pair of concentrated load, separated by a distance, d . So, in this case d is the depth of this element, and these would then be treated same in the same way as this Dirac delta function for this traction in y direction.

So, in this case, it would be traction along x directions of opposite signs. So, we will have $\delta x - x_7$, multiplied by $y - y_7$, $-\delta x - x_8$, $y - y_8$. So, that is the variation and using this variation it can be evaluated.

So, once we have this traction along x along this particular edge 7, 8, then this expression can be evaluated in exactly the same fashion as this. So, instead of traction, we will have these Dirac delta functions and then the integrals one-dimensional integrals can be evaluated.

(Refer Slide Time: 31:00)

Finite Elements of Two and Three Dimensions
Approximations for 2-D Elastic Continuum
Finite Elements for Two-Dimensional Domains

Example Problem-12

- ▶ The equations of equilibrium for element 1 is established as:
 $K_1 u_1 = f_1 + Q_1$
- ▶ The contributions from individual elements may now be assembled into a global system of equations describing the whole domain.
- ▶ Since a total of 8 nodes with 2 degrees of freedom per node are used in the discretisation, there will be 16 (8 nodes \times 2 DOFs per node) equations to describe the complete domain.
- ▶ The key is to find an appropriate placement in the global template for each coefficient in the element equilibrium equation, which necessitates mapping of local degrees of freedom in an element definition to the global ordering of primary variables.

$$\begin{pmatrix} u_i \\ v_i \\ u_j \\ v_j \\ u_k \\ v_k \\ u_l \\ v_l \end{pmatrix}^{(1)}$$

\longleftrightarrow

$$\begin{pmatrix} u_2 \\ v_2 \\ u_5 \\ v_5 \\ u_4 \\ v_4 \\ u_1 \\ v_1 \end{pmatrix}$$

Manish Shrikhande manish.shrikhande@eeq.iitr.ac.in
Finite Element Method and Computational Structural Dynamics
Department of Earthquake Engineering Indian Institute of Technology Roorkee

So, once we have this evaluation, now similar equations the problem element number 1, now for the element number 1, the equilibrium equations look like stiffness matrix for element 1 multiplied by the nodal coordinates of for element nodal, variables for primary variables of our element number 1 is equal to the element force vector for node of element number 1, and the secondary variables reactions for element number 1. And these are then placed into appropriate locations of the global stiffness, global equation, equation for the entire system.

So, contributions from individual elements may then be assembled into global system. So, we are done with first element. And these element level equations are plugged into the placeholders for the global system of equations. Since, we have a total of 8 nodes with two degrees of freedom. So, there will be 16 equations in this particular problem, right. So, there are 8 number of nodes. So, total 8 number of nodes, each node has 2 degrees of freedom u_j and v_j . So, 2 degrees of freedom at each node.

So, there are total of 16 variables that we are looking at, so we need to have provision for 16 by 16 stiffness matrix, and 16 number of variable, 16 number of element level element forces, nodal equivalent forces, nodal equivalent forces and the secondary variable tractions.

Now, how does the assembly take place? The key is to find appropriate placement and that is how the entire assembly process is, one different elements contribute the contribution of different elements goes in along the degrees of freedom that they share or the nodes that they share with the adjacent elements. And one, the way it is done, it is automatically done in the program.

Again using very powerful data structuring techniques. If you are interested, you may look at the text book very excellent and easy to read text book on finite element programming by Hinton and Owen which beautifully explains this process of assembly, how it is done automatically in a computer program and how to prepare a data structure suitable data structure to facilitate that.

So, essentially what it involves is mapping of element level degrees of freedom we know there are two degrees two displacements at each node. So, $u_i, v_i, u_j, v_j, u_k, v_k, u_l, v_l$, we know that this i, j, k, l , are numbered in counter clockwise direction. In the global mesh, node i may refer to some particular node according to the mesh design. So, in this particular example, i^{th} node corresponds to node 2, j^{th} node corresponds to node 5, k^{th} node corresponds to node 4, and l^{th} node corresponds to node 1.

So, appropriately what happens here is whatever is the coefficient or terms corresponding to u_i , they would go to corresponding terms of node 2 or x displacement of node 2. So, in the global system of equations that corresponds to element 3 because node 1 will have u and v , and node 2 will have u and v . So, u_i would go to third variable

(Refer Slide Time: 36:08)

And then, once we are done with element number 2, we are we will go to element number 3, and add super impose element number 3 equations on this and then finally,

And once we impose that boundary condition then the rest of the system of equations can be solved for unknowns, unknown variables.

Example Problem-13

For the current example problem, the solution for unknown nodal displacements may be calculated as:

$$\begin{pmatrix} u_1 \\ v_1 \\ u_2 \\ v_2 \\ u_3 \\ v_3 \\ u_4 \\ v_4 \\ u_5 \\ v_5 \\ u_6 \\ v_6 \\ u_7 \\ v_7 \\ u_8 \\ v_8 \end{pmatrix} = \begin{pmatrix} 0 \\ 0 \\ 0 \\ 0 \\ 0 \\ 1.9322 \times 10^{-5} \\ -1.7578 \times 10^{-5} \\ -6.5006 \times 10^{-7} \\ -1.6748 \times 10^{-5} \\ -1.8684 \times 10^{-5} \\ -1.9049 \times 10^{-5} \\ 4.7338 \times 10^{-5} \\ -7.2203 \times 10^{-5} \\ -1.3838 \times 10^{-5} \\ -6.7086 \times 10^{-5} \end{pmatrix} \text{ m}$$

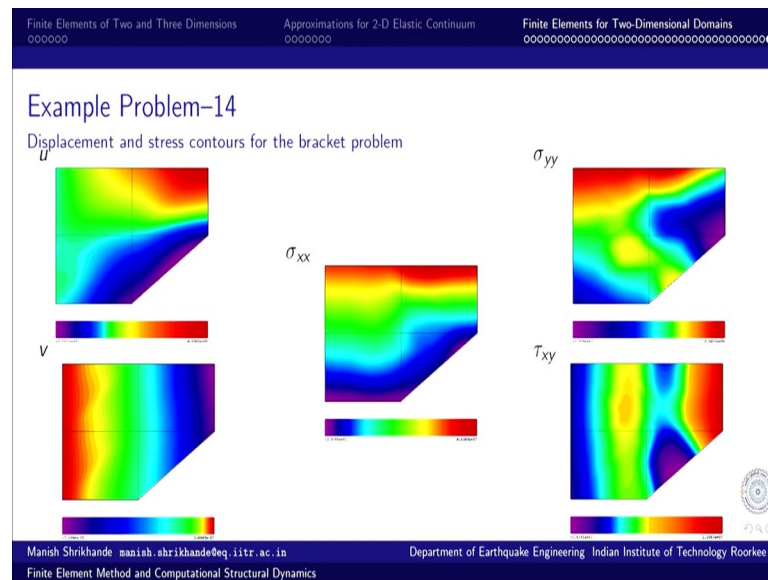
Manish Shrikhande manish.shrikhande@eeq.iitr.ac.in
Finite Element Method and Computational Structural Dynamics

Department of Earthquake Engineering Indian Institute of Technology Roorkee

So, this is what the solution of primary variables looks like. And if I plot, I mean once I do this then I can go back to interpolation model then I have the complete variation of displacement all through the domain with respect to x and y . The displacement is of course, continuous across the element boundaries. So, I can have the deflected shape of the bracket which is what this looks like.

This is often referred to as wire frame diagram of the deflected shape. And once I have, once I extract these and go back to the element level equations, I can derive the secondary variables the stresses and so.

(Refer Slide Time: 39:02)



This is the displacement u and v , all through throughout the as a function of x and y . This is σ_{xx} . You can see very beautifully the constant stress, I mean the tension at the top, and compression at the bottom, and very beautifully my color-coded contours can be seen. And this is σ_{yy} , and this is τ_{xy} shear stress. And these are just the representative of how the typical results of a finite element stress would look like, a finite element result look like.

Of course, these are what we call as smooth contours, in raw contours as we compute do not look so continuous and beautiful. They would be there would be jagged lines. Because the continuity of stresses is not guaranteed, continuity of displacement is of course, guaranteed along across the element boundaries, but not of these stresses because the derivatives are going to be discontinuous.

And should, therefore, the stresses should also be discontinuous over the across the boundary. And that level of discontinuity is what provides us an estimate of the error in the solution or adequacy of or whether we need to refine the mesh or we need to look at further refinement in the finite element results. But that is a another topic, error analysis. We will see if time permits.

So, we stop here. In the next lecture, we begin with a new topic called Mapped Elements or Distorted Elements.

Thank you.