

Mechanics of Material
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Displacement due to uniaxial loading, temperature and bending
Lecture – 51
Bending Equation

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$$u = -\frac{d\Delta}{dx} (y - y_0) e_x + \Delta(x) e_y$$

$$\Delta = \Delta(x) u_x$$

$$u_x = \theta (y - y_0)$$

$$u_x = -\frac{d\Delta}{dx} (y - y_0)$$

$$\tan(\theta) = \frac{d\Delta}{dx}$$

$$\tan \theta \approx \theta$$

Plane section normal to the neutral axis remains plane and normal to the neutral axis

Neutral axis: The line about which the stress (σ_x) is zero

Now, let us proceed, we assume a displacement field, u to be minus $d\Delta$ by $d x$ delta y by $d x$ to y minus y naught e_x plus delta of y as a function of x e_y ok.

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The image shows a whiteboard with the following handwritten content:

$$\underline{u} = -\frac{d\Delta y}{dx}(y-y_0)\underline{e}_x + D_y(y)\underline{e}_y$$

$$\text{Grad}(\underline{u}) = \underline{H} = \begin{pmatrix} \frac{\partial u_x}{\partial x} & \frac{\partial u_x}{\partial y} & \frac{\partial u_x}{\partial z} \\ \frac{\partial u_y}{\partial x} & \frac{\partial u_y}{\partial y} & \frac{\partial u_y}{\partial z} \\ \frac{\partial u_z}{\partial x} & \frac{\partial u_z}{\partial y} & \frac{\partial u_z}{\partial z} \end{pmatrix} = \begin{pmatrix} -\frac{d^2\Delta y}{dx^2}(y-y_0) & -\frac{d\Delta y}{dx} & 0 \\ \frac{d\Delta y}{dx} & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

$$\underline{\underline{\epsilon}} = \frac{1}{2}(\underline{H} + \underline{H}^T) = \begin{pmatrix} -\frac{d^2\Delta y}{dx^2}(y-y_0) & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix} \quad \epsilon_{xx} = -\frac{d^2\Delta y}{dx^2}(y-y_0)$$

Now, I want to compute gradient of \underline{u} . So, that I can compute the strain, gradient of \underline{u} is \underline{H} which is given by $\frac{\partial u_i}{\partial x_j}$. If you recollect we saw this \underline{H} we are looking at strains $\frac{\partial u_x}{\partial x}$; $\frac{\partial u_x}{\partial y}$; $\frac{\partial u_x}{\partial z}$; $\frac{\partial u_y}{\partial x}$; $\frac{\partial u_y}{\partial y}$; $\frac{\partial u_y}{\partial z}$; $\frac{\partial u_z}{\partial x}$; $\frac{\partial u_z}{\partial y}$; $\frac{\partial u_z}{\partial z}$, this for the assumed displacement field would be $-\frac{d^2\Delta y}{dx^2}(y-y_0)$; $-\frac{d\Delta y}{dx}$; 0 ; $\frac{d\Delta y}{dx}$; 0 ; 0 ; 0 ; 0 ; 0 , ok.

Now, the linearized strain as we saw before going by $\frac{1}{2}(\underline{H} + \underline{H}^T)$ which in this case happens to be given by this matrix, $\begin{pmatrix} -\frac{d^2\Delta y}{dx^2}(y-y_0) & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$. Here, a few comments are to be made, the first comment is you have only ϵ_{xx} strain as $-\frac{d^2\Delta y}{dx^2}(y-y_0)$ and there is no shear strain or other components of the normal strain in effect what you are saying is there is no Poisson effect in this problem that is if you consider Poisson effect where there is an actual strain will be a lateral contraction that is operating in the other 2 directions which are ignored as a first approximation for the problem.

Basically what we add was we add only axial strain, we ignore the Poisson effect which means you are ignoring the lateral contraction that will happen when there is an actual strain, that we have ignored as a first approximation.

The second thing is, since we assume plane instructions remain plane and normal to the neutral axis, there is no shear strain and there is no angle change in this (Refer Time:

03:18) system, there is no angle change in this (Refer Time: 03:21) system. So, since you have the plane being perpendicular even after deformation to a particular line which was initially it was perpendicular to; that is basically because of our assumption of the deformation field for the beams.

So, the strain is given by this expression, then we appeal to 1 dimensional consultation where in σ_{xx} would be related to ϵ_{xx} through the relation $E \epsilon_{xx}$ and ϵ_{xx} . This is 1 dimensional consultation because we are ignoring we pass on it is effects.

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The image shows a handwritten derivation on a whiteboard. At the top, it defines the strain tensor $\underline{\underline{\epsilon}}$ as the symmetric part of the displacement gradient $\underline{\underline{H}}$:

$$\underline{\underline{\epsilon}} = \frac{1}{2} \left[\underline{\underline{H}} + \underline{\underline{H}}^T \right]$$

The displacement gradient $\underline{\underline{H}}$ is given by the partial derivatives of the displacement components u_x, u_y, u_z with respect to the spatial coordinates x, y, z :

$$\underline{\underline{H}} = \begin{pmatrix} \frac{\partial u_x}{\partial x} & \frac{\partial u_x}{\partial y} & \frac{\partial u_x}{\partial z} \\ \frac{\partial u_y}{\partial x} & \frac{\partial u_y}{\partial y} & \frac{\partial u_y}{\partial z} \\ \frac{\partial u_z}{\partial x} & \frac{\partial u_z}{\partial y} & \frac{\partial u_z}{\partial z} \end{pmatrix}$$

For a beam under bending, the displacement field is assumed to be $u_x = -y \frac{d^2 \Delta y}{dx^2} x$, $u_y = \Delta y$, and $u_z = 0$. Substituting these into the strain tensor, the only non-zero component is ϵ_{xx} :

$$\epsilon_{xx} = -\frac{d^2 \Delta y}{dx^2} y$$

The stress σ_{xx} is then related to the strain ϵ_{xx} through Hooke's law for a linear elastic material:

$$\sigma_{xx} = E \epsilon_{xx}$$

Substituting the expression for ϵ_{xx} into Hooke's law, the stress is:

$$\sigma_{xx} = -E \frac{d^2 \Delta y}{dx^2} y$$

The derivation also shows the stress tensor $\underline{\underline{\sigma}}$ as a diagonal matrix with σ_{xx} in the top-left corner and zeros elsewhere:

$$\underline{\underline{\sigma}} = \begin{pmatrix} -\frac{d^2 \Delta y}{dx^2} y & 0 & 0 \\ 0 & 0 & 0 \\ 0 & 0 & 0 \end{pmatrix}$$

Finally, the stress σ_{xx} is identified as the bending moment M per unit length L :

$$\sigma_{xx} = \frac{M}{L}$$

So, now then σ_{xx} is given by minus E times $d^2 \Delta y / dx^2$ into y minus y naught. Now, what we will do is we will go back and plug this expression for σ_{xx} the equilibrium equation that we are, that is we will plug σ_{xx} in in this expression, in this expression, in this expression and then use this equilibrium equations this, this and this to find what is the variation of σ_{xx} , basically unknown function the unknown function here is $d^2 \Delta y / dx^2$. So, we want to estimate this, so we will see how to estimate this.

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The image shows a whiteboard with the following handwritten content:

$$\sigma_{xx} = -E \frac{d^2 y}{dx^2} (y - y_0)$$

$$P = \int \sigma_{xx} da_x = \int E \frac{d^2 y}{dx^2} (y - y_0) da_x = 0$$

$$= \frac{d^2 y}{dx^2} \int E (y - y_0) da_x = 0 \Rightarrow \int E (y - y_0) da_x = 0$$

Assume, Beam is Homogeneous, i.e. E is a constant \rightarrow No Spatial Variation
 $\Rightarrow E$ is same in tension or Compression.

$$\int E (y - y_0) da_x = E \int (y - y_0) da_x = 0 \Rightarrow \int (y - y_0) da_x = 0 \Rightarrow y_0 \int da_x = \int y da_x$$

$$y_0 = \frac{\int y da_x}{\int da_x} = y_{\text{centroid}} \Rightarrow y_0 = y_{\text{centroid}}$$

So, we saw that σ_{xx} is given by $-E \frac{d^2 y}{dx^2} (y - y_0)$. Now, the actual force P is $\sigma_{xx} da_x$; that will be $E \frac{d^2 y}{dx^2} (y - y_0) da_x$, integrated over da_x . This da_x is nothing but $dy dz$ as you have seen many a times now, this as we equated to 0. Now, the integration with respect to y and z is independent of x . So, I can pull out $\frac{d^2 y}{dx^2}$ outside the integration. So, this will be $E \frac{d^2 y}{dx^2} \int (y - y_0) da_x$ is equal to 0.

Since $\frac{d^2 y}{dx^2}$ cannot be 0, identically this will imply that $\int E (y - y_0) da_x$ has to be equated to 0. Now, let us assume beam is homogeneous, that is E is a constant implying no special variation. So, I am not considering the case where the beam is made of 2 different materials for example, concrete and steel or wood and steel, wood and aluminum and so on. So, I am considering a case where E is the same throughout the beam and this also means that E is same in both tension and compression. This also means that E is same in tension and compression.

We will see why I am saying this shortly. Now, what happens now is, now this as I can pull out E since E is a constant. So, $\int E (y - y_0) da_x$ will give me now $E \int (y - y_0) da_x = 0$. This, implies $\int (y - y_0) da_x$ has to be 0. Since, we saw that E cannot be 0 because E has to be greater

than 0 in material parameters when I estimated the restriction of material around is we saw that.

Now, this will imply y_{naught} integral $d a x$ must be equal to integral y times $d a x$ or y_{naught} is given by integral y times $d a x$ divided by integral $d a x$. What is this integral $y d a x$ divided by integral $d a x$ is nothing but this centroid of the cross section, the y component of the centroid of the cross section; y of the centroid of the cross section. So, basically y_{naught} is identified as this implies y_{naught} is equal to y centroid of the cross section. Now, we have found what y_{naught} is or we have found expression for y_{naught} , next let us go ahead and find $d^2 \delta y$ by $d x^2$.

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$$M_z = - \int \sigma_{xx} y da_x = - \int \sigma_{xx} (y - y_0) da_x \quad (\because P=0, \int \sigma_{xx} y_0 da_x = y_0 \int \sigma_{xx} da_x = 0)$$

$$M_z = \int E \frac{d^2 \delta y}{dx^2} (y - y_0) da_x = + \frac{d^2 \delta y}{dx^2} \int E (y - y_0)^2 da_x$$

$$\frac{d^2 \delta y}{dx^2} = \frac{M_z(x)}{\int E (y - y_0)^2 da_x} = \frac{M_z(x)}{E I_{zz}} \quad ; \quad I_{zz} = \int (y - y_0)^2 da_x$$

Homogeneous Beam

$$\frac{E d^2 \delta y}{dx^2} = \frac{M_z}{I_{zz}} \quad \sigma_{xx} = -E \frac{d^2 \delta y}{dx^2}$$

Moment of Inertia about z direction at the neutral axis.

Now, I know that M_z moment is given by integral $\sigma_{xx} y d a x$ of x ; this I am claiming the same as $\sigma_{xx} y_{naught} d a x$, since P is 0, integral $\sigma_{xx} y_{naught} d a x$ would be y_{naught} into $\int \sigma_{xx} d a x$, so that is equal to 0 because just now we show that $\int \sigma_{xx} d a x$ is 0. So, I can rewrite the moment equation as this. Now, I substitute for σ_{xx} from the previous equation that is it is minus E times $d^2 \delta y$ by $d x^2$ into y_{naught} into $y_{naught} d a x$.

So, this is nothing but again $d a x$ is $d y d z$ this is δy is a function of x alone. So, I can move that outside the integral, this would be $d^2 \delta y$ by $d x^2$; integral E times y_{naught} whole square $d a x$. In other words, this is nothing but from

here I get d^2y by dx^2 to be M_z , so function of x divided by integral a of x ; E times y minus y naught square d a x . We saw that if the beam is homogenous, E will be a constant and hence you get M_z of x divided by E times y minus y naught whole square d a x , this is for a homogeneous beam assumption.

So now, what is this? This is nothing but M_z of x divided by E times I_z , where I_z is given by integral y minus y naught whole square d a x a of x , this is moment of inertia about z axis and neutral axis, about z ed neutral axis direction at the neutral axis, but I said direction of the neutral axis.

So, we get the equation minus E times d^2y by dx^2 is M_z by I_z . We saw before that σ_x by y minus y naught is again minus E times d^2y by dx^2 .

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The image shows a digital whiteboard with the following text and equation:

Beam bending equation:
$$-\frac{\sigma_x}{(y - y_0)} = \frac{M_z}{I_z} = \frac{E d^2y}{dx^2}$$

So, combining these 2 equations you get the beam bending equation, which is σ_x by y minus y naught equal to M_z by I_z ; E into d^2y by dx^2 . So, we have the beam bending equation given by this equation in here.